Mechanical Activation of Piezo1 but not $Na_V1.2$ Channels by Ultrasound

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ABSTRACT

Ultrasound can modulate the electrical activity of the brain and other excitable tissues but the mechanisms underlying this effect are not understood, either at the molecular level or in terms of the physical modality through which ultrasound exerts its effects. Here we report an experimental system that allows for stable patch-clamp recording in the presence of ultrasound at 43 MHz, a frequency known to stimulate neural activity. We describe the effects of ultrasound on two ion channels proposed to be involved in the response of excitable cells to ultrasound: the mechanosensitive Piezo1 channel and the voltage-gated sodium channel $Na_V1.2$. Our patch-clamp recordings, together with finite-element simulations of acoustic field parameters indicate that Piezo1 channels can be activated by ultrasound through cell membrane stress caused by acoustic streaming. $Na_V1.2$ channels were not affected through this mechanism, but their activation and inactivation rates could be accelerated by ultrasound-induced heating.

Key Words: Mechanotransduction, Piezo Channel, Sodium Channel, Electrophysiology, Membrane Biophysics, Ultrasound Neuromodulation, Ultrasound Neurostimulation, Acoustic Streaming, Membrane Stress

INTRODUCTION

It has long been known that ultrasound (US) can modulate electrical activity in excitable tissues (Fry, et al. 1958, Gavrilov, et al. 1996, Harvey 1929). In recent years, several research groups have investigated this effect with the motivation of developing US neuromodulation as a tool for treating mental and neurological disorders (Bystritsky, et al. 2011, Kim, et al. 2014, Kim, et al. 2015, King, et al. 2013, King, et al. 2014, Lee, et al. 2015, Lee, et al. 2016, Legon, et al. 2014, Mehic, et al. 2014, Menz, et al. 2013, Min, et al. 2011, Tufail, et al. 2010, Tyler, et al. 2008, Ye, et al. 2016, Yoo, et al. 2011, Younan, et al. 2013). A distinct advantage of US in this context is its ability to be focused deep within tissue with excellent spatial resolution and to function without surgical implants or genetic manipulation. These possibilities motivate the investigation of the mechanistic basis of US effects on excitable tissues.

Despite increased interest in recent years, models and hypotheses for the mechanism of US effects on excitability (Krasovitski, et al. 2011, Plaksin, et al. 2016, Plaksin, et al. 2013, Prieto, et al. 2013, Sassaroli and Vykhodtseva 2016, Tyler 2011) have been much more abundant than definite evidence in support of any particular mechanism. Complicating the picture is the essentially polymodal nature of US interaction with biological tissue (O'Brien 2007). In addition to changes in density and particle displacement related to the primary acoustic pressure, there may also be effects related to acoustic cavitation, temperature changes due to acoustic energy absorption, and second-order effects of radiation force and acoustic streaming. In different clinical and experimental contexts, and for different sets of US stimulus parameters, different modalities may be of primary importance. Clarity could be provided by a system for

studying the behavior of the basic "units" of biological electrical activity—individual ion channels—in the presence of US. The gold standard for this type of investigation is patch-clamp recording, which allows for detailed characterization of ion channel kinetics and membrane voltage dynamics. However, published accounts (Tyler, et al. 2008) and our own experience indicate that the gigaOhm seals required for patch-clamp recording are unstable in presence of US, and are irreversibly damaged at relatively low US intensities. This instability, however, appears to be frequency dependent. Here we report an experimental system that allows stable patch-clamp recording in the presence of US at 43 MHz, at intensities known to modulate neural activity in the salamander retina *in vitro* (Menz, et al. 2013).

Using this system, we investigated the effects of US on heterologously expressed mechanosensitive Piezo1 channels and on $Na_V1.2$ voltage-gated sodium channels. The choice of these two channels was guided by the hypothesis that US modulates action potential firing *in vivo* by causing membrane stress, thereby affecting the activity of mechanosensitive ion channels. Piezo channels are one of the few eukaryotic channels known to be directly activated (as opposed to modulated) by membrane stress (Coste, et al. 2010, Coste, et al. 2012), and since their discovery in 2010 have come to be regarded as the principal mechanoreceptor channel in mammalian cells (Volkers, et al. 2015, Wu, et al. 2016). In the context of US neuromodulation, it is notable that Piezo channels are expressed in the brain at the mRNA level (Coste, et al. 2010). Voltage-gated sodium channels (Na_V channels) are a strong candidate for the molecular target of US neuromodulation because of their central role in action potential generation, and because

they can be modulated by membrane stress (Beyder, et al. 2010, Morris and Juranka 2007, Wang, et al. 2009).

MATERIALS AND METHODS

Experimental set-up. The experimental set-up for simultaneous US stimulation and patch-clamp recording of cultured cells is illustrated in Figure 1A. The set-up is based on a modified Axioskop-2 microscope (Zeiss Microscopes) with a 40x W N-Achroplan objective. The condenser was removed and replaced with a custom-built housing for the 43-MHz transducer used in these experiments. US is transmitted from below the cells in the direction perpendicular to the bottom of the experimental chamber. Cells are illuminated with a ring of white LED lights mounted on the transducer housing, allowing visualization of the cells through the epifluoresence pathway with an image of sufficient quality for patch-clamp recording. The transducer is coupled to the bottom of the chamber by a drop of distilled water held by a small rubber O-ring attached to the tip of the transducer with silicone grease. The experimental chamber containing the cells was made from the lid of a 35-mm cell culture dish with a ~1 cm hole in the center. A thin film of plasma-treated polystyrene (Goodfellow USA, Coraopolis, PA) with cultured cells adhering to it was sealed onto the bottom of the chamber using silicone grease. Polystyrene films were used as the cell substrate instead of standard tissue culture dishes in order maximize US transmission and minimize heating due to the absorption of acoustic energy. The thickness of the film was 25 microns, except for the results shown in Figure 5D.

Electrophysiology. Voltage-clamp recordings were performed using an Axopatch 200B amplifier with a Digidata 1440 A digitizer and pClamp 10.4 software (Molecular Devices, Sunnyvale, CA). Currents were sampled at 100 kHz and filtered at 5 kHz for Piezo1 channels or 10 kHz for $Na_V1.2$ channels. Patch-clamp pipettes were pulled from thin-walled glass using a Sutter Instruments P-87 puller and had resistances between 2 and 7 MOhm when filled with the internal solution. The internal and external solutions were, respectively, in mM: 125 CsCl, 1 $MgCl_2$, 1 $CaCl_2$, 4 Na_2ATP, 0.4 Na_2GTP, 5 EGTA, pH 7.3 (CsOH) and 127 NaCl, 3 KCl, 1 $MgCl_2$, 10 HEPES, 10 glucose, pH 7.3 (NaOH) for the experiments on Piezo1 channels; and 140 KCl, 2 $MgCl_2$, 1 $CaCl_2$, 11 EGTA, 10 HEPES, pH 7.2 (NMDG) and 140 NaCl, 3 KCl, 2 $MgCl_2$, 1 $CaCl_2$, 10 HEPES, pH 7.2 (NMDG) for the experiments on $Na_V1.2$ channels. Series resistance compensation was not used due to the small size of the currents in these experiments. For experiments on $Na_V1.2$ channels, data were rejected if the magnitude of the maximum series resistance error ($\Delta V_{\max}$), calculated according to

(Eq. 1) $$\Delta V_{\max} = I_{\max} R_a$$

where $I_{\max}$ is the peak current and R_a is the access resistance, was more than 10 mV. For experiments on $Na_V1.2$ channels, pipettes were coated with dental wax to reduce pipette capacitance.

<u>Elevated cells.</u> In some experiments, cells were detached and elevated above the bottom of the chamber to examine the effect of height on the response to US. For these experiments, the height of the cell was determined by reading the display on the controller for the micromanipulator (Sutter Instruments MPC-325), or (in experiments using an Eppendorf

Patchman micromanipulator lacking this feature) by counting the number of rotations of the focus adjustment knob required to bring the cell into focus starting from the focal plane at the bottom of the chamber, and then converting this value into height using the calibration factor provided by the manufacturer. In contrast to cells at the bottom of the chamber, the gigaOhm seals of elevated cells were sometimes damaged by US stimulation, as evidenced by a sudden and irreversible drop of the apparent membrane resistance to <200 MΩ. This effect, which is most likely a result of acoustic streaming, became more frequent with increased height above the bottom of the chamber. The height at which the frequency of this effect made experiments very challenging was ~400 microns, so experiments were limited to the height range 0-400 microns.

