# IP$_3$R activity increases frequency of RyR-mediated sparks by elevating dyadic Ca$^{2+}$

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## Abstract

Calcium (Ca$^{2+}$) plays a critical role in the excitation contraction coupling (ECC) process that governs the contraction of cardiomyocytes during each heartbeat. While ryanodine receptors (RyRs) are the primary Ca$^{2+}$ channels responsible for mediating cell-wide Ca$^{2+}$ transients during ECC, Ca$^{2+}$ release via inositol 1,4,5-trisphosphate (IP$_3$) receptors (IP$_3$Rs) have been reported to elicit ECC-modulating effects. Recent studies suggest that the proximal localization of IP$_3$Rs at dyads grants their ability to modify the occurrence of Ca$^{2+}$ sparks (elementary Ca$^{2+}$ release events that constitute ECC-associated Ca$^{2+}$ transients) which may underlie the modulatory effects on ECC. Here, we aim to uncover the mechanism by which IP$_3$Rs affect Ca$^{2+}$ spark dynamics. To this end, we developed a mathematical model of the dyad that incorporates IP$_3$Rs to reveal their impact on local Ca$^{2+}$ handling and corresponding Ca$^{2+}$ spark formation. Consistent with published experimental data, our model predicts that the propensity for Ca$^{2+}$ spark formation increases with IP$_3$R activity. Our simulations support the hypothesis that IP$_3$R activity elevates Ca$^{2+}$ within the dyad, sensitizing proximal RyRs for future release. However, this lowers Ca$^{2+}$ in the JSR available for release and thus results in Ca$^{2+}$ sparks with the same duration but lower amplitudes.

## Keywords
Ca$^{2+}$ microdomains, Ca$^{2+}$ sparks, calcium, cardiomyocyte, IP$_3$R, RyR

## Abbreviations
Ca$^{2+}$, calcium; ECC, excitation contraction coupling; RyR, ryanodine receptor; IP$_3$, inositol 1,4,5-trisphosphate; IP$_3$R, IP$_3$ receptor; IP$_3$R-1, type 1 IP$_3$R; IP$_3$R-2, type 2 IP$_3$R; LTCC, L-type Ca$^{2+}$ channel; SR, sarcoplasmic reticulum; JSR, junctional SR; NSR, network SR; CICR, Ca$^{2+}$-induced Ca$^{2+}$ release; GPCR, G protein-coupled receptor; ET-1, endothelin-1; IICR, IP$_3$-induced Ca$^{2+}$ release; CaM, calmodulin; TnC,

troponin C; CSQ, calsequestrin; SERCA, sarco/endoplasmic reticulum ATPase; 1D, 1-dimensional; PSF, point spread function

# 1   Introduction

Underpinning the heart's pumping action is the concerted contraction and relaxation of individual cardiomyocytes, governed by the excitation-contraction coupling (ECC) process. In ventricular cardiomyocytes, ECC is initiated at the arrival of an action potential (AP) that permits a small calcium ($Ca^{2+}$) influx through voltage-gated L-type $Ca^{2+}$ channels (LTCCs) into $10 - 15$ nm wide microdomains delimited by T-tubules and the junctional cisternae of the sarcoplasmic reticulum (SR) (**Figure 1**). The $Ca^{2+}$ influx into these microdomains (henceforth dyads) induces a larger $Ca^{2+}$ release from the SR via resident ryanodine receptors (RyRs). This $Ca^{2+}$-induced $Ca^{2+}$ release (CICR) raises local $Ca^{2+}$ concentration ($[Ca^{2+}]$) in the dyad, producing elementary $Ca^{2+}$ release events that underlie ECC known as $Ca^{2+}$ sparks (1,2). By virtue of the distribution of T-tubules at 1.8 µm intervals across the cardiomyocyte that penetrate throughout the cell volume forming dyads, the synchronous evocation of $Ca^{2+}$ sparks at dyads by an AP facilitates the transient global rise in cytosolic $Ca^{2+}$ levels. This allows sufficient $Ca^{2+}$ to bind to troponin C (TnC) in myofilaments, thereby enabling the cross-bridge cycle that contracts the cardiomyocyte (3).

