

A minimal robophysical model of quadriflagellate self-propulsion

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Abstract.

Locomotion at the microscale is remarkably sophisticated. Microorganisms have evolved diverse strategies to move within highly viscous environments, using deformable, propulsion-generating appendages such as cilia and flagella to drive helical or undulatory motion. In single-celled algae, these appendages can be arranged in different ways around an approximately 10 μm cell body, and coordinated in distinct temporal patterns. Inspired by the observation that some quadriflagellates (bearing four flagella) have an outwardly similar morphology and flagellar beat pattern, yet swim at different speeds, this study seeks to determine whether variations in swimming performance could arise solely from differences in swimming gait. Robotics approaches are particularly suited to such investigations, where the phase relationships between appendages can be readily manipulated. Here, we developed autonomous, algae-inspired robophysical models that can self-propel in a viscous fluid. These macroscopic robots (length and width = 8.5 cm, height = 2 cm) have four independently actuated ‘flagella’ that oscillate back and forth under low-Reynolds number conditions ($\text{Re} \sim \mathcal{O}(10^{-1})$). We tested the swimming performance of these robot models with appendages arranged in one of two distinct configurations, and coordinated in one of three distinct gaits. The gaits, namely the pronk, the trot, and the gallop, correspond to gaits adopted by distinct microalgal species. When the appendages are inserted perpendicularly around a central ‘body’, the robot achieved a net performance of 0.15 – 0.63 body lengths per cycle, with the trot gait being the fastest. Robotic swimming performance was found to be comparable to that of the algal microswimmers across all gaits. By creating a minimal robot that can successfully reproduce cilia-inspired drag-based swimming, our work paves the way for the design of next-generation devices that have the capacity to autonomously navigate aqueous environments.

¹ *Keywords:* swimming, cell motility, robot, bio-inspired, gait, quadrupedal locomotion,
² algae

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3 1. Introduction

4 The capacity for self-generated movement is a distinguishing feature of most living
5 organisms. In the macroscopic world, locomotion is typically associated with inertia [1].
6 On the other hand, movement at the microscopic scale is subject to low Reynolds number
7 physics, and cannot take advantage of inertial coasting. Without motility, a bacterium
8 can only coast a minuscule distance an order of magnitude below the Ångström scale
9 [2]. Over billions of years of evolution, microorganisms have become adept at swimming,
10 evolving distinct mechanisms for powering and maintaining self-movement through a
11 fluid, often achieving speeds of several tens of body lengths per second. This active
12 motility confers a significant survival advantage, allowing microbes to navigate freely
13 towards regions or locations where nutrients or resources are more plentiful [3]. Depending
14 on the arrangement and number of locomotor appendages, single cells can execute
15 swimming gaits that are surprisingly reminiscent of animals. For example, the model
16 biflagellate alga *Chlamydomonas* actuates two equal-length flagella in a breaststroke [4],
17 while quadriflagellate algae (single cells with four-flagella) exhibit distinctive quadrupedal
18 gaits such as the trot or the gallop [5] (Fig. 1A,B).

19 In recent years, advances have been made in understanding the biomechanics of
20 microswimming. Here, the Reynolds number is small, $Re = UL/\nu$, where L is a typical
21 lengthscale of the swimmer, U a typical velocity scale, and ν is the kinematic viscosity
22 of the fluid. Equally important is the oscillatory Reynolds number $Re^{\text{osc}} = L^2\omega/\nu$ [6],
23 where ω the typical stroke frequency (which sets a tip velocity of ωL). When both are
24 small, flows are then governed by the Stokes equations: $0 = \nabla p - \mu \nabla^2 \mathbf{v}$ and $\nabla \cdot \mathbf{v} = 0$
25 (where \mathbf{v} and p are the flow and pressure fields), and have no explicit time-dependence.
26 Microorganisms are able to break time-reversal symmetry using non-reciprocal strokes or
27 body deformations, often involving whip-like appendages called cilia and flagella [2, 7].
28 While bacteria make use of rigid helical flagella [8], eukaryotes actuate motile cilia which
29 produce asymmetric waves of propulsion [9, 10]. For a microorganism oscillating a $10 \mu\text{m}$
30 flagellum at 50 Hz, $Re \sim 10^{-3}$, and $Re^{\text{osc}} \sim 10^{-2}$. One further asymmetry is required for
31 forward propulsion [11] - this is shape asymmetry, which is ensured by the slender aspect
32 ratio of all cilia and flagella (about 100). A rod sweeping through a fluid in the direction
33 perpendicular to the axis of the rod experiences approximately twice the drag compared
34 to when it is moved in the parallel direction [12]. Organisms across all scales have been
35 found to exploit this basic anisotropy in frictional forces for locomotion [13, 14, 15].