Ultrasound. US stimuli were produced by a custom-made focused 43-MHz transducer with a quartz lens. The transducer was calibrated using the method described previously(Prieto, et al. 2013). The focal region of the US beam produced by this transducer is approximately cylindrical with a 90-micron diameter and 500-micron height. Outside of the focal region, the sound beam diverges rapidly, such that the sound reflected from the microscope objective above the cells does not interfere significantly with the sound in the focal region; we therefore do not expect standing waves to be significant in our experimental set-up. This idea was confirmed by finite-element simulations of the acoustic pressure field. The transducer was excited by 43-MHz sine-wave pulses from a Hewlett-Packard 8116A function generator, amplified by an ENI 403LA 37-dB amplifier (ENI, Inc., Rochester, NY). The function generator was controlled by an Agilent 33220A function generator (Agilent Technologies, Palo Alto, CA), which was triggered by 5-V pulses from the digitizer. The focal

region was localized along the *z*-axis to the surface of the polystyrene film using a pulse-echo method as described previously (Prieto, et al. 2013). The focal region was localized in the *x*-*y* plane by observing the acoustic streaming in response to US using Molecular Probes 1-micron fluorescent beads (ThermoFisher Scientific). A small volume (1-10 microliters) of bead suspension, diluted 1:100 in the external solution, was added to the chamber. In response to US, a circulatory acoustic streaming flow occurs, such that the beads appear to converge on the focal zone. Using the controls intended for centering the condenser, the focal zone was positioned at the center of the microscope's field of view. The solution in the chamber was then replaced with fresh solution and the cell chosen for recording was placed in the center of the focal region by adjusting the microscope stage.

Data analysis. Data were analyzed and curve fitting was performed using Igor Pro (Wavemetrics, Lake Oswego, OR). Ion channel activation and inactivation rates were determined from single exponential fits to the time course of the current rise and decay. All current records analyzed, with the exception of one of the experiments on Piezo1 channels, are the average of at least three trials.

Experimental design and statistical analysis. Two tailed t-tests were used to establish statistical significance. Significance was defined as $P < 0.05$. All mean results are reported and displayed as mean ± standard deviation (SD). For the data in Figure 2H-I, one recording with an anomalous current-voltage relationship was excluded based on a Dixon's Q-test (Barnett and Lewis 1994).

Finite-element simulations. Finite-element simulations of the acoustic pressure, acoustic streaming, and acoustic heating in the experimental chamber during our experiments were performed using COMSOL (COMSOL Inc., Palo Alto, CA). The simulation domain had radially symmetrical geometry, and was 6 mm in the axial direction by between 1 and 5 mm in the axial direction (depending on the property simulated), with a 940-micron diameter by 100-micron height arc on the lower axial boundary of the simulation domain representing the quartz lens of the transducer. The distance from the transducer lens to the bottom surface of the polystyrene was 4.2 mm, and the distance from the polystyrene to the upper boundary of the simulation was 2 mm, corresponding to the focal distance of the microscope objective. For each simulation, an appropriate mesh was determined by iteratively refining the mesh, starting from the default "Very Fine" mesh setting, until changes in the mesh had no effect on the result of the simulation. The default linear or non-linear solver was used for each simulation, except for the acoustic streaming simulation where a segregated Incomplete LU solver was used. Due to the radially symmetrical geometry, a symmetrical boundary condition was applied at the inner radial boundary in all simulations. All other boundary conditions are detailed in the descriptions of the individual simulations.

Material properties for water and polystyrene used in the simulations, along with sources for these values, are listed in Table 1. The attenuation coefficient for polystyrene was determined by measuring the echoes from the top and bottom surfaces of a 1.6 mm thick sample of the polystyrene used as the cell culture substrate and simulating the sound propagation and echo in COMSOL. Briefly, the amplitude of the reflected pulse from the bottom surface of the sample (the surface closest to the transducer) was used to determine

the reflection coefficient between water and polystyrene. The amplitude of the reflected pulse from the top surface was compared with that expected based on this reflection coefficient in the absence of attenuation to determine the attenuation coefficient. We measured the attenuation coefficient at 15, 20, 25, and 30 MHz and extrapolated to 43 MHz assuming linear attenuation with frequency.

<u>Simulation of the acoustic pressure field.</u> The steady acoustic pressure field was determined by solving the acoustic wave equation (Pierce 1994)

(Eq. 2)
$$\nabla \cdot \left(-\frac{1}{\rho}\left(\nabla p_{US}\right)\right) - \frac{k^2 p_{US}}{\rho} = 0$$

where ρ is the resting density, p_{US} is the complex acoustic pressure, and k is the complex wavenumber defined as

$$k = \frac{2\pi f}{c} - i\alpha$$

with frequency f, speed of sound c, and attenuation coefficient α. Equation 2 was solved with a 150-micron perfectly matched layer on the outer radial boundary and hard reflecting boundary conditions on all remaining boundaries. An oscillating pressure of 50 kPa or 67 kPa was applied at the boundary corresponding to the transducer lens. The pressure amplitude on the lens was chosen so that the simulated pressure in the focal region, in a simulation performed with the polystyrene film replaced by water, matched the pressures used in the experiments, as determined from the transducer calibration. The radial domain size was 1 mm.

Simulation of acoustic streaming. Acoustic streaming at steady-state was simulated by solving the Navier-Stokes Equations for an incompressible fluid with a volume force term due to acoustic radiation force (Nyborg 2008):

(Eq. 3a) $$\rho(\mathbf{u}\cdot\nabla)\mathbf{u} = \nabla\cdot(-p+\mu\nabla\mathbf{u}) + F_V$$

(Eq. 3b) $$\rho\nabla\cdot(\mathbf{u}) = 0$$

where $\mathbf{u}$ is the acoustic streaming velocity, p is the hydrostatic pressure, μ is viscosity, and F_V is the radiation force per unit volume:

(Eq. 4) $$F_V = \frac{2\alpha I}{c}$$

with the US intensity I defined as:

(Eq. 5) $$I = \frac{\mathrm{Re}\left[p_{US}^{*}\mathbf{v}_{US}\right]}{2}$$

where p^{*}_{US} is the complex conjugate of the complex acoustic pressure and $\boldsymbol{v}_{US}$ is the complex acoustic velocity, obtained by

(Eq. 6) $$\mathbf{v}_{US} = \frac{-\nabla p_{US}}{i\rho 2\pi f}$$

No slip (zero-velocity) boundary conditions were applied at the upper and lower axial boundaries, at the outer radial boundary and at the upper and lower surfaces of the polystyrene film. We also included the effect of radiation force on the polystyrene. The polystyrene buckles slightly under the influence of radiation force, which determines the position of the no-slip boundary. The polystyrene was modeled as an elastic material with geometric nonlinearity included due to the thinness and softness of the material. The displacement was determined by solving:

(Eq. 7) $$\nabla \cdot \sigma + F_V = 0$$

where σ is the stress tensor and F_V is the radiation force. The components of the stress tensor are a function of the strain (gradient of the displacement) and of Young's modulus (E) and Poisson's ratio (ν) (Landau and Lifshitz 1986).

The radial domain size for the simulated displacement was 5 mm and the displacement was fixed at zero at the outer radial boundary. The displacement of the polystyrene film was fully coupled to the movement of the fluid using the Fluid-Structure Interaction multiphysics module of COMSOL. Since boundary-layer effects can have very significant impact on fluid dynamics simulations, we performed simulations in which the distance to the radial boundary was changed between 1 and 5 mm. These simulations indicated that the presence of this arbitrary (non-physical) boundary does not affect the simulated streaming within the focal region. We also performed simulations in which the distance to upper axial boundary (a physically meaningful boundary corresponding to the surface of the microscope objective) was changed, to determine whether the position of this boundary (which was changed during experiments on elevated cells due to adjustment of the microscope focus) affects acoustic streaming. These simulations indicated that the position of the upper axial boundary does not affect the streaming in the focal region.

