

1 Mechanics of *Hydra* Detachment from Substrates:
2 The Role of Substrate Rigidity and Starvation

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4 Running title: Mechanics of *Hydra* detachment

5 Abstract

6 *Hydra* is a fresh water hydrozoan living as a solitary polyp with a sedentary
7 feeder lifestyle attached to a substrate. In times of food shortage they are
8 reported to detach from their substrate and move either by drifting or ‘somer-
9 saulting’. The attachment to the substrate is usually by the basal-body which
10 secretes a mucosal adhesive. The mechanical strength of the adhesion of *Hydra*
11 has not been quantified so far. Here, we measure the force required to detach
12 *Hydra vulgaris* and *Hydra magnipapillata* from a surface and the role of phys-
13 ical and physiological factors. In order to do this, we have developed a flow
14 chamber with a calibrated jet of water. We find *H. vulgaris* adhering to a hard
15 substrate - a glass cover slip- requires more force to detach it as compared to
16 a soft substrate- polyacrylamide gel. While *H. vulgaris* after one week of star-
17 vation detaches with very similar values of stress, *H. magnipapillata* detaches
18 more readily when starved. These results suggest that the strength of adhesion
19 is strongly affected by the stiffness of the substrate, while nutritional status
20 dependence of detachment force appears to be species dependent. Given that
21 *Hydra* detachment is required during locomotion, our measurements on the one
22 hand suggest the magnitude of forces the animal must exert to detach itself.
23 Additionally, our results suggest active detachment of the base might be re-
24 quired for *Hydra* to achieve movement, and only a small contribution coming
25 from weakening adhesion.

26 Keywords

27 *Hydra*, detachment, mechanics, shear stress.

28 1 Introduction

29 Aquatic life forms ranging from single-celled to multi-cellular have evolved a
30 variety of strategies to remain static through adhesion to substrates. The specific
31 mechanism by which they achieve this adhesion ranges from suckers and
32 nanometer scale spatulae to biological adhesives (reviewed by Gorb (2008)).
33 Amongst aquatic animals the adhesion of the mussel *Mytilus edulis* has been
34 particularly well studied (reviewed by (Waite, 2002)). The mussel shells attach
35 to rocky substrates with byssal threads with multiple proteins contributing differing
36 mechanical properties (Lin et al., 2007), of which the amino acid 3,4-
37 dihydroxy-L-phenylalanine (dopa) is considered a vital component (Lee et al.,
38 2006).

39 *Hydra* on the other hand, are fresh water dwelling Hydrozoans of phylum
40 Cnidaria that live as solitary polyps, typically found attached to substrates like
41 stems, branches or leaves under-water. Renewed interest in *Hydra* is due to
42 its regenerative ability, along with genome sequence and the evolutionary relatedness
43 of the regenerative pathways to vertebrates (Fujisawa, 2006; Watanabe
44 et al., 2009). In their natural environment, *Hydra* are subject to gentle flows
45 and so far their movement has been attributed to passive drifting. Wagner has
46 noted that the resistance of an attached *Hydra* to water flows might be an adaptation
47 to the diverse environmental conditions (still and flowing water) that it is
48 exposed to and its inability to actively swim, once suspended in water (Wagner,
49 1905). The movement of *Hydra* that are already attached to a substrate was
50 observed to occur by ‘somersaulting’- animals attach their tentacles to a new
51 position, detach the basal disk with body contraction, straighten and reattach
52 the basal disk at a new position near the hypostome (Wagner, 1905). Annandale
53 reported *Hydra vulgaris* from Indian samples to be observed to be actively
54 ‘crawling’, which involved similar ‘somersaulting’ motion (Annandale, 1911).
55 The movement was thought to enable the individual to leave unfavourable environments.
56 While it is known that muscles that are ectodermal and longitudinal drive contraction,
57 while endodermal circular muscles drive extension of *Hydra*,
58 the biomechanics of *Hydra* detachment has yet to be examined.

