Sperm chemotaxis is driven by the slope of the chemoattractant concentration field

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Abstract

Spermatozoa are attracted to their conspecific female gamete by diffusive molecules released from the egg investments, a process called chemotaxis. The decapeptide speract induces metabolic and permeability changes in Strongylocentrotus purpuratus sea urchin sperm. In spite of decades since speract purification from S. purpuratus egg investments, sperm chemotaxis has not been demonstrated in this species. By studying how the stimulus function, which spermatozoa experience during the accumulation of bound chemoattractants throughout their trajectory, influences both their motility response and their internal Ca\(^{2+}\) oscillations, we were able to show, for the first time, that S. purpuratus spermatozoa exhibit chemotaxis under sufficiently steep speract concentration gradients. We demonstrate that this process arises through frequency entrainment of the coupled metabolic oscillators.
Introduction

Broadcast spawning organisms, such as marine invertebrates, release their gametes into open sea, where they are often subject to extensive dilution that reduces the probability of gamete encounter (Lotterhos, 2010). In many marine organisms, female gametes release diffusible molecules that attract homologous spermatozoa (Lillie, 1913, Miller, 1985, Suzuki, 1995). Propelled by their beating flagella, spermatozoa detect and respond to chemoattractant concentration gradients by steering their swimming trajectory toward the gradient source: the egg. Though it was in bracken ferns where sperm chemotaxis was first identified (Pfeffer, 1884), sea urchins are currently the best-characterized model system for studying sperm chemotaxis at a molecular level (Alvarez et al., 2012, Cook et al., 1994, Darszon et al., 2008, Strunker et al., 2015, Wood et al., 2015).

The sea urchin egg is surrounded by an extracellular matrix which contains short sperm-activating peptides (SAPs), that modulate sperm motility through altering intracellular Ca\(^{2+}\) concentration ([Ca\(^{2+}\)]\(_i\)) and other signaling intermediates (Darszon et al., 2008, Suzuki, 1995). The probability of sperm-egg encounter is enhanced by the prompt transduction of the biochemical signals triggered by SAPs into the appropriate reorientation of the sperm trajectory.

The decapeptide speract is one of best characterized members of the SAP family due to its powerful stimulating effect on metabolism, permeability and motility in *S. purpuratus* and *L. pictus* spermatozoa. In 2003, it was shown that the binding of speract to its receptor located in the flagellar plasma membrane, triggers a train of [Ca\(^{2+}\)]\(_i\) increases in immobilized *S. purpuratus* spermatozoa (Wood et al., 2003). This calcium signal was thought to regulate the activity of dynein motor proteins in the flagellum, thus might modulate the swimming path of sperm (Brokaw, 1979).
A direct link between [Ca\(^{2+}\)]\(_i\) signaling and sperm motility was established through the use of optochemical techniques to rapidly, and non-turbulently, expose swimming sea urchin spermatozoa to their conspecific attractant in a well-controlled experimental regime (Wood et al., 2005, Bohmer et al., 2005). Currently, it is well established that the transient [Ca\(^{2+}\)]\(_i\) increases triggered by chemoattractants produce a sequence of turns and straight swimming episodes (turn-and-run), where each turning event results from the rapid increase in the [Ca\(^{2+}\)]\(_i\) (Wood et al., 2005, Bohmer et al., 2005, Shiba et al., 2008, Alvarez et al., 2012). The turn-and-run seems to be a general requirement for sperm chemotaxis, however it is not sufficient on its own to produce a chemotactic response (Guerrero et al., 2010a, Strunker et al., 2015, Wood et al., 2007, Wood et al., 2005).

Our current understanding of chemotaxis, suggests that sperm first sample the chemoattractant concentration gradient by swimming in periodic paths (either circular 2D or helical 3D). During the sampling phase, the accumulation of bound chemoattractants triggers [Ca\(^{2+}\)]\(_i\) transients that control the waveform of the flagellar beat. In this way, the alternate periods of asymmetrical (turn) and symmetrical (run) flagellar beating give rise to a looping swimming pattern that guides up to the source of the chemoattractant gradient.

Friedrich and Jülicher proposed a generic theory that captures the essence of sperm navigation following periodic paths in a non-homogeneous chemoattractant field, where the sampling of a periodic concentration stimulus \(s(t)\) is translated by intracellular signaling \(i(t)\) into the periodic modulation of the swimming path curvature \(k(t)\) (Friedrich and Jülicher, 2007, Friedrich and Jülicher, 2008). As result, the periodic swimming path drifts in a direction that depends on the internal dynamics of the
signaling system. In this theory, the latency of the intracellular signaling (the \([\text{Ca}^{2+}]_i\) signal), expressed as the phase shift between \(s(t)\) and \(k(t)\), is a crucial determinant of the directed looping of the swimming trajectory up the chemical concentration field. This theory also predicts that chemotaxis is a robust property of the system that does not require fine-tuning of parameters (if the signaling system is adaptive) (Friedrich and Jülicher, 2008, Friedrich and Jülicher, 2009). In other words, there is a large range of parameters for which sperm chemotaxis is a robust outcome, providing an effective way for sampling the local chemoattractant concentration field and detecting the direction of the concentration gradient (Kaupp et al., 2003, Kashikar et al., 2012, Pichlo et al., 2014, Friedrich and Jülicher, 2008, Friedrich and Jülicher, 2007).

Even though the conceptual framework of Friedrich & Jülicher provides insights into the mechanism governing sperm chemotaxis, it does not explore the scenario whereby chemoattractants trigger an autonomous \([\text{Ca}^{2+}]_i\) oscillator operating in the absence of a periodic stimulus. The existence of an autonomous \([\text{Ca}^{2+}]_i\) oscillator triggered by chemoattractants (Wood et al., 2003, Espinal et al., 2011, Aguilera et al., 2012) suggests that sperm chemotaxis might operate in a dynamical space where two autonomous oscillators, namely the stimulus function and the internal \(\text{Ca}^{2+}\) oscillator, reach frequency entrainment (Pikovsky et al., 2001).

In spite of 30 years of research since speract’s isolation from \(S. \text{purpuratus}\) oocytes (Hansbrough and Garbers, 1981, Suzuki, 1995), chemotaxis of \(S. \text{purpuratus}\) sperm towards this peptide has not yet been demonstrated (Cook et al., 1994, Darszon et al., 2008, Guerrero et al., 2010b, Kaupp, 2012, Miller, 1985, Wood et al., 2015). A comparison between individual \(L. \text{pictus}\) and \(S. \text{purpuratus}\) sperm responses to a specific chemoattractant concentration gradient generated by photoactivating caged speract (CS) revealed that only \(L. \text{pictus}\) spermatozoa exhibit chemotaxis under these conditions.
conditions (Guerrero et al., 2010a). In that study, *L. pictus* spermatozoa experience 
[Ca\(^{2+}\)]\(_i\) fluctuations and pronounced turns while swimming in descending speract 
gradients, that result in spermatozoa reorienting their swimming behavior along the 
positive chemoattractant concentration gradient. In contrast, *S. purpuratus* spermatozoa 
experience similar trains of [Ca\(^{2+}\)]\(_i\) fluctuations that in turn drive them to relocate, but 
with no preference towards the center of the chemoattractant gradient (Guerrero et al., 
2010a).

The precise triggering of [Ca\(^{2+}\)]\(_i\) fluctuations associated to sperm turning events 
towards the chemoattractant gradient, as well as the turn-and-run behavior seem to be 
general requirements for sperm chemotaxis in marine invertebrates (Bohmer et al., 
2005, Guerrero et al., 2010a, Jikeli et al., 2015, Kashikar et al., 2012, Shiba et al., 2008, 
Strunk et al., 2015, Wood et al., 2015), and could be important features of sperm 
chemotaxis in general.

In the present work, we investigate whether *S. purpuratus* spermatozoa can undergo 
chemotaxis. Particularly, we examined whether there is a physical limit to the sampling 
of the chemoattractant concentration gradient needed for detection that has, to date, 
prevented observation and characterization of their chemotactic response. We report 
that *S. purpuratus* spermatozoa are chemotactic only when exposed to much steeper 
speract concentration gradients than those previously employed. Furthermore, we 
explored the coupling between the recruitment of speract molecules during the sperm 
voyage, the triggered signaling cascade and the internal Ca\(^{2+}\) oscillator; and demonstrate 
that sperm chemotaxis arises through coupled metabolic oscillators.
Results

Chemotaxis refers to the directed movement of an organism or a cell in a chemical gradient. The first step in a chemotactic response is the sampling of a chemoattractant field, from where in a given time \( T \), statistical fluctuations limit the precision with which the searcher can determine the concentration of the chemoattractant (Berg and Purcell, 1977, Vergassola et al., 2007, Dusenbery, 2011).