<u>Simulation of acoustic heating.</u> Acoustic heating was simulated by solving the heat transfer equation (Hand 1998):

(Eq. 8) $$\rho C_p \frac{\partial T}{\partial t} + \rho C_p \mathbf{u} \cdot \nabla T = \nabla \cdot \left(K_T \nabla T \right) + Q$$

where C_p is the specific heat capacity (at constant pressure), T is temperature, $\mathbf{u}$ is the streaming velocity, K_T is the thermal conductivity, and the heat source Q is determined by the US intensity and attenuation coefficient according to $Q = 2I\alpha$. A thermally insulating boundary condition was applied on all boundaries except the inner radial boundary. The radial domain size was 1 mm.

Note on heating measurements. We anticipated that absorption of acoustic energy by polystyrene could cause significant heating in our experiments, which we minimized by using thin (25-50 micron) polystyrene films as the cell-culture substrate rather than standard polystyrene cell-culture dishes. Despite these precautions, a small but significant temperature change occurs during our experiments, although the precise value of the temperature change is difficult to quantify because reliable thermal measurements in the presence of highly focused ultrasound fields are not technically feasible. The peak temperature rise in these experiments is expected to be localized to an area approximately equal in diameter to the focal volume (90 microns), which determines the spatial resolution that would be required for an accurate measurement. Solid-state temperature measurement devices such as thermocouples are available on this scale, but these will reflect and distort the sound field and will therefore not provide an accurate measurement. We attempted to measure temperature changes based on changes in the holding current for an open patch-clamp pipette tip (Shapiro, et al. 2012, Yao, et al. 2009) since this method has the appropriate spatial resolution, but in preliminary attempts we discovered that US would sometimes cause an apparent decrease in the temperature measured using this method, indicating that US can affect open-tip holding current through mechanisms other than heating. Instead, we estimated

the heating using simulations, combined with measurement of the US attenuation in polystyrene, as described above.

Video imaging. Supplementary Movie 1 was acquired using a Hamamatsu Orca Flash 4 microscope camera with Micromanager software (Edelstein, et al. 2010) at 100 frames per second with a 10-ms exposure time.

Cell culture and transfection. Rat $Na_V1.2$ channels were stably expressed in CHO cells. The CHO/$Na_V1.2$ cell line was provided by William Catterall (University of Washington). Mouse Piezo1 channels were transiently expressed in CHO or HEK cells by Lipofectamine transfection according to the manufacturer's protocol. Piezo1 in the pIRES2-eGFP vector was provided by Ardem Patapoutian (Scripps Research Institute, La Jolla, CA). Successfully transfected cells were identified by eGFP fluorescence and were recorded from 2-3 days after transfection. Cultured cells were maintained in an incubator at 37 C with 5% CO_2. The culture media was DMEM/F12 (1:1, high glucose, with sodium bicarbonate and L-glutamine) with 10% FBS and 100 μg/mL penicillin/streptomycin for CHO cells and DMEM (high glucose, with sodium pyruvate and L-glutamine) with 10% FBS and 100 μg/mL penicillin/streptomycin for HEK cells. For transiently transfected cells, antibiotics were omitted and FBS was reduced to 5%. For the stably transfected $Na_V1.2$ cell line, 200 μg/mL G418 was included to maintain expression. For recording, cells were trypsinized and dissociated and plated on ~2 cm square pieces of polystyrene film (25 or 50 microns thick) which had previously been treated with atmospheric plasma for 20 minutes in a Herrick plasma cleaner at the high RF setting to create a positively charged surface for cell

adhesion. Polystyrene films with cultured cells were placed in 6-well plates, covered with the standard media, and maintained in the incubator until used (1-3 days after plating). For experiments where the goal was to compare the effects of US on adhered versus detached cells, cells that had been plated at least 24 hours previously and appeared flat and elongated were considered fully adhered. Fully adhered cells could be detached from the polystyrene substrate without damaging the seal by a slow, very gentle back-and-forth movement of the patch-pipette. For experiments on detached cells where it was not essential that the cells be fully adhered at the start of the experiment, we used the following procedure to obtain loosely adhered that cells that could be easily detached without damaging the seal: first, cells were trypsinized and dissociated and a small volume (25-100 microliters) of cell suspension was added to the experimental chamber; the chamber was then placed in the incubator, inside a larger dish with a damp KimWipe to prevent dehydration, for 1-4 hours before recording.

RESULTS

We measured the effect of US on heterologously expressed ion channels using a focused 43-MHz transducer that has previously been shown to potentiate action potential firing in the salamander retina (Menz, et al. 2013), using the experimental set-up illustrated in Figure 1A. US is transmitted from below the experimental chamber through a thin (25 or 50 micron) polystyrene film. We measured the effects of US under three different whole-cell recording conditions (Figure 1B): first, with the cells expressing the channels adhered to the polystyrene film at the bottom of the recording chamber (*top*); second, with the cells

detached from the polystrene substrate and elevated above its surface by between 60 and 390 microns (*bottom left*); and finally with the cells detached from the substrate but not elevated above its surface ("ground level," *bottom right*). We chose to examine these different conditions because US can affect biological systems through several different modalities and, as examined in more detail below, the strength of these modalities, absolutely and relative to one another, can vary widely depending on the location within the experimental chamber. Specifically, in our experimental set-up the spatial distributions of acoustic heating and acoustic streaming have opposite dependences on height above the surface of the polystyrene film: heating is much more prominent for cells adhered to the film or very close to it than for cells elevated above it, while acoustic streaming is absent at the surface of the film but becomes increases with height. The strength of a third potential modality, acoustic radiation force, decreases with height but varies much less dramatically than heating or streaming over the range examined in these experiments (0-390 microns). Thus, examining the effects of US as a function of height above the bottom of the chamber can help distinguish which modalities are responsible for US effects on ion channels.

As seen in Figure 2, US can activate Piezo1 channels when the cell expressing the channels is elevated above the bottom of the chamber, but not when the cell is at "ground level," and the Piezo1 current in response to US increases with height. Figure 2A shows an example of the current at -80 mV in response to US at 50 W/cm^2 in a Piezo1-transfected HEK cell elevated above the bottom of the chamber by 330 microns (*black current trace*) and at "ground level" (0 microns, *gray current trace*), where there is no apparent response to ultrasound. This phenomenon was reproducible, with a detectable current response to US at 50 or 90 W/cm^2, at a height of 60 microns or higher, in 26 out of 32 cells tested, but

never at ground level, whether or not the cell was the adhered to the bottom of the chamber. The effect of height on the US-activated current was reversible (Figure 2B-C), indicating that the sensitivity to US in elevated cells is not due to irreversible damage to the cell rendering Piezo channels more responsive to US (as might occur, for example, if US were to disrupt the cytoskeleton, reducing resting membrane tension and thereby making a previously inactivated population of Piezo channels available for activation by US). The dependence of the peak US-activated current on height is summarized in Figure 2D for twelve different cells for which recordings at multiple heights were obtained. The effect of height on the current amplitude is summarized for all cells in Figure 2E, which shows the distribution of current amplitudes for cells below (*triangles*) or above (*circles*) 50 microns. The US-activated current inactivates with a time constant on the order of that previously reported for heterologously expressed Piezo1 channels (~13 ms (Coste, et al. 2010, Coste, et al. 2012)), although there was significant variability in time constant among cells tested (Figure 2F). Inactivation rates and peak current values are weakly voltage-dependent, as is characteristic of Piezo channels (Figure 2G-I).

We considered whether the US-activated current might be due to endogenous mechanosensitive channels. HEK cells express an endogenous mechanosensitive current that can be eliminated by knocking out Piezo1, but this current exhibits very slow inactivation (time constant >>100 ms) (Dubin, et al. 2017). To confirm that the currents in response to US are due to heterologously expressed Piezo1 channels, we compared currents in Piezo1-transfected cells with those in control cells of the same batch transfected with eGFP in the same expression vector. In nine out nine cells transfected with eGFP alone, there was no apparent US-activated current at 50 or 90 W/cm^2 at a height of

220 or 390 microns (Figure 2E, *red stars*) compared with seven out of seven cells concurrently transfected with Piezo showing a significant (30 – 500 pA) response to US under the same conditions. Based on these results, and on the relatively fast inactivation time constants of the currents (Figure 2F) compared to those reported for endogenous channels, we conclude that the US-activated current is primarily due to heterologously expressed Piezo1 channels, although we cannot entirely rule out some small contribution from endogenous channels.

