

Pairwise hybrid incompatibilities dominate during allopatric speciation for a simple genotype-phenotype map of embryonic spatial patterning

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ABSTRACT Understanding the origin of species is as Darwin called it “that mystery of mysteries”. Yet still, how the processes of evolution give rise to non-interbreeding species is not well understood. In an empirical search for a genetic basis, transcription factor DNA binding has commonly been identified as being an important factor in the development of reproductive isolation. This is supported by computational and theoretical models based on the biophysics of transcription factor DNA binding that provide a mechanistic basis of such incompatibilities between allopatrically evolving populations. However, gene transcription mediated by such binding events occurs within the context of larger gene regulatory networks, so the question remains how important are such pair-wise interactions compared to higher order interactions in determining incompatibilities. Orr calculated that as the order of interaction increases there are more pathways for an incompatibility to occur. Here, we show, using simulations based on a simple biophysical genotype phenotype map of spatial patterning in development, that biophysics provides a stronger constraint, leading to pair-wise incompatibilities arising more quickly and being more numerous than higher order incompatibilities, when there is stabilising selection on each allopatric lineage. In addition, we show that incompatibilities arise more quickly for smaller populations and in a manner supporting previous conclusions from models of hybrid incompatibility based solely on transcription factor DNA binding.

KEYWORDS speciation, Dobzhansky-Muller incompatibilities, sequence entropy, population size, co-evolution, genotype phenotype map

1 Introduction

2 The detailed genetic mechanisms by which non-interbreeding
3 species arise is still largely not understood. Darwin called it that
4 “mystery of mysteries” Darwin (1859); he struggled to under-
5 stand how natural selection could give rise to hybrid inviability
6 or infertility. In a modern setting, Darwin’s conundrum was, if
7 a hybrid incompatibility were due to a single locus, how could
8 two species, fixed for *AA* and *aa*, respectively, evolve from a
9 common ancestor, since at least one of these species would need
10 the inviable genotype *Aa* to give rise to offspring. A solution
11 to this problem was conceived independently by Dobzhansky,

12 Muller and Bateson Dobzhansky (1936); Muller (1942); Bateson
13 (1909), by which neither population need pass through a bot-
14 tleneck. If instead incompatibilities arise due to non-linear or
15 epistatic fitness interactions between different loci, it is possi-
16 ble for example, that two lines that are geographically isolated
17 from each other, evolve independently from a common ancestor
18 *ab* (allopatric evolution), fix the allelic combinations *aB* and *Ab*
19 respectively, yet the hybrid genotype *AB* is inviable.

20 The work of Orr provided a framework to understand how
21 incompatibilities might arise in allopatry Orr (1995); Orr and
22 Turelli (2001), when populations are small, $\mu N \ll 1$, (where μ
23 is the mutation rate for a typical loci and N the effective popu-
24 lation size) and essentially monomorphic. Orr suggested that
25 as two lines fix independent substitutions from a common an-
26 cestor, any combination of alleles that may arise in hybrids that
27 have not been ‘tested’ or explicitly fixed by the process of evolu-
28 tion, could potentially cause an incompatibility. Assuming an

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1 infinite number of loci and that back-substitutions or multiple
2 substitutions at the same loci are not possible, Orr showed that
3 the number of incompatibilities involving n -loci increases as
4 $\sim K^n$, for K substitutions separating the two lines and assuming
5 $n \ll K$. In particular, pair-wise incompatibilities ($n = 2$) rise as
6 K^2 . This has been likened to a “snowball” effect as the number of
7 incompatibilities rises rapidly with the number of substitutions
8 that separate the lines. As an outcome of this framework Orr
9 also hypothesised that more complex DMIs, which involve more
10 than 2-loci should be easier to evolve as the fraction viable paths
11 between common ancestor genotype and the genotypes of the
12 two present day species is larger Orr (1995) for a fixed number
13 of incompatible genotypes. It is, however, an open question,
14 whether these predictions would hold for more realistic fitness
15 landscapes.

16 To address the question of how incompatibilities develop
17 in more realistic fitness landscapes, Tulchinsky et al Tulchin-
18 sky et al. (2014b,a) developed sequence-based simulations that
19 investigated the mechanistic basis of the evolution of hybrid in-
20 compatibilities and the effects of pleiotropy. Khatri & Goldstein
21 Khatri and Goldstein (2015a,b) used a similar model, based on
22 a simple biophysically motivated genotype-phenotype map of
23 a single transcription factor binding to DNA; they found that
24 the nature and rapidity of the growth of DMIs between two
25 allopatric populations strongly depended on the product of the
26 effective population size and the scale of fitness of the stabil-
27 ising landscape; for populations that are small compared to
28 the inverse of the scale of fitness, DMIs arise more quickly. This
29 arises as there are exponentially more sequences that bind poorly
30 than well, which means drift pushes common ancestors on av-
31 erage closer to incompatible regions and so less substitutions
32 are needed for incompatibilities to arise in hybrids. In this low
33 population size limit DMIs arise quadratically with divergence
34 time, in agreement with Orr’s predictions, however, the underly-
35 ing mechanism for this power law is very different to Orr’s. For
36 populations that are larger than the inverse of this scale of fitness,
37 DMIs take longer to arise as common ancestors have a smaller
38 drift load and require more substitutions for hybrids to reach
39 incompatible regions. However, in addition they arise more
40 slowly, because within a discrete stabilising landscape, as the
41 population size increases the scaled fitness differences becomes
42 larger and more deleterious, giving rise to slower divergence.
43 This is in line with predictions of the nearly neutral theory ??,
44 where the number of nearly neutral mutations decreases as the
45 population size increases. In this case the growth of DMIs has
46 a characteristic non-power law form, with negative curvature
47 on a log-log plot, indicating that the hybrid binding energies
48 change diffusively. However, real gene regulatory systems are
49 more complex than a single TF binding to DNA, so again the
50 question arises do these predictions hold for more complex gene
51 regulatory systems with more realistic fitness landscape?

52 Although there has been much progress in understanding
53 evolution in terms of selection, mutation and genetic drift, much
54 of this work has been reliant on phenomenological fitness land-
55 scapes, which encompass in a heuristic manner smoothness,
56 epistasis and neutrality Higgs and Derrida (1992); Kauffman
57 and Levin (1987). In recent years, the question of the structure
58 of real fitness landscapes has gained prominence, where the
59 redundancy of the mapping from genotype to phenotype can
60 give rise to non-trivial properties of the evolution of pheno-
61 types Fontana (2002); Khatri et al. (2009); Hayden et al. (2011);
62 Goldstein (2011); Schaper and Louis (2014); Greenbury et al.

(2014); Manrubia and Cuesta (2015); Greenbury et al. (2016). In
63 particular, Khatri et al Khatri et al. (2009) introduced a simple
64 genotype-phenotype map for spatial gene expression regulation
65 in development, from which emerged a number of non-trivial
66 features such as a balance between selection and sequence en-
67 tropy deciding the course of evolution at small population sizes
68 and a partitioning of the effective phenotypic landscape into neu-
69 tral and selective degrees of freedom; none of these emergent
70 phenomena could be predicted on the basis of purely phenotypic
71 considerations. In this paper, we will use a slightly modified
72 version of the spatial patterning model in Khatri et al. (2009),
73 that has explicit sequence representation of each loci, to examine
74 the Dobzhansky-Muller mode of evolution of incompatibilities
75 in allopatry as a function of population size and under stabilis-
76 ing selection in each lineage. Our results show that population
77 size interplays with fitness and sequence entropic effects in a
78 way mirrored by previous results on single TF-DNA binding. In
79 addition, we find that, unlike Orr’s prediction, pair-wise inter-
80 actions between loci dominate the growth of DMIs, suggesting
81 that biophysics provides a stronger constraint on their evolu-
82 tion than the simple combinatorics of pathways between the
83 common ancestor and present day lineages.

85 Materials and Methods

86 GP map

87 The genotype-phenotype map we use was described in detail by
88 Khatri et al. Khatri et al. (2009), so here we summarise its basic
89 elements and recapitulate the main results in the Supporting
90 Information, since the methodology used is slightly different.
91 The evolutionary task set for the gene regulation module is to
92 turn an exponentially decaying morphogen gradient across a
93 field of cells in an embryo (M), into a sharp step function profile
94 of a downstream transcription factor TF with its transition at
95 the mid-point of the embryo, as shown in Fig.1. The model for
96 this gene regulation is essentially the model of the genotype-
97 phenotype map, where the binding affinity of the different pro-
98 tein species to themselves or to the regulatory region of TF is de-
99 termined by matching binary sequences. More specifically, gene
100 regulation is controlled by two non-overlapping binding sites,
101 the promoter P and an adjacent binding site B , together with
102 two protein species, the morphogen M and RNA Polymerase R ,
103 where the length of these sequences are ℓ_b . The various binding
104 energies are proportional to the Hamming distance between the
105 relevant sequences, where for protein-DNA binding the cost of
106 a mismatch is ϵ_b and for protein-protein interactions ϵ_g . So for
107 example, $E_{MP} = \epsilon_b \rho(\mathbf{b}_M, \mathbf{b}_P)$ is the binding energy between the
108 morphogen and the promoter, where ρ is the Hamming distance
109 and $\mathbf{b}_M, \mathbf{b}_P$ are the morphogen binding sequence and promoter
110 binding sequence respectively.