59 In its sedentary mode *Hydra sp.* is attached by its basal-body also called
60 the basal disk or ‘foot’ to substrates by a glandular secretion (Brien, 1960).
61 Like other parts of Hydra, the animal can also regenerate the basal disk when
62 amputated (Amimoto et al., 2006). Histologically the cells of the disk consist
63 of the inner endoderm and outer ectoderm. The cells are glandular, conical
64 in shape and filled with granules (Bode et al., 1986). The cells secrete large
65 amounts of mucus needed for the attachment of the animal to substrates. The
66 basal disk was thought for long to be a closed structure but more recently a
67 pore-like structure in the disk called the aboral pore (Shimizu et al., 2007) has
68 been found. While the histology of the foot is understood, the nature of the
69 mucus as a bioadhesive which works under water could be interesting both from
70 a fundamental perspective of adhesives, as well as applications in biocompatible
71 materials (Waite, 2002). Recently, the glue from *Hydra magnipapillata* has been
72 characterised and found to be based on glycans and glycoproteins (Rodrigues

73 et al., 2016a). Additional gene-expression analysis has revealed 21 transcripts to
74 be expressed in the basal disk alone (Rodrigues et al., 2016b). While remaining
75 attached is important for *Hydra*, the ability to detach and move is equally
76 important. It would appear addressing the biomechanics of detachment, could
77 connect the mechanics of muscle-generated forces with the biochemistry of ‘glue’
78 attachment.

79 Measuring the force for detachment of larger animals such as *Mytilus sp.* has
80 involved mechanical spring-based instruments (Bell and Gosline, 1996; Denny,
81 1987), while sea anemone detachment has been measured using force trans-
82 ducers (Koehl, 1977). Such instruments however do not mimic the naturally
83 occurring flows that aquatic animals are likely to experience and measurements
84 could suffer from artefacts from mechanical contact. Flow chambers address
85 some of these shortcomings and have been used to study biofouling using tur-
86 bulent (Schultz et al., 2000) or pumped flows coupled to inline flow meters (Neal
87 et al., 1996). Flow tanks that have been described for whole organism studies
88 (Vogel and LaBarbera, 1978) and parallel-plate flow chambers used to studying
89 leukocyte adhesion (Chen and Springer, 1999) are successful means to measure
90 detachment dynamics of samples ranging from cells to whole organisms.

91 We have chosen to characterise the biomechanics of *Hydra sp.* detachment
92 with calibrated fluid flows, from which we estimate the drag force required to
93 displace the animal from a substrate. We use this device to measure the flow rate
94 required to detach the *Hydra* from substrates of different stiffness and proceed
95 to examine the role that starvation and substrate stiffness plays in the stress
96 required to detach the individuals.

97 2 Results

98 2.1 *Hydra* sizes

99 In order to estimate the forces exerted by flow, we needed to morphologically
100 characterise the *Hydra* we used in the study. Two species were chosen due to the
101 differences in sizes and availability, namely *H. vulgaris* and *H. magnipapillata*.
102 While qualitatively in a dissection microscope *H. vulgaris* was seen to be shorter
103 than *H. magnipapillata*, their widths appeared comparable (Figure 1). This
104 was confirmed by the estimate of the mean cross-sectional diameter of the foot
105 of *Hydra vulgaris* to be 0.279 mm for (Figure 1(a)-(d)) and 0.342 mm for *H.*
106 *magnipapillata* (Figure 1(e)-(h)). Given the base of the hydra is approximately
107 circular, we estimated the mean base area of *H. vulgaris* and *H. magnipapillata*
108 to be 0.24 and 0.37 mm^2 respectively. These values show small differences
109 in base-areas, while the length of *H. magnipapillata* is approximately two-fold
110 greater than *H. vulgaris*, varying between individuals. For our measurements,
111 the flow was directed to the base of the anim and is not expected to influence
112 our measurement as a result. As a next step, we needed to calibrate the flow
113 flow chamber.