A relatively small change in position has a profound effect on the response of Piezo1 channels to US. How can a change in position as small as 100 microns render previously unresponsive channels responsive to US? This distance is equivalent to only about three wavelengths at 43 MHz, and is very small relative to the US attenuation coefficient in water (46 neper/m) (www.ondacorp.com/images/Liquids.pdf). To help explain the change in the sensitivity of Piezo1 channels to US in elevated cells, we performed finite-element simulations of three different modalities by which US may affect ion channel activity: acoustic radiation force, acoustic streaming, and acoustic heating (Figures 3 and 4). Radiation force and acoustic streaming are hypothesized to activate Piezo channels by causing cell membrane stress, while heating typically accelerates ion channel kinetics and can therefore activate ion channels if the temperature rise is sufficiently large or the channel kinetics are especially sensitive to heat. The effects of US on all relevant aspects of the experimental set-up, including the solution in the experimental chamber, the polystyrene film at the bottom of the chamber, and the water coupling the chamber to the transducer, were simulated. Although there is some uncertainty in the absolute values of the simulated field parameters due to estimates of the material properties input to the

simulation (see references in Table 1 for the range of material properties for various forms of polystyrene and related polymers), our conclusions are based primarily on the spatial distributions of the field parameters, which are not affected by small changes in material properties, rather than their absolute values.

The simulated spatial profile of radiation force in response to ultrasound at 90 W/cm^2 is shown in Figure 3A for the liquid media (*top panel*) and for the polystyrene film (*middle panel*). Standing waves could create variations in US intensity with spatial period equal to the wavelength, possibly explaining the variation in response to US with height, but the simulation indicates that there is no significant effect of standing waves on the US intensity distribution within the chamber, as there is no apparent periodic variation in radiation force, which is proportional to the US intensity. Based on the simulation, cells at ground level experience a radiation force per unit volume of ~6 N/m^3, but this is apparently not effective, in and of itself, in activating Piezo1 channels. More importantly, the spatial profile of radiation force is incompatible with direct activation of Piezo1 channels by radiation force, since radiation force decreases with height above the bottom of the chamber (Figure 4A, *thin solid line*), while the response of Piezo1 channels to US increases (Figure 2D). We also simulated the displacement of the polystyrene film in response to radiation force (Figure 3A, *bottom panel*), because a large displacement of the film could affect the spatial distribution of the acoustic field parameters, and a gradient in the displacement could cause stress in cells adhered to the film. The peak displacement is very small (0.3 microns) indicating that displacement of the polystyrene film is not a significant factor in our experiments.

In contrast, the simulated spatial profile of the axial component of the acoustic streaming velocity correlates well with the dependence of the US-activated Piezo1 current on height (Figure 3B). As typically seen with acoustic streaming (Duck 1998), high-velocity streaming in the direction of propagation occurs within the US beam, along with a slower, circulatory return flow outside the beam, as indicated by the regions with negative velocity. (The circulatory flow also has a radial velocity component, much smaller than the axial component (9.4 mm peak radial velocity amplitude), not shown in Figure 3B). Critically, there is a thin stagnant layer, directly above the surface of the polystyrene film, in which no streaming motion occurs (Figure 3B, *bottom panel*). This stagnant layer is the result of the zero-velocity boundary condition at the rigid surface of the polystyrene film. Outside the stagnant layer, the axial streaming velocity increases steeply with height and then levels off over the range probed in the Piezo1 experiments (Figure 4A, *dashed line*). This result explains why Piezo1 channels in cells at the bottom of the chamber are immune to US stimulation, despite the significant radiation force, and why the Piezo1 current in response to US increases with height. We conclude that forces acting on the cell membrane as a result of acoustic streaming are sufficient to activate Piezo1 channels. These forces are caused by relative motion between the cell, which is held in place by the patch-pipette, and the surrounding fluid. A cell suspended in a moving fluid will experience a drag force made up of normal and shear components (King 2002), resulting in membrane stress. Consistent with this idea, we observed that the membranes of elevated cells are distorted and apparently displaced in the axial direction during US stimulation (Supplementary Movie 1). Thus, radiation force is critical for the response of Piezo1 channels to US, insofar as it

drives acoustic streaming, but the boundary conditions that regulate acoustic streaming are equally critical.

Although the spatial profile of acoustic streaming velocity appears sufficient to explain our results on Piezo1 channels, we considered whether hydrodynamic pressure caused by streaming might also contribute to the effects of US on Piezo1 channels. The simulated spatial distribution of the hydrodynamic pressure (Figure 3C) indicates that there is a significant negative hydrodynamic pressure within the chamber at the surface of the polystyrene film. However, the amplitude of this hydrodynamic pressure decreases with height (Figure 3A, *dotted line*), which is again inconsistent with the dependence of US-activated Piezo1 current on height. We therefore conclude that, relative to the streaming velocity, hydrodynamic pressure is not an important factor in the activation of Piezo1 channels by US.

Finally, the simulated spatial profile of acoustic heating indicates that temperature changes due to US absorption cannot explain the effects of US on Piezo1 channels. The simulated spatial profile of the temperature rise at the end of a 1-s US stimulus at 90 W/cm^2 (during which the temperature change reaches an approximate steady state) shows that US-induced temperature rise is greatest at center of polystyrene film (Figure 3D), as expected based on the higher ultrasound attenuation coefficient and lower specific heat of polystyrene relative to water. In contrast to the effect of height on the response of Piezo1 channels to US, the temperature rise steeply declines with height (Figure 4A, *thick solid line*). We therefore conclude that heating does not play a significant role in the response of Piezo1 channels to US. This conclusion is consistent with the observation that mechanically

activated Piezo current have unexceptional temperature sensitivity, as reflected in their Q10 value of ~2.8 for inactivation (A. Patapoutian and B. Coste, personal communication).

To further support our conclusion that the activation of Piezo1 channels by US correlates with acoustic streaming, we compared three sets of recordings in which we measured US-activated currents from the same cell across a wide range of heights (Figure 4B). The peak Piezo1 current in response to US increases from its minimum value at zero height to a maximum value at 390 microns, the highest level tested, where it appears to approach a plateau. At this level and above, patch-clamp seals were frequently irreversibly damaged by US, presumably due to acoustic streaming, so we were unable to confirm the presence of a plateau in the Piezo current response. Nonetheless, the height dependence of the US response in Piezo-transfected cells over the range tested clearly resembles the axial profile of the acoustic streaming velocity, which rises from zero at the polystyrene surface to a plateau value at above 300 microns, rather than that of the radiation force, which remains relatively constant between zero and 400 microns, or the hydrodynamic pressure, which declines as the velocity increases, or the temperature change, which rapidly declines from its peak value at the zero microns to a small fraction of the peak at 200 microns (Figure 4A). This result confirms that acoustic streaming is critical for activating Piezo channels with US in our experiments.