36 Despite the adoption of cilia and flagella as a common propulsion mechanism, the
37 microscale locomotion strategies of microorganisms have diversified significantly across
38 different phyla [16]. It is not well-understood why different gaits exist nor how they are
39 coordinated. For centuries, locomotor gaits have been studied in the context of terrestrial
40 animals, where the sequences of relative movement sustained by subsets of limbs or
41 legs have fascinated researchers. In vertebrates, gaits are thought to be generated by
42 central pattern generators (CPGs) [17]. But how can orderly, deterministic appendage
43 coordination occur in single cells in the absence of nervous control [16, 18]? Recent
44 theoretical and experimental work have show that dynamic gait selection, at least in
45 flagellates, appears to be an active and species-dependent process driven by intracellular

46 and mechanical coupling [18, 19]. Notably, distinct quadriflagellates can self-propel
47 at different speeds despite an apparently identical arrangement of flagella around the
48 cell body [5, 20]. Since the ancestral form of the green algal lineage may have
49 been a unicell with four flagella [21], there is much incentive to understand the precise
50 mechanisms of appendage coordination in such systems.

51 In the quest to address these open questions of movement control, extant organisms
52 can provide only a limited parameter space of possibilities in terms of size, shape,
53 beat frequency etc, often making it challenging to investigate certain configurations or
54 physical regimes. Theoretical and computational approaches have been instrumental in
55 shaping our understanding of active propulsion [12, 22], but these can be computationally
56 expensive or reliant on simplifying assumptions. Meanwhile robophysical modelling has
57 emerged as a powerful and versatile technique for elucidating organismal behaviour by
58 engineering customised configurations that can be easily tested in controlled laboratory
59 settings [23, 24, 25]. The revolution in robophysical modelling has been driven in
60 part by cheap electronics (motors, microcontrollers), and increasingly accessible control
61 technologies that can complement theoretical modelling to provide real biological insights
62 [24, 26]. However, trying to model cell movement is a significant conceptual challenge
63 when working *at the microscale*. Even though increasingly controllable micro- and nano-
64 devices have been fabricated to mimic the locomotive behaviours of biological swimmers
65 [27, 28], these are overwhelmingly driven by external magnetic, electric or chemical fields.
66 For instance magnetic fields are often unable to deliver the fine spatial control, required
67 to independently actuate individual artificial cilia in a given array or network. Thus
68 detailed investigations of the effect of gait on microswimming has mostly been restricted
69 to theoretical microswimmers [29, 30], occasionally in artificial or colloidal swimmers [31],
70 but seldom in microrobots [32].

71 The intrinsic limits of device manufacture at small scales severely undermines the
72 suitability of microrobots as realistic models of cell motility. To understand the influence
73 of gait on self-propulsion at low-Reynolds number, our goal is to build a dynamically-
74 scaled robophysical model which is truly *self-powered*, where the movement of individual
75 locomotor appendages can be prescribed and controlled independently. In contrast to
76 traditional 'microrobots', the large size allows us to explore and take advantage of
77 increasingly sophisticated electronics and control architectures [33, 34]. We can readily
78 reprogram these "roboflagellates" to execute specific swimming gaits, making them
79 uniquely suited to testing theories of bio-inspired and autonomous locomotion at low-
80 Reynolds number. This paper is organised as follows. We first identified and measured
81 the relative swimming performance of three species of quadriflagellate algae that exhibit
82 near-identical morphology but distinct swimming speeds. Next we built a centimetre-
83 sized robot which can self-propel in high-viscosity fluid when mimicking the asymmetric
84 beat pattern of the algal flagella, verifying that low-Reynolds number kinematics are
85 recapitulated. By arranging the robotic flagella in one of two possible configurations
86 (parallel or perpendicular) relative to a central "cell body", we imposed and tested three
87 distinct flagellar actuation patterns (gaits) that occur naturally in the algal flagellates,
88 namely the pronk, the trot, and the gallop. In each case, we compared the hydrodynamic
89 swimming performance of the robot to that of the corresponding algal species. Finally,