Like RyRs, inositol 1,4,5-trisphosphate ($IP_3$) receptors ($IP_3Rs$) are another family of $Ca^{2+}$-regulated $Ca^{2+}$ channels that reside on the SR of cardiomyocytes (4). $IP_3Rs$ are positively regulated by their ligand, $IP_3$ (5), which is produced following phospholipase C activation and phosphatidylinositol 4,5-bisphosphate hydrolysis downstream of G protein-coupled receptors (GPCRs) as well as certain receptor growth factor receptors (6). Indeed, ventricular cardiomyocytes stimulated by $G_q$ coupled GPCR agonists, such as endothelin-1 (ET-1), lead to $IP_3$-induced $Ca^{2+}$ release (IICR) via $IP_3Rs$, which are shown to promote ECC-modulating effects such as arrhythmia and positive inotropy (7–12).

Despite lower expression levels (13) and $Ca^{2+}$ conductance (5) relative to RyRs, $IP_3Rs$ elicit these ECC-modulating effects by their localisation to functionally relevant $Ca^{2+}$ signalling sites in the cell (14). A notable example is the colocalization of $IP_3Rs$ and RyRs at dyads (8,15). It has been recently shown that stimulating the activity of $IP_3Rs$ significantly increases the frequency of dyadic $Ca^{2+}$ spark events (15). In this regard, IICR is hypothesised to elevate dyadic $Ca^{2+}$, thereby priming and recruiting otherwise "silent" RyRs towards activation (8,11,14,15). The resulting increase in $Ca^{2+}$ spark propensity may then lead to ECC-modulating effects observed (11,16).

Here, we employed computational modelling to simulate the effects of IICR in the dyad. We developed a 1D spatial model of a dyad containing RyRs and type 2 $IP_3Rs$. Using this model, we varied the number of $IP_3Rs$ in the dyad and simulated its effect on the local $Ca^{2+}$ dynamics as well as the properties of $Ca^{2+}$ $Ca^{2+}$ sparks generated. Our model predicts that increasing $IP_3R$ activity increases the baseline dyadic $[Ca^{2+}]$ at the expense of that in the JSR. This elevation of dyadic $[Ca^{2+}]$ then sensitises colocalized RyRs to be more active, consequently increasing the propensity of $Ca^{2+}$ spark formation. The decrease in JSR $[Ca^{2+}]$ thus resulted in $Ca^{2+}$ sparks with lower amplitudes but a similar duration.

# 2   Methods

## 2.1   Model Formulation

We model the spatiotemporal evolution of $Ca^{2+}$ as a system of partial differential equations (PDEs) at three interconnected compartments: cytosol, junctional SR (JSR), and network SR (NSR). The spatiotemporal evolution of $[Ca^{2+}]$ in these compartments is described by the variables $[Ca^{2+}]_c$, $[Ca^{2+}]_{JSR}$, and $[Ca^{2+}]_{NSR}$ respectively. These are shown in order in the equations below:

$$\frac{\partial [Ca^{2+}]_c}{\partial t} = \mathcal{D}_c \nabla^2 [Ca^{2+}]_c + \sum_{i=1}^{4} J_{B_i} + J_{RyR} + J_{IP_3R} - J_{SERCA} + J_{Leak}$$

$$\frac{\partial [Ca^{2+}]_{JSR}}{\partial t} = \mathcal{D}_{JSR} \nabla^2 [Ca^{2+}]_{JSR} + J_{B_5} - \left( J_{RyR} + J_{IP_3R} \right) + J_{Refill}$$

$$\frac{\partial [Ca^{2+}]_{NSR}}{\partial t} = \mathcal{D}_{NSR} \nabla^2 [Ca^{2+}]_{NSR} + J_{SERCA} - J_{Refill} - J_{Leak}$$

where $\mathcal{D}_c$, $\mathcal{D}_{JSR}$, and $\mathcal{D}_{NSR}$ represent the diffusivity of Ca²⁺ in the cytosol, JSR, and NSR compartments, respectively. $J_{RyR}$ and $J_{IP_3R}$ correspond to the Ca²⁺ release fluxes by open RyRs and IP₃Rs respectively. $J_{SERCA}$ corresponds to the Ca²⁺ uptake flux by sarco/endoplasmic reticulum Ca²⁺ ATPase (SERCA). $J_{Refill}$ corresponds to the Ca²⁺ refill flux from the NSR into the JSR compartment. $J_{B_i}$ corresponds to the flux of Ca²⁺ binding to mobile and immobile buffer species $i$.