Assuming that the number of ligands that collide with a single receptor per unit of time is proportional to radius of the receptor \( s \), the concentration of the attractant \( c \) and the diffusion constant of the attractant \( D \), the total number of collisions between \( N_R \) receptors in a time interval \( T \) is \( x = N_R s D c T \). The statistical fluctuation on the number of ligand-receptor collisions \( x \) follows a Poisson distribution, and therefore the expected value \( E[x] \) and the variance \( Var[x] \) are the same. In this scenario, the coefficient of variation is \( \frac{SD[x]}{E[x]} = \frac{\sqrt{Var[x]}}{E[x]} = \frac{\sqrt{E[x]}}{E[x]} \). This coefficient can be interpreted as the smallest fractional error attainable in the determination of the concentration of the attractant (see Theory section) (Berg and Purcell, 1977, Dusenbery, 2011, Vergassola et al., 2007).

\[
 u_1 = (N_R s D c T)^{-1/2} \quad (1)
\]

Uncertainty in the determination of the presence of chemoattractant molecules

The flagella of sea urchin spermatozoa possess a large number of high affinity SAP receptors (\( 10^4 \) - \( 10^6 \) depending on species and author) (Nishigaki and Darszon, 2000, Shimomura and Garbers, 1986, Smith and Garbers, 1982, Kaupp et al., 2008). This great number of SAP receptors provides spermatozoa with a fine-tuned sensory system able to respond to a wide range of SAP concentrations (\( 10^{12} \) - \( 10^6 \) M). That these receptors are distributed along the entire flagellum increases the probability of capturing even a few arriving molecules (Kaupp et al., 2003, Kashikar et al., 2012).
To understand the theoretical sampling capabilities of sea urchin spermatozoa, given a specific speract concentration (in a homogeneous concentration field), and flagellar receptor availability, we first computed the smallest fractional error attainable in the detection of a chemoattractant $u_l$ at different concentrations, considering distinct receptor numbers. A theoretical $S. purpuratus$ spermatozoon possessing $N_r = 2 \times 10^4$ receptors, will sample a medium containing 10 nM of speract in time $T = 0.5$ s with an uncertainty of $u_l = 0.006$ (Figure 1 – figure supplement 1a and Figure 1 – table supplement 1). Here, the time scale $T$ was considered as the time needed for a $S. purpuratus$ spermatozoon to traverse a distance equivalent to approximately half a circumference of its circular trajectory when confined to the water-glass boundary, in a homogeneous field of speract. The least fractional error attainable in the determination of the concentration of the attractant has no upper boundary, however $u_l$ values greater than 1 indicates that stochastic fluctuations dominate over the signal (Berg and Purcell, 1977). For simplicity, $u_l$ was considered to operate in the range $[0, 1]$. In other words, our model spermatozoon, when exposed to a chemoattractant concentration of 10 nM, will be able to determine the presence of speract molecules with 0.6% of uncertainty in 0.5 s.

If the spermatozoon finds itself in a sampling space of 1 pM of speract, this level of uncertainty increases to 60% over the same interval. It is likely that at speract concentrations below picomolar, an $S. purpuratus$ spermatozoon will require relatively extended sampling times to determine the presence of speract molecules. For example, for a spermatozoon sampling a 1 pM concentration field of speract for 60 seconds, the uncertainty of receptor occupancy drops to 6% (Figure 1 – figure supplement 1b).

An important difference between $S. purpuratus$ and $L. pictus$ spermatozoa is their respective receptor density; the former possesses 3.15x fewer receptors than the latter (2
x $10^4$ versus $6.3 \times 10^4$ receptors per cell, respectively) (**Figure 1 – table supplement 1**)

(Nishigaki and Darszon, 2000, Nishigaki et al., 2001). Note that if *S. purpuratus* spermatozoa were to possess a receptor density equivalent to *L. pictus*, they would be able to sample a speract concentration field of 1 pM within 0.5 seconds with an uncertainty of $u_1 = 0.4$ (40%) (**Figure 1 – figure supplement 1a**). Given the differences in sensitivities between these two species, the lack of chemotactic responses in *S. purpuratus* spermatozoa could possibly be explained, at least partially, by their lesser density of speract receptors, compared to *L. pictus*.

**Uncertainty in the determination of sperm position in a non-homogeneous chemoattractant concentration field**

To detect a chemoattractant concentration gradient, the signal at two given points, sampled throughout a time interval must be greater than the noise, which in this case would arise from spontaneous stochastic activation of one or more components of the chemoattractant signaling pathway (**Figure 1a**) (Berg and Purcell, 1977, Vergassola et al., 2007, Wood et al., 2015, Dusenbery, 2011). Information regarding the direction of the chemoattractant gradient results from the magnitude difference of the signal at two sampled positions (for further explanation see **Theory** section). The smallest fractional error in determination of the chemoattractant gradient direction due to the occupation state of receptors is:

$$u_2 = v^{-1}T^{-3/2}(N_R S D)^{-1/2}c^{1/2}e^{-1}$$  \(2\)

where $v$ is the swimming velocity and $e = dc/dr$ is the chemoattractant concentration gradient (**Figure 1a**) (Berg and Purcell, 1977, Dusenbery, 2011, Vergassola et al., 2007). A locality requirement must also be met: the change in concentration across the distance spanned during $T$ must be small, compared to the concentration itself.
The slope of the chemoattractant concentration gradient $\varepsilon$ directly impinges on the ability to reliably determine the source of the attractant (Figure 1b). At high concentrations of speract ($10^{-8}$ to $10^{-6}$ M) the change in receptor occupancy experienced by *S. purpuratus* spermatozoa, at two given distinct locations, allows the reliable assessment of the chemical gradient, when $u_2 < 0.05$ (Figure 1b) (As for $u_1$, $u_2$ was considered to operate in the range [0, 1]) (Vergassola et al., 2007, Berg and Purcell, 1977). However, at low concentrations of speract (below of $10^{-9}$ M), where other parameters are identical, stochastic fluctuations tend to dominate over the signal (Figure 1b). Under these low-concentration regimes, the slope of the chemoattractant gradient is determinant; shallow gradients ($\varepsilon < 10^{-12}$ M $\mu$m$^{-1}$) increase the uncertainty for detection of gradient polarity. In contrast, steeper chemoattractant gradients, *i.e.* $\varepsilon \geq 10^{-11}$ M $\mu$m$^{-1}$, would allow *S. purpuratus* spermatozoa to determine the orientation of the chemical gradient dependably (Figure 1b).

Previously it was shown that *L. pictus* spermatozoa experience chemotaxis in response to exposure to a gradient of speract (Guerrero et al., 2010a). Notably, under the same experimental regime *S. purpuratus* spermatozoa undergo a motility response to the presence of speract, yet do not demonstrate chemotaxis. It seems reasonable to speculate that the lack of chemotactic responses in *S. purpuratus* spermatozoa is due to their reduced receptor density relative to *L. pictus* spermatozoa. The chemoreception model predicts that, in contrast to *L. pictus*, *S. purpuratus* spermatozoa might determine the direction of the attractant gradient in the range of $10^{-10}$ to $10^{-9}$ M (Figure 1b, orange and yellow lines with $u_2 < 5\%$), and thus can only determine the orientation of the chemical gradient in steep, but not shallow concentration gradients (Figure 1b, yellow and cyan lines with $u_2 < 5\%$). If the latter holds to be true, then the model predicts that
S. purpuratus spermatozoa should be able to experience chemotaxis when exposed to relatively steep speract gradients, which have not been tested experimentally to date.

The ratio of uncertainties for reliable assessment of the direction of a speract gradient between S. purpuratus and L. pictus spermatozoa can be expressed as:

\[
\frac{u_2 \text{purpuratus}}{u_2 \text{pictus}} = S \frac{\varepsilon_\text{pictus}}{\varepsilon_\text{purpuratus}},
\]

\[
S = \left( \frac{v_\text{purpuratus}}{v_\text{pictus}} \right)^{-1} \left( \frac{T_\text{purpuratus}}{T_\text{pictus}} \right)^{-3/2} \left( \frac{N_R_\text{purpuratus}}{N_R_\text{pictus}} \right)^{-1/2}
\]

(3, 4)

with \(S\) being a factor that scales the slope of the speract gradient to a regime that allows S. purpuratus spermatozoa to detect the local direction of the chemical gradient. In other words, the model predicts that S. purpuratus spermatozoa should undergo chemotaxis in a speract gradient approximately three times steeper than the gradient that drives chemotaxis in L. pictus spermatozoa, with \(\varepsilon_\text{purpuratus} = 3.25 \varepsilon_\text{pictus}\).

In summary, the chemoreception model suggests that S. purpuratus spermatozoa detect chemoattractant gradients with less sensitivity than those of L. pictus. It also predicts that S. purpuratus spermatozoa may detect chemoattractant gradients in the \(10^{-9}\) M regime with sufficient certainty only if the difference in sampling concentration is greater than \(10^{-11}\) M \(\mu\)m\(^{-1}\) (steep concentration gradients) (Figure 1b). Given this prediction, we designed and implemented an experimental condition where S. purpuratus spermatozoa should experience chemotaxis.

S. purpuratus spermatozoa accumulate at steep speract concentration gradients

Our experimental setup is designed to generate specific concentration gradients by focusing a brief (200 ms) flash UV light along an optical fiber, through the objective, and into a field of swimming S. purpuratus spermatozoa containing CS at 10 nM in artificial sea water (Tatsu et al., 2002, Guerrero et al., 2010a). To test experimentally whether S. purpuratus undergo chemotaxis, as predicted from the chemoreception
model, we varied the slope of the chemoattractant gradient by separately employing four optical fibers of distinct diameters that could be arranged into five different configurations ($f1$, $f2$, $f3$, $f4$, $f5$).

Each configuration produces a characteristic pattern of UV illumination within the imaging field (Figure 2). The UV intensity was measured at the back focal plane of the objective for each fiber configuration (Figure 2a and Figure 2 – table supplement 1). The spatial derivative of the imaged UV light profile was computed for use as a proxy for the slope of the speract concentration gradient (Figure 2b). By examining these UV irradiation patterns, we calculated that, at the gradient peak, the highest concentration of speract released through photo-liberation from CS is generated by the $f5$ fiber, followed by $f4 > f3 > f2 > f1$. The steepest UV irradiation gradients are those generated by the $f2$, $f3$ and $f5$ fibers (Figure 2b).