90 we discuss the relevance of these results for understanding how functional differences in
91 swimming performance may arise from morphologically similar structures, and highlight
92 the implications of this from an eco-evolutionary perspective.

93 **2. Methods**

94 *2.1. Microalgal culturing and imaging*

95 Three species of algae (*P. parkeae*, *P. tetraarhynchus* and *C. carteria*) were cultured
96 axenically according to previously published protocols [5, 18]. Free-swimming individuals
97 were tracked in open microfluidic chambers using a high-speed camera (Phantom Vision
98 Research). Brightfield imaging was conducted with 40x or 60x objectives using standard
99 inverted microscopes (Leica DMI8 and Nikon T2000-U) under white light illumination.
100 Free-swimming trajectories were obtained from high-speed videos in which single cells
101 crossed the focal plane, with the use of the open source software TrackMate (Fiji) [35].
102 Ten cells per species were used to determine the performance of each swimming gait
103 (Supplementary Video 1). Tracks in which cells performed transient gaits, tumbles, or
104 changed directions were not used in this analysis. The body length of each cell was
105 measured along the long axis (anterior-posterior) of the organism. An average body
106 length of $13.95 \pm 2.05 \mu\text{m}$, $12.54 \pm 0.65 \mu\text{m}$, and $12.82 \pm 0.72 \mu\text{m}$ was found for *P. parkeae*,
107 *P. tetraarhynchus* and *C. carteria* respectively.

108 *2.2. A self-powered roboflagellate*

109 We designed a dynamically-scaled robot to ensure that the robotic model can be self-
110 powered and does not require external fields - all controllers and servos are fully self-
111 contained (Fig.1C). We performed robotic experiments in a highly viscous fluid (mineral
112 oil, McMaster, 1000 cSt, product no. 1401K75) to match the low Reynolds number
113 regime experienced by the algae (Figure 1C). A subset of trials were conducted in
114 glycerin (vegetable glycerin, Blue Water Chem Group, product no. B07FQWDTH7)
115 of comparable viscosity to the mineral oil, to enable better visualisation and tracking
116 of appendage movement. Each robot consisted of a 3D printed body (length and width
117 = 8.5 cm, height = 2 cm) attached to four flagella that are independently actuated
118 by waterproof servo motors (Savox, product no. SW0250MG). Each appendage was
119 oriented such that the stroke lies in the plane perpendicular to the body (Figure 1D).
120 Foam (FOAMULAR Insulating Sheathing (IS) XPS Insulation) was attached on the
121 robot body to achieve neutral buoyancy and allow it to swim untethered. Commanded
122 appendage positions were achieved using a micro controller (Photon, Particle, part ID:
123 PHOTONH) that allowed us to actuate our robot with the use of Wi-Fi. Our micro-
124 controller and each motor were connected via a IOT Servo Shield (Actuonix, part ID:
125 IOT-SHIELD-PHOTON), a circuit board specific to our micro controller. Four LEDs
126 were placed on the 3D printed body to facilitate tracking. The robot was powered with
127 three lithium ion polymer batteries (3.7 V, 2500 mAh), each powering directly the micro
128 controller, the motors, and any attached LEDs. With this micro controller, the robots
129 were able to achieve self-propulsion.