111 Given an exponential morphogen concentration profile
112 $[M](x, \alpha)$ as a function of the position of embryonic cells, x ,
113 and a fixed concentration $[R]$ of RNAP, in each cell, we fol-
114 low Shea and Ackers Shea and Ackers (1985) to calculate the
115 TF concentration profile $[TF](x)$, assuming transcription is pro-
116 portional to the probability of RNAP being bound to the pro-
117 moter p_{RP} . This probability can be calculated using the canon-
118 ical ensemble of statistical mechanics, where details can be
119 found in Khatri et al. (2009). Since the TF cannot affect its
120 own transcription, the steady state concentration profile is
121 then simply proportional to the probability of RNAP being
122 bound to the promoter: $[TF](x) \propto p_{RP}(G, R, M(x, \alpha))$. where
123 $G = [\mathbf{b}_R, \mathbf{b}_M, \mathbf{b}_P, \mathbf{b}_B, \mathbf{g}_R, \mathbf{g}_M]$ is the genome and is a function of

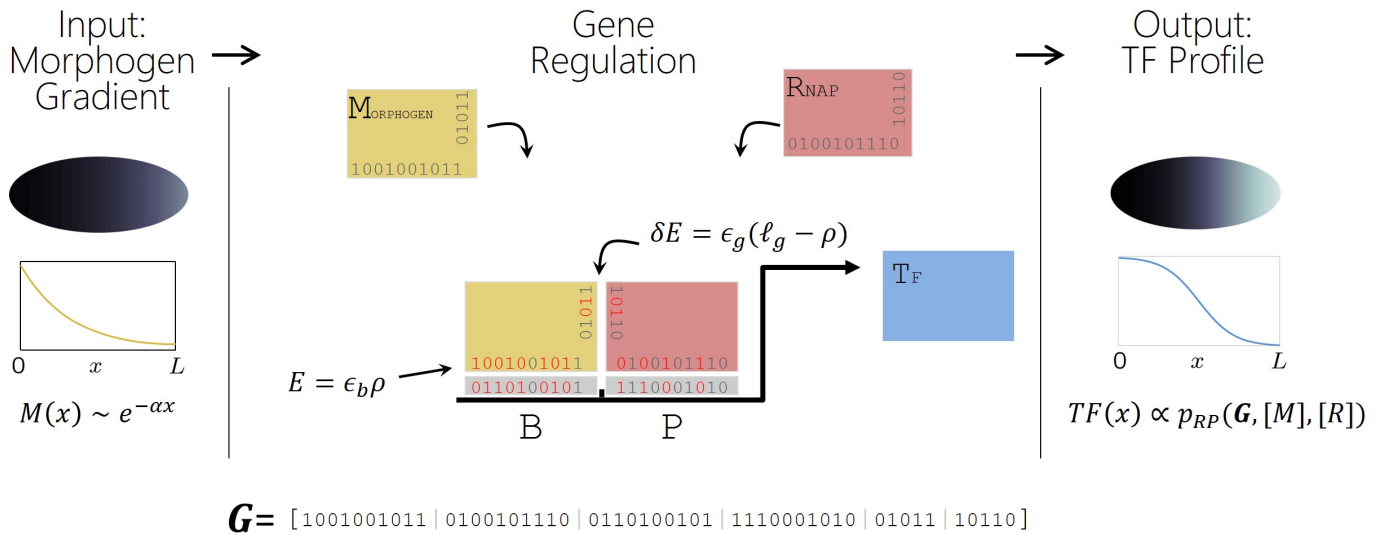


Figure 1 An overview of the genotype-phenotype map. The gene regulatory module has input a morphogen gradient $[M](x)$ across a 1-dimensional embryo of length L and outputs a transcription factor $TF(x)$. Gene regulation of TF using a morphogen and RNAP (R) is controlled in a bottom-up manner, by binding to its regulatory region consisting of a promoter P and adjacent binding site B ; E represents binding free energies of proteins to one of the two binding sites of the regulatory region of the transcription factor T , δE are glue free energies to aid in co-operative binding of paired protein complexes. Each energy is calculated by the number of mismatches ρ (Hamming distance), shown in red, between relevant binary sequences, together with mismatch energies ϵ_b and ϵ_g for binding and glue energies respectively. Transcription of T is controlled by the probability of RNAP being bound to the P , p_{RP} .

1 all the binding and glue sequences from proteins and DNA. The
2 proportionality constant will be given by the ratio of the rate of
3 transcription and translation to the rate of degradation of TF ,
4 which is not important in our study, since we are only interested
5 in the shape or contrast of $[TF](x)$ that can be achieved.

6 Monte Carlo Scheme

7 We use a kinetic Monte Carlo scheme to simulate a Wright-Fisher
8 evolutionary process for the genome G and α on each lineage, as
9 detailed in [Khatri and Goldstein \(2015b\)](#). In particular, the rate
10 of fixation of one-step mutants are calculated based on Kimura's
11 probability of fixation [Kimura \(1962\)](#), where we are assuming
12 a regime of small effective population size ($n\mu_0 N_e \ll 1$, where
13 $n = 50$ is the number of coding base-pairs in the genome and
14 μ_0 is the base-pair mutation rate and N_e is effective size of the
15 haploid asexual population). Here, we determine the goodness
16 of the spatial gene regulation, from the resulting concentration
17 profile $[TF](x)$ by mapping to a Malthusian fitness by use of a
18 functional that promotes expression of the TF in the anterior
19 half, whilst penalising expression in the posterior half, with
20 truncation selection below critical a value W^* . The functional
21 \mathcal{W} is:

$$\mathcal{W}[[TF](x)] = \frac{\int_0^{L/2} [TF](x) dx - \int_{L/2}^L [TF](x) dx}{\frac{L}{2} \max_x \{[TF](x)\}}. \quad (1)$$

22 We then map this to a Malthusian fitness as follows:

$$F = \begin{cases} \kappa_F \ln(\mathcal{W}) & \text{if } \mathcal{W} > W^* \\ -\infty & \text{if } \mathcal{W} < W^* \end{cases} \quad (2)$$

23 where κ_F is the strength of selection for the trait represented by
24 the spatial patterning process. Note that although here the exact
25 form of the fitness is slightly different to the one used in [Khatri
26 et al. \(2009\)](#), the qualitative behaviour is the same as shown in
27 the Supporting Information.

28 Speciation Simulations

29 Starting from a random initial genome and fixed population size
30 N_e , the spatial patterning simulation is run for 100,000 substi-
31 tutions, in order to effectively equilibrate the system (typically
32 10,000 substitutions are required to adapt to an ensemble of fit
33 states). For each replicate allopatric speciation run, two simula-
34 tions are performed from the same common ancestor, each using
35 the same fitness function Eqn.2. For each simulation pair, the
36 common ancestor is drawn from the equilibrium distribution
37 taking the end state of the sequences of the long simulation and
38 running it further for a further 100 substitutions. Although, non-
39 neutral substitutions can fix on each line, since each line has the
40 same environment we consider this situation to be neutral with
41 respect to selection on each line.

42 For simplicity, we assume that our genome is composed of
43 4 loci: 1) RNAP (R), 2) Morphogen (M), 3) Regulatory region
44 of TF (T) and 4) the morphogen gradient steepness α , and we
45 form hybrids between the two lines assuming independent re-
46 assortment of these loci and complete linkage within each loci.
47 We define a hybrid genotype by a 4 digit string where each digit
48 corresponds to one of the loci defined above and takes one of
49 two values, which correspond to the allele from the 1st line or
50 2nd line; for example, the hybrid rMTa corresponds to R loci
51 having an allele from the 1st line, M loci from the 2nd loci, T
52 loci from the 2nd loci and α loci the allele from the 1st loci. Note
53 that the underlying sequence of each hybrid changes as different

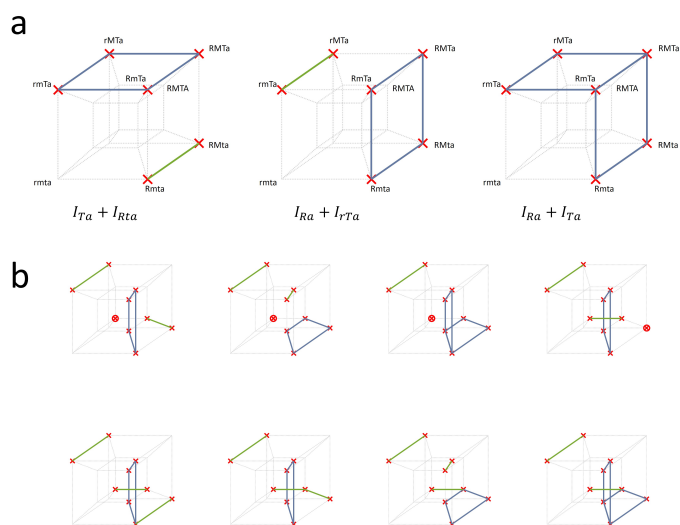


Figure 2 Decomposition of hybrid DMIs on a Boolean hyper 4-cube. Each point on the 4-cube represent each possible hybrid genotype across 4 loci, including the genotype of each parental lineage, where each red cross represents an incompatible hybrid genotype. As shown in a) the pattern of DMIs can be explained by different combinations of fundamental types of DMIs, where a blue square or face identify a subspace of genotypes that correspond to a 2-point DMI and a single green edge of line corresponds to a 3-point DMI. b) A more complicated pattern of hybrid DMIs and their decomposition into fundamental types, where a red dot corresponds to a single isolated 4-point DMI.