114 2.2 Characterizing the laminar range of the flow-chamber

115 The flow chamber setup consists of a syringe pump connected by tubing to
116 a trough filled with medium, in which *Hydra* is submerged and subjected to
117 the flows (Figure 2(a)). In order to characterise the nature of the fluid flow
118 in the chamber, we needed to ensure the forces are generated due to laminar
119 flows. Flowing safranin-stained water into the apparatus in the absence of any
120 obstacle, turbulence was observed beyond the pipe exit (E) at a certain point
121 of turbulence (T) onset (Figure 2(b)). This distance from the pipe exit (E)
122 at which turbulence (T) due to eddies is observed was measured for multiple
123 flow rates. The plot of distance (T) as a function of flow rate (Q) shows that
124 even for the fastest flow-rates, the eddies begin 10 mm from the pipe exit (E)
125 (Figure 2(c)). As a result, we placed the *Hydra* at 5 mm from E in subsequent
126 experiments, to ensure that the force experienced is due to laminar flows alone.

127 In order to independently confirm the consistency of this data, we also esti-
128 mated the Reynolds number for each as a function of flow rate (Q) using:

$$Re = \frac{\rho \cdot d \cdot V}{\eta} \quad (1)$$

129 where, Re is Reynolds Number, ρ ($kg \cdot m^{-3}$) is the density of fluid, d (m)
130 is the diameter of the pipe V ($m \cdot s^{-1}$) is the fluid velocity and η ($N \cdot s \cdot m^{-2}$)
131 is the dynamic viscosity. In order to estimate Re in terms of Q , we used the
132 relation $Q = V \times A$, where Q ($m^3 \cdot s^{-1}$) is the volume rate of the fluid and A
133 (m^2) is the area of cross section of the jet of fluid. The resulting values of Re
134 (Table 1) appeared to converge on a mean value of $Re = 1.3243(\pm 0.8785) \cdot 10^3$.

135 In order to relate the onset of turbulence with the distance from the pipe
136 exit (x), we rewrite Equation 1 in the simple form by using the relation $V = Q/A$
137 and substituting all constant values in one lumped constant c_1 as follows:

$$x = c_1/Q \quad (2)$$

138 This equation is qualitatively comparable to the experimental estimates of x as
139 a function of Q , i.e. an inverse relation (Figure 2(c)). The constant c is given
140 by:

$$c_1 = \frac{Re \cdot A \cdot \eta}{\rho} \quad (3)$$

141 We substitute known values into this Equation 3 and estimate c_1 to be 1.13
142 cm^4/s , given: $\eta = 1.002 \cdot 10^{-3} N \cdot s/m^2$ and $\rho = 10^3 kg/m^3$ and the cross-
143 sectional area of the flow $a_f = 5.0265 \cdot 10^{-7} m^2$, since the jet diameter = 0.8
144 mm. When we compare the fit of Equation 2 to our data of turbulence onset
145 distance (T) with Q (Figure 2(c)), the parameter $c_1^{fit} = 1.06 cm^4/s$. This
146 validates our approach from first principles.

147 In addition, the onset of turbulence over the entire range of flow rates remains
148 $x \geq 1$ cm. Thus, we confirm that the force experienced by objects at a distance
149 of less than ~ 1 cm can be treated as being due to a laminar flow for all values
150 of Q .

151 **2.3 Drag force experienced by *Hydra***

152 In order to estimate the force at which *Hydra* detaches from the substrate, we
153 need to relate the volume flow rate with force. To this end, the drag-force (F_{drag})
154 exerted by the fluid flow on *Hydra* was estimated from the drag-equation:

$$F_{drag} = C_d \cdot \rho \cdot V^2 \cdot A/2 \quad (4)$$

155 where C_d is the drag coefficient, ρ is the density of the fluid, V is the velocity
156 of flow and A is the projected area of the body in the path of fluid flow. We
157 can relate the linear velocity (V) of the fluid with the volume flow rate (Q) as
158 follows:

$$V = Q/a_f \quad (5)$$

159 where a_f is the cross-sectional area of the flow. On substituting V (Equation
160 5) in the expression for drag force (Equation 4) gives us:

$$F_{drag} = \frac{C_d \cdot \rho \cdot Q^2 \cdot A}{2 \cdot a_f^2} \quad (6)$$

161 .
162 Thus the detachment force of *Hydra* can be calculated using Equation 6. The
163 drag coefficient (C_d) is dimensionless constant and depends on properties of the
164 object and fluid such as shape and Reynolds number. For our calculations,
165 choice of C_d was made by approximating the shape of *Hydra* to a cylinder with
166 its long axis normal to the direction of flow. Based on the length:width ratio of
167 *Hydra*, the drag coefficient of 0.68 was chosen, based on standard results for a
168 cylinder with length to the diameter ratio of 2:1 (Stoecker, 2004). For the sake
169 of simplicity, we assumed the shape to be constant for the *Hydra* across all the
170 conditions.