We also considered whether, and through what modality, US could affect voltage-dependent $Na_V1.2$ sodium channels. Since US can potentiate action potential firing and Na_V channels are central to action potential firing, and since the kinetics of least some subtypes of Na_V channels can be modulated by membrane stretch (Beyder, et al. 2010, Morris and Juranka 2007, Wang, et al. 2009), it seemed possible that these channels might also

respond to US. We therefore measured the effects of US on heterologously expressed $Na_V1.2$ channels in cells at ground level and in elevated cells. A representative example of the effect of US at 43 MHz and 90 W/cm^2 on $Na_V1.2$ channels in an adhered cell is shown in Figure 5A. The current in response to a step from a holding potential of -100 mV to -10 mV is shown with (*black current trace*) and without (*gray current trace*) US stimulation. There was no significant effect of US on the amplitude of the $Na_V1.2$ current in adhered cells (-379 ± 159, control, versus -382 ± 150 pA, US; mean ± SD, N =6)P = 0.87, paired t-test)). However, the macroscopic rates of activation and inactivation both increase in adhered cells in the US condition, from 2420 ± 530 s^{-1} to 3280 ± 820 s^{-1} (activation) and from 600 ± 140 s^{-1} to 740 ± 150 (inactivation) (mean ± SD, N = 6). The effect of US was significant for both activation (paired t-test, P = 0.006) and inactivation (paired t-test, P = 2×10^{-5}). These accelerated kinetics are consistent with either a thermal or a radiation force mechanism. Based on our finite-element simulations, the US-induced temperature rise is 0.8 C at the surface of the polystyrene film at the start of voltage step. For a Q10 of 3, which is typical of ion channel kinetics and similar to that reported for Na_V currents (Frankenhaeuser and Moore 1963, Hodgkin, et al. 1952), this would produce a 1.09-fold acceleration of channel gating kinetics, slightly lower than the experimental result (1.23 ± 0.06). However, the exact value of the temperature rise in the simulation depends on the specific heat and thermal conductivity of the polystyrene film, for which we lack exact values. The effect of US on $Na_V1.2$ kinetics was not dependent on the size of the voltage step used to activate the channels (data not shown), indicating that the effect of US is orthogonal to the effect of voltage, consistent with a temperature effect. However, membrane stretch produced by applying pressure to cell-attached membrane patches also increases the macroscopic rates

of activation and inactivation in Na_V channels (although the stretch sensitivity of the specific subtype studied here has not to our knowledge been tested). Thus, it is difficult to tell based on these results alone whether the effects of US on Na_V channels are due to acoustic heating, or membrane stress caused by radiation force or acoustic streaming, or some combination of these.

To determine the physical basis of the effect of US on Na_V channel kinetics, we measured the dependence of this effect on height above the bottom of the experimental chamber and, applying the rationale already established for Piezo channels, compared this dependence with the simulated axial profiles of radiation force, acoustic streaming velocity, hydrodynamic pressure, and steady-state temperature change (Figure 4A). As for adhered cells, there was no significant effect of US on the current amplitude in either detached cells at ground level, or in detached and elevated cells (Figure 5B). In contrast to the effect of US on Piezo1 channels, however, the effect of US on the kinetics of $Na_V1.2$ channels declines with height (Figure 5C), and the steep decline matches the axial profile of the US-induced temperature change (Figure 4A, *thick solid line*). Thus, unlike for Piezo1 channels, elevating the cell expressing the channels does not reveal a previously suppressed mechanical response to US in $Na_V1.2$ channels. This result indicates that the response of $Na_V1.2$ channels to US is exclusively or predominantly due to heating. As a final test of this conclusion, we compared the time course of the effect of US on the $Na_V1.2$ inactivation rate with the simulated time course of the temperature change at the surface of the polystyrene film. Figure 5D shows the ratio of the $Na_V1.2$ inactivation rate during the US stimulus to the inactivation rate in the control condition, as a function of the time between the start of the US stimulus and the start of the channel-activating voltage step (Δt). The simulated time

course of the US-induced temperature change is plotted on the right axis for comparison. As expected, the time courses are very similar. We therefore conclude that the acceleration of channel gating by US is predominantly due to heating, with little, if any, contribution from membrane stress, either due to radiation force directly or to acoustic streaming. Since this heating is dependent on US absorption by the polystyrene film, the response of $Na_V1.2$ channels to US in these experiments is of limited relevance in the context of US effects on brain and excitable tissues *in vivo*. The absence of a mechanical effect, despite the hypothesized mechanical basis of US neuromodulation and the sensitivity of Na_V channels to membrane stretch in other contexts, is the most important result of our experiments on Na_V channels.

DISCUSSION

Comparison with previous US/electrophysiology experiments. Here we report stable patch-clamp recordings in the presence of ultrasound. Previously, Tyler and coworkers reported that patch-clamp seals were unstable in the presence of US, and therefore concluded that patch-clamp recording was not a practical technique for investigating the bio-effects of US (Tyler, et al. 2008). The 43-MHz frequency and 90 W/cm^2 intensity used in our experiments are both much higher than those used by Tyler and co-workers (440 kHz and 2.9 W/cm^2), suggesting that relatively high-frequency US may be critical for compatibility with patch-clamp recording. In addition, the US beam in our experiments is very tightly focused, which may contribute to our success at combining US stimulation with patch-clamp recording. Finally, another potentially important factor is the absence of significant standing wave effects, which were minimized by the tightly focused US field.

Kubanek and coworkers reported on effects of pulsed US at 10 MHz on ion channel currents recorded using two-electrode voltage-clamp (TEVC) in *Xenopus* oocytes (Kubanek, et al. 2016). They observed a small increase in $Na_V1.5$ conductance in response to US but were unable to resolve the peak currents or kinetics due to the limited clamp speed of TEVC. They did not differentiate between possible sources of US effects; however, they did measure a significant temperature rise that could explain the effects they observed and that is consistent with our findings on $Na_V1.2$ channels.

The radiation force hypothesis of US effects on ion channels. The goal of these experiments was to better understand the mechanisms by which US affects the activity of excitable tissues, and to test the hypothesis that US affects ion channel activity by a mechanical effect on cell membranes. While this idea has been proposed by several groups, a correlation between the amount of acoustic radiation force, the amount of membrane stress, or other relevant parameter, and the strength of the response to US has not been demonstrated previously. In fact, US neuromodulation of mouse motor cortex *in vivo* appears to be more effective at lower frequencies (King, et al. 2013, Tufail, et al. 2010, Ye, et al. 2016), the opposite of the expectation for a radiation force mechanism. However, it is not yet clear to what extent confounds such as differences in beam profile and absorption by the skull distort the frequency response in *in vivo* experiments. US neuromodulation of the salamander retina *in vitro* is more effective at higher frequency, as expected for a radiation force-based mechanism (Menz, et al. 2016). In heart tissue, Dalecki et al. demonstrated that aortic pressure in the frog heart is modulated by US, and by varying both frequency and beam width, determined that the effect of US was proportional to radiation force, providing

compelling evidence for a radiation force mechanism for US modulation of electrical activity (Dalecki, et al. 1997).

Our results on Piezo1 channels are consistent with the general idea that radiation force is involved in the effect of US on electrical activity, insofar as acoustic streaming represents the response of liquid media to radiation force, and demonstrate that US can activate mechanosensitive ion channels. However, critical issues must be resolved to apply these results to *in vivo* neuromodulation experiments. First, the 43-MHz frequency used in our experiments is much higher than the typical frequency range for *in vivo* experiments (~0.2-3 MHz). As already mentioned, there is an apparent discrepancy between the frequency dependence of *in vivo* and *in vitro* US neuromodulation experiments. We cannot rule out the possibility that effects based on cavitation that are not present in our experiments are important at lower frequencies. In addition, methods to evaluate acoustic streaming effects in brain tissue will need to be developed. Relatively simple viscoelastic models of US-tissue interactions do not account for acoustic streaming. US-induced tissue displacement, produced through the same radiation force mechanism that drives acoustic streaming, is the closest analogue to acoustic streaming in these models (Sarvazyan, et al. 2010). Dual-phase models of tissue as a fluid-saturated porous solid will be more appropriate for modeling effects of acoustic streaming.

Comparison with in vitro neurostimulation experiments at 43 MHz. There is a very interesting parallel between our experiments and those of Menz and Baccus, who used the same focused 43-MHz transducer to stimulate neural activity in the salamander retina *in vitro*, while recording the activity of retinal ganglion cells using a multi-electrode array

(Menz, et al. 2013). They found that retinal ganglion cells do not respond directly to US, but show increased activity due to US stimulation of cells in more superficial layers. A key feature of their experiments in the context of the present results is that the ganglion cell layer is directly above the rigid surface of the multi-electrode array, which therefore enforces a no-slip boundary condition at its surface, just as the polystyrene substrate does in our experiments. Thus ganglion cells may be unresponsive to US in these experiments not for reasons related to their biophysical properties, but instead because they are located in a stagnant layer adjacent to the surface of the array. Consistent with this idea, they found that the membranes of ganglion cells undergo much less displacement than those of neurons in more superficial layers in response to US stimulation (Menz, et al. 2013). Together with the experiments reported here, these results motivate further investigation of the neuromodulatory effects of higher frequency US and of the role of acoustic streaming in these effects.