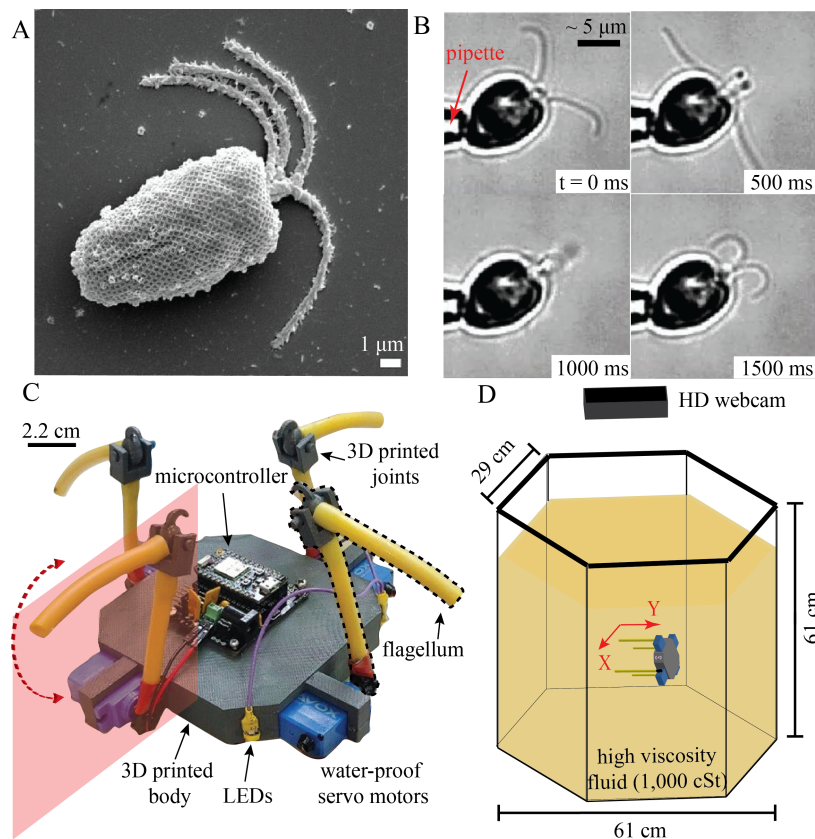


Figure 1: **Design and fabrication of a dynamically-scaled robophysical model of a microswimmer with four flagella.** (A) Scanning electron microscope (SEM) image of *Pyramimonas gelidicola* [36]. (B) Snapshots of *Pyramimonas parkeae* held by a pipette. (C) Robophysical model of quadriflagellate algae. (D) Experimental set-up. The arena is a hexagonal tank filled with mineral oil of high viscosity, to match the low Reynolds number regime experienced by the algae.

130 2.3. Actuation of robotic flagella

131 Each robotic appendage comprised a two-link flagellum (length = 6.5 cm, diameter =
132 3.1 mm, polypropylene-based thermoplastic elastomer (TPE)) connected by a 3D printed
133 joint allowing each flagellum to passively bend and break drag symmetry (Figure 1C).
134 Inspired by the flagellar beating waveform of the organisms, we implemented a simple
135 two-link flagellum in the robot that was able to break time-reversal symmetry (Figure 2).
136 An irreversible stroke pattern was achieved with the use of 3D printed hinges between
137 the two flagella segments. Instead of actively prescribing the shape of the flagella over
138 a beat cycle, symmetry breaking was achieved passively. No external control such as
139 magnetic fields were used, our robot was completely open loop. Each gait maintained a
140 constant phase difference between adjacent flagella set by prescribed joint angles of the
141 proximal segment (Figure 3, Supplementary Video 2). Each gait was uploaded to the
142 microcontroller via Wi-Fi, allowing the controllers to actuate the motors offline. Unless
143 otherwise specified, all gaits were prescribed with a flagellar beat frequency of 0.14 Hz.
144 The applied torque was constant throughout each beat cycle (3.5 kg/0.34 Nm operating
145 at 4.8 V).

146 For the movement of the robot in mineral oil (kinematic viscosity $\mu/\rho = 10 \text{ cm}^2/\text{s}$),

147 the Reynolds number (Re) 0.14 ($L = 3.8$ cm, $U = 0.38$ cm/s), while the oscillatory
148 Reynolds number Re^{osc} was 0.20 ($L = 3.8$ cm, $\omega = 0.14$ Hz). For the experiments
149 conducted in glycerin (kinematic viscosity $\mu/\rho = 11.83$ cm²/s), $Re = 0.23$ ($L = 6.89$ cm,
150 $U = 0.40$ cm/s), and $Re^{\text{osc}} = 0.55$ ($L = 6.89$ cm, $\omega = 0.14$ Hz).