171 The calculated estimates of force with increasing flow rate was fit to a func-
172 tion of the form:

$$F = c_2 \cdot Q^2 \quad (7)$$

173 where Q is the volume flow rate (cm^3/s) and c_2 is a constant to compare the
174 drag force dependence on flow rate between the species and the drag coefficients
175 assumed.

176 **2.4 Detachment force of *H. vulgaris* and *H. magnipapillata***

177 In order to measure detachment forces, individual *Hydra* were allowed to attach
178 to a glass coverslip (coated or uncoated) in the incubator. At the time of
179 the measurement, the coverslip was removed from the incubator and placed
180 at the same distance from the pipe exit ($x=0.5$ cm). The flow rate (Q) was
181 gradually increased until the *Hydra* detaches. Measurements were repeated for
182 ~ 10 individuals of *H. vulgaris* and *H. magnipapillata*. By starving one set of
183 animals, we addressed the effect of nutritional state. We also compared the
184 effect of changing substrate stiffness on *H. vulgaris*. For each experimental

185 condition, the flow rate at which *Hydra* detached was used to calculate the force
186 of detachment (Equation 7). Using the estimate of the area of the basal disk, we
187 thus estimate the shear stress of detachment. *H. vulgaris* detaches from glass
188 substrates over a wide range of shear stresses: 0.6 to 1.8 MPa, for starved and
189 fed samples (Figure 3(a)). *H. magnipapillata* detachment was only recorded in a
190 few cases, since in most of the cases, as they did not detach within the maximal
191 limit of ≤ 60 ml/min of the instrument (Figure 3(a)). Fed and unfed animals
192 of *H. vulgaris* show a very minor difference in detachment shear stresses. The
193 mean shear stress of detachment from a soft substrate of 5% polyacrylamide
194 is lower than that measured on glass. However, more measurements will be
195 required for statistical significance. However the shear stress required to detach
196 *H. magnipapillata* is higher than that for *H. vulgaris* in the same nutritional
197 state (Student's paired t-test with 95% confidence interval). We therefore find
198 the detachment shear stress to be in the range of 10^3 N/m², which is two orders
199 of magnitude smaller than in molluscs (Denny, 1987). This suggests the *Hydra*
200 adhere to their substrate with a weak glue in a substrate-stiffness dependent
201 manner.

202 3 Conclusions and Discussion

203 Here for the first time we have quantitatively characterised the mechanics of
204 *Hydra* detachment from an attached position on a substrate. We use a simple
205 flow device which generates known amounts of shear stress through flows gen-
206 erated in the aqueous. We can show that smaller shear stresses are required to
207 detach *H. vulgaris* from soft as compared to hard substrates. We find *H. vulgaris*
208 and *H. magnipapillata* show little difference between fed and unfed (starved for
209 one-week) conditions.

210 The steady attachment of sedentary aquatic animals can occur by multiple
211 mechanisms, but the strength of the attachment to the substratum is related
212 to its behaviour as well as the flows in which it lives. Typically free flowing
213 streams with a gentle flow are reported to have flow speeds in the range of 0.5
214 m/s to 3 m/s. We can make an order of magnitude estimate of the shear stress
215 (S) based on the fluid drag-force (F_{drag}) from Equation 4 which simplifies to
216 $S = \rho \cdot v^2 / 2 \cdot r$ using the C_d of 0.68 based on the 2:1 ratio of length to the diameter
217 (Stoecker, 2004) of *Hydra* and assuming the *Hydra* can be treated as cylinders,
218 so the projected half-area of the a cylinder is affected by drag. To estimated
219 the projected areas of *H. vulgaris* and *H. magnipapillata* the height is measured
220 to be approximately 3 and 5 mm respectively and radius (r) is the same as in
221 Section 2.1. As a result *H. vulgaris* will be expected to experience shear stresses
222 between 1.8×10^3 and 6.6×10^4 N/m², while *H. magnipapillata* is expected
223 to experience between 2.5×10^3 and 8.9×10^4 N/m². Given that we measure
224 detachment shear stresses for both species ranging between 6×10^2 to 3×10^3
225 N/m² (Figure 3(a)), it would suggest normal flows that the animal is likely to
226 experience, would be sufficient to detach the animal from the substrate. This
227 would suggest, that in addition to active motion, *Hydra* can also be passively

228 detached from its substrate. This is corroborated by observations of drifting
229 animals found in their natural habitat. Careful observations in still and moving
230 streams combined with measurements of flow-rates *in situ* could be potentially
231 used to test this prediction.