28 between loci, for example, using a Fourier series on a Boolean
 29 hypercube Weinberger (1991); Neher and Shraiman (2011); how-
 30 ever, such decompositions suffers from a problem that the
 31 terms involving interactions between 3 or more loci are diffi-
 32 cult to interpret physically. A more explicit representation of the
 33 form $F(\mathbf{g}) = F_0 + \sum_i F_i(g_i) + \sum_{ij} F_{ij}(g_i, g_j) + \sum_{ijk} F_{ijk}(g_i, g_j, g_k) +$
 34 $F_{1234}(g_1, g_2, g_3, g_4)$ is not possible as the problem is hugely un-
 35 derdetermined; given that there are only $2^4 - 2 = 14$ possible
 36 hybrids (not including the well-adapted genotypes of line 1 and
 37 line 2) and a total of $I_{max} = 3^L + 1 - 2^{L+1} = 50$, for $L = 4$
 38 loci¹, different interactions it is not possible to determine un-
 39 ambiguously which interactions are responsible for a particular
 40 set of DMIs across all hybrids. The approach we take is look
 41 at patterns of DMIs across all hybrid genotypes and construct
 42 the most parsimonious fundamental interactions that could give
 43 rise to this pattern of DMIs across the hybrid genotypes. Here
 44 we define parsimonious to be, all other things being equal, the
 45 minimum number of possible interactions needed to explain the
 46 observed pattern of DMIs, which from a Bayesian perspective
 47 would have the smallest Occam factors MacKay (2007).

48 If we consider a pair-wise incompatibility say $T \leftrightarrow a$,
 49 which we denote I_{Ta} , any hybrid-genotype that contains the
 50 alleles Ta must be a DMI; this defines a 2D subspace (or
 51 face) of 4 DMIs across loci R and M on a 4D Boolean cube
 52 $\{rmtA, rMTa, RmTa, RMTa\}$. Similarly, a 3-point DMI I_{mtA}
 53 defines an edge of 2 DMIs along the 1D subspace of loci 4:
 54 $\{rmtA, RmtA\}$. Each 4-point DMI takes up a single point in
 55 the hybrid-genotype space. The approach is then to find all
 56 combinations of possible DMIs that could explain the pattern
 57 of incompatibilities across the hybrids, but consider only those
 58 with the least number of total DMIs. Of this subset of minimal
 59 DMIs, we assume there is no a priori reason each subset should
 60 not be given equal weighting and so then our measure of the
 61 number of DMIs of each type is the average number of 2-point,
 62 3-point and 4-point DMIs. For example, in Fig.2a, the hybrids
 63 may have a pattern of DMIs as shown by the red crosses; the
 64 most parsimonious representation are 3 equally weighted possi-
 65 bilities: 1) $I_{Ta} + I_{Rta}$, 2) $I_{Ra} + I_{rTa}$, 3) $I_{Ta} + I_{Ra}$. If this pattern of
 66 hybrid incompatibilities were to arise it would be decomposed
 67 into $n_{Ta} = 2/3$, $n_{Ra} = 2/3$, $n_{Rta} = 1/3$, $n_{rTa} = 1/3$, so the total
 68 number of 2-point DMIs $n_2 = 4/3$ and 3-point DMIs $n_3 = 2/3$,
 69 corresponding to a total number of $n = n_2 + n_3 = 2$ DMIs,
 70 which is the minimum number of DMIs needed to explain the
 71 pattern of hybrid incompatibilities shown in Fig.2a. In Fig.2b
 72 is shown a more complicated pattern of hybrid DMIs and the
 73 parsimonious decomposition into k -point DMIs.

1 substitutions are accepted in each line; the notation only refers
 2 to alleles fixed at any point in time. Under these assumption we
 3 can then look at how the average fitness and average number of
 4 incompatibilities varies as the number of substitutions that
 5 separate two lines increases. As substitutions in α are quite com-
 6 mon due to a small curvature of its landscape, we only count
 7 substitutions in α when it makes a transition between the two
 8 basins of attraction mentioned above and discussed in Khatri
 9 *et al.* (2009).

10 The value chosen for W^* is somewhat arbitrary, the larger
 11 W^* the more quickly incompatibilities will arise; here we choose
 12 $W^* = 0.2$ ($F^* = \kappa_F \ln W^* \approx -0.0016$), which gives a good
 13 balance between the number of incompatibilities on a given
 14 timescale and realism for their actual inviability; arguably
 15 $W^* = 0$, but this requires quite lengthy simulations in order
 16 for incompatibilities to arise.

17 **Decomposing DMIs**

18 Examining the number of DMIs that arise across different hy-
 19 brid genotypes is not the most fundamental representation of
 20 the incompatibilities, since for example, if a particular hy-
 21 brid $RmTa$, where the 1st and 3rd loci take alleles from line
 22 2 and the 2nd and 4th take alleles from line 1, is found to
 23 be incompatible, it is not known whether this incompatibil-
 24 ity arises from a pairwise interaction between R and m or be-
 25 tween m and T , or a triplet interaction between R, m and T or
 26 all of these concurrently. Ideally we would like to be able to
 27 decompose fitness contributions into n -point interactions be-

¹ The total number of n -point DMIs is $(2^n - 2) \binom{L}{n}$, as there are $\binom{L}{n}$ combinations of n loci amongst L total loci and then considering a binary choice of alleles across both lines, there are a total of 2^n allelic combinations or states, 2 of which are the fit allelic combinations where all alleles come from one lineage or the other giving $2^n - 2$. For example, between each pair of loci there are $2^2 - 2 = 2$ mismatching combinations of alleles (e.g. rM and Rm) that could give DMIs and $\binom{L}{2} = L(L-1)/2 = 6$ pairwise interactions. A similar argument would give a total of 24 3-point DMIs as there are $2^3 - 2 = 6$ mismatching combinations of alleles at 3 loci (e.g., excluding rmt and RMT) and $\binom{L}{3} = 4$ 3-point interactions and similarly, $14 \binom{L}{4} = 14$ for 4-point interactions. In total, the max number of DMIs is $I_{max} = \sum_{n=2}^L (2^n - 2) \binom{L}{n} = 3^L + 1 - 2^{L+1}$, which for $L = 4$ loci is $I_{max} = 50$.

1 Results

2 *Evolutionary properties of genotype-phenotype map on each* 3 *lineage*

4 The properties of this genotype-phenotype map have been pre-
5 viously explored [Khatri et al. \(2009\)](#). As the some small details of
6 the evolutionary dynamics have changed here, we show in the
7 Supporting Information that the basic pertinent findings from
8 [Khatri et al. \(2009\)](#) are recapitulated. An important property of
9 this genotype-phenotype map is that only a single mechanism
10 of patterning is found, which is that RNAP (R) binds with inter-
11 mediate affinity to the promoter (P), but through a high affinity
12 protein-protein interaction with the morphogen (M), the mor-
13 phogen binds to the first binding site (B) only above a critical
14 morphogen concentration, thereby giving a spatial switch once
15 the morphogen falls below this concentration; evolution then
16 fine tunes the relationship between the binding energies (E 's),
17 glue energies (δE 's) and the steepness of the morphogen gradi-
18 ent α to turn off transcription at the mid-point of the embryo.
19 Despite a single global solution there are many different combi-
20 nations of the binding and glue energies and α that give good
21 patterning, and for each of these many underlying genotypes
22 (G). A further key property is that despite this redundancy some
23 energy phenotypes such as E_{MB} , δE_{RM} and E_{RP} are under strong
24 stabilising selection, whilst the remaining energy phenotypes,
25 E_{RB} , E_{MP} , δE_{RR} and δE_{MM} are almost neutral and under very
26 weak stabilising selection (see Supporting Information and [Kha-
27 tri et al. \(2009\)](#)). Finally, at large population sizes it is found that
28 the evolutionary dynamics exhibits what is known as quenched-
29 disorder in statistical physics, where energy phenotypes are that
30 less constrained take different *random* values between indepen-
31 dent evolutionary runs, which indicates there is an underlying
32 roughness to the fitness landscape [Khatri et al. \(2009\)](#).

33 *Different hybrids genotypes have different growth rates of* 34 *DMIs*

35 In [Fig.3](#), we plot a typical time series, for scaled population sizes
36 of $2N\kappa_F = 1$ and $2N\kappa_F = 10$, of how fitness of two different
37 hybrids (Rmta a & b and RmtA c & d) changes over time t
38 separating a pair of lines together with the threshold in fitness
39 (dashed black line) used later to count DMIs. We see that the
40 fitness of hybrids generally decreases in a stochastic fashion;
41 when the fitness of a hybrid drops below the threshold F^* , a
42 DMI arises and is indicated by the disappearance of the fitness
43 line ($F = -\infty$) for that hybrid and so as can be seen from [Fig.3](#),
44 at any given time only a subset of all possible hybrids might
45 be incompatible. We also see that as the fitness of hybrids is
46 stochastic, DMIs that arise do not stay, as one might expect
47 within the Orr framework [Orr \(1995\)](#); [Orr and Turelli \(2001\)](#). A
48 further observation is that for smaller scaled population sizes the
49 common ancestor fitness is lower and incompatibilities appear
50 to arise more quickly; this is a consistent with previous studies of
51 a more simple genotype-phenotype map of a transcription factor
52 binding a single binding site [Khatri and Goldstein \(2015a,b\)](#),
53 where the smaller populations have a larger genetic drift load
54 and so common ancestors are more likely to be closer to an
55 inviable binding threshold.