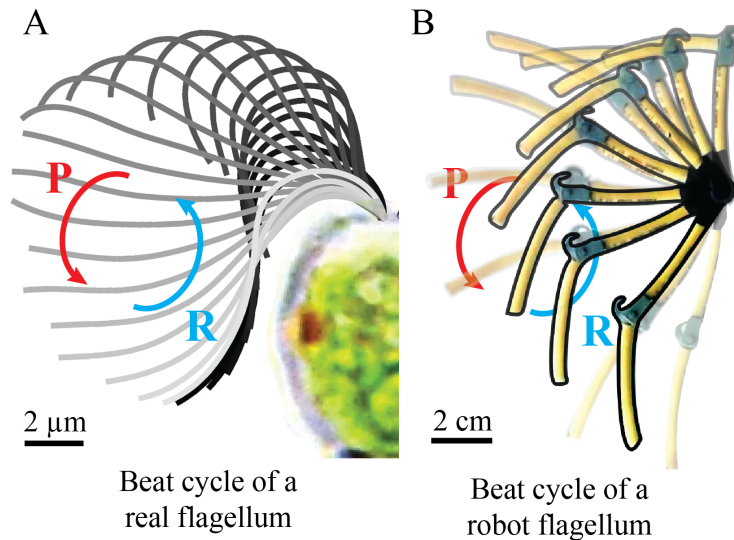


Figure 2: **Breaking time-reversal symmetry with a hinged two-link bio-inspired flagellum.** One beat cycle of an (A) algal flagellum compared to a (B) robot flagellum. P: power stroke, R: recovery stroke. Each robot flagellum segment has a length of 6.5 cm and diameter of 3.1 mm. Asymmetric beat patterns are achieved via a 3D printed joint. The movement patterns of the algal flagellum was measured in water, for the robot this was visualized in a high-viscosity fluid (glycerin).

151 2.4. Prescribing the swimming gait in the roboflagellate

152 We imposed three distinct gaits observed in quadriflagellate algae – the pronk, the trot,
153 and the gallop. The different coordination patterns were achieved by prescribing the
154 phase differences between adjacent appendages. The resulting gait sequences can be
155 confirmed for an immobilised robot body, where the distance from each flagellum tip
156 to the cell body was used as proxy for phase. In the pronk gait, all four appendages
157 move simultaneously, without any phase difference ($\varphi = 0^\circ$) between adjacent flagella
158 (Figure 3A). The trot gait is defined by alternating pairs of flagella each of which is
159 generating a pattern analogous to a breaststroke, with a phase difference of half a gait
160 cycle ($\varphi=180^\circ$) (Figure 3B). In the gallop gait, each appendage moves with a phase
161 difference of a quarter-gait cycle relative to its neighbour (Figure 3C). The directionality
162 (clockwise or counter-clockwise) of the gallop gait is determined by the phase difference
163 (φ) between the first appendage (m_1) and an adjacent appendage (m_2 or m_4). We tested
164 the gallop gait in both a clockwise ($\varphi=60^\circ$ between m_1 and m_2) and counter-clockwise
165 ($\varphi=180^\circ$ between m_1 and m_2) direction.

166 *2.5. Motion tracking*

167 Due to the of opaqueness of the oil, we attached LEDs to the robot's body to enable
168 motion tracking. Further LEDs were attached to the flagella. All LEDs were digitized
169 using custom MATLAB algorithms. We approximated the center of geometry of the robot
170 by averaging the position of the LEDs over time. Then, we used the tracks to determine
171 the distance traversed by the robot in units of body lengths per beat cycle. A total of
172 9 trials were taken per gait, for each robot configuration. A trial was terminated either
173 when the robot touches a boundary, or if the LEDs were no longer visible as the robot
174 sediments over time (this is due to the foam trapping fluid and increasing in mass). Thus,
175 each trial comprised 6 – 10 cycles per gait. To visualize and confirm movement of the
176 flagella during active swimming, we used glycerin as an alternative high viscosity fluid.
177 However, because glycerin is not a dielectric fluid, wi-fi connectivity was interrupted and
178 the circuits were negatively affected. To resolve this, we substituted our micro controller
179 (Pro Trinket, Adafruit, product ID: 2000) and sealed the circuits with a gasket and a 3D
180 printed cap.