232 The measurements we report here are made on two kinds of artificial substrates-
233 glass and 5% polyacrylamide gel. While glass is very stiff with a Young's modu-
234 lus of 72.9 MPa, the gel used has a reported stiffness of ~ 8 kPa (Tse and Engler,
235 2010). Our measurements suggest the *Hydra vulgaris* is less firmly attached on
236 a soft substrate as opposed to a hard substrate. The comparison with the bulk
237 modulus of elasticity of freshwater aquatic plant leaves, which ranges between
238 1 and 10 MPa (Touchette et al., 2014), would suggest our measurements cover
239 the range of stiffness that *Hydra* could be expected to encounter when attached
240 to leaves. While on the one hand the differences in detachment are less than an
241 order of magnitude and subject to large variations, it would be interesting in
242 future to systematically vary substrate stiffness and examine the role it plays in
243 movement of the animal. Additionally, the mechanical properties of the specific
244 plants and other objects to which *Hydra* is naturally found attached to, could
245 also determine whether there is any role at all for substrate stiffness.

246 In our experiments, we have used inert substrates during the detachment
247 measurements, in order to study the role of mechanical properties in the ab-
248 sence of any material properties. However, it could have been possible that
249 the chemical nature of the 'glue' might also be modulated during detachment,
250 such as by hydrolysis by some enzyme produced by the animal itself. However
251 a recent study that investigated the glue concluded that active attachment is
252 more likely to be the primary method by which *Hydra* achieves detachment
253 (Rodrigues et al., 2016a). While our data corroborates this by an independent
254 means, since we do not observe a clear starvation dependent weakening of the
255 bond, it would be useful in future to use motion-capture to carefully test these
256 theories of 'somersaulting' movement of *Hydra* to capture the entire cycle of
257 movement (Figure 3(b)) and the mechanics involved.

258 The force required to detach *Hydra* is two orders of magnitude smaller than
259 the detachment stress of $1.2 \cdot 10^5$ N/m² reported for the well studied sedentary
260 marine mussel *Mytilus sp.* (Denny, 1987). We hypothesise that the difference
261 in habitat of *Hydra sp.* which mostly inhabits ponds and slow-flowing streams
262 means that the detachment stresses do not need to be as high as those observed
263 in *Mytilus* mussels, typically found attached to inter-tidal rocks subject to con-
264 stant wave action (Bell and Gosline, 1996). Based on reports by Annandale and
265 others, it is reasonable to assume this weaker adhesion of *Hydra* is an adaptation
266 to the forces generated by currents it usually experiences in its natural habitat
267 and for the mode of motility that it adopts.

268 The measurement setup, while simple, provides useful initial answers to a
269 mechanical approach to animal behaviour. Potentially in future higher flow rate
270 methods would require taking into consideration the turbulent regime (Schultz
271 et al., 2000). In addition the starvation conditions in the native environment
272 that trigger 'somersaulting' movement are not clearly reported. In our work
273 we have empirically chosen a week of starvation. In future a controlled study

274 on the factors and duration of nutrient withdrawal, combined with mechanics
275 could help us better understand the triggers that govern the decision of *Hydra*
276 to move.

277 In conclusion we have characterised the shear stress of detachment of two
278 species of *Hydra* and find them to range between 0.14 to ≥ 0.8 kPa. We have
279 shown for the species examined, the detachment stress is independent of the
280 nutritional state (i.e. fed as compared to starved for one week) and only weakly
281 dependent on substrate-stiffness. Additionally we find *H. magnipapillata* is de-
282 tached at a higher stress value as compared to *H. vulgaris*. It leads us to
283 hypothesise that the active detachment of *Hydra sp.* is likely to be driven by
284 active muscle contractions. This work sets the stage for a more comprehensive
285 study of the mechanics of locomotion by this organism.