56 To examine the overall trends in the number of incompatibil-
57 ities for each hybrid type, as a function of divergence time μt ,
58 we plot the average number of DMIs in [Fig.4](#), where we have
59 averaged over pairs of complementary genotypes (e.g. RmtA
60 and rMTa), which behave the same on average as each line is
61 exploring the same fitness landscape independently (not shown).

62 As denoted in the figure, we distinguish hybrid-genotypes based
63 on the types of potential mismatch: R -type is characterised by a
64 mismatch of the R loci with M and T loci, M -type is a mismatch
65 of M with R and T loci, T -type is a mismatch of T with R and
66 M and α -type is a mismatch of the α loci with the rest of the loci.
67 In [Fig.4](#), we first note that each hybrid-genotype behaves in a
68 different way in a population size dependent manner. In general,
69 there is an initial growth in the average number of DMIs as the
70 time separating lines increases, followed by a slowing down of
71 the growth or a plateau. It is also clear that as the population size
72 is increased DMIs take longer to arise, which is also consistent
73 with previous work on the dynamics of incompatibilities due to
74 transcription factor DNA binding [Khatri and Goldstein \(2015b\)](#),
75 where this is caused by a slowing of the substitution rate as in a
76 stabilising discrete fitness landscape more deleterious changes
77 are needed for any evolutionary change. We also see that for
78 small population sizes the initial growth of DMIs is power law
79 and approximately quadratic, as predicted by [Orr \(1995\)](#);
80 [Orr and Turelli \(2001\)](#), but as argued previously [Khatri and Gold-
81 stein \(2015b\)](#) the underlying mechanism is very different. On the
82 other hand for large population sizes there is no clear power law,
83 which is again consistent with previous simulations [Khatri and
84 Goldstein \(2015b\)](#) and also theoretical calculations [Khatri and
85 Goldstein \(2015a\)](#) that predict that as common ancestor popula-
86 tions are further away, the growth of DMIs follows a diffusive
87 law, which has a characteristic negative curvature on a log-log
88 plot.

89 However, in addition to these general trends, we see that
90 the growth of DMIs is different for different hybrids; for small
91 population sizes ($2N\kappa_F = 0.1$ & $2N\kappa_F = 1$), M -type, T -type
92 and R -type dominate the growth of DMIs, in this order and
93 with only a small difference between them, whilst α -type arise
94 far more slowly; we might expect this since substitutions in α
95 only tend to shift the pattern away from the mid-point of the
96 embryo, which with the model of fitness defined in [Eqn.2](#) only
97 moderately affects fitness. As the population size increases, and
98 at small times, we see that initially M -type DMIs arise more
99 slowly relative to T -types and R -types, but at longer times there
100 is a cross-over where M -type DMIs dominate R -type; although
101 the simulations do not run out to sufficiently long times for
102 the largest population sizes, this cross-over appears to move to
103 longer times at increasing population sizes.

104 How can we understand this general behaviour? It is clear
105 from [Eqns. 1&2](#) that good patterning or fitness is only dependent
106 on the binding and glue energy phenotypes (as well as α) and so
107 this in general requires co-evolution of the relevant sequences
108 to maintain these energies within certain constraints; e.g. the
109 binding energy of R to the promoter of T , E_{RP} mustn't be too
110 strong and so on each line the sequences will co-evolve to main-
111 tain this constraint. From previous work [Khatri and Goldstein
112 \(2015a,b\)](#) we expect that it is not only the population size that
113 determines how quickly hybrid incompatibilities arise, but its
114 product with the strength of selection κ_β maintaining a particu-
115 lar interaction; so the energy E_{MB} is under very strong selection
116 for strong binding of the morphogen to the 1st binding site and
117 we would expect incompatibilities due to this interaction to arise
118 more slowly when $2N\kappa_{MB} \gg 1$.

119 *Decomposition of DMIs*

120 As discussed in the model section the DMIs shown in [Fig.4](#)
121 will have contributions from many different fundamental in-
122 compatibility types, which can be 2-point, 3-point and 4-point

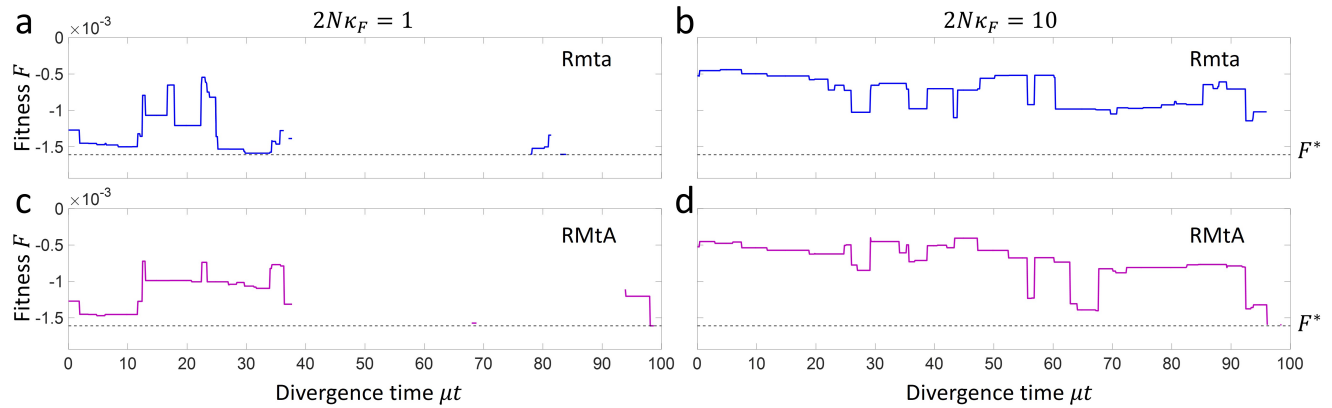


Figure 3 Plot of the times series of two hybrids Rmta (a & b) and RMtA (c & d) at population sizes of $2N\kappa_F = 1$ (a & c) and $2N\kappa_F = 10$ (b & d).

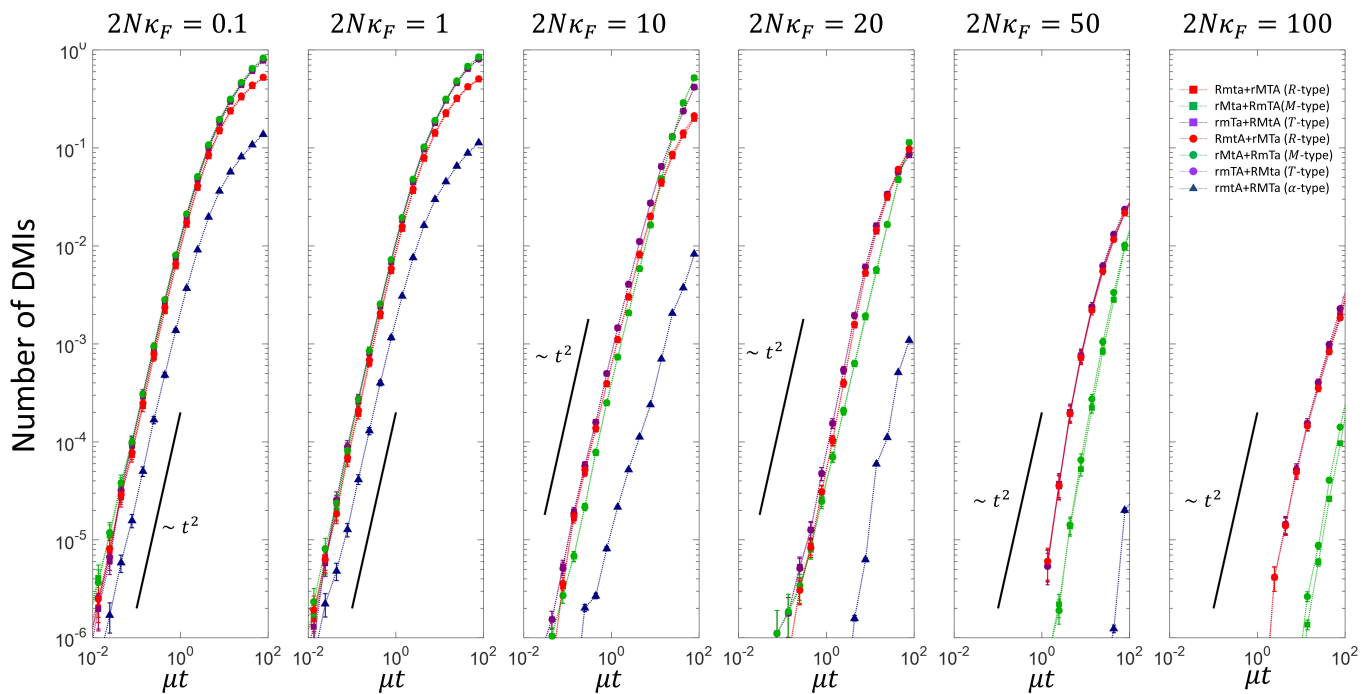


Figure 4 Plot of the number of DMIs for each hybrid genotype since divergence for different scaled population sizes. Complementary hybrid genotypes are summed over, since their behaviour should on average be the same, as there is the same stabilising selective pressure on each lineage.

in nature. Using the method described above to decompose DMIs into fundamental types, we plot the total number of each type of DMI versus divergence time in Fig.5, where the panels correspond to different scaled effective population sizes from $2N\kappa_F = 0.1$ to $2N\kappa_F = 100$. We see clearly that pair-wise DMIs are dominant at all population sizes and divergence times, though the difference is diminished at larger population sizes. In addition, for small population sizes, the relative proportion of 2-point to 3- and 4-point DMIs passes through a minimum at intermediate divergence times. These results show that contrary to the prediction of Orr that higher order DMIs should be easier to evolve, higher order DMIs evolve more slowly and are in smaller number compared to pair-wise DMIs.