The reaction diffusion of Ca²⁺ buffers are described by

$$\frac{\partial [CaB]_i}{\partial t} = \mathcal{D}_{B_i} \nabla^2 [CaB_i] - J_{B_i}$$

where $[CaB]_i$ corresponds to the concentration of Ca²⁺-bound buffer species $i$, with $i \in \{1,2,3,4,5\}$ representing buffers ATP, calmodulin (CaM), Fluo-4, troponin C (TnC), and calsequestrin (CSQ) respectively. $\mathcal{D}_{B_i}$ corresponds to the diffusivity of Ca²⁺-bound buffer species $i$. Immobile buffers TnC and CSQ have $\mathcal{D}_{B_4} = \mathcal{D}_{B_5} = 0$.

## 2.2   Calcium Fluxes

The flux for each buffer species $i$ is given by

$$J_{B_i} = k_{off}[CaB_i] - k_{on}[Ca^{2+}]\left([B_i^{Tot}] - [CaB_i]\right)$$

where $[B_i^{Tot}]$ corresponds to the total concentration of buffer species $i$. $k_{on}$ and $k_{off}$ corresponds to the forward and backward reaction rates of buffer species $i$ with Ca²⁺ respectively.

The refill flux from the NSR to the JSR compartment is given by

$$J_{refill} = g_{refill}\left([Ca^{2+}]_{NSR} - [Ca^{2+}]_{JSR}\right)$$

where $g_{refill}$ is the refill flux rate adjusted to achieve a realistic JSR refill time constant of ~130 ms (17–19) in simulations where the number of IP₃Rs is 10 (see **Figure 1**). We assume the average number of IP₃Rs in a cluster to be 10.

The release fluxes from RyRs and IP₃Rs are given by

$$J_{RyR} = n_{RyR} g_{RyR}\left([Ca^{2+}]_{JSR} - [Ca^{2+}]_c\right)$$

$$J_{IP_3R} = n_{IP_3R} g_{IP_3R}\left([Ca^{2+}]_{JSR} - [Ca^{2+}]_c\right)$$

where $n_{RyR}$ and $n_{IP_3R}$ correspond to the number of open RyRs and IP₃Rs, respectively, whereas $g_{RyR}$ and $g_{IP_3R}$ correspond to the flux rate of RyR and IP₃R release, respectively. The value of $g_{RyR}$ was adjusted to yield a characteristic Ca²⁺ spark profile in the simulation condition where only RyRs are present in the dyad. $g_{IP_3R}$ is less than $g_{RyR}$ by a factor of 2.85 as the Ca²⁺ conductance of IP₃Rs is estimated to be lower than that of RyRs by that amount (5).

Fluxes due to SERCA uptake activity were directly adapted from (20), which takes the form

$$J_{SERCA} = 2v_{cycle}A_p$$

where $v_{cycle}$ is the cycling rate per SERCA molecule and $A_p$ is the cytosolic concentration of SERCA homogenously spread throughout the bulk cytosolic region. The complete expression of each term is provided in Supplementary Materials.

An SR leak flux was also introduced to maintain the cytosolic $Ca^{2+}$ background concentration of 0.1 µM. We use the same formulation as the SERCA model to balance $J_{SERCA}$ such that $[Ca^{2+}]_c$ does not fall below 0.1 µM. Therefore, the SR leak flux is expressed as,

$$J_{Leak} = J_{SERCA}([Ca^{2+}]_c = 0.1)$$

All parameter values are listed in Supplementary Materials.

## 2.3 Calcium Channels

Two types of $Ca^{2+}$ channels are stochastically simulated in the model – RyRs and IP$_3$Rs. The gating of each RyR is directly adapted from the 2-state RyR model developed in (21); IP$_3$Rs are modelled after the 6-state Siekmann model that incorporates non-steady state kinetics developed and used in (22,23). Mathematical expressions of the IP$_3$R model used in (22,23) were parameterised specifically to fit the open probability curve of type 1 IP$_3$Rs (IP$_3$R-1). IP$_3$R-1 have different channel activities for the same range of $[Ca^{2+}]$ compared to type 2 IP$_3$Rs (IP$_3$R-2) (compare **Figure 8**C and **Figure 8**D in Supplementary Materials), the isoform most expressed in cardiomyocytes (5). Therefore, we varied coefficients and exponents of the IP$_3$R model in (22,23) to fit the open probability curve of IP$_3$R-2 data from (24). This allowed us to obtain an IP$_3$R model specific to IP$_3$R-2. This modification is essential due to the higher levels of $Ca^{2+}$ required for channel inhibition, which thus allow IP$_3$R-2 to remain active for longer in the conditions of the dyad. Details are described in Supplementary Materials. In simulations containing IP$_3$R-2s, their gating behaviour were computed with a fixed IP$_3$ concentration of 0.15 µM, similar to that used in (22).