286 4 Materials and Methods

287 4.1 Growth and handling of *Hydra*

288 *Hydra vulgaris Ind-Pune* (Figure 1(a)-(d)) and (Reddy et al., 2011) and *Hydra*
289 *magnipapillata* (Figure 1(e)-(h)) were obtained from ARI (Pune, India). They
290 were maintained in ~ 200 ml of ‘M’ solution containing 0.1 mM KCl, 1 mM
291 NaCl, 1 mM CaCl₂ · 2H₂O, 1 mM Tris (pH 8) and 0.1 mM MgSO₄ · 7H₂O in
292 water (Sugiyama and Fujisawa, 1977). The animals were maintained at 18 °C
293 in an incubator with a lamp with timer kept on for $\sim 12h$ to artificially induce
294 day-night cycles (Thermo Scientific, USA) and the beaker cleaned on a daily
295 basis. *Hydra* were fed two to three hatched *Artemia sp.* (brine shrimp) every
296 two days that were grown in a 0.6% saline solution and washed and filtered in
297 tap water before being used as feed.

298 4.2 Imaging and microscopy

299 Individual *H. vulgaris* and *H. magnipapillata* animals placed in a petri dish with
300 a drop of ‘M’ solution to prevent desiccation and imaged using a Leica dissection
301 microscope S8 APO illuminated with a Leica (24 DC) LED control unit and
302 equipped with an EC 3 camera (Leica Microsystems, Germany). The onset of
303 turbulence was recorded using autofocus and automatic exposure settings in
304 video mode on a Canon EOS 1000D camera (Canon Inc., Japan).

305 4.3 Flow chamber and detaching *Hydra*

306 A syringe pump (PhD Ultra, Harvard Apparatus, USA) with a 20 ml plastic
307 syringe (BD Biosciences, India) was connected to polycarbonate tubing of inner
308 diameter (I.D.) 3 mm (BioRad, USA) and further connected with an adaptor
309 (BioRad, USA) to a tube of I.D. 0.8 mm, taped to the bottom of a glass trough
310 (Figure 2(a)). For *Hydra* detachment experiments, the animals were allowed to
311 attach to a glass coverslip (MicroAid, Pune, India) such that the animal was
312 at a distance of 0.5 cm from the tube outlet, in line with the fluid flow. Flow

313 experiments typically involved increasing the flow rate from 10 ml/min with
314 increments of ~ 2 ml/min until the animal detached due to the force generated.
315 Those experiments in which the *Hydra* was either attached to the substrate with
316 its tentacles, or did not attach at all, or failed to detach at all flow rates were
317 ignored in analysis of detachment shear stress.

318 **4.4 Modulating substrate stiffness**

319 For experiments to measure the effect of substrate stiffness, a 0.75 mm thick 5%
320 polyacrylamide gel was prepared using a 5 ml solution of 5% acrylamide and
321 0.22% bisacrylamide, 25 μ l APS (1/200 volume) and 2.5 μ l TEMED (1/2000
322 volume) (all reagents Sigma-Aldrich, Mumbai, India) and curing for 15 min
323 between two plates layered with water in a standard polyacrylamide gel elec-
324 trophoresis (PAGE) setup (BioRad, USA). This gel of stiffness ~ 8 kPa (Tse
325 and Engler, 2010), was layered on the coverslip and the *Hydra* was allowed to
326 attach to the gel. The remainder of the experiment was performed in a manner
327 similar to the experiments for detaching *Hydra* from glass coverslips.

328 **4.5 Data analysis**

329 Images of *Hydra* and the onset of turbulence were processed using ImageJ
330 (Schneider et al., 2012). Fitting data to functions was performed using the
331 non-linear fitting tool (*nlinfit*) in MATLAB (Mathworks Inc., MA, USA), as
332 was all plotting. Statistical testing of mean shear stresses was performed by
333 comparing the means by a Student's t-test.