As mentioned in the introduction the Orr model also predicts that n -point DMIs should increase as $\sim t^n$. Here, we find that for small population sizes 2-point, 3-point and 4-point DMIs all increase as a power law at small times, indicated by a straight line on a log-log plot, with a larger exponent for 3-point and 4-point DMIs. To more quantitatively assess the exponent, we fit the data for $2N\kappa_F \leq 20$ using the phenomenological equation:

$$I(t) = \frac{I_0 t}{T + t} \left(1 - \exp(-t/\tau)^{\gamma-1}\right) \quad (3)$$

which has the asymptotic form of $I(t) \sim t^\gamma$ for $t \ll \tau$ and $t \ll T$ and an opposite limit of $I(t \rightarrow \infty) = I_0$. We see that for the total number of DMIs and for 2-point, 3-point and 4-point DMIs, this form fits the data well at intermediate and small population sizes. We tabulate the power law exponent derived from these fits in Table 1. We see that the total number of DMIs and 2-point DMIs have a power law exponent close to $\gamma = 2$, which is consistent with the Orr model and with Fig.4 which shows a similar power law, further showing that 2-point DMIs are dominant in determining the growth of hybrid incompatibilities. However, the higher order incompatibilities do not quite follow the Orr prediction, where n -point DMIs have an exponent $\gamma = n$, although 4-point DMIs have a larger exponent than 3-point DMIs; 3-point DMIs have an exponent that varies between $\gamma = 2$ and $\gamma = 3$, while 4-point DMIs have an exponent between $\gamma = 3$ to $\gamma = 3.5$. In all these cases an alternative model for the power law behaviour, as argued in Khatri and Goldstein (2015b), is that at small population sizes, where genetic drift is dominant and there is a large drift load, common ancestor populations are poised at the incompatibility boundary (truncation selection threshold) and the growth of DMIs at short times is due to how likely a few critical substitutions are to arrive very quickly, which is given by a Poisson process; if the critical number of substitutions is K^* then for short times we would expect $P_I(t) \sim (\mu t)^{K^*}$ and so given that at least n substitutions are needed for a n -point incompatibility, we would expect $K^* \geq n$. It is possible the inconsistency here could be resolved by more accurate measurement of the power law, by exploring simulations at even shorter times, using a larger number of replicate simulations (here there are 10^6 replicate simulations at short times).

At larger population sizes ($2N\kappa_F \geq 50$), Fig.5 shows that there is no clear power law and instead there is a negative curvature in the growth of DMIs on a log-log plot. This is consistent with a model of DMI growth where high fitness corresponds to high binding affinity, so that the common ancestor distribution is peaked away from the inviability boundary, which would arise with large populations that have a small drift load, meaning that hybrids diffuse to the boundary to give rise to DMIs; one such analytically tractable model was investigated in Khatri and Goldstein (2015a) and predicted that the number of DMIs is a

$2N\kappa_F$	0.1	1	10	20
Total	1.94 ± 0.05	1.99 ± 0.05	1.95 ± 0.05	2.00 ± 0.12
2-point	1.93 ± 0.05	1.98 ± 0.06	1.84 ± 0.06	1.97 ± 0.13
3-point	2.66 ± 0.13	2.81 ± 0.17	2.76 ± 0.24	2.14 ± 0.15
4-point	3.17 ± 0.19	3.43 ± 0.41	3.14 ± 0.24	2.93 ± 0.19

Table 1 Table of values of the exponent γ characterising the power law of growth of DMIs at short times and small scaled population sizes.

complementary error function, which has an asymptotic form

$I(t) \sim \frac{\sqrt{4\mu t}}{K^*} e^{-(K^*)^2/4\mu t}$, which due to the essential singularity as $t \rightarrow 0$ has the property of negative curvature on a log-log plot. However, this form does not fit the simulation data well (not shown). Given the multidimensional nature of this spatial patterning model, it is possible that we need to consider the analogous result to the effective one-dimensional diffusion studied in Khatri and Goldstein (2015a), which results in a multidimensional generalisation of the error function Brown (1963), where in n dimensions $\text{erf}_n(x) = \Gamma(x^2, n/2) / \Gamma(n/2)$, so that the number of DMIs has asymptotic form $I(t) \sim \left(\frac{4\mu t}{K^*}\right)^{1-\frac{n}{2}} e^{-(K^*)^2/4\mu t}$; this again, however, does not fit the data in Fig.5 well. A functional form that is a good fit for the data at large population sizes is

$$I(t) = \frac{\sqrt{4(\mu t)^\beta}}{K^*} e^{-(K^*)^2/4(\mu t)^\beta}, \quad (4)$$

which arises when considering fractional Brownian processes with exponent β ; normal diffusion or Brownian motion arises when $\beta = 1$, while $\beta < 1$ corresponds to subdiffusive behaviour, while $\beta > 1$ is superdiffusive. It is clear by examining the exponent β in Table 2 from fits of the data in Fig.5 at large population size ($2N\kappa_F \geq 50$) that the DMIs arise as a result of a subdiffusive process, where 2-point, 3-point and 4-point DMIs have an exponent $\beta \approx 1/3$ for $2N\kappa_F = 50$ and $\beta \approx 1/4$ for $2N\kappa_F = 100$. The most likely mechanism that would give rise to subdiffusive behaviour is a broad spectrum of times between substitutions; even though in the simulations the kinetic Monte carlo scheme is based on a Poisson process for a given genotypic state G , the distribution of rates could vary significantly as populations explore the fitness landscape. This would be consistent with the results in Khatri et al. (2009), which reveal the underlying fitness landscape of this spatial patterning genotype-phenotype map to be rough, which could lead to broad distribution of substitution rates in each lineage and effective subdiffusive behaviour of the hybrids Bertin and Bouchaud (2003). Finally as expected the average number of substitutions needed at large population sizes is large, irrespective of the value of n , with values of K^* ranging from 6 to 9.

2-point DMIs In Fig.6 we have plotted the number of 2-point DMIs of each type, where for example, I_{mt} is a 2-point DMI caused by an incompatibility between the M locus and T locus. For each type of 2-point incompatibility there are 2 binding energy traits that could contribute. So increases in I_{mt} could be due to an incompatibility in a hybrid in E_{MB} or E_{MP} ; in this case, as the binding energy E_{MP} is almost neutral Khatri et al. (2009), we would expect incompatibilities to arise predominantly

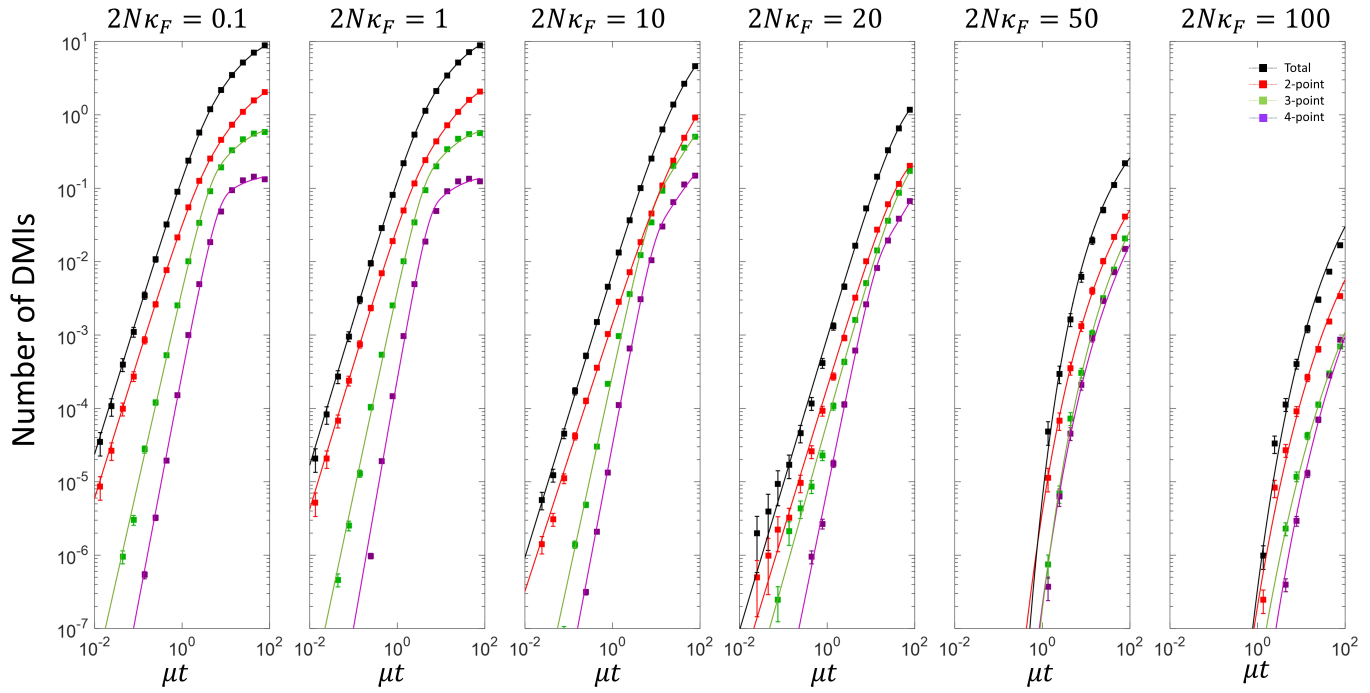


Figure 5 Plot of the total number of DMIs vs divergence time, together with their decomposition into the total number of 2–point, 3–point, 4–point DMIs, for various scaled populations sizes. For $2N\kappa_F \leq 20$ the solid lines correspond to fits of the simulation data to Eqn.3, while for $2N\kappa_F \geq 50$ correspond to fits to Eqn.4.