## 2.4 Model Geometry

The dyad and its surrounding cytoplasmic space are represented on a 1-dimensional (1D) simulation domain of size 8 µm. The 1D simulation domain reflects the portion of a typical experimental confocal line scan where a dyad is located. The buffering of $Ca^{2+}$ by mobile buffers ATP, CaM, and Fluo-4 take place throughout this domain. The domain consists of 200 elements of size 0.04 µm, with the center nine elements (0.36 µm long) representing the dyadic region where RyRs and IP$_3$Rs are placed as shown in **Figure 1**. Elements outside this region represent the bulk cytosol where $Ca^{2+}$ is subject to additional buffering by TnC and uptake into the NSR by SERCA. In all simulations, the number of RyRs in each element is fixed at 15, consistent with the average number of RyRs in a cluster as determined by super resolution microscopy techniques in healthy cardiomyocytes (25–27). However, similar data on IP$_3$R clusters are not yet available. Therefore, to investigate the effect of enhanced IP$_3$R activity, the number of colocalized IP$_3$Rs in their element is varied between 0, 5, 10, and 20, corresponding to situations where there are no, low, intermediate, and high levels of IP$_3$R activity relative to the number of RyRs. The JSR compartment is designated the same location and the same number of elements as the dyadic region. Open RyRs and IP$_3$Rs thus result in a $Ca^{2+}$ flux from the JSR into the dyad that is driven by the difference in $[Ca^{2+}]$ between the two compartments. $Ca^{2+}$ in the JSR is subject to buffering by CSQ and refill from the NSR compartment. The non-junctional regions of the NSR compartment are homogenously distributed with SERCA that pumps $Ca^{2+}$ from the bulk cytosol into the SR. SR leak fluxes are likewise present along non-junctional regions of the NSR compartment and leaks $Ca^{2+}$ into the bulk cytosolic region to maintain a baseline cytosolic $[Ca^{2+}]$ of 0.1 µM. No-flux conditions were imposed on the boundaries of the simulation domain.

**Figure 1. Schematic diagram of the compartments, fluxes, and arrangement of Ca²⁺-handling proteins in the dyad and its 1D representation in the model.** *The dyad is represented on an 8-μm, 1D computational domain with three compartments: cytosol, JSR and NSR. The center nine elements of the cytosolic compartment represent the dyadic region where RyRs and IP₃Rs are located while the remaining elements represent the bulk cytosol where Ca²⁺ is additionally subject to $J_{SERCA}$, $J_{Leak}$ and $J_{B_4}$. Ca²⁺ from the JSR is released in the dyadic region via open RyRs and IP₃Rs and diffuses along the cytosolic compartment, reacting with buffers before eventually being pumped back into the NSR which refills the JSR.*

## 2.5    Numerical Methods and Implementation

The system of PDEs were discretised using the forward time centered space finite difference scheme, similar to (28). The resulting system of ODEs was solved using the explicit Euler method with adaptive time stepping capped at a maximum of $1 \times 10^{-4}$ ms and a regular spatial resolution of 0.04 μm. Stochastic IP₃R and RyR gating states were solved using a hybrid Gillespie method as described in (29). The time at which any one receptor changes state determines the time step forward for which the system is solved (adaptive time stepping). Each simulation was repeated 200 times. In all simulations, the model was run for 1000 ms to ensure the system achieves steady state before any results were recorded. All codes and computations were implemented in MATLAB (The MathWorks Inc., Natick, Massachusetts).