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405 **6 Tables**

Flow rate, Q (ml/min)	Reynolds Number, Re
10	264.863
20	529.726
30	794.589
40	1059.45
50	1324.32
60	1589.18
70	1854.04
80	2118.9
90	2383.77
100	2648.63

Table 1: Flow is laminar within the experimental range of flow rates.

406 **7 Figures**

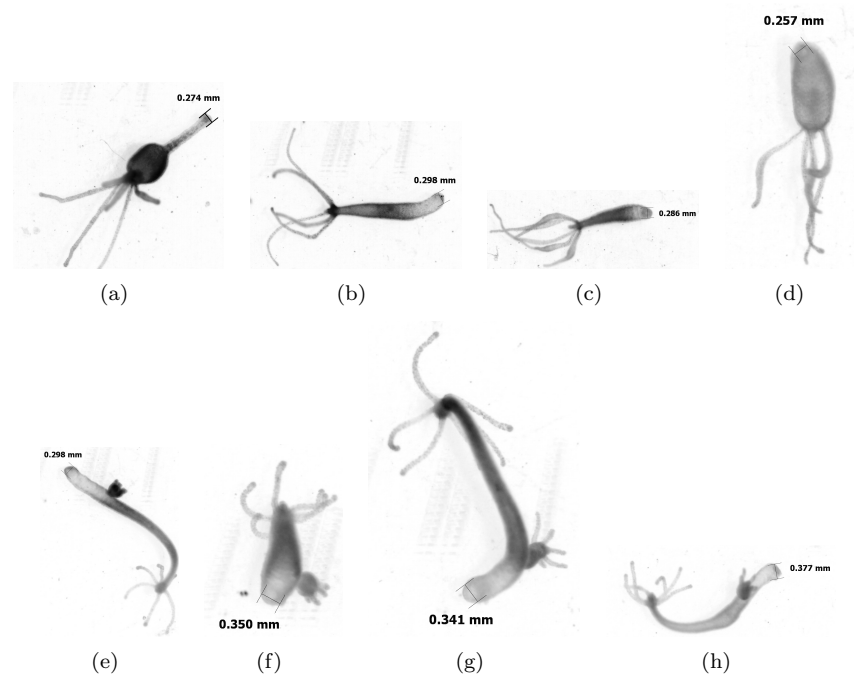


Figure 1: **Estimating the size of *Hydra*:** Images of live animals of (a)-(d) *Hydra vulgaris* and (e)-(h) *Hydra magnipapillata* under a dissection microscope in a drop of “M” solution. The scale bar in every image indicates the base diameter.