Irrespective of the optical fiber used, the photo-activation of caged speract triggers the stereotypical $\text{Ca}^{2+}$-dependent motility responses of $S. \text{purpuratus}$ spermatozoa (Figure 2c, Movies 1, 2, 5, 6 and 7). To determine whether these changes lead to sperm accumulation, we developed an algorithm which automatically scores the number of spermatozoa at any of four defined concentric regions (R1, R2, R3, and R4) relative to the center of the speract concentration gradient (Figure 2 – figure supplement 1).

Photo-liberation of speract through either $f2$ or $f3$ fiber, but not through $f1$, $f4$ or $f5$ fibers, lead to the accumulation of $S. \text{purpuratus}$ spermatozoa towards the center of the speract gradient (zones R1 and R2) within the first 5-10 seconds after UV irradiation (Figure 3, Figure 3 - figure supplement 1, Figure 3 - figure supplement 2, Movies 1, and 2).

Interestingly, for $f2$, the number of spermatozoa increases in R1 and R2 and decreases in R3 and R4, indicating that cells from the R3 and R4 regions most probably
relocate towards R1 and R2 (Figure 3a, Figure 3 - figure supplement 2 and Movie 1). In the case of the $f3$ gradient, the number of spermatozoa increases in the R1, R2 and R3 regions and decreases in the R4 region, suggesting that spermatozoa in R4 and possibly outside of the imaged field are entering the other regions (Figure 3a and Movie 2). The maximum sperm accumulation (about two-fold) occurs in R1 for the $f3$ gradient (Figure 3a and Movie 2).

In the case of exposure to the $f5$ speract concentration gradient the number of spermatozoa showed a tendency to increase in R2, R3 and R4, although it was only statistically significant in R4 (Figure 3a, Figure 3 - figure supplement 2 and Movie 6). Gradients $f1$, $f4$ and negative controls (Low [Ca$^{2+}$]$_i$ or High [K$^+$]$_e$) did not show increased sperm number in any region (Figure 3a, Figure 3 - figure supplement 1, Figure 3 - figure supplement 2, Movies 3, 4, 5 and 6).

We also evaluated the corresponding [Ca$^{2+}$]$_i$ changes across the imaging field for each imposed speract concentration gradient. Gradients $f2$ to $f5$ increase [Ca$^{2+}$]$_i$ in spermatozoa at least two-fold, while the increase for the $f1$ gradient was modest (Figure 3b). Interestingly, [Ca$^{2+}$]$_i$ levels rose highest upon exposure to the $f4$ gradient, even though the number of spermatozoa did not increase significantly (Figure 3a), which underlines the notion that elevated [Ca$^{2+}$]$_i$ levels are necessary, but not sufficient to drive the accumulation of spermatozoa (Figure 3), as previously suggested (Kaupp et al., 2003, Wood et al., 2005, Bohmer et al., 2005, Guerrero et al., 2010a, Alvarez et al., 2012).

In summary, S. purpuratus spermatozoa accumulate towards the center of the speract gradients generated by the $f2$- and $f3$-fibers, which are the two optical fibers that generate UV light profiles with steeper slopes compared to the $f1$ and $f4$ fibers (Figure 2b). Notably, use of fibers $f4$ and $f5$ photo-releases higher concentrations of speract (by
providing higher UV energies than other fibers) (Figure 2a, Figure 2 table supplement 1), yet they trigger neither the maximum accumulation of S. purpuratus spermatozoa at the center of the chemoattractant field, nor the most elevated [Ca$^{2+}$], response.

S. purpuratus spermatozoa undergo chemotaxis upon exposure to steep speract gradients

The spatial derivative of the UV profiles shown in Figure 2b indicates that the steeper light gradients generated from UV irradiation are those of f2, f3 and f5, which are assumed to generate the most pronounced speract gradients of similar form. This assumption is strictly only valid at the instant of UV exposure, as subsequently the speract gradient dissipates over time with a diffusion constant of $D \approx 240 \mu$m$^2$ s$^{-1}$.

We further sought to understand how the stimulus function, which S. purpuratus spermatozoa experience during the accumulation of bound speract throughout their trajectory, influences their motility response. For this purpose, we computed the spatio-temporal dynamics of the speract gradient for f1, f2, f3, f4 and f5 fibers (Figure 4a, b and Figure 4 - figure supplement 1) and analyzed the trajectories of spermatozoa swimming in these five distinct speract gradient configurations (Figure 4c and Figure 4 - figure supplement 2). From these trajectories, we derived the chemotaxis indices (LECI, (Yoshida et al., 2002)) (Figure 4d and Figure 5), and computed the stimulus function of individual spermatozoa in response to each of the five speract gradient forms (Figure 4e and Figure 4 - figure supplement 2).

Fibers f2, f3 and f5 triggered chemotaxis in S. purpuratus spermatozoa (Figure 4, Figure 4 - figure supplement 2 and Figure 5). Under such conditions, only those spermatozoa located in the regions R3 and R4 at the moment of speract uncaging...
underwent chemotaxis (Figure 5; see Figure 4 and Figure 4 - supplement figure 2 for
single sperm reorienting towards the center of speract gradients generated by the f2 and
f3 fibers, respectively). For both, f2 and f3 conditions, sperm chemotaxis initiated within
the first 3 s after speract exposure, and lasted less than 10 s (Figure 5, Figure 3 - figure
supplement 1, Movies 1 and 2).

Figure 4b shows that, for the f2 condition, the slope of the speract gradient within
R2 and regions is barely altered during the initial 3 seconds of speract exposure (gray
shading). Spermatozoa in that region sample a pseudo-static speract gradient over a
short (2-3 s) period, which is apparently sufficient to permit detection of the direction of
the gradient. In contrast, spermatozoa located at other regions (R1, R3, R4) experience a
speract gradient with a slope that changes during the sampling time (Figure 4a, 4b).

In summary, f2, f3 and f5 optical fibers generate speract concentration gradients that
drive the accumulation of S. purpuratus spermatozoa. Of note, the f4 fiber generates a
speract gradient that activates motility responses of S. purpuratus spermatozoa but not
chemotaxis (Movie 6), although it was previously shown that it triggers the
accumulation of L. pictus spermatozoa (Guerrero et al., 2010a).

The model of chemoreception presented in the previous section (equations (3, 4))
predicts a scaling rule for chemotactic responses between S. purpuratus and L. pictus
spermatozoa of S > 3.25. Moreover, the UV-irradiation profiles shown in figure 2b
indicate that f2, f3, and f5 fibers generate steeper speract gradients than f4 fiber.

To reliably determine the direction of the chemoattractant concentration gradient, the
signal difference dc between two sampled positions dr must be greater than the noise
(Figure 1a). In order to test the prediction of the chemoreception model, we computed
the local relative slope ξ detected by single spermatozoa exposed to a given speract
concentration gradient, with ξ = c^{-1/2} dc/dr (Figure 4e).
We found that, in agreement with the chemoreception model, the maximum relative slope $\xi_{\text{max}} = \text{Max}(\xi_1, \xi_2, \xi_3, \ldots, \xi_n)$ experienced by *S. purpuratus* spermatozoa, when exposed to $f_2$ and $f_3$ speract gradients, were up to 2-3 times greater than that experienced when exposed to the $f_4$-generated speract gradient (*Figure 6a*). In addition, they were also up to 2-3 times greater than the relative slope experienced by *L. pictus* spermatozoa when exposed to $f_4$ speract gradient (*Figure 6a*), hence supporting the predicted scaling rule for the detection of the speract concentration gradient between *L. pictus* and *S. purpuratus* spermatozoa.

The slope of the speract gradient is the critical determinant for the strength of coupling between the stimulus function and the internal $\text{Ca}^{2+}$ oscillator

To test the idea that the slope of the speract gradient regulates the coupling between the stimulus function and the internal $\text{Ca}^{2+}$ oscillator triggered by speract, we made use of a generic model for coupled phase oscillators (Pikovsky et al., 2001). In its simplest form, the model describes two phase oscillators of intrinsic frequencies $\omega_1$ and $\omega_2$ coupled with a strength $\gamma$ through the antisymmetric function of their phase difference $\phi = \phi_1 - \phi_2$. The time evolution of $\phi$ then follows an Adler equation $d\phi/dt = \Delta \omega - 2\gamma \sin(\phi)$, which is the leading order description for weakly-coupled non-linear oscillators. In the present case, the two coupled oscillators are the internal $\text{Ca}^{2+}$ oscillator and the oscillations in the stimulus function induced in spermatozoa swimming across a speract gradient. The former occurs even for immotile cells, for which there are no stimulus oscillations under a spatially uniform speract field (*Figure 6 - supplement 1*, and *Movie 8*); while the later exists under two tested negative controls: cells swimming in Low $\text{Ca}^{2+}$ and in High $K^+$ artificial sea water, both of which inhibit $\text{Ca}^{2+}$ oscillations (see *Figure 3, Movie 3 and 4*).
There are two immediate predictions from the Adler model: first, there is a minimum coupling strength necessary for the two oscillators to synchronize \((\gamma_{\text{min}} = \Delta \omega / 2)\). For weaker coupling \((i.e. \gamma < \gamma_{\text{min}})\), the two oscillators run with independent frequencies and, hence, the phase difference increases monotonically with time; second, and within the synchronous region \((i.e. \gamma > \gamma_{\text{min}})\), the phase difference between the oscillators is constant and it does not take any arbitrary value, but rather follows a simple relation to the coupling strength \((\phi_{\text{sync}} = \arcsin(\Delta \omega / 2\gamma))\). Figure 6b shows the two regions in the parameter space given by \(\Delta \omega\) and \(\gamma\). The boundary between these two regions corresponds to the condition \(\gamma = \gamma_{\text{min}}\) and it delimits what is known as an Arnold’s tongue.