# 3 Results

## 3.1 1D model reproduces calcium spark dynamics

The first column of **Figure 2**A illustrates the typical time evolution of $Ca^{2+}$ in different compartments of the model during a $Ca^{2+}$ spark in RyR-only simulations i.e., no $IP_3Rs$. To replicate CICR during ECC arising following $Ca^{2+}$ influx via LTCCs, $Ca^{2+}$ sparks were initiated by introducing a 2-ms $Ca^{2+}$ flux, reaching $\approx$ 30 µM, to elements in the dyadic region where RyRs are placed at the 1000 ms time point. This influx can be observed by the initial rise in $[Ca^{2+}]$ (with no variance) at the center of the dyadic region. The resultant opening of RyRs occurs rapidly and releases a greater amount of $Ca^{2+}$ from the JSR, thus providing a temporary positive feedback mechanism for the opening of other RyRs via CICR. RyRs open shortly after the initiation trigger and peaked at $\approx$ 13 RyRs for $\approx$ 9 ms before closing completely after $\approx$ 16 ms, consistent with simulation results from (21) whereby RyRs terminate after $\approx$ 20 ms of activity. During this time, dyadic $[Ca^{2+}]$ increased to $\approx$ 300 µM on average and declined back to $\approx$ 0.1 µM due to diffusion and chelation by buffers in the cytosol. While in the JSR, $[Ca^{2+}]$ declines and reaches its nadir at $\approx$ 300 µM $\approx$ 13 ms after the initiation trigger, during which point RyRs have already begun closing. These results support the induction decay mechanism of $Ca^{2+}$ spark termination proposed by (30) whereby the decay of the $Ca^{2+}$ flux through RyRs due to JSR depletion eventually impedes the inter-RyR regenerative CICR during a $Ca^{2+}$ spark, resulting in its termination. $[Ca^{2+}]$ in the JSR is then gradually replenished by $Ca^{2+}$ in the NSR at a time constant of $\approx$ 130 ms in the 10 $IP_3R$ simulation, consistent with experimental data (18,19). Together, our model depicting the 1D representation of the dyad is thus capable of reproducing $Ca^{2+}$ spark dynamics in reasonable agreement to that reported in other modelling and experimental studies (17–19,21).

**Figure 2. Ca²⁺ dynamics associated with LTCC-initiated Ca²⁺ sparks for different number of colocalized IP₃Rs. A:** *Time evolution of [Ca²⁺] in the dyad and JSR and the corresponding activity of RyRs and IP₃Rs during a Ca²⁺ spark. Solid lines and surrounding shaded region show mean and 95% confidence intervals respectively.* **B:** *Amplitude and FDHM of Ca²⁺ sparks generated based on [Ca²⁺] in the dyad.*

## 3.2 Increased IP$_3$R-2 activity decreases Ca$^{2+}$ spark amplitude and Ca$^{2+}$ stores

**Figure 2**A illustrates the effect of incorporating an increasing number of IP$_3$R-2s at the dyadic region. Despite varying the number of IP$_3$Rs in the dyad, our model is able to robustly simulate Ca$^{2+}$ spark events. Due to the significantly faster state transition rates and higher activity at lower [Ca$^{2+}$] of IP$_3$R-2s relative to RyRs, incorporating IP$_3$Rs in the dyad essentially causes a continuous Ca$^{2+}$ "leak" from the JSR into the dyad. Hence, a 1 s wait time was allocated to allow the system to equilibrate to steady state before simulating any Ca$^{2+}$ release events. This amount of time was sufficient for the system to equilibrate as triggering Ca$^{2+}$ sparks in simulations with longer wait times did not alter the resultant steady state [Ca$^{2+}$]. After equilibration, the average baseline dyadic [Ca$^{2+}$] rose above (inset of first row in **Figure 2**A) while that in the JSR fell below (second row of **Figure 2**A) the model's initial conditions of 0.1 µM and 1 mM respectively. Moreover, the magnitude of these changes increases with the number of IP$_3$Rs present in the dyad. We thus attribute these effects to the increased average number of open IP$_3$Rs (fourth row of **Figure 2**A). Altogether, our results suggest that enhancing IP$_3$R activity elevates dyadic [Ca$^{2+}$] at the expense of [Ca$^{2+}$] in the JSR.

To test the effect of IP$_3$R colocalization on Ca$^{2+}$ spark dynamics, we initiated Ca$^{2+}$ sparks in simulations where IP$_3$Rs are present by introducing a Ca$^{2+}$ flux into RyR-containing elements as described earlier. Generated Ca$^{2+}$ sparks have amplitudes that decrease with the increasing number of colocalized IP$_3$Rs (**Figure 2**B). This correlates well with the lower JSR [Ca$^{2+}$] available for release at steady state. However, the duration of these Ca$^{2+}$-triggered Ca$^{2+}$ sparks, measured by its full duration at half maximum (FDHM), is not significantly different (**Figure 2**B). This result can also be indirectly inferred from the time to complete closing of RyRs and time to nadir of [Ca$^{2+}$] in the JSR that are not significantly altered with increasing IP$_3$R activity. Mechanistically, the elevated dyadic [Ca$^{2+}$] together with the lower JSR [Ca$^{2+}$] at steady state collectively results in RyR Ca$^{2+}$ release fluxes that sustain inter-RyR CICR while depleting the JSR such that the Ca$^{2+}$ spark duration remains unchanged. In all cases, the occurrence of Ca$^{2+}$ sparks coincide with the transient opening of RyRs while the average IP$_3$R activity remained relatively constant throughout the simulation. This suggests that RyRs, and not IP$_3$Rs, are primarily responsible for the manifestation of Ca$^{2+}$ sparks, which is consistent with experimental results that show an almost complete loss of Ca$^{2+}$ spark events when RyRs are inhibited (15). Our model also successfully reproduced the experimental observation that [Ca$^{2+}$] in the JSR decreases to the same level after a Ca$^{2+}$ spark event regardless of its initial concentration (19), further bolstering our confidence of this model in simulating Ca$^{2+}$ sparks.