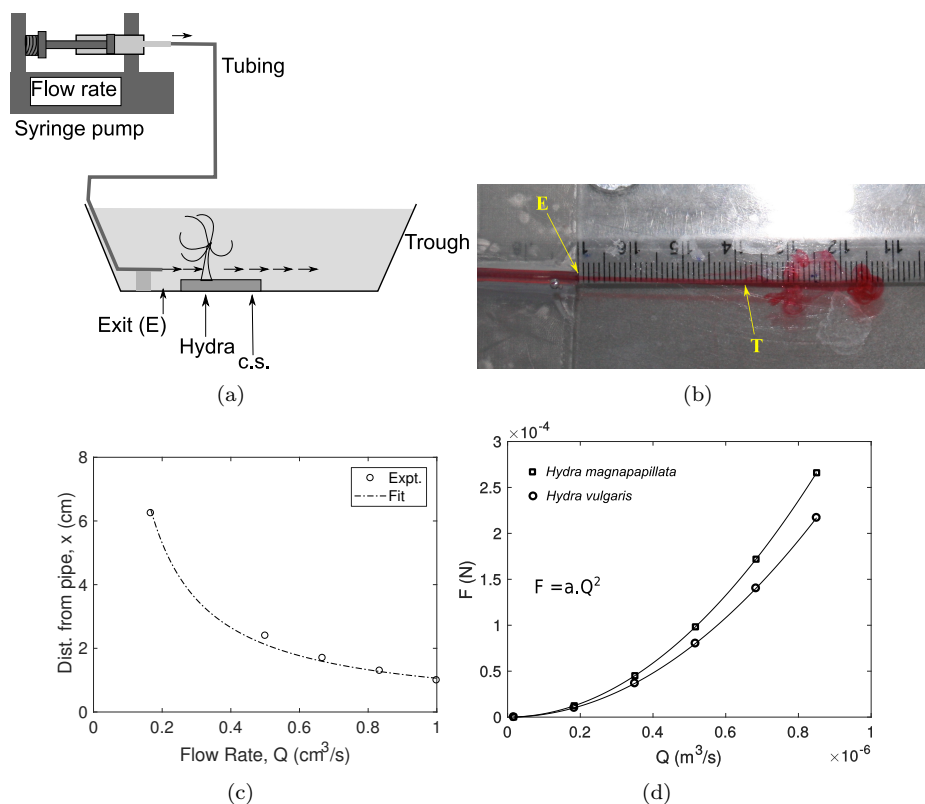


Figure 2: **Detachment of Hydra by laminar flow set up.** (a) A schematic representation of the experimental set up with the *Hydra* placed on a coverslip (c.s.) at a fixed distance from the pipe exit (E) submerged in buffer in a glass trough. The arrows indicate the direction of flow of water from the syringe pump. (b) The view from the top of the distance of the onset of turbulence (T) flow from the pipe exit (E) estimated by flowing safranin containing water. (c) The distance from E at which turbulence is seen is plotted against the flow rate (circle) and fit (dashed line) by the equation $x = c_1/Q$ (Equation 2). (d) The drag force experienced by Hydra due to flow is calculated (Equation 4) and fit by Equation 7. The fit parameter for *Hydra vulgaris* (circle) is $c_2 = 3 \cdot 10^8 \text{ kg} \cdot \text{m}^{-5}$ and for *Hydra magnapapillata* (square) is $c_2 = 3.68 \cdot 10^8 \text{ kg} \cdot \text{m}^{-5}$.

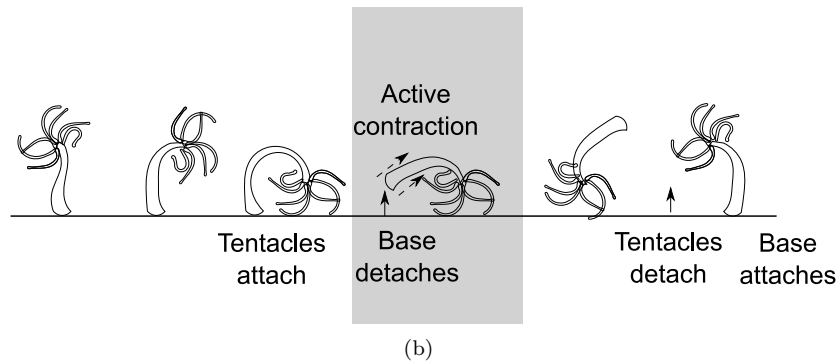
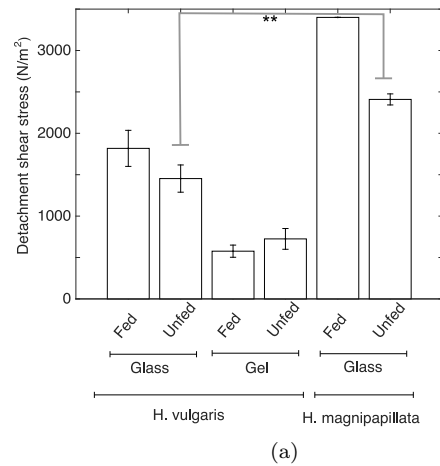


Figure 3: **Detachment shear stress:** (a) The mean shear stress (\pm S.E.) of detachment for *H. vulgaris* and *H. magnipapillata* under fed and unfed conditions for individual hydra attached to glass (stiffness $7.29 \cdot 10^7$ kPa). *H. vulgaris* detachment was also tested on a 5% polyacrylamide gel (stiffness ~ 8 kPa). The mean shear stress was compared using a pairwise, two-sided t-test with $\alpha = 0.05$ (**). (b) The schematic depicts the role of active detachment of the base from the substrate during ‘somersaulting’ movement by *Hydra*. The detachment shear stress measurements presented here could suggest the amount of force required to be generated by active contraction of *Hydra*.