We measure the difference in intrinsic frequency by looking at the instantaneous frequency of the internal Ca\(^{2+}\) oscillator just before and after the speract gradient is established. The range of measured \(\Delta \omega\) is shown in the same figure as a band of accessible conditions in our experiments (mean of \(\Delta \omega\), black line; mean ± standard deviation, green dashed lines). If the driving coupling force between the oscillators is the maximum slope of the speract gradient, \(i.e. \gamma = \zeta_{\text{max}}\), we would expect to find a minimum slope \((\zeta_{\text{max}}^{\text{min}})\) below which no synchrony is observed.

This is indeed the case as clearly shown in Figure 6a, 6c and 6d (magenta line). Moreover, and for cells for which synchronization occurs, the measured phase difference is constrained by the predicted functional form of \(\phi_{\text{sync}} = \phi_{\text{sync}}(\Delta \omega, \gamma)\) as can be verified from the collapsed data shown in Figure 6c, and 6d within the theoretical estimates.

Altogether, the excellent agreement of this simple model of coupled phase oscillators with our data, points to the slope of the speract gradient as the driving force behind the observed synchronous oscillations and, as a result, for the chemotactic ability of sea
urchin spermatozoa.
Discussion

Marine spermatozoa, together with many motile microorganisms, explore their environment via helical swimming paths, whereupon encountering a surface these helices collapse to circular trajectories. The intrinsic periodicity of either swimming behavior commonly results in the periodic sampling of the cells’ chemical environment with direct implications for their ability to accurately perform chemotaxis.

A strict requirement for sperm chemotaxis is the presence of extracellular Ca\(^{2+}\). For chemotaxis to occur, the timing of the Ca\(^{2+}\) transients (*i.e.* the intracellular Ca\(^{2+}\) oscillations) triggered by the chemoattractants must also be kept in phase with the polarity of the chemoattractant concentration field, which in this, and other studies, is referred as the stimulus function (Bohmer et al., 2005, Guerrero et al., 2010a, Kaupp et al., 2008, Friedrich and Jülicher, 2008). This requisite coupling ensures that the turning events start at the descending phase of the chemoattractant concentration field; otherwise spermatozoa are driven away by Ca\(^{2+}\)-dependent motility adjustments. The periodic sampling of chemoattractants by the sperm flagellum continuously feeds back to the signaling pathway governing the intracellular Ca\(^{2+}\) oscillator, hence providing a potential coupling mechanism for sperm chemotaxis. Indirect evidence for the existence of a feedback loop operating between the stimulus function and the Ca\(^{2+}\) oscillator triggered by chemoattractants has been found in *L. pictus, A. punctulata* and *Ciona intestinalis* species, whose spermatozoa show robust chemotactic responses towards their conspecific chemoattractants (Guerrero et al., 2010a, Shiba et al., 2008, Jikeli et al., 2015, Bohmer et al., 2005).

For almost three decades, chemotaxis had not been observed for the widely studied *S. purpuratus* species under diverse experimental conditions, raising doubts about their capabilities to ‘sense’ and respond to the spatial cues provided by the speract
concentration gradients. To tackle whether *S. purpuratus* spermatozoa are able to ‘sense’ a chemoattractant concentration gradient, we use a model of chemoreception developed by Berg and Purcell, which considers the minimal requirements needed for a single searcher (*i.e.* a sperm cell) to gather sufficient information to determine the orientation of a non-uniform concentration field. By considering the difference between *L. pictus* and *S. purpuratus* spermatozoa in terms of the number of chemoattractant receptors, receptor pocket size, sampling time, swimming velocity, sampling distance, and the local mean and slope of the chemoattractant concentration field, we predicted that *S. purpuratus* should be able to detect the polarity of a speract concentration field. The model predicts that speract gradient necessary to guide *S. purpuratus* spermatozoa would be up to three times steeper than the gradient that drives chemotactic responses on *L. pictus* spermatozoa. We tested this prediction experimentally by exposing *S. purpuratus* spermatozoa to various defined speract concentration gradients. We show that *S. purpuratus* spermatozoa exhibit chemotactic responses but, as predicted by the chemoreception model, only if the speract concentration gradients are sufficiently steep (*i.e.* speract gradients that are at least three times steeper than the speract concentration gradient that drives chemotaxis in *L. pictus* spermatozoa). The shallower speract gradients previously tested are therefore unable to generate any chemotactic response in *S. purpuratus* spermatozoa.

To investigate further the molecular mechanism involved in sperm chemotaxis, we measured both the stimulus function and the triggered-internal Ca\(^{2+}\) oscillations for up to five hundred *S. purpuratus* spermatozoa exposed to five distinct speract concentration gradients. We demonstrate that the slope of the chemoattractant concentration field is a major determinant for sperm chemotaxis in *S. purpuratus*, and might be an uncovered feature of sperm chemotaxis in general. A steep slope of the
speract gradient entrains the frequencies of the stimulus function and the internal Ca\textsuperscript{2+} oscillator triggered by the periodic sampling of a non-uniform speract concentration field. We assessed the transition boundary of the coupling term (the slope of the speract concentration field) for the two oscillators to synchronize, and found it to be very close to the boundary where *S. purpuratus* start to experience chemotaxis. The agreement of our data with a model of weakly coupled phase oscillators, points to the slope of the speract gradient as the driving force behind the observed synchronous oscillations and, as a result, for the chemotactic ability of sea urchin spermatozoa.

One can further hypothesize about the evolutionary origin of the described differences in sensitivity to chemoattractant concentration gradients between *S. purpuratus* and *L. pictus* spermatozoa given the significant differences between their ecological reproduction niches. The turbulent environment where sea urchin reproduce directly impinges on the dispersion rates of small molecules such as speract, hence, imposing ecological pressure against sperm chemotaxis. For instance, the reproduction success of *L. pictus*, *S. purpuratus* and Abalone species has been shown to peak at particular hydrodynamic shearing values (Zimmer and Riffell, 2011, Riffell and Zimmer, 2007, Mead and Denny, 1995, Hussain et al., 2017). We might ask what are the typical values of the chemoattractant gradients encountered by the different species in their natural habitats. The correct scale to consider when discussing the small-scale distribution of chemicals in the ocean is the Batchelor scale, \( l_B = (\eta D^2 / \zeta)^{1/4} \), where \( \eta \) is kinematic viscosity, \( D \) the molecular diffusivity and \( \zeta \) is the turbulent dissipation rate (Batchelor, 2006, Aref et al., 2014). Turbulence stirs dissolved chemicals in the ocean, stretching and folding them into sheets and filaments at length scales down to the Batchelor scale: below \( l_B \) molecular diffusion dominates and chemical gradients are smoothened out.
S. purpuratus is primarily found in the low intertidal zone. The purple sea urchin thrives amid strong wave action and areas with churning aerated water. These more energetic zones, including tidal channels and breaking waves, generate relatively high levels of turbulence ($\zeta \sim 10^{-4} \text{ m}^2\text{s}^{-3}$) which lead to relatively small values of $l_B$ and, hence, to steep gradients (i.e. $l/l_B$). L. pictus, on the contrary, is mostly found at the edge of or inside kelp beds, well below the low tide mark where the levels of turbulence are much more moderate ($\zeta \sim 10^{-6} \text{ m}^2\text{s}^{-3}$) (Jiménez, 1997, Thorpe, 2007). This difference in turbulent kinetic energy dissipation rate has a significant effect on the largest chemical gradients available in a particular habitat for sperm chemotaxis. The ratio of $l_B$ for the different habitats scales as $l_B_{\text{purpuratus}}/l_B_{\text{pictus}} \sim (\zeta_{\text{pictus}}/\zeta_{\text{purpuratus}})^{1/4} \sim 3$, which fits considerably well with the relative sensitivity to speract of the two species. Furthermore, we have shown that S. purpuratus spermatozoa experience chemotaxis toward steeper speract gradients than those that guide L. pictus spermatozoa, which is also compatible with the distinct chemoattractant gradients they might naturally encounter during their voyage searching for the egg.

The chemoattractant concentration gradients generated in the present study were near-instantaneously set up by the photo-release of speract in still water. Further experimental studies are needed to assess the chemotactic ability of sea urchin spermatozoa to more realistic chemoattractant gradients (as those shaped, for instance, by hydrodynamic forces in their natural environment) and to shed light into the mechanisms governing chemotaxis and their ecological implications.
Materials and Methods

Materials

Undiluted *S. purpuratus* or *L. pictus* spermatozoa (Pamanes S. A. de C. V., Ensenada, Mexico and Marinus Scientific, LLC. Newport Beach, CA, USA respectively) were obtained by intracoelomic injection of 0.5 M KCl and stored on ice until used within a day. Artificial seawater (ASW) was 950 to 1050 mOsm and contained (in mM): 486 NaCl, 10 KCl, 10 CaCl$_2$, 26 MgCl$_2$, 30 MgSO$_4$, 2.5 NaHCO$_3$, 10 HEPES and 1 EDTA (pH 7.8). For experiments with *L. pictus* spermatozoa, slightly acidified ASW (pH 7.4) was used to reduce the number of spermatozoa experiencing spontaneous acrosome reaction. Low Ca$^{2+}$ ASW was ASW but pH 7.0 and with 1 mM CaCl$_2$, and Ca$^{2+}$-free ASW was ASW with no added CaCl$_2$. [Ser5; nitrobenzyl-Gly6]speract, referred to throughout the text as caged speract (CS), was prepared as previously described (Tatsu et al., 2002). Fluo-4-AM and pluronic F-127 were from Molecular Probes, Inc. (Eugene, OR, USA). PolyHEME [poly(2-hydroxyethylmethacrylate)] was from Sigma-Aldrich (Toluca, Edo de Mexico, Mexico).