Absence of a significant mechanical effect on $Na_V1.2$ channels. Although we cannot entirely rule out the idea that some small part of the response of $Na_V1.2$ channels to US is due to mechanical effects, most of the response is clearly due to heating. Why then don't $Na_V1.2$ channels respond to US when subjected to an acoustic streaming velocity sufficient to activate Piezo channels (Figure 5C), despite their reported mechanosensitivity? A simple explanation would be that more membrane tension is required to affect $Na_V1.2$ kinetics than is required to activate Piezo1 channels. This explanation is difficult to rule out. The tension required for half-maximal activation of Piezo1 channels is 2.7 mN/m in cell-attached patches (Lewis and Grandl 2015), but the tension required to modulate Na_V

kinetics is not known. Instead, quantitative comparison of the effects of tension on Na_V and Piezo channels must be based on the pressures applied to membrane patches, which is problematic since membrane tension depends on both the applied pressure and the patch's radius of curvature (Sokabe, et al. 1991), causing considerable variation in the tension produced by a given pressure step (Moe and Blount 2005). Comparison of pressure values without information on patch curvature is therefore extremely approximate, and the observation that pressures required to activate Piezo channels and modulate Na_V channels in on-cell patches are similar (on the order of -10 to -40 mmHg) (Bae, et al. 2013, Beyder, et al. 2010, Coste, et al. 2012, Lewis and Grandl 2015, Morris and Juranka 2007)) is not conclusive.

Another potential explanation for the lack of a mechanically-mediated response to US in $Na_V1.2$ channels is that Na_V channel mechanosensitivity was demonstrated in cell-attached membrane patches, and the mechanical environment of the ion channels in the membrane patch may be different from that of the remaining channels in the cell membrane (Suchyna, et al. 2009, Ursell, et al. 2011). A third alternative is that Na_V channels and Piezo1 channels may respond to different forms of membrane stress through different mechanisms, and the fluid stress from acoustic streaming may act through one of these mechanical pathways but not the other. Broadly speaking, mechanosensitive ion channels are thought to respond to membrane stress through lipid bilayer tension (*e.g.*, reconstituted Piezo channels (Syeda, et al. 2016), bacterial Msc channels (Sukharev, et al. 1994), and mammalian K2P channels (Brohawn, et al. 2014)) or interactions with the membrane-associated cytoskeleton (*e.g.*, NOMPC channels (Zhang, et al. 2015)). Which of these mechanisms is most important for Na_V channels is not known.

CONCLUSIONS

The experiments reported here provide insight into the molecular and biophysical basis of US-induced bio-effects in the nervous system and other tissues. Our results are consistent with the idea that mechanical forces associated with US can activate mechanosensitive ion channels, but a subtle and key point is that radiation force alone is not sufficient. Instead, we conclude that acoustic streaming produced by radiation force is required to activate Piezo1 channels in our experiments, and this leads us to propose that acoustic streaming or radiation force-induced tissue displacement may be required for the neuromodulatory effects of US seen in various *in vivo* and *in vitro* preparations. For both streaming and tissue displacement, boundary layer effects will suppress the response near the boundary. Thus cultured cells adhering to a rigid surface will not provide an appropriate model system for understanding the effects of US on ion channels. Fortunately, the approach described here can be readily applied to other ion channels, providing a platform for investigating the endogenous response of ion channels to US or developing genetically engineered US-sensitive actuators for "sonogenetics."

SUPPLEMENTARY DATA

Supplementary Movie 1. Cell membrane distortion in response to ultrasound.

A CHO cell at the end of a patch pipette at a height of 170 microns above the bottom of the chamber is exposed to US at 50 W/cm^2 for 200 ms. The direction of US propagation is normal to the field of view. The image was acquired at 100 fps and playback is at 7 fps. The length of the video is 500 ms.

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FIGURE CAPTIONS

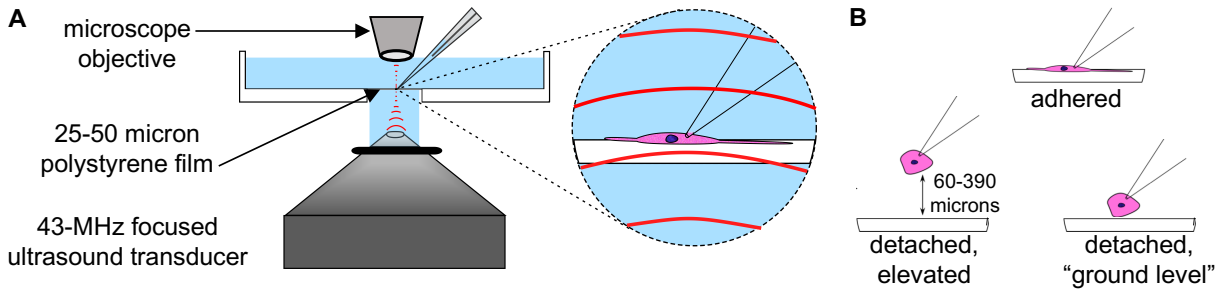

Figure 1. Experimental Set-Up. A. Diagram of experimental set-up. Focused US at 43 MHz is propagated from below to cells adhering to a thin (25- or 50-micron) polystyrene film, oriented perpendicular to the direction of US propagation. See Experimental Procedures for details. **B.** Cartoon illustrating the three different whole-cell recording configurations used in our experiments. Cells were either adhered to the polystyrene film (*top*), or detached and elevated above it (*bottom left*), or detached but not elevated ("ground level," *bottom right)*. The cartoon is not to scale.