## 3.3 IP$_3$Rs increase propensity for spontaneous Ca$^{2+}$ sparks

By virtue of elevating dyadic [Ca$^{2+}$], IP$_3$Rs may play a role in enhancing the formation of Ca$^{2+}$ sparks. Indeed, cardiomyocytes treated with G$_q$ agonists or IP$_3$ significantly increased the number of spontaneous Ca$^{2+}$ spark events (7,8,31). This was attributed to IICR, but the underlying mechanism that led to this observation is not fully resolved. To test whether the colocalization of IP$_3$Rs in the dyad is a potential cause for the increase in spontaneous Ca$^{2+}$ spark events, we performed simulations in the absence of LTCC initiation – all Ca$^{2+}$ sparks generated occur spontaneously. After a 1 s wait time for system equilibration, the simulation was allowed to run for a further 2 s from which our results were obtained.

We recorded the number of Ca$^{2+}$ spark events generated from these simulations and their associated properties (amplitude and FDHM). We find that the percentage of simulations with at least 1 Ca$^{2+}$ spark event increases with the number of IP$_3$Rs (**Figure 3**A). Consistent with triggered Ca$^{2+}$ sparks, the average amplitudes of spontaneously generated Ca$^{2+}$ sparks also decrease (**Figure 3**B) with an increase in the number of colocalized IP$_3$Rs while their FDHM remain unchanged (**Figure 3**C). We hypothesise that the increase in spontaneously generated Ca$^{2+}$ sparks is due to the sensitization of RyRs by an

elevated dyadic [Ca$^{2+}$]. To verify that RyRs are more active due to their sensitization by IICR, we also recorded the number of RyR openings that did not develop into Ca$^{2+}$ sparks (an example of detecting these events is shown in **Figure** 7 of Supplementary Materials). Expectedly, the average number of spontaneous RyR openings that do not lead to the formation of Ca$^{2+}$ spark events also increased with the number of colocalized IP$_3$Rs (**Figure 3**D), reinforcing that RyRs are indeed more active in the presence of greater IP$_3$R activity. The increased number of spontaneous RyR openings thus raises the probability for Ca$^{2+}$ spark formation and contributes to the decreased JSR [Ca$^{2+}$] at steady state to some degree. Altogether, our results support the notion that IICR via colocalized IP$_3$Rs elevates dyadic [Ca$^{2+}$] sensitizing RyRs to promote Ca$^{2+}$ spark formation.

*Figure 3. The number of spontaneous Ca$^{2+}$ sparks increase with IP$_3$R activity. A: Percentage of simulations where more than 1 Ca$^{2+}$ spark event spontaneously occurred and the average number of Ca$^{2+}$ spark events per simulation. B: Amplitude of spontaneous Ca$^{2+}$ spark events. C: FDHM of spontaneous Ca$^{2+}$ spark events. D: Average number of spontaneous RyR opening events that do not lead to a Ca$^{2+}$ spark.*

## 4  Discussion

While the increased activity of IP$_3$Rs in the cardiomyocyte is shown to elicit ECC-modulating effects (7–12), the mechanistic basis of this action has not been demonstrated. However, evidence increasingly suggests that IICR can modulate ECC by the localisation of IP$_3$Rs at functionally important

Ca$^{2+}$ signalling sites (14), a quintessential example of which are dyads (8,15). Specifically, the expression of IP$_3$Rs at dyads is hypothesised to sensitise native RyRs towards activation via IICR (8,14,15), thus increasing the propensity for the formation of Ca$^{2+}$ sparks – elementary Ca$^{2+}$ release events underlying ECC-associated Ca$^{2+}$ transients. Using a 1D model of the dyad that includes both RyRs and IP$_3$R-2s, we set out to test this hypothesis while uncovering the underlying local Ca$^{2+}$ dynamics in the dyad.