Loading of Ca$^{2+}$-fluorescent indicator into spermatozoa

Undiluted spermatozoa were suspended in 10 volumes of low Ca$^{2+}$ ASW containing 0.2% pluronic F-127 plus 20 µM of fluo-4-AM and incubated for 2.5 h at 14 °C. Spermatozoa were stored in the dark and on ice until use.

Imaging of fluorescent swimming spermatozoa

The cover slips were briefly immersed into a 0.1% wt/vol solution of poly-HEME in ethanol, hot-air blow-dried to rapidly evaporate the solvent and mounted on reusable chambers fitting a TC-202 Bipolar temperature controller (Medical Systems Corp.). The temperature plate was mounted on a microscope stage (Eclipse TE-300; Nikon) and
maintained at a constant 15 ºC. Aliquots of labeled sperm were diluted in ASW and transferred to an imaging chamber (final concentration ~2x10^5 cells ml^{-1}). Epifluorescence images were collected with a Nikon Plan Fluor 40x 1.3 NA oil-immersion objective using the chroma filter set (ex HQ470/40x; DC, 505DCXR; em, HQ510LP) and recorded on a DV887 iXon EMCCD Andor camera (Andor Bioimaging, NC). Stroboscopic fluorescence illumination was supplied by a Cyan LED no. LXHL-LE5C (Lumileds Lighting LLC, San Jose, USA) synchronized to the exposure output signal of the iXon camera (2 ms illumination per individual exposure). Images were collected with AndorIQ 1.8 software (Andor Bioimaging, NC) at 30.80 fps in full-chip mode (observation field of 200 x 200 µm).

*Image processing*

The background fluorescence was removed by generating an average pixel intensity time-projection image from the first 94 frames before uncaging, which was then subtracted from each frame of the image stack by using the Image calculator tool of ImageJ 1.49u (Schneider et al., 2012). For Figure 2c, the maximum pixel intensity time projections were created every 3 s from background-subtracted images before and after the UV flash.

*Quantitation of global changes of spermatozoa number and [Ca^{2+}],*

To study the dynamics of overall sperm motility and [Ca^{2+}], signals trigger by the distinct speract gradients, we developed an algorithm that provides an efficient approach to automatically detect the head of every spermatozoa in every frame of a given video-microscopy (C/C++, OpenCV 2.4, Qt-creator 2.4.2). Fluorescence microscopy images generated as described previously were used. The following steps summarize the work-flow of the algorithm (Figure 2 - figure supplement 1):

1. Segment regions of interest from background: This step consists of thresholding
each image (frame) of the video (x, y, t) to segment the zones of interest (remove noise and atypical values). Our strategy includes performing an automatic selection of threshold value for each Gaussian blurred image ($I_G$) ($\sigma = 3.5 \, \mu m$) considering the mean value ($M_i$) and the standard deviation ($S_i$) of the image $I_G$.

The threshold value is defined by: $T_i = M_i + 6S_i$.

2. Compute the connected components: The connected components labeling is used to detect connected regions in the image (a digital continuous path exists between all pairs of points in the same component - the sperm heads). This heuristic consists of visiting each pixel of the image and creating exterior boundaries using pixel neighbors, accordingly to a specific type of connectivity.

3. Measure sperm head fluorescence. For each region of interest, identify the centroid in the fluorescence channel (sperm head) and measure the mean value.

4. Compute the relative positions of the sperm heads within the imaging field, and assign them to either R1, R2, R3 or R4 concentric regions around the centroid of the UV flash intensity distribution. The radii of R1, R2, R3 or R4, were 25, 50, 75 and 100 $\mu m$, respectively.

5. Repeat steps 1 to 4 in a frame-wise basis.

Step 1 of the algorithm filters out shot noise and atypical values; step 2 divides the images in N connected components on the location of the sperm heads; step 3 quantitates sperm head fluorescence, and finally step 4 computes the relative sperm position on the imaging field. A similar approach has been recently used to identify replication centers of adenoviruses in fluorescence microscopy images (Garces et al., 2016).

We automatically analyzed up to 267 videos of $S. \ purpuratus$ spermatozoa, each containing tens of swimming cells, exposed to five distinct speract concentration
gradients.

Computing the dynamics of speract concentration gradients

The dynamics of the chemoattractant gradient was computed using the Green’s function of the diffusion equation:

\[ c(r, t) = \frac{C_0}{\sqrt{4\pi D t}} e^{-r^2/4D} + C_b \]  

(5)

This equation for the concentration tells us that the profile has the form of a Gaussian. The width of the Gaussian is \( \sigma = \sqrt{4D(t + t_0)} \), and hence it increases as the square root of the time. \( C_b \) is the basal concentration of the chemoattractant, \( D \) is the molecular diffusivity.

The speract concentration gradients were generated via the photolysis of 10 nM caged speract (CS) with a 200 ms UV pulse delivered through each of four different optical fibers with internal diameters of 200 µm, 600 µm, 2 mm and 4 mm (at two different positions). Light intensity was normalized dividing each point by the sum of all points of light intensity for each fiber and multiplying it by the fiber potency (measured at the back focal plane of the objective) in miliwatts (mW) (Figure 2 - table supplement 1). Each spatial distribution of instantaneously-generated speract concentration gradient was computed by fitting their corresponding normalized spatial distribution of UV light (Residual standard error: \( 2.738 \times 10^{-5} \) on 97 degrees of freedom), considering an uncaging efficiency of 5%, as reported (Tatsu et al., 2002).

The diffusion coefficient of a molecule is related to its size by the Stokes-Einstein equation:

\[ D = \frac{kT}{6\pi\eta R_h} \]  

(6)

where \( k \) is Boltzmann’s constant, \( T \) is the temperature, \( \eta \) is the viscosity of the solvent, and \( R_h \) is the hydrodynamic radius (Lakowicz, 2006). The hydrodynamic radius...
\( R_h \) of speract was calculated by modeling the molecules in terms of equivalent hydrodynamic spheres.

\[
R_h = \left( \frac{3MV}{4\pi} \right)^{1/3}
\]

(7)

where \( M \) is the molecular weight, and \( \bar{\rho} \) is the specific gravity (Lakowicz, 2006). The volume of an equivalent spherical particle is \( V_e = \frac{4}{3}\pi R_h^3 \). Equations 6 and 7 show that the radius and diffusion coefficient are weakly dependent on the molecular weight.

The diffusion coefficient of speract has not been measured experimental, nonetheless it can be estimated following equations 6 and 7. The diffusion coefficient of a similar chemoattractant molecule, resact (with fourteen amino acids), has been reported, \( D_{resact} = 239 \pm 7 \, \mu m^2 s^{-1} \) (Kashikar et al., 2012). If we consider that speract is a decapeptide, the 1.4 fold difference in molecular weight between speract and resact would imply a \((1.4)^{1/3}\) fold difference in their diffusion coefficients, which is close to the experimental error reported (Kashikar et al., 2012). For the sake of simplicity, the spatio-temporal dynamics of the distinct instantaneously generated speract gradients was modeled considering a speract diffusion coefficient of \( D_{speract} = 240 \, \mu m^2 s^{-1} \).

The hydrodynamic radius of speract \( (R_h = 8.8 \, \text{Å}) \) was computed with equation (6), with \( D_{speract} = 240 \, \mu m^2 s^{-1} \), \( k = 1.38 \times 10^{-23} \, \text{J K}^{-1} \), \( T = 288.15 \, \text{K} \) and \( \eta = 0.001 \, \text{N s /m}^2 \).

**Computing \([Ca^{2+}]_i\) dynamics, chemotactic behavior, and the stimulus function of single spermatozoa**

Spermatozoa were tracked semi-automatically by following the head centroid with the MtrackJ plugin (Meijering et al., 2012) of ImageJ 1.49u. Single cell \([Ca^{2+}]_i\) signals were computed from the mean value of a 5 x 5 pixels region, centered at each sperm head along the time.

Chemotactic behavior was quantified using the linear equation chemotaxis index \((LECI)\) employing different temporal windows before (control) and after uncaging.
LECI is defined as the negative value of the slope \( \text{LECI} = -l \) of a least square linear regression \( r(t) = lt + r_0 \) (Yoshida et al., 2002) where \( r(t) \) is the sperm head distance to the center of the speract concentration gradient, \( t \) is the time after UV pulse measured in seconds, with \( t \in [0 - 3; 3.2 - 6.2; 6.2 - 9.2] \) seconds. Positive \( \text{LECI} \) indicate movement towards the chemoattractant source. The center of the speract gradient was identified as the centroid of the UV flash intensity distribution.

The head position of each spermatozoa was used to compute the local concentration of speract at \( r(x, y) \) over each frame. The stimulus function of single spermatozoa \( s = c(r, t) \) was computed by solving equation (5), considering both their swimming trajectories, and the spatio-temporal evolution of a given speract concentration gradient. The profiles of UV light were used to compute the initial conditions at \( c(r, t_o) \).

The phase- and temporal-shifts between time derivative of the stimulus function \( ds/dt \) and the internal \( \text{Ca}^{2+} \) oscillator triggered by speract, were computed from their normalized cross-correlation function.