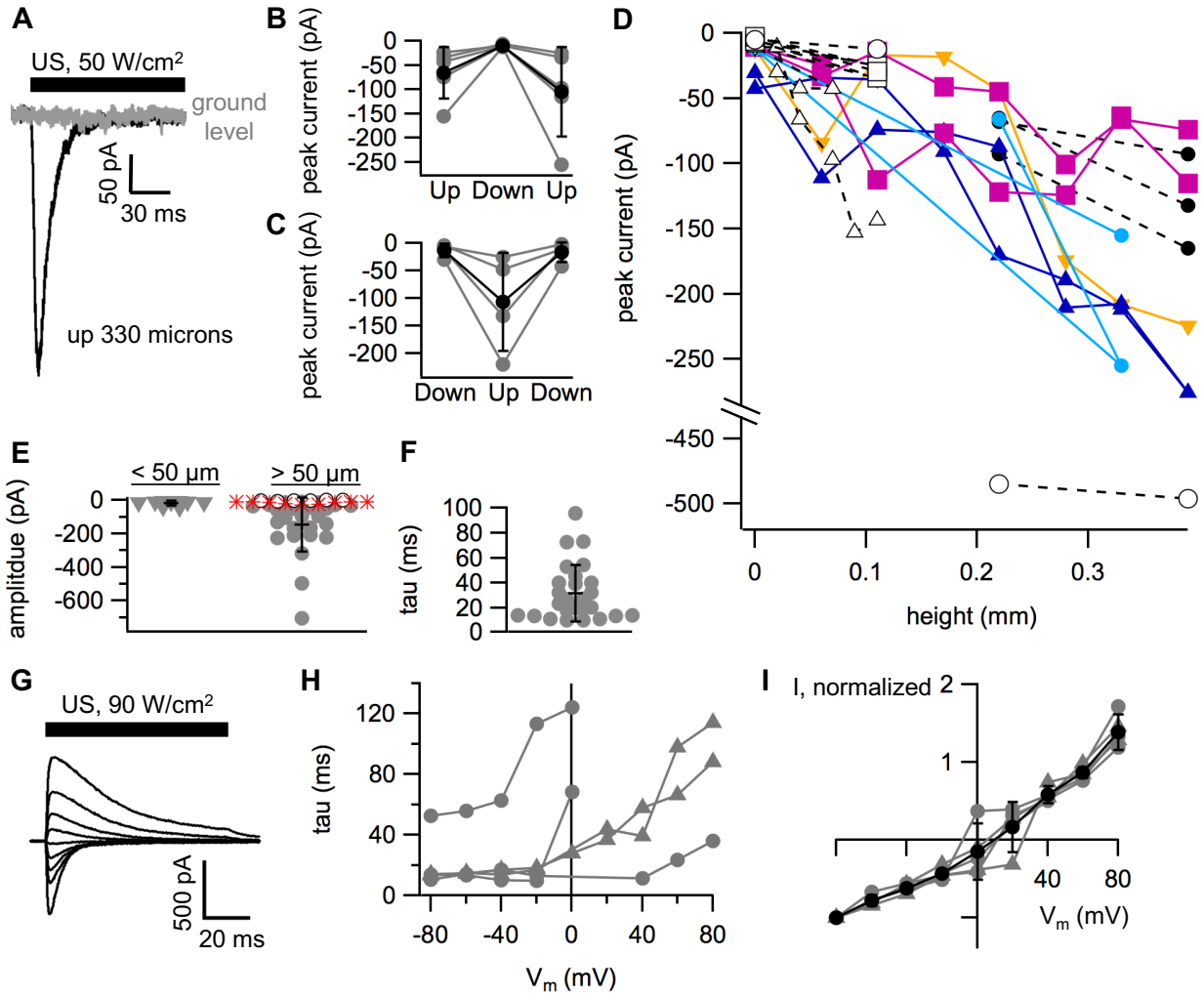

Figure 2. Activation of Piezo1 channels by ultrasound. A. Example currents at -80 mV in response to a 200-ms, 50 W/cm² US application in a HEK cell expressing Piezo1 channels, recorded while the cell is at the bottom of the recording chamber (*gray current trace*, "ground level") and after detaching the cell from film and elevating it by 330 microns (*black current trace*). **B-C.** Reversibility of the effect of height on US-activated currents. Cells were sequentially exposed to US at 50 or 90 W/cm² at an elevated height (70-390 microns, "Up"), then at the bottom of the chamber ("Down"), then back to the original height (*B*); or at the bottom of the chamber, then at an elevated height, then back to the bottom of the

chamber (*C*). Individual cells are shown as gray circles and mean ± SD values as black circles. N=5 in *B* and N=4 in *C*. **D.** Peak US-activated current as a function of height for N=12 cells. Each group of connected symbols represents one cell tested at various heights. US was either at 50 W/cm^2 (*small symbols*) or 90 W/cm^2 (*large symbols*). The different symbol shapes, colors, and line styles are intended to help distinguish individual cells and have no additional symbolic meaning. **E.** Distribution of current amplitudes at -80 mV for cells examined at heights > 50 microns (*circles*, N = 31) or < 50 microns (*triangles*, N = 9, same cells as at > 50 microns; data at lower heights not available in all cells). Open circles indicate Piezo1-transfected cells with no detectable Piezo current (N = 6), and red stars indicate control cells transfected with eGFP only, all of which showed no detectable US-activated current (N = 9). For this analysis, data at the highest height available in each of the two categories was used for cells tested at multiple heights. One cell was excluded from this analysis because the height was not quantified. The mean and standard deviation for cells with detectable Piezo1 current are indicated (-16 ± 14 pA, < 50 microns (N = 9); -145 ± 163 pA, > 50 microns (N = 25)) The difference in current amplitude between groups was statistically significant (for all cells: unpaired t-test, unequal variance, P = 0.0006, N = 25; for cells in both groups: paired t-test, P = 0.004, N = 9). **F.** Distribution of inactivation time constants (*tau*) for US-activated currents at -80 mV for N = 26 cells. The mean and standard deviation are indicated (33.1 ± 23.5 ms) **G.** Example currents at a series of voltages from -80 to +80 mV (in 20 mV steps) in a CHO cell expressing Piezo1 channels in response to a 200-ms, 90=W/cm^2 US application. The cell was 170 microns above the bottom of the chamber. **H.** Inactivation time constant for Piezo1 currents expressed in CHO cells (*gray triangles*) or HEK cells (*gray circles*) in response to US at 90 W/cm^2 as a function of

membrane voltage (V_m) for N=6 different cells. **I.** Peak US-activated current as a function of membrane voltage (V_m) for the cells in *H*, normalized to the current at -80 mV, along with the mean ± SD values (*black circles*).

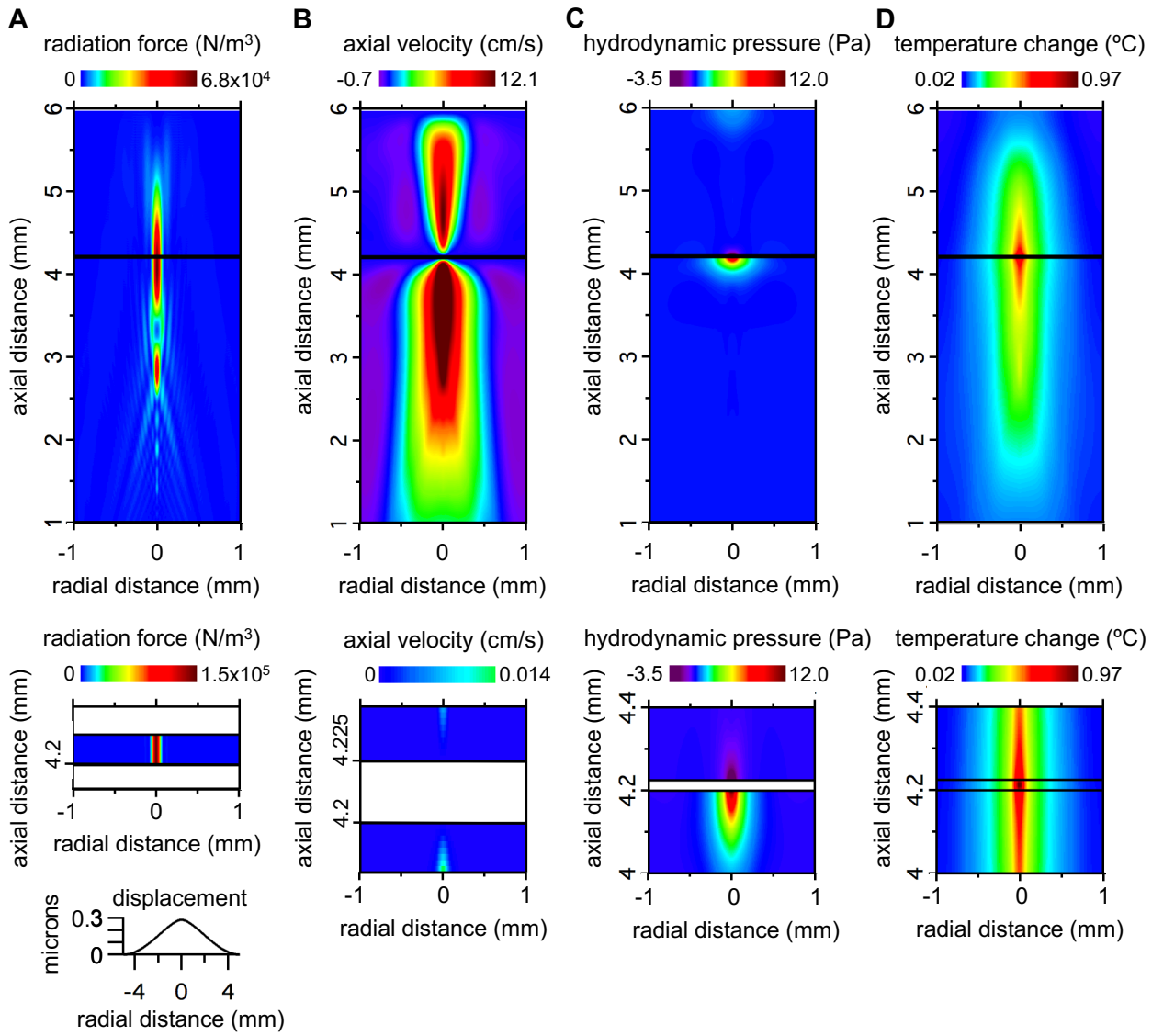

Figure 3. Simulated spatial profiles of acoustic field parameters. A. Acoustic radiation force in the liquid media (*top panel*) and in the polystyrene film (*middle panel*). The bottom panel shows the simulated displacement of the surface of the polystyrene film in response to acoustic radiation force. **B.** Axial component of the acoustic streaming velocity. The bottom panel shows the streaming velocity on an expanded scale, highlighting relevant features in the vicinity of the polystyrene film. **C.** Hydrodynamic pressure due to acoustic

streaming (shown on an expanded scale in the bottom panel). **D.** Steady-state temperature change (shown on an expanded scale in the bottom panel).