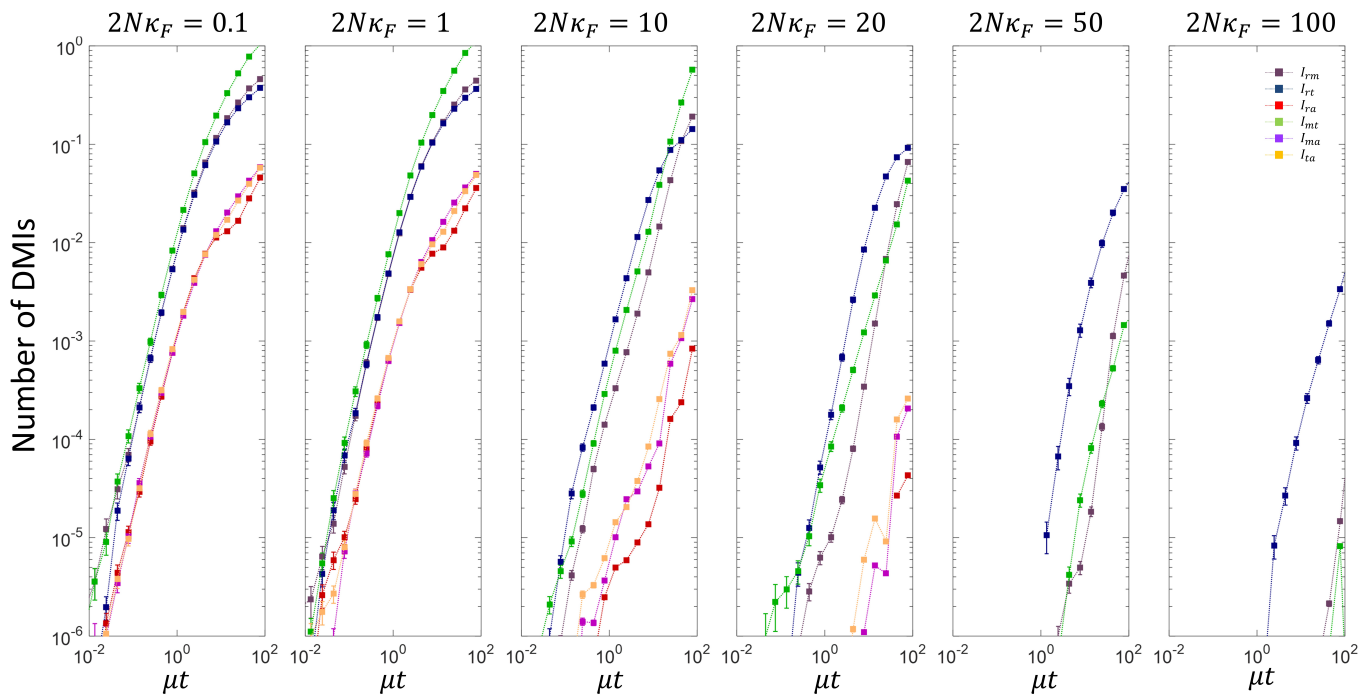


Figure 6 Plot of the spectrum 2–point DMIs vs divergence time for different scaled population sizes.

$2N\kappa_F$		50	100
Total	β	0.47 ± 0.03	0.33 ± 0.02
	K^*	6.58 ± 0.12	7.37 ± 0.26
2-point	β	0.32 ± 0.01	0.25 ± 0.01
	K^*	6.71 ± 0.07	7.51 ± 0.18
3-point	β	0.33 ± 0.01	0.22 ± 0.01
	K^*	7.53 ± 0.10	8.09 ± 0.27
4-point	β	0.31 ± 0.02	0.25 ± 0.02
	K^*	7.50 ± 0.28	8.66 ± 0.39

Table 2 Table of values of the parameters characterising the sub-diffusive growth of DMIs for large scaled population sizes; $\beta = 1$ corresponds to normal diffusive motion, $\beta < 1$ to sub-diffusion and $\beta > 1$ super-diffusion, while K^* corresponds roughly to the number of substitutions required to reach the inviable region.

1 from E_{MB} . Similarly, we expect I_{rt} to be dominated by E_{RP} and
2 not E_{RB} and I_{rm} dominated by δE_{RM} and not δE_{MM} or δE_{RR} .
3 When it comes to incompatibilities involving the α -locus, there
4 is no clear phenotypic trait that can be identified and as we will
5 see the analysis of these DMIs will not be so clear.

6 Examining Fig.6, we see at small population sizes $2N\kappa_F \leq 1$,
7 that all the DMIs grow approximately quadratically at short
8 times with a saturating form at long times, as seen in Fig.5. In
9 addition, we see that when population sizes are small, the num-
10 ber of DMIs for a particular pair wise interaction correlates with
11 the strength of selection on that trait; for example, the binding
12 energy trait E_{MB} , which has the strongest selective constraint,
13 gives rise to the most number of DMIs (I_{mt}) at all times; the next
14 most critical energy trait in terms of selective constraint is δE_{RM} ,
15 which has the next highest number of DMIs (I_{rm}). Although,
16 this observation would appear to be intuitive, interpreting it in
17 light the results of a simple model of transcription factor DNA
18 binding [Khatri and Goldstein \(2015b\)](#) are not straightforward;
19 the results of this work suggest that for small population sizes,
20 the rate that incompatibilities arise decreases with increasing
21 strength of selection and decreasing sequence length, so it is
22 possible these two effects could confound each other. Here for
23 example, E_{MB} has a stronger selective constraint compared to
24 δE_{RM} , but a longer sequence length, as each DNA-protein inter-
25 action interface has 10 binary digits versus each protein-protein
26 interaction interface that has 5. This sequence length effect arises,
27 simply as a result of the fact that the overall substitution rate of a
28 binding energy is proportional to the number of nucleotides that
29 code for it, and so whilst each substitution will have same phe-
30 notypic effect, the hybrid binding energy changes more quickly
31 for longer sequences, as substitutions are more frequent. In ad-
32 dition, it is not obvious how to directly map this single binding
33 interface model onto the multidimensional situation here; one
34 possibility is that the effect of the fitness threshold for truncation
35 selection impacts on each pair-wise interaction differently, for
36 example, where the *average* number of substitutions, that affect
37 the binding energy trait E_{MB} , needed to give rise to an incom-
38 patibility I_{mt} is smaller than for other binding energy traits.

39 However, as the scaled population size increases, we see that

40 the time for I_{mt} incompatibilities to arise sharply increases, while
41 the time for I_{rm} increases less rapidly and I_{rt} even less rapidly.
42 This is consistent with the simple model of transcription factor
43 DNA binding [Khatri and Goldstein \(2015b\)](#) and as observed
44 with the hybrid DMIs in Fig.4, as E_{MB} , which contributes most
45 to I_{mt} is under the greatest selection pressure and so as the pop-
46 ulation size changes these should change most rapidly. We see
47 that for large population sizes, it is not the phenotypic traits
48 under the strongest selection that give rise to significant DMIs
49 at short times, but those under a weaker selective constraint;
50 traits under weaker selection will be affected more by the se-
51 quence entropic pressure for poorer binding affinities and so the
52 common ancestor is more likely to be closer to the inviability
53 boundary. However, if a trait is effectively neutral, i.e. that
54 selection is sufficiently weak that for no values that the trait can
55 take do incompatibilities arise then these will not give rise to
56 incompatibilities; the energy traits E_{RB} , E_{MP} , δE_{RR} and δE_{MM}
57 have this property, as is evident by examining their marginal
58 distribution functions which follow the neutral expectation [Kha-
59 tri et al. \(2009\)](#). The 2-point incompatibilities involving the α
60 locus are more difficult to interpret, since there is no clear trait
61 in the patterning model associated with an interaction solely
62 between the α locus and R , M , or T loci; if α was resolved into a
63 sequence for a protease and it's interaction with a 3rd sequence
64 of the M loci, in addition to the binding and glue sequences, then
65 the value of α itself would be a trait determined by a pair-wise
66 interaction between this 3rd sequence of M and the protease loci,
67 but the current model does not include this feature. The most
68 identifiable phenotype associated with α is the position of the
69 mid-point of the embryo, but this trait involves a co-evolution
70 of E_{MB} and α and so represents a 3-point interaction between M ,
71 T and α loci, which will be discussed below. It is likely that the
72 2-point DMIs involving α are therefore false and a consequence
73 of the parsimonious DMI decomposition method used, which
74 assumes an equal prior on all possible DMIs that have a mini-
75 mum number of DMI types. This could be rectified by having
76 a zero-prior on all pair-wise incompatibilities involving the α
77 locus; however, here we have not implemented this as 2-point
78 DMIs involving α are typically an order of magnitude smaller
79 than the other DMIs.