**Analysis of speract induced \( \text{Ca}^{2+} \) transients with immobilized spermatozoa.**

Imaging chambers were prepared by coating cover slips with 50 \( \mu \text{g/ml} \) poly-D-lysine, shaking off excess, and allowing to air-dry. Coated cover slips were then attached to imaging chambers. Fluo-4 labeled spermatozoa were diluted 1:40 in ASW, immediately placed into the chambers, and left for 2 min, after which unattached sperm were removed by washing with ASW. The chambers were then filled with 0.5 ml of ASW containing 500 nM of caged speract, and mounted in a TC-202 Bipolar temperature controller (Medical Systems Corp.). Images were collected with Andor iQ 1.7 software (Andor Bioimaging, NC) at 90 fps in full-chip mode, binning 4x4 (observation field of 200 \( \mu \text{m} \times 200 \mu \text{m} \)). The imaging setup was the same as that used for swimming spermatozoa. The caged speract was photo-released with a 200 ms UV
pulse delivered through an optical fiber (4 mm internal diameter) coupled to a Xenon UV lamp (UVICO, Rapp Opto Electronic). The optical fiber was mounted on a “defocused” configuration to minimized the generation of UV light heterogeneities.

Images were processed off-line using ImageJ 1.4.5s. Overlapping spermatozoa and any incompletely adhered cells, which moved during the experiment, were ignored. Fluorescence measurements in individual sperm were made by manually drawing a region of interest around the flagella with the line tool of ImageJ.

Programs were written in R statistical software.

**Theory**

*Constraints over the detection of chemoattractant concentration fields.*

Strategies for chemoattractant sampling based on local concentration gradients require the concentration to be high enough to ensure that the measured average difference at two nearby locations is larger than typical fluctuations (Wood et al., 2015). As stochastic fluctuations limits the precision with which a microorganism can, in a given time $T$, determine the concentration of a surrounding chemoattractant (Berg and Purcell, 1977, Vergassola et al., 2007), the rate of encounter between a chemoattractant $c$ diffusing with effective diffusivity $D$ and its receptor, with a binding site of effective radius $s$, is given by:

$$J(r) = 4\pi s D c(r)$$

(8)

For reliable assessment of the local concentration $c(r)$, a spermatozoon provided with $N_R$ number of chemoattractant receptors, collects detection events over time $T_{int}$. The average number of detection events will then be $J(r)N_R T_{int}$. Typical stochastic fluctuations are of the order of the square root of the mean (Berg and Purcell, 1977). Then the condition for the signal to emerge from the noise reads:
\[ \sqrt{N_R s D c T_{\text{int}}} \ll 1 \]  

(9)

Then, it follows that the spermatozoon will experience an uncertainty \( u_1 = (N_R s D c T_{\text{int}})^{-1/2} \) (eq. (1)), for measuring a variation on the local concentration of chemoattractant (Figure 1 - figure supplement 1).

To reliably measure a concentration gradient of chemoattractant, the difference in counts collected by the spermatozoon in the interval measurement must be above the noise level (Figure 1a). The corresponding conditions can be stated as (Berg and Purcell, 1977, Vergassola et al., 2007):

\[
\left( v T_{\text{int}} \frac{dc}{dr} \right) N_R s D T_{\text{int}} \gg \sqrt{N_R s D c T_{\text{int}}} 
\]

(10)

\[
v T_{\text{int}} \frac{d \log c}{dr} \ll 1
\]

(11)

Here, \( v \) is the swimming velocity and \( dc/dr \) is the concentration gradient of the chemoattractant. The first inequality indicates that the signal-to-noise ratio for the difference in the number of hits experienced by the swimming spermatozoon across the integration time \( T_{\text{int}} \) be larger than unity. The second inequality is the requirement of locality, i.e. the change in concentration across the distance spanned during \( T_{\text{int}} \) needs to be small compared to the local concentration itself.

The information to determine the direction of the chemoattractant gradient results from the magnitude difference of the signal at two sampled positions. Then the condition for the signal to emerge from the noise on a chemoattractant field reads:

\[
v^{-1} T^{-3/2} (N_R s D)^{-1/2} c^{1/2} \epsilon^{-1} \ll 1
\]

(12)

where \( \epsilon = dc/dr \) is the chemoattractant concentration gradient. Then, it follows that the least fractional error of the chemoattractant gradient direction due to the state of occupation of receptors is \( u_2 = v^{-1} T^{-3/2} (N_R s D)^{-1/2} c^{1/2} \epsilon^{-1} \) (Figure 1, eq. (2)).
Assuming an exponentially decaying concentration, reliable integration time $T_{\text{int}}$ scales as $\exp(r/3\lambda)$ (Vergassola et al., 2007).

Statistical analyses

Data are presented for individual spermatozoa ($n$) collected from up to three sea urchins. All statistical tests were performed using R software (R Core Team, 2016). The significance level was set at 95%.

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Author contributions

A.G., A.D. and I.T. conceived the project; A.G. and V.J.S. performed the experiments; H.R.G., A.G., I.T., V.J.S. and M.V. analyzed the data; A.G., I.T. and H.R.G performed the mathematical model calculations and wrote the corresponding section; A.D., A.G., H.R.G. and I.T. participated in the design and drafting of the manuscript J.C., C.D.W. and C.B. provide feedback for conceptualization and drafting of the manuscript. All authors approved the final version of the article.

Competing interests

The authors declared that no competing interests exist.
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Figure 1. Physics of chemoreception. a. Determining the direction of the chemoattractant gradient requires that the signal difference $dc$ between two sampled positions $dr$ must be greater than the noise. b-d. The uncertainty in the determination of the chemoattractant gradient direction, $u_2$, plotted against the slope of the gradient, $dc/dr$, in log-log scale, for different chemoattractant concentrations, (b) S. purpuratus, (c) L. pictus, and (d) A. punctulata spermatozoa (See Figure 1 - table supplement 1 for the list of parameter values taken in consideration for panels b-d).
Figure 1 - figure supplement 1. Uncertainty for the detection of a homogeneous chemoattractant concentration field. The uncertainty for determining the chemoattractant concentration $u_1$ versus the receptor number $N_R$ for different concentrations ($10^{-6}$ to $10^{-12}$ M), for a sampling interval of $T = 0.5$ s (a), or a $T = 60$ s (b). Vertical lines indicate the number of receptors of Sp = *S. purpuratus* (dark green), Lp = *L. pictus* (dark blue) and Ap = *A. punctulata* (black) spermatozoa. Parameters: sampling interval $T = 0.5$ s (a) and 60 s (b), speract diffusion coefficient $D = 239 \, \mu m^2 s^{-1}$, hydrodynamic radius of speract (as proxy of receptor pocket size) $R_h = 0.79$ nm (see figure 1 - figure supplement 2 and figure 1 - table supplement 1).
<table>
<thead>
<tr>
<th>Spermatozoan species</th>
<th><em>S. purpuratus</em></th>
<th><em>L. pictus</em></th>
<th><em>A. punctulata</em></th>
</tr>
</thead>
<tbody>
<tr>
<td><em>N</em>&lt;sub&gt;R&lt;/sub&gt; [per cell]</td>
<td>2 x 10&lt;sup&gt;4b&lt;/sup&gt;</td>
<td>6.3 x 10&lt;sup&gt;4b&lt;/sup&gt;</td>
<td>3 x 10&lt;sup&gt;3&lt;/sup&gt;</td>
</tr>
<tr>
<td><em>D</em> [cm&lt;sup&gt;2&lt;/sup&gt;s&lt;sup&gt;-1&lt;/sup&gt;]</td>
<td>2.4 x 10&lt;sup&gt;-6c&lt;/sup&gt;</td>
<td>2.4 x 10&lt;sup&gt;-6c&lt;/sup&gt;</td>
<td>2.4 x 10&lt;sup&gt;-6c&lt;/sup&gt;</td>
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<tr>
<td><em>s</em> [cm]</td>
<td>1.65 x 10&lt;sup&gt;-7c&lt;/sup&gt;</td>
<td>1.65 x 10&lt;sup&gt;-7c&lt;/sup&gt;</td>
<td>1.65 x 10&lt;sup&gt;-7c&lt;/sup&gt;</td>
</tr>
<tr>
<td><em>T</em> [s]</td>
<td>0.39 ± 0.08&lt;sup&gt;a&lt;/sup&gt;</td>
<td>0.52 ± 0.22&lt;sup&gt;a&lt;/sup&gt;</td>
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<tr>
<td><em>v</em> [cms&lt;sup&gt;-1&lt;/sup&gt;]</td>
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<td>144 ± 36 x 10&lt;sup&gt;-4a&lt;/sup&gt;</td>
<td>200 x 10&lt;sup&gt;-4&lt;/sup&gt;</td>
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<tr>
<td><em>dr</em> [cm]</td>
<td>28 ± 6 x 10&lt;sup&gt;-4a&lt;/sup&gt;</td>
<td>46 ± 14 x 10&lt;sup&gt;-4a&lt;/sup&gt;</td>
<td>60 x 10&lt;sup&gt;-4&lt;/sup&gt;</td>
</tr>
</tbody>
</table>