181 **3. Results and Discussion**

182 *3.1. The trot is the fastest gait in the algae*

183 We identified the quadriflagellates as an ideal study group owing to their morphological
184 diversity (in size, shape, aspect-ratio), and abundance in marine, terrestrial as well as
185 freshwater habitats. A key trait distinguishing quadriflagellate genera is the arrangement
186 or insertion of flagella around the anterior of the cell [21, 37]. Here we take advantage
187 of this diversity to compare the swimming behaviour of three species (*Pyramimonas*
188 *tetrarhynchus*, *Pyramimonas parkeae*, and *Carteria crucifera*) that employ three distinct
189 gaits - respectively the pronk, the trot, and the gallop (Supplementary Video 1). We
190 conjecture that inter-species differences in quadriflagellate swimming performance can
191 be attributed to differences in gait alone - where the same basic stroke is applied to
192 ensembles of appendages but according to distinct phase relationships.

193 Two of these algae belong to the genus *Pyramimonas*, a Prasinophyte algae belonging
194 to an early diverging class which is thought to have given rise to the core Chlorophyte
195 algae, comprising species with two, four, eight, or up to sixteen flagella [38, 5]. Four
196 flagella of identical length and beat pattern emerge from a deep anterior groove or pit
197 in the cell body. The third species, *C. crucifera*, is a Volvocalean flagellate that is closely
198 related to the model biflagellate *Chlamydomonas*. Despite this phylogenetic divergence,
199 all three species are similar in body size and flagellar morphology (approx $\sim 20 \mu\text{m}$ in
200 length and $\sim 15 \mu\text{m}$ in width) and appear obovoid to cordate in side profile [39, 40].

201 In all three cases, cells swim smoothly flagella-first (puller-type) at speeds of $\mathcal{O}(100)$
202 $\mu\text{m}/\text{s}$. The translational motion is coupled to an axial rotation to produce swimming
203 along helical trajectories [41]. Abrupt gait transitions can occur either spontaneously
204 or when triggered by mechanical contact, during which the flagella are directed to the
205 front of the cell in a so-called shock-response [42]. Cells can also reversibly stop and start
206 swimming, when all or some of the flagella transiently cease to beat [18].