## 4.1    IP$_3$R-mediated Ca$^{2+}$ release prime RyRs for release

A notable finding in our simulations is that the probability of spontaneous Ca$^{2+}$ spark events increase with the number of colocalized IP$_3$Rs. We were also able to uncover the mechanism by which this occurs through our recording of [Ca$^{2+}$] evolution with time at different compartments of the dyad. IP$_3$Rs demonstrate vastly different gating behaviour from RyRs. While RyRs are almost always in their closed state at baseline [Ca$^{2+}$], IP$_3$Rs are shown to be more highly active, rapidly fluctuating between their open and closed state throughout the simulation time course (fourth row of **Figure 2**A). Consequently, the activity of IP$_3$Rs in the dyad is akin to introducing a continuous leak from the JSR into the dyad. Increasing the number of IP$_3$Rs increases the magnitude of this "leak", as can be seen from a lower [Ca$^{2+}$] in the JSR, due to an increased number of open IP$_3$Rs on average. This SR [Ca$^{2+}$]-modulating "leak" due to IICR is consistent with that proposed by (32). On the other hand, however, the baseline [Ca$^{2+}$] in the dyad is increased due to IICR (inset in first row of **Figure 2**A). The elevation in dyadic [Ca$^{2+}$], in turn, increases the open probability of colocalized RyRs which is reflected by both the increased number of spontaneous Ca$^{2+}$ sparks (**Figure 3**A) and RyR openings that do not initiate a Ca$^{2+}$ spark (**Figure 3**C). Altogether, our results support the notion that IICR via IP$_3$Rs expressed in dyads increases the propensity for RyR-mediated Ca$^{2+}$ spark formation by increasing dyadic [Ca$^{2+}$].

Our findings have important implications about the wider role of IP$_3$Rs in cardiomyocytes. As we show that IICR increases the probability of Ca$^{2+}$ spark events by raising dyadic [Ca$^{2+}$], this mechanism may be useful in activating RyR clusters that are usually "silent" during ECC. This can potentially explain the enhanced Ca$^{2+}$ transient amplitude in cardiomyocytes where IP$_3$Rs are stimulated (7–11). Indeed, in a recent study in which a dyadic Ca$^{2+}$ reporter was employed IP3R activation was found to result in an increase in the number of dyads recruited during ECC (15). In diseased cardiomyocytes, the greater expression of IP$_3$Rs (8,12) suggest a compensatory mechanism for the increased decoupling of RyRs from LTCCs due to T-tubule degradation (33,34). But IICR in dyads could also contribute to increased spontaneous Ca$^{2+}$ release events in cardiomyocytes where IP$_3$Rs are stimulated, which can have arrhythmogenic consequences (7–11).

## 4.2    Model limitations and implications

We developed a 1D spatial model of a dyad that reproduced all major characteristics of a Ca$^{2+}$ spark. This enabled its utilization in conducting a qualitative investigation into the influence of IP$_3$R activity on the dynamics of Ca$^{2+}$ sparks in the dyad. While computationally less expensive, the reduced order of our model from 3D to 1D requires simplifying assumptions that presents several limitations, which we discuss below.

### 4.2.1    Arbitrary RyR and IP$_3$R Placement

In our model, we chose to fix the number of RyRs in a cluster as 15 based on recent estimates obtained from super resolution imaging data (25–27). Since similar data on IP$_3$R clusters is unavailable in the literature, the number of IP$_3$Rs in a cluster is varied to test the effect of increased IP$_3$R activity on the same RyR cluster. These clusters are then arbitrarily placed in elements of the dyadic region as shown in **Figure 1**. The 1D nature of our model precludes our ability to place each RyR in its own element in 3D space such that it can detect Ca$^{2+}$ that has diffused from other RyRs in the cluster. RyRs and IP$_3$Rs

that belong to the same cluster are placed in one element such that all $Ca^{2+}$ channels in that element are assumed to detect the same dyadic $[Ca^{2+}]$. Similar assumptions have also been employed in previous modelling studies simulating $Ca^{2+}$ sparks (20,35,36). While we acknowledge that developing models of higher dimensions permits one to incorporate the spatial arrangement of individual RyRs in the dyad, which influences $Ca^{2+}$ spark fidelity (17,27), our reduced-order model is sufficient for our purposes of illustrating the general effect of increasing IP$_3$R activity on $Ca^{2+}$ spark dynamics and derive an underlying mechanism for its increased occurrence.