80 **3-point & 4-point DMIs** In Fig.7, we have plotted the 3-point
81 DMIs as a function of divergence time μt , where the panels
82 from left to right represent increasing scaled population size.
83 We see that for small population sizes, 3-point DMIs between
84 the R , M and T loci dominate at all times and in particular that
85 the different types of DMIs of this type are all roughly equal,
86 $I_{Rmt} \approx I_{rMt} \approx I_{rMT}$. In addition, we see that all other DMIs arise
87 more slowly and each of the 9 other types of 3-point DMIs are
88 all again approximately equal. However, at larger population
89 sizes this degeneracy is lifted amongst the different types of
90 DMIs and different 3-point DMIs grow at different rates. How
91 can we understand this general behaviour?

92 The patterning solution found in these evolutionary simu-
93 lations involves the morphogen binding strongly to the first
94 binding site recruiting RNAP to bind to the promoter to turn
95 on transcription, through a high affinity interaction between the
96 morphogen and RNAP; the spatial position along the length of
97 the embryo where the transcription switches from on to off is
98 controlled by an interaction with the steepness of the morphogen
99 gradient α . Given this, incompatibilities between R , M and T
100 loci could arise through a 3-point interaction where the R loci
101 interacts with the parts of the T and M loci coding for E_{RP} and

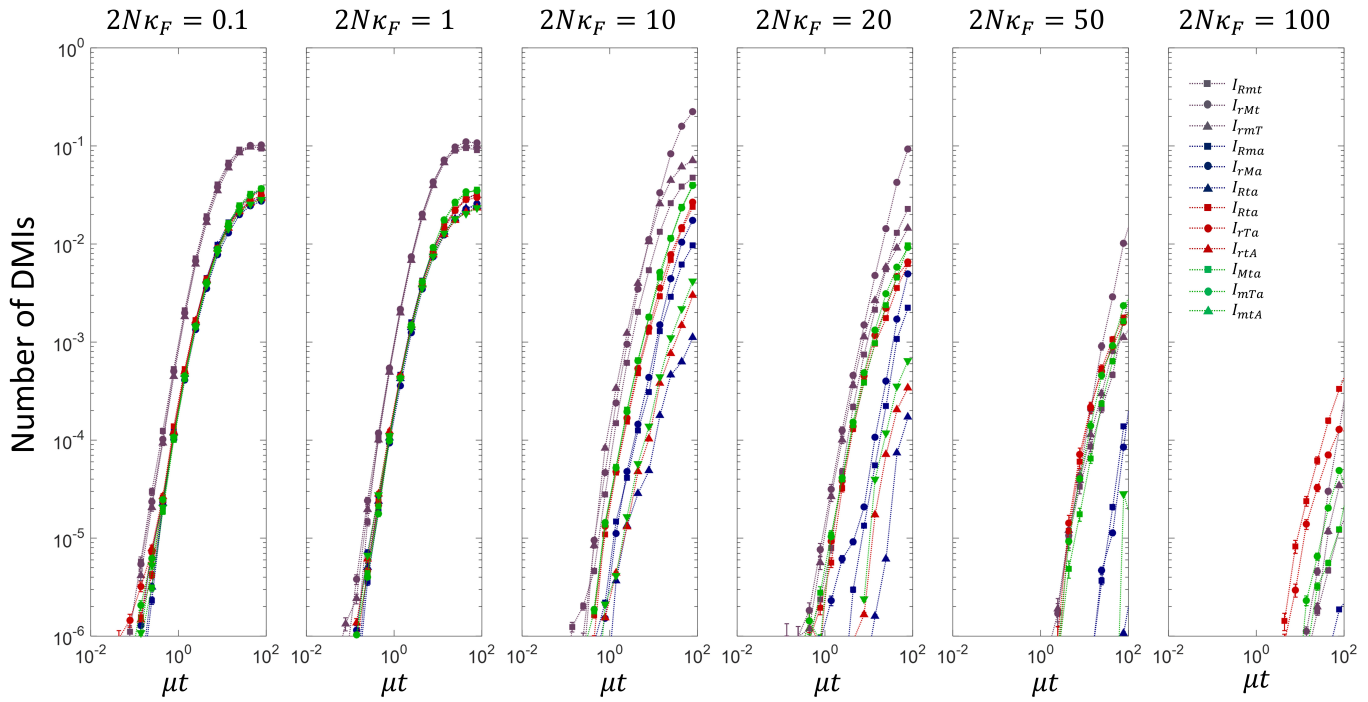


Figure 7 Plot of the spectrum 3–point DMIs vs divergence time for different scaled population sizes.

1 δE_{RM} , or where M loci interacts with the parts of the T and R loci coding for E_{MB} and δE_{RM} . So in analogy to 2–point DMIs, where a pair of loci give rise a single phenotypic binding energy trait, whose value contributes to fitness, here the triplet of loci, R , M and T , give rise to two binding energy traits, which together contribute to fitness. These two traits will co-evolve to maintain good fitness, balanced by the constraints of sequence entropy on the underlying 3 loci which will diminish at large population sizes. On the other hand 3–point incompatibilities between say M , T and α could arise due to an interaction of the E_{MB} binding energy trait with α ; in this model this is subject to a sequence entropy constraint between only two loci. This is true for all the 3–point interactions that involve the α loci. Qualitatively, this then explains the behaviour at low population sizes, as sequence entropy dominates fitness, meaning that the behaviour of the different 3–point DMIs will be dominated by their underlying sequence entropy constraints.

18 The sequence entropy constraints for the 3–point interactions involving the α loci is straightforward and given by a binomial degeneracy function $\Omega(E) = \binom{\ell}{E/\epsilon}$, so that the sequence entropy function $S(E) = \ln(\Omega)$ is approximately quadratic in E , where here E represents one of the binding energies that interacts with α . However, for the other 3–point interactions that don't involve α , but involve the R , M and T loci, the sequence entropy constraint will be related to a degeneracy function $\Omega(E, \delta E) = \Omega(E)\Omega(\delta E)$, where the joint number of sequences that give E and δE is a product, since these energy traits are coded by different sequences, even though they come from the same loci (the joint number of sequences $\Omega(E_{MB}, E_{MP}) \neq \Omega(E_{MB})\Omega(E_{MP})$ since the protein binding sequence of the morphogen that determines E_{MB} and E_{MP} is the same in this case). Given that the joint number of sequences that give E and δE is a product of two binomial coefficients, the sequence entropy function will approximately be a sum of two quadratic terms

35 $S(E, \delta E) = -\frac{2}{\ell_b} (E/\epsilon_b - \ell_b/2)^2 - \frac{2}{\ell_g} (\delta E/\epsilon_g - \ell_g/2)^2$. At small population sizes, where genetic drift dominates selection, we expect the distribution of common ancestors to be distributed such that they are poised at the incompatibility boundary for E and δE ; incompatibilities then arise when substitutions arise that take hybrids across the boundary.

41 Given that a 3–point DMI between the R , M and T genetic loci corresponds to co-evolution of a pair of binding energy traits, instead of a single binding energy trait for 2–point and 3–point DMIs that involve α , means the fraction of substitutions that lead to incompatibilities versus those that keep the hybrids compatible/fit becomes larger when going from one to two dimensions. This then explains why 3–point DMIs between the R , M and T loci gives rise to incompatibilities more quickly than those involving the α loci, as seen in Fig.7. Also since at small population sizes the only influence that fitness will have is in defining the region of incompatibility for the traits of interest, we see that for each type of DMI there are very little differences in the rate of growth of DMIs.

54 4–point DMIs correspond to an interaction where all four loci require a particular combination of alleles for good patterning. As previously noted they are of much smaller number compared to 2– and 3–point DMIs, so here, we do not examine these DMIs in detail. However, we note that the 4–point DMIs shown in Fig.8, show a similar pattern as found with 3–point DMIs, where for small scaled population sizes the DMIs tend to cluster, which suggests, as found for 2– and 3–point DMIs, this is due to sequence entropy constraints dominating the growth of DMIs; on the other hand, at large scaled population sizes this degeneracy is lifted and each hybrid has a different growth rate of DMIs, depending on their particular contribution to fitness and how that balances against the sequence entropy constraint.