Figure 1 - table supplement 1. Parameters of the chemoattractant sampling model for each species. Note that the main differences between species are the number of receptors *N*<sub>R</sub> and the velocity *v* of the spermatozoon. *D* diffusion coefficient; *s* hydrodynamic radius of speract (as proxy of chemoattractant receptor's pocket radius); *T* sampling time (time to swim half circumference in the boundary close to the water-glass interface); *dr* = sampling interval (circumference diameter). <sup>a</sup>Measured in this study (mean ± sd; *N* = 3 sea urchins; *n* = 495 (*S. purpuratus*), 56 (*L. pictus*)) spermatozoa. <sup>b</sup>(Nishigaki et al., 2001, Nishigaki and Darszon, 2000). <sup>c</sup>Calculated in this study (see section: Computing the dynamics of speract concentration gradients in Materials and Methods, and Figure 1 - figure supplement 2).
**Figure 1 - figure supplement 2. Speract and resact modeling.** Five speract and resact structures, respectively, were modeled using PEP-FOLD (Maupetit et al., 2009, Thevenet et al., 2012, Maupetit et al., 2010), and used to estimate an equivalent hydrodynamic radius by computing the length of the mayor axis of each model: a. Speract (16.18 ± 0.41 Å). b. Resact (18.97 ± 1.79 Å). The estimates of the mayor axis of the speract model were similar to their corresponding hydrodynamic diameter estimates \( d = 2R_h = 2(8) \text{ Å} = 1.6 \text{ nm} \) computed with the Stokes-Einstein equation \( R_h = 0.88 \text{ nm} \) (see section: *Computing the dynamics of speract concentration gradients* in **Materials and Methods**). Models were visualized with chimera 1.11.2 (Pettersen et al., 2004).
Figure 2. Screening of speract concentration gradients. a. Radial profile of the UV light scattered at the glass-liquid interface for each optical fiber (f1-f5). b. Derivatives of radial distribution for each optical fiber. c. Spatial distribution of the UV flash energy (left), and typical motility and [Ca^{2+}], responses of spermatozoa exposed to different concentration gradients of speract (right). F-F₀ time projections, showing spermatozoa head fluorescence each 3 s before and after 200 ms UV photoactivation of 10 nM caged speract in artificial sea water. The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca^{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca^{2+}]. Scale bars of 50 µm.
Figure 2 - figure supplement 1. Automatic segmentation of swimming spermatozoa. a. Working-flow of the algorithm: Video microscopy images were background subtracted by removing the temporal average intensity projection (static fraction) of the first un-stimulated frames (93 frames = 3 seconds) from the whole video (25 seconds). The resulting images were convolved with a low-pass spatial frequency filter to eliminate typical detector shot-noise. The resulting images were thresholded to generate arrays of regions of interest (ROIs), a heuristic search for connected components is then applied to label single ROIs and to assign the corresponding pixels to unique spermatozoa. Scale bar of 15 µm. b. The positions of the sperm heads within the imaging field are assigned to either R1, R2, R3 or R4 concentric regions around the centroid of the UV flash intensity distribution. Each ROI was also used to compute the sperm head fluorescence from the raw video microscopy images (as the mean value of the ROI). Scale bar of 50 µm.
<table>
<thead>
<tr>
<th>Physical diameter (mm)</th>
<th>UV power at the back focal plane of the objective (mW)*</th>
</tr>
</thead>
<tbody>
<tr>
<td>$f_1$ 0.2</td>
<td>0.07</td>
</tr>
<tr>
<td>$f_2$ 0.6</td>
<td>1.25</td>
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<tr>
<td>$f_3$ 2</td>
<td>4.7</td>
</tr>
<tr>
<td>$f_4$ 4</td>
<td>7.8</td>
</tr>
<tr>
<td>$f_5$ 4</td>
<td>9.46</td>
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</tbody>
</table>

Figure 2 - Table supplement 1. Physical diameter of the optical fibers, and UV light power measured at the back focal plane of a 40×/1.3NA oil-immersion objective.

*Typically, there is a 20% loss of light power due to scattering within the optics.
Figure 3. Motility and \([\text{Ca}^{2+}]_i\) responses of *S. purpuratus* spermatozoa exposed to specific concentration gradients of speract. 

**a.** Fold change in sperm number, defined as the number of spermatozoa at the peak of the response (6 s) relative to the mean number before speract stimulation (0-3 s) (see figure 3 - figure supplement 1). 

**b.** Relative changes in \([\text{Ca}^{2+}]_i\) experienced by spermatozoa at the peak response (6 s) after speract stimulation. Negative controls for spermatozoa chemotaxis are artificial sea water with nominal \(\text{Ca}^{2+}\) (Low \(\text{Ca}^{2+}\)); and artificial sea water with 40 mM of \(\text{K}^+\) (High \(\text{K}^+\)). Both experimental conditions prevent chemotactic responses by inhibiting the \(\text{Ca}^{2+}\) membrane permeability alterations triggered by speract; the former disrupts the \(\text{Ca}^{2+}\) electrochemical gradient, and the later disrupt the \(\text{K}^+\) electrochemical gradient required as electromotive force needed for the opening of \(\text{Ca}^{2+}\) channels. Number of experiments (n) on the top of each experimental condition. Each box contains 50\% of single events, the inner lines indicate the median and the error bars delimit the 95\% outliers. The notch display a confidence interval around the median (median +/- 1.57 x IQR\(^{1/2}\)). If two boxes notches do not overlap, there is evidence that their medians differ (95\% confidence). *Statistical significance, \(p < 0.05\); Multiple comparison test after Kruskal
Wallis. Studies at the population level are presented in **figure 3 - figure supplement 1**, while studies considering the absolute number of spermatozoa at two periods: 3s and 6s (3 seconds before and after speract exposure, respectively) are presented in **figure 3 - figure supplement 2**.
Figure 3 - figure supplement 1. Sperm response to speract photo-release, collated data from individual experiments. Sperm motility responses to different speract concentration gradients ($f_1$, $f_2$, $f_3$, $f_4$, $f_5$) at R1 (a), R2 (b), R3 (c) and R4 (d) concentric regions. Negative controls for sperm chemotaxis are artificial sea water with nominal 0 Ca$^{2+}$ (Low Ca$^{2+}$); and artificial sea water with 40 mM of K$^+$ (High K$^+$). Each time trace, represent the mean sperm density from up to 20 video microscopy experiments (raw data distributions at two periods 3s and 6s are presented in figure 3 - figure supplement 2). Note that peak responses occurred at second 6, three seconds after speract exposure (indicated as vertical dotted lines).
Figure 3 - figure supplement 2. Absolute number of sperm cells, three seconds before and three seconds after speract photo-release through either f1, f2, f3, f4, f5 optical fibers. Absolute number of spermatozoa scored at either R1, R2, R3 or R4 concentric regions around the centroid of the UV flash intensity distribution. Each box contains 50% of single events, the inner lines indicate the median and the error bars delimit the 95% outliers. *Statistical significance, p < 0.05; Mann-Whitney test.
Figure 4. Steep speract gradients attract S. purpuratus spermatozoa. a. Dynamics of the f2 speract gradient. The blue dashed line \((t_0 = 0 \text{ s})\) corresponds to a Gaussian distribution fitted to the UV light profile, and illustrates the putative shape of the instantaneously-generated speract concentration gradient. Solid black lines illustrate the temporal evolution of the speract concentration field after \(t = 1, 2, 3, \ldots, 20 \text{ seconds}\). b. Temporal changes in \(f2\) speract field computed radially (each 10 µm) from the center of the gradient. A region where the concentration of speract barely changes is shown in gray. c. Characteristic motility changes of a S. purpuratus spermatozoon exposed to the \(f2\) speract gradient. Solid lines illustrate its swimming trajectory 3 s before (gray) and 6 s after (black) speract exposure. d. Spermatozoon head distance to the source of the speract gradient versus time, calculated from black trajectory in c. LECI index is computed from the least squared linear regression of the head distance to the source, a negative slope (as shown) indicates chemotaxis (positive LECI). e. Stimulus function computed from the swimming behavior of the spermatozoon in c, considering the dynamics of a.
Figure 4 - figure supplement 1. Modeling of the dynamics of speract gradient based on the UV light profile of distinct optical fibers. The radial profiles of the UV light scattered at the glass-liquid interface of f1, f3, f4, f5 optical fibers are shown in gray. The speract gradient was generated as in figure 4, but with the corresponding f1 (a), f3 (c), f4 (e) and f5 (g) optical fibers. Left panels - The dynamics of the speract gradient computed as is in figure 4. The blue dashed line (t0 = 0) corresponds to a
Gaussian distribution fitted to the UV light profile, and illustrates the putative shape of
the instantaneously generated speract gradient. Solid black lines illustrate the shape of
the speract gradient after $t = 1, 2, 3, \ldots, 20$ seconds. Right panels - Simulated temporal
changes in speract concentration gradients of $f_1$ (a), $f_3$ (c), $f_4$ (e) and $f_5$ (g) at each 10
radial $\mu$m point from the center of the concentration gradient.
Figure 4 - figure supplement 2. Characteristic motility changes of a *S. purpuratus* spermatozoon exposed to *f3* and *f4* speract gradients (chemotactic vs non-chemotactic response). Panels a, b and c show single cell responses to the *f3* speract gradient (chemotactic); and panels d, e and f to the *f4* speract gradient (non-chemotactic). a, d. Solid lines illustrate the spermatozoon swimming trajectory 3 s before (gray) and 6 s after (black) speract gradient exposure. b, e. Spermatozoon head distance to the source of the speract gradient versus time calculated from a and d, respectively. b, e. Stimulus function computed from a and d, considering the spatio-temporal dynamics of speract computed for the *f3* and *f4* gradients, respectively.
Figure 5. *S. purpuratus* spermatozoa selectively experience chemotaxis towards speract gradients. Chemotactic behavior quantified using the linear equation chemotaxis index (LECI) before (a), as control experiment, and after exposure to either $f1, f2, f3, f4$ or $f5$ speract gradients (b, c). The color scale indicates the mean LECI value calculated from tens of *S. purpuratus* spermatozoa located in either R1, R2, R3, or R4 regions prior speract photo-release (3 s after UV light irradiation (b), or 6 s after UV light exposure (c)). Positive LECIs (shown in red) indicate movement towards the chemoattractant source. Note that spermatozoa that were in R3 and R4 distal regions prior speract gradient exposure experience chemotaxis only in response to $f2, f3$, and $f5$ speract gradients (b, c). Numerals indicate the number of spermatozoa analyzed in each experimental condition. *Statistical significance, $p < 0.05$; Student's t-test.
Figure 6. The slope of the speract gradient generates a frequency-locking phenomenon between the stimulus function and the internal Ca\textsuperscript{2+} oscillator triaggered by speract. a. Maximum relative slopes ($\xi_{\text{max}}$) experienced by *S. purpuratus* (Sp) when exposed to $f_1, f_2, f_3, f_4, f_5$ speract gradients. The maximum relative slopes experienced by *L. pictus* spermatozoa (Lp) towards $f_4$ experimental regime are also shown. Note that $\xi_{\text{max}}$ for $f_2, f_3$, and $f_5$, are up to 2-3 times greater than in $f_4$, regardless of the species. b. Arnold’s tongue indicating the difference in intrinsic frequency of the internal Ca\textsuperscript{2+} oscillator of *S. purpuratus* spermatozoa, just before and after the speract gradient exposure. c. Phase difference between the time derivative of the stimulus function and the internal Ca\textsuperscript{2+} oscillator of *S. purpuratus* spermatozoa, obtained by computing the cross-correlation function between both time series (Figure 6 - figure
supplement 2). **d.** Phase difference between the time derivative of the stimulus function and the internal Ca\(^{2+}\) oscillator of *S. purpuratus* spermatozoa expressed in temporal delays. **b-d.** Gray points represent the collapsed data of all *f*\(_1\), *f*\(_2\), *f*\(_3\), *f*\(_4\), *f*\(_5\) experimental regimes. Red, black and blue points indicate chemotactic spermatozoa (LECI > 1), located in R3, and R4 regions just before the speract gradient is established under *f*\(_2\), *f*\(_3\) and *f*\(_5\) experimental regimes, respectively. Magenta lines represent the transition boundary (*γ*\(_{\text{min}}\) = \(\xi^{\text{max}}\) \(\approx\) 2.7 x 10\(^{-4}\)) below which synchrony is not observed, obtained from the theoretical estimates (black curves) of panels **c** and **d**. Green lines indicate confidence intervals.
Figure 6 - figure supplement 1. Speract induces Ca\textsuperscript{2+} oscillations in immobilized S. purpuratus spermatozoa. Spermatozoa were immobilized on cover slips coated with poly-D-lysine (see Materials and Methods), and ASW containing 500 nM caged speract added. Recordings were performed 3s before and during 6s after 0.2s of UV irradiation. f4 optical fiber was used for the UV light path, to generate the speract concentration gradient. The optical fiber was mounted on a “defocused” configuration to minimize UV light heterogeneities producing photoactivated speract gradients. Time traces indicate the [Ca\textsuperscript{2+}], of selected spermatozoa of Movie 8. Note that the photo-release of speract induces a train of [Ca\textsuperscript{2+}] episodes in immobilized spermatozoa, hence provides evidences for the presence of an internal Ca\textsuperscript{2+} oscillator triggered by speract.
Figure 6 - figure supplement 2. Cross-correlation analysis of \([\text{Ca}^{2+}]_i\) and stimulus function derivative (ds) signals. Representative examples of \([\text{Ca}^{2+}]_i\) and the derivate of the stimulus function (ds) were plotted and then analyzed by cross-correlation analysis (CCF). Examples of a pair of spermatozoa for the two principal chemotactic gradients \((f_2\ and\ f_3)\) are shown. a, b. Representative examples of two spermatozoa in \(f_2\) gradient. c, d. Representative examples of two spermatozoa in \(f_3\) gradient.
Movie 1. Typical motility and Ca\textsuperscript{2+} responses of *S. purpuratus* spermatozoa towards *f*\textsubscript{2}-generated speract concentration gradient. Spermatozoa swimming in artificial sea water containing 10 nM caged speract, 3 s before and 5 s after 200 ms UV irradiation. An optical fiber of 0.6 mm internal diameter was used for the UV light path to generate the speract concentration gradient. Real time: 31 frames s\textsuperscript{-1}, 40x /1.3NA oil-immersion objective. Note that spermatozoa located at R2, R3 and R4 regions prior speract exposure swim up the speract concentration gradient, towards the center of the imaging field. The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca\textsuperscript{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca\textsuperscript{2+}]. Six *S. purpuratus* spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 2. Typical motility and Ca\textsuperscript{2+} responses of \textit{S. purpuratus} spermatozoa towards \textit{f3}-generated speract concentration gradient. An optical fiber of 2 mm internal diameter was used for the UV light path to generate the speract concentration gradient. Other imaging conditions were set up as for Movie 1. Note that spermatozoa located at R2, R3 and R4 regions prior to speract exposure swim up the speract concentration gradient, towards the center of the imaging field. The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca\textsuperscript{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca\textsuperscript{2+}]. Six \textit{S. purpuratus} spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 3. Chemotaxis of *S. purpuratus* spermatozoa requires extracellular calcium.