Axial distances are relative to the surface of the transducer and radial distances are relative to the center of the US beam. Results are shown for US at 90 W/cm^2 and 25-micron polystyrene film. Similar results were obtained for US at 50 W/cm^2 and 50-micron polystyrene (not shown). Changing the US intensity or the thickness of the polystyrene film changed the amplitude of the acoustic field parameters but did not substantially change their spatial profile.

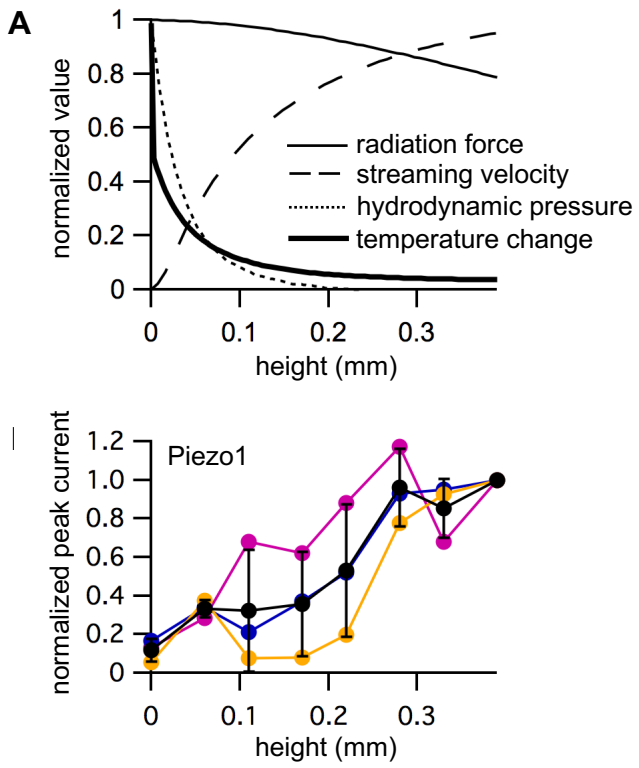

Figure 4. Activation of Piezo1 channels by ultrasound depends on acoustic streaming velocity. A. Normalized axial values of the simulated acoustic radiation force (*thin solid line*), acoustic streaming velocity in the axial direction (*dashed line*), hydrodynamic pressure amplitude (*dotted line*), and steady-state temperature change (*thick solid line*) as a function of height above the bottom of the experimental chamber. **B.** Normalized peak Piezo1 current at -80 mV in detached cells in response to US as a function of height above the bottom of the experimental chamber, normalized to the peak current at a height of 390 microns for N = 3 cells (*colored circles, corresponding to the colors in Figure 2D*), along with the mean (± SD) values (*black circles*).

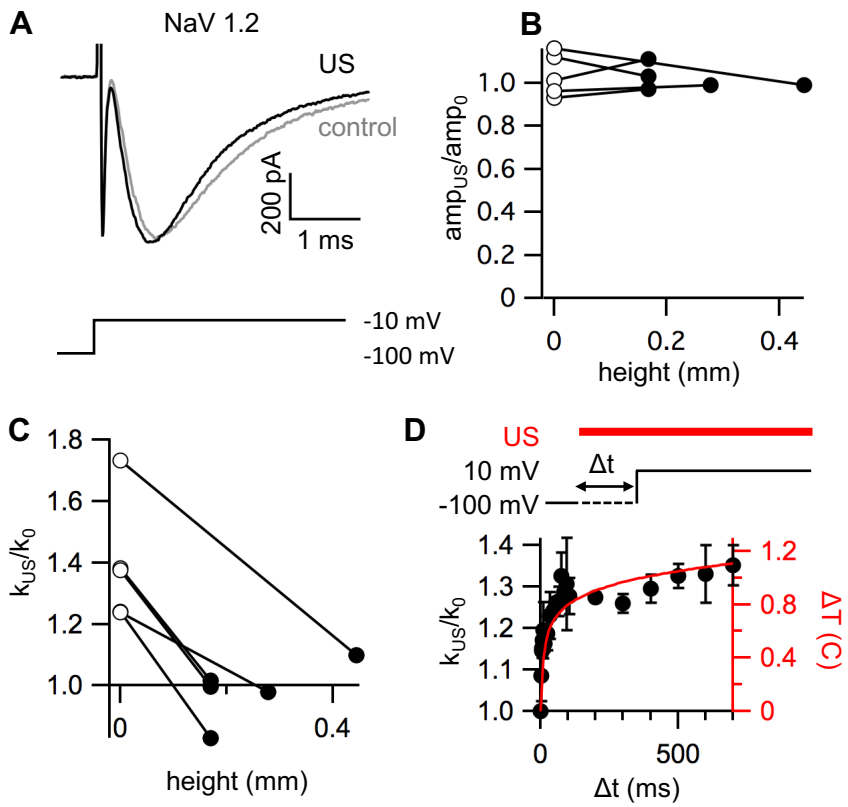

Figure 5. Effects of ultrasound on $Na_V1.2$ Channels. A. Example currents for $Na_V1.2$ channels (in adhered cells at ground level) in response to a voltage step from -100 to -10 mV with (*black current trace*) and without (*grey current trace*) stimulation. The 90-W/cm^2, 1-s US stimulus was applied starting 800 ms before the voltage step. **B.** Ratio of the $Na_V1.2$ current amplitude at -10 mV in the US condition to the amplitude in the control condition (amp_{US}/amp_0) for five detached cells at zero height (*open circles*) and at various heights above the bottom of the experimental chamber (*black circles*), for the same stimulus protocol used in *A*. There was no significant effect of height on the amplitude ratio (paired t-test, P = 0.71). **C.** Ratio of the $Na_V1.2$ inactivation rate at -10 mV in the US condition to the inactivation rate in the control condition (k_{US}/k_0) in five detached cells at zero height (*open circles*) and at various heights above the bottom of the experimental chamber (*black circles*), for the same stimulus protocol used in part *A*. The difference was statistically

significant (paired t-test, P = 0.001, N = 5) **D.** Comparison of the time courses of simulated acoustic heating and effects of US on $Na_V1.2$ channel inactivation. The mean (± SD) ratio of the $Na_V1.2$ inactivation rate at -10 mV in the US condition to the inactivation rate in the control condition (k_{US}/k_0) is shown as a function of the US exposure time (Δt) prior to the voltage step for N = 4 different cells at each interval, along with the simulated temperature change as a function of the time relative to the start of the US stimulus (*red line*, right axis). The experiments and simulation were done with 50-micron polystyrene.

Table 1. Values of Material Properties Used in Finite-Element Simulations.

	water	**polystyrene**
density (kg/m^3)	1000[a]	1040[b]
speed of sound (m/s)	1500[a]	2300[b]
attenuation coefficient (neper/m)	46[c]	160[d]
heat capacity (J/kg·K)	4180[a]	1200[e]
thermal conductivity (W/m·K)	0.6[a]	0.1[e]
viscosity (Pa·s)	0.001[a]	(not applicable)
Young's modulus (Pa)	(not applicable)	10^9 [f]
Poisson's ratio	(not applicable)	0.4[f]

Sources for material properties are as follows: [a]standard value; [b]based on typical acoustic properties of plastics (www.ondacorp.com/images/Plastics.pdf); [c] www.ondacorp.com/images/Liquids.pdf; [d]measured (see Experimental Procedures, *Finite Element Simulations*); [e]based on typical thermal properties of plastics (www.professionalplastics.com/professionalplastics/ThermalPropertiesofPlasticMaterials.pdf); [f]based on typical mechanical properties of plastics (www.professionalplastics.com/…/MechanicalPropertiesofPlastics.pdf).