A minimal robophysical model of quadriflagellate self-propulsion

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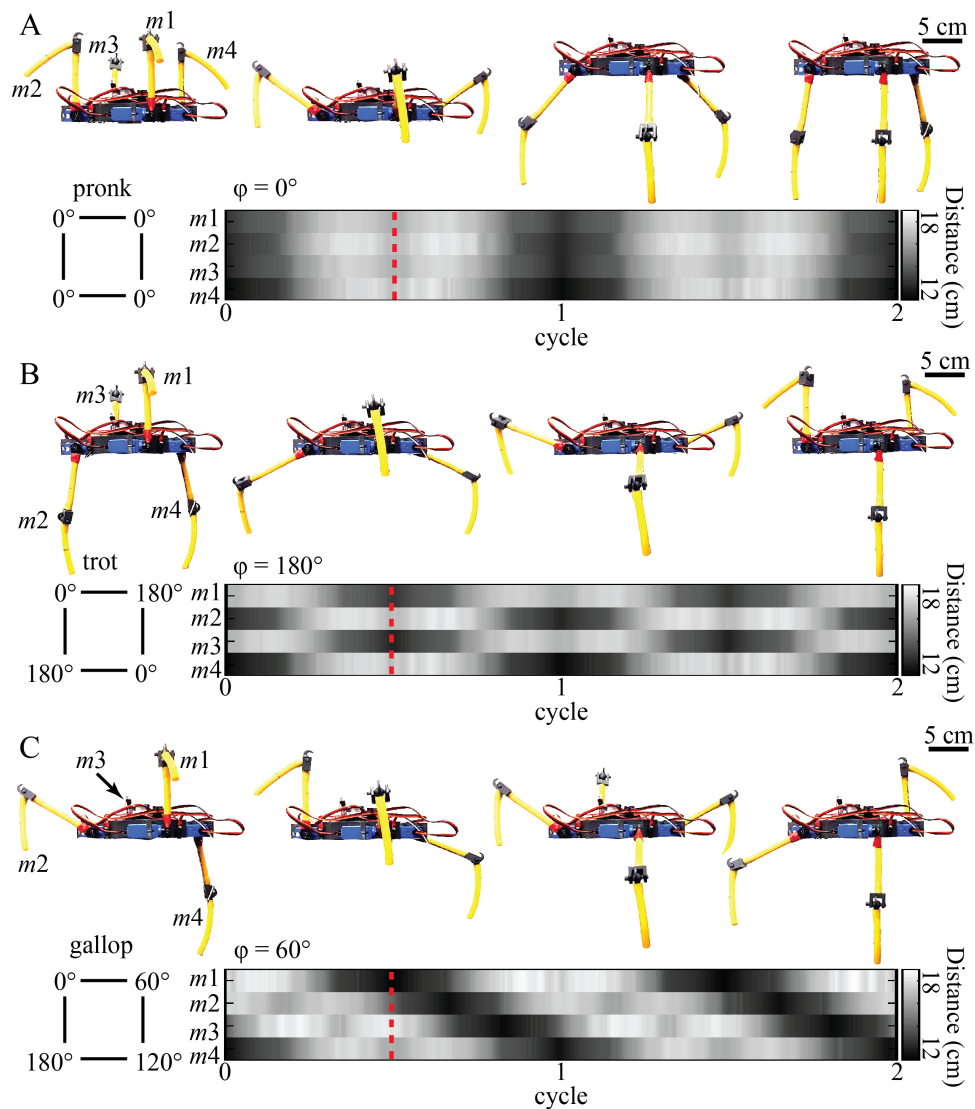


Figure 3: Quadriflagellate gaits prescribed to the robot. Distance from the center of geometry of the robot to the tip of each flagella was used as a proxy for the phase between adjacent flagella, labelled $m1-4$. (A) The pronk gait: zero phase difference ($\varphi=0^\circ$) between adjacent flagella. (B) The trot gait: alternating pairs of flagella with a phase difference of half a gait cycle ($\varphi=180^\circ$). (C) The gallop gait: adjacent flagella with a phase difference of a quarter of a gait cycle ($\varphi=60^\circ$). Snapshots of the robot showing the flagella configurations during each gait over half a gait cycle. The dashed red line delineates half a gait cycle from the start of the recording. [Note to visualise the gaits fully the robot was not placed in fluid.]

207 In all cases, free-swimming trajectories are superhelical, where small-scale swirls
 208 at the scale of single-cells are produced by the periodic flagellar oscillations. Three
 209 representative tracks, projected onto the focal plane, are shown in Figure 4 (A,C,E).
 210 Using the large-scale tracks, we estimated for each of the three gaits the displacement
 211 per cycle, including the cumulative displacement as a function of phase during the beat
 212 cycle (Figure 4G) as well as the mean forward progress per complete cycle (Figure 4H).
 213 Measured swimming speeds were $126 \pm 24 \mu\text{m/s}$ for the pronk, $408 \pm 46 \mu\text{m/s}$ for the trot,
 214 and $127 \pm 25 \mu\text{m/s}$ for the gallop. Our results show that the trot gait is the fastest gait
 215 in the microalgae. Meanwhile the pronk and gallop gaits lead to comparable propulsion