Spermatozoa swimming in artificial sea water with nominal calcium containing 10 nM caged speract 3 s before and 5 s after exposure to 200 ms UV light. Nominal calcium disrupts the electrochemical gradient required for Ca^{2+} influx, hence blocking the triggering of the internal Ca^{2+} oscillator by speract. The f2 fiber (0.6 mm diameter) was used to uncage speract in this control. Other imaging conditions were set up as for Movie 1. Note that spermatozoa re-located after speract uncaging but they fail to experience the Ca^{2+}-driven motility alteration triggered by speract. As a consequence they fail to experience chemotaxis (compare with Movie 1). The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca^{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca^{2+}]. Six *S. purpuratus* spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 4. Disrupting the K⁺ electrochemical gradient blocks chemotaxis of *S. purpuratus* spermatozoa. Cells were swimming in artificial sea water containing 40 mM of KCl, and 10 nM caged speract 3 s before and 5 s after exposure to 200 ms UV light. High K⁺ in the ASW blocks the hyperpolarization required for opening Ca²⁺ channels, and hence prevents the triggering of the internal Ca²⁺ oscillator by speract exposure. The f2 fiber (0.6 mm diameter) was used to uncage speract in this control. Other imaging conditions were set up as for Movie 1. Note that spermatozoa re-located after speract uncaging but they fail to experience the Ca²⁺-driven motility alteration triggered by speract, and thus they fail to experience chemotaxis (compare with Movie 1). The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca²⁺ indicator, showing maximum (red) and minimum (blue) relative [Ca²⁺]. Six *S. purpuratus* spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 5. Typical motility and Ca$^{2+}$ responses of *S. purpuratus* spermatozoa towards *f*1-generated speract concentration gradient. An optical fiber of 0.2 mm internal diameter was used for the UV light path to generate the speract concentration gradient. Other imaging conditions were set up as for Movie 1. Note that some spermatozoa re-located after speract uncaging but they fail to experience chemotaxis (compare with Movie 1). The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca$^{2+}$ indicator, showing maximum (red) and minimum (blue) relative [Ca$^{2+}$]. Six *S. purpuratus* spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 6. Typical motility and Ca\(^{2+}\) responses of \textit{S. purpuratus} spermatozoa towards \textit{f4}-generated speract concentration gradient. An optical fiber of 4 mm internal diameter was used for the UV light path to generate the speract concentration gradient. Other imaging conditions were set up as for Movie 1. Note that spermatozoa re-located after speract uncaging but they fail to experience chemotaxis (compare with Movie 1). The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca\(^{2+}\) indicator, showing maximum (red) and minimum (blue) relative \([\text{Ca}^{2+}]_i\). Six \textit{S. purpuratus} spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 7. Typical motility and Ca\textsuperscript{2+} responses of *S. purpuratus* spermatozoa towards *f*\textsubscript{5}-generated speract concentration gradient. An optical fiber of 4 mm internal diameter was used for the UV light path to generate the speract concentration gradient. Other imaging conditions were set up as for Movie 1. Note that spermatozoa located at R2, R3 and R4 regions prior to speract exposure swim up the speract concentration gradient, towards the center of the imaging field. The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca\textsuperscript{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca\textsuperscript{2+}]. Six *S. purpuratus* spermatozoa were manually tracked for visualization purposes. Scale bar of 50 µm.
Movie 8. Photo-release of caged speract induces Ca\textsuperscript{2+} oscillations in immobilized *S. purpuratus* spermatozoa. Spermatozoa were immobilized, by coating the cover slip with a thin layer of poly-D-lysine, in artificial sea water containing 500 nM caged speract, 3 s before and during 6 s after 0.2 s of UV irradiation. The f/4 optical fiber was used for the UV light path to generate the speract concentration gradient. The optical fiber was mounted on a “defocused” configuration to minimize the generation of UV light heterogeneities. 93 frames s\textsuperscript{-1}, 40x/1.3NA oil-immersion objective, 4x4 binning.

The pseudo-color scale represents the relative fluorescence of fluo-4, a Ca\textsuperscript{2+} indicator, showing maximum (red) and minimum (blue) relative [Ca\textsuperscript{2+}]\textsubscript{i}. The brightness and contrast scale was adjusted for better visualization of [Ca\textsuperscript{2+}]\textsubscript{i} transients in the sperm flagella (as a consequence some heads look artificially oversaturated, however no fluorescence saturation was observed in the raw data).
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