### 4.2.2 Visualisation of $Ca^{2+}$ Spark Fluorescence

$Ca^{2+}$ spark characteristics obtained from experiments are derived from the fluorescence measurement of indicator dyes. To corroborate experimental observations with modelling results, many mathematical models incorporate the reaction kinetics of the indicator dye and simulate the resulting fluorescence measurement. Although we included the reaction kinetics of $Ca^{2+}$ with the indicator dye Fluo-4 in our model, we could not reliably corroborate its fluorescence with experimental observations, which show that $Ca^{2+}$ spark amplitudes do not change in the presence of IICR (15). Contrastingly, our modelling results show progressively decreasing $Ca^{2+}$ spark amplitudes with increasing IP$_3$R activity, which we attribute to the decreased JSR $[Ca^{2+}]$ available for release.

Whether this decrease in $Ca^{2+}$ spark amplitude, measured in terms of $[Ca^{2+}]$, is noticeably reflected in the fluorescence equivalent measurement cannot be directly confirmed from our model. This is because we find that the rise in dyadic $[Ca^{2+}]$ during a $Ca^{2+}$ spark saturates the indicator dye, resulting in a plateau of the fluorescence trace (see **Figure 8**A in Supplementary Materials). The saturation of Fluo-4 and the consequent plateau of the fluorescence trace thus prohibits us from conclusively inferring the amplitude and FDHM of the $Ca^{2+}$ spark. We partly attribute this saturation to the diffusion of $Ca^{2+}$ that is restricted to one dimension due to our model design. Additionally, the 1D setup of our model limited our ability to convolve the simulated fluorescence with a point spread function (PSF) of the objective lens of a confocal microscope. Hence the 1D model precludes a realistic visualisation of $Ca^{2+}$ sparks as they would be experimentally observed. We acknowledge this as a major limitation of our model. Nevertheless, simulated spark properties in in our model (**Figure 2**A) is consistent with many previous experimental and model findings. Hence, we base our findings on $[Ca^{2+}]$ instead.

Besides our model limitations, we also consider the discrepancy between our modelling results and experimental data to be partially attributable to reaction kinetics between the dye and dyadic $Ca^{2+}$. The dye does not fully provide a measure of dyadic $Ca^{2+}$ due to the dissociation constant between $Ca^{2+}$ and the indicator dye. Therefore, our model results that show decreasing $Ca^{2+}$ spark amplitude with increased IICR activity may not be detectable experimentally. Furthermore, the placement of the line scan in imaging experiments may not overlap precisely on the dyad and the resolution of the microscope which is $\approx 0.4$ um is roughly similar to the size of a dyad (37,38).

## 5   Conclusions

By incorporating the behaviour of both RyRs and IP$_3$Rs in our 1D model of the dyad, we show that IP$_3$R activity elevates dyadic $[Ca^{2+}]$, which sensitizes proximal RyRs towards activation. The colocalization of IP$_3$Rs with RyRs in the dyad thus increases the propensity for RyR-mediated $Ca^{2+}$ sparks which potentially underlies the ECC-modulating effects seen in ventricular cardiomyocytes treated with G$_q$ agonists. In this regard, further work is needed to link our findings of IP$_3$R-influenced $Ca^{2+}$ spark formation to multiscale whole-cell cardiomyocyte models incorporating IP$_3$ signalling (39) and $Ca^{2+}$ cycling (40,41) to elucidate its overall impact on global cytosolic $Ca^{2+}$ transient dynamics and ECC (42,43).

## Acknowledgements

This research was supported in part by the Australian Government through the Australian Research Council Discovery Projects funding scheme (project DP170101358) to EJC and VR, the Australian Research Council Centre of Excellence in Convergent Bio-Nano Science and Technology (project CE140100036) to EJC, the KU Leuven Global PhD Partnerships with The University of Melbourne Grant (GPUM/21/036) to VR and HLR, and The University of Melbourne's Research Computing Services and the Petascale Campus Initiative. HLR wishes to acknowledge financial support from the Research Foundation Flanders (FWO) through Project Grant G08861N and Odysseus programme Grant 90663. JC would like to thank Dr Pengxing Cao for the useful discussions on IP$_3$R modelling. S.T.J. is supported by the Australian Research Council (Project No. DE200100988).

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