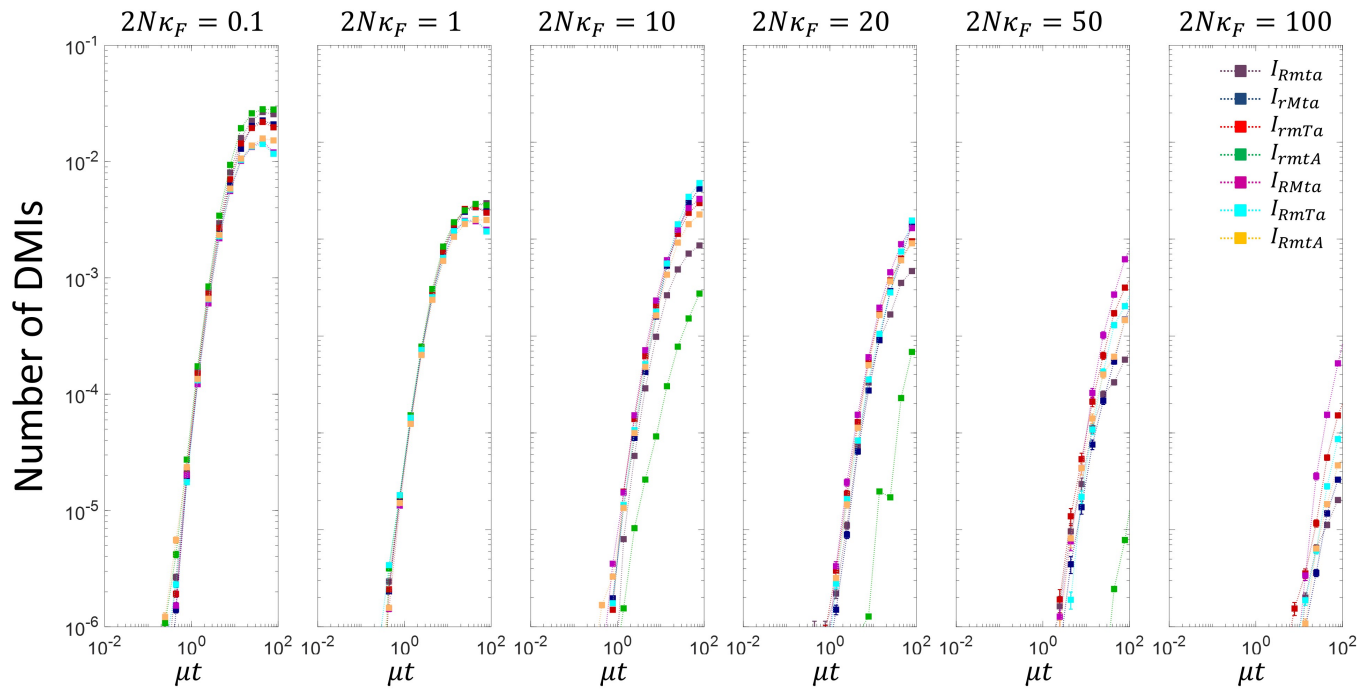


Figure 8 Plot of the spectrum 4–point DMIs vs divergence time for different scaled population sizes.

1 Discussion

2 There is still very little understood about the underlying genetic
 3 basis that gives rise to reproductive isolation between lineages.
 4 Gene expression divergence is thought to be a strong determi-
 5 nant of the differences between species King and Wilson (1975);
 6 Wolf *et al.* (2010); Wray (2007); Abzhanov *et al.* (2006); Wittkopp
 7 *et al.* (2008) with growing body of evidence for their direct role
 8 in speciation, particularly through transcription factors Ting
 9 *et al.* (1998); Brideau *et al.* (2006); Mack and Nachman (2016).
 10 Here building on previous works which modelled the mechanis-
 11 tic basis and growth of DMIs in models of transcription factor
 12 DNA binding Tulchinsky *et al.* (2014b,a); Khatri and Goldstein
 13 (2015a,b), we have investigated the growth of DMIs for a simple
 14 genotype-phenotype map of gene regulation for spatial pattern-
 15 ing in embryonic development, previously studied in Khatri
 16 *et al.* (2009). Our results in this more complicated gene regula-
 17 tory system confirm the basic conclusions from simple models
 18 of transcription factor binding Khatri and Goldstein (2015a,b)
 19 that 1) as the population size decreases below the inverse of
 20 the characteristic scale of fitness ($2N\kappa_F \ll 1$) incompatibilities
 21 arise more quickly, 2) they grow in this regime as a quadratic
 22 power law with divergence time ($P_I \sim (\mu t)^2$) and 3) for large
 23 scaled population sizes incompatibilities arise more slowly with
 24 a characteristic negative curvature on a log-log plot indicative of
 25 a diffusive process. We note that although we find a quadratic
 26 growth of DMIs with divergence time (only at small scaled popu-
 27 lation sizes), which is as predicted by Orr’s framework Orr
 28 (1995), the underlying reason is likely to be very different in
 29 these models and arises as the common ancestor is likely to be
 30 close to the inviable region that gives non-functional binding
 31 Khatri and Goldstein (2015b).

32 In the case of simple models of transcription factor DNA
 33 binding Khatri and Goldstein (2015a,b), smaller diverging popu-
 34 lations, or traits under weaker selection, were found to develop
 35 incompatibilities more quickly, as their common ancestor is al-

36 ready less well adapted due to sequence entropic pressures
 37 dominating fitness at smaller population sizes. Here we see that
 38 this basic principle that incompatibilities arise more quickly due
 39 to a higher drift load of the common ancestor remains valid
 40 for a more complicated gene regulatory system. Although this
 41 question requires greater empirical attention, there is direct and
 42 indirect evidence that smaller populations develop incompati-
 43 bilities more quickly; for example, the greater species diversity
 44 in smaller habitats, such as Hawaii Mayr (1970), the island of
 45 Cuba Glor *et al.* (2004) and East African Great Lakes Santos and
 46 Salzburger (2012); Owen *et al.* (1990), contrasted with the much
 47 slower speciation rate for animals with large ranges or popu-
 48 lation sizes Mayr (1970, 1954); Rubinoff and Rubinoff (1971);
 49 Cooper and Penny (1997). In addition, there is more direct evi-
 50 dence from the net rates of diversification Coyne and Orr (2004)
 51 inferred from phylogenetic trees Nee (2001); Barraclough and
 52 Nee (2001), which support this population size trend.

53 The results of this model have also revealed a number of other
 54 emergent properties for the growth of hybrid incompatibilities,
 55 not obtainable by simply modelling transcription factor DNA
 56 binding. For example, for small populations we find clustering
 57 of the different behaviours of growth of 3–point DMIs, which
 58 can be explained by the different sequence entropy constraints
 59 on different binding energies. Also we found that although the
 60 growth of DMIs at large population sizes has a characteristic
 61 negative curvature on a log-log plot, predicted theoretically by
 62 Khatri *et al.* Khatri and Goldstein (2015a), indicating that hybrid
 63 traits randomly diffuse, a simple model of diffusion does not
 64 fit the simulation data well; instead a model of sub-diffusion
 65 that would arise if there are a number of kinetic traps giving a
 66 broad distribution of substitution times does fit the data well.
 67 This is consistent with the finding that the genotype-phenotype
 68 map has a rough fitness landscape, which is only revealed at
 69 sufficiently large population sizes Khatri *et al.* (2009).

70 However, most importantly we find that pair-wise or 2–point

DMIs dominate compared higher order DMIs (3- and 4- point in this model with 4 loci). This is in contrast to Orr's theoretical argument that the fraction of viable paths from the common ancestor to the current day species increases as we consider higher order DMIs Orr (1995). This argument partly rests on the assumption that the number of inviable genotype remains fixed as a larger number of loci are considered, which would seem a very strong assumption. In the same paper Orr also argues that since there are $\binom{L}{n}$ possible n -point DMIs (the number of combinations of n loci amongst L loci), so as long as $n < L/2$, we would expect an increase in the number of DMIs as n increases; for $L = 4$ as in this paper, this would suggest 2-point DMIs are most numerous, but Orr's calculation in fact undercounts the number of DMIs, which as we show above increases as $(2^n - 2)\binom{L}{n}$, in which case 3-point DMIs, would a priori be more numerous. Our results would then suggest, at least in this simple, but still relatively complex model, that biophysical constraints provide a stronger constraint on the relative number of DMIs of different orders than a purely combinatorial argument would suggest. Although recent results of Weinreich et al Weinreich et al. (2013), would seem to contradict our conclusions, their finding of extensive complex epistasis relates to higher order interactions between sites within a single loci, coding for protein stability or enzymatic activity, whereas our work relates to epistasis between multiple loci.

Of course there is an inherent simplicity with our gene regulatory module for spatial patterning, which requires only two proteins to bind to a regulatory region to turn on transcription; a key direction to investigate would be the effect of multiple transcription factors binding to enhancer regions to control gene expression Bintu et al. (2005); Spitz and Furlong (2012); Levo and Segal (2014), where there could be a large scope for complex epistasis across many loci coding for a large number of transcription factors. However, as our results show, despite the possibility and a prior expectation of a larger number of triplet interactions, pair-wise interactions dominate; for complex transcriptional control, if pair-wise interactions between proteins, and proteins and DNA dominate, for example in determining the binding affinity of transcriptional complexes, then our conclusions would hold.

Overall, our results point to a basic principle, where developmental system drift or cryptic variation True and Haag (2001); Haag (2014); Gavin-Smyth and Matute (2013), play a key role in speciation; basic body plans or phenotypes are conserved, but co-evolution of the components and loci of complicated gene regulatory networks can change differently in different lineages, giving incompatibilities that grow in allopatry in a manner that is controlled by the drift load of the common ancestor, which is turn determined by a balance between selection pushing populations towards phenotypes of higher fitness and genetic drift pushing them towards phenotypes that are more numerous (higher sequence entropy). In particular, although in principle more complicated regulation should give rise to more complex patterns of epistasis Orr (1995), our findings suggest that more simple, pair-wise, incompatibilities dominate the development of reproductive isolation between allopatric lineages under stabilising selection.

Supporting Information

Acknowledgments

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