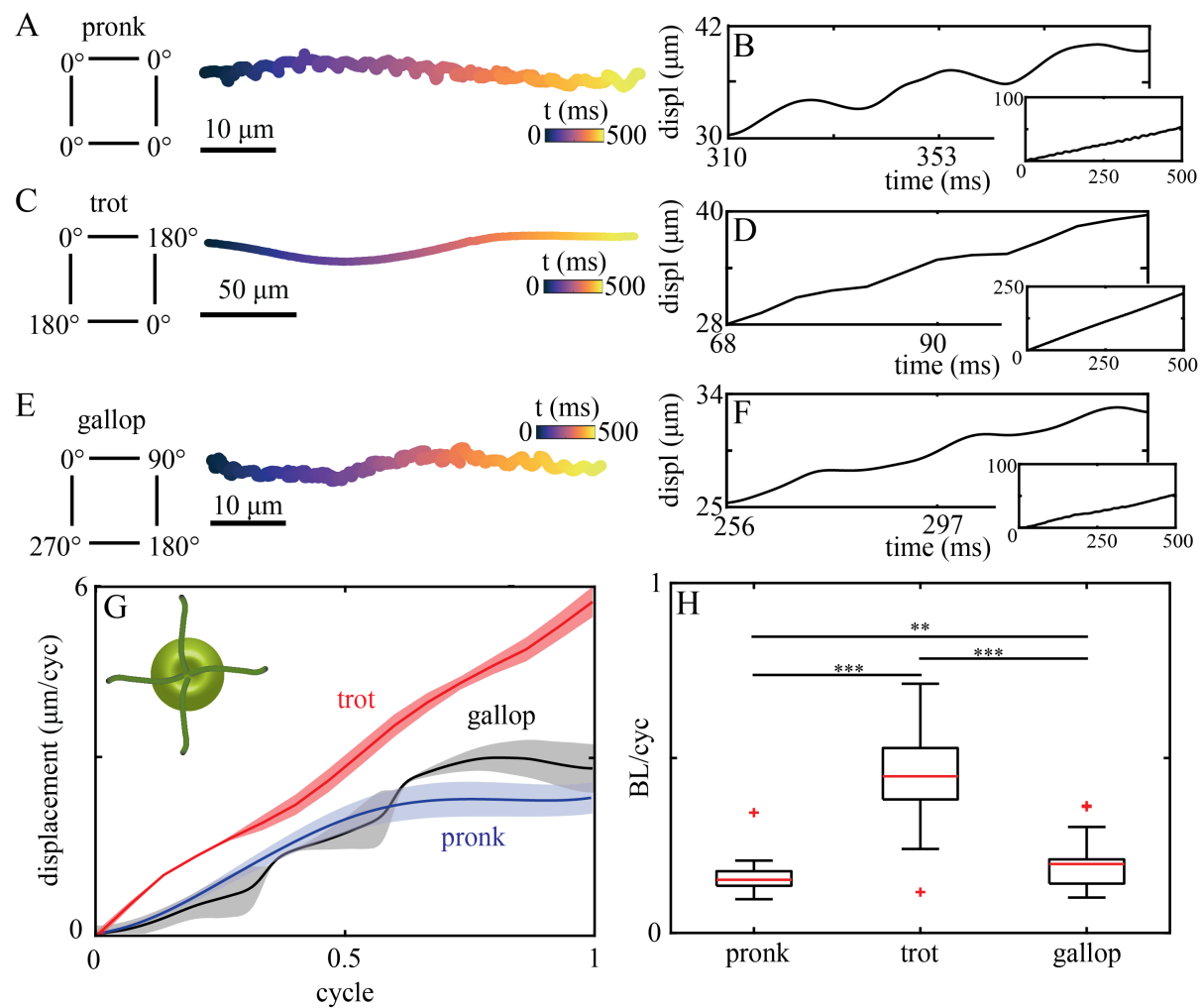


Figure 4: **Gaits, kinematics, and hydrodynamic performance of quadriflagellate algae.** All experiments were conducted in culture media - which had the same viscosity as water. For the pronking gait of *Pyramimonas tetrahynchus*: (A) a sample (cell-centroid) trajectory colored by time, and (B) forward displacement over time for three cycles. Inset shows forward displacement over time of trajectory. For the trotting gait of *Pyramimonas parkeae*: (C) a sample (cell-centroid) trajectory colored by time, and (D) forward displacement over time for three cycles. Inset shows forward displacement over time of trajectory. For the galloping gait of *Carteria crucifera*: (E) a sample (cell-centroid) trajectory colored by time, and (F) forward displacement over time for three cycles. Inset shows forward displacement over time of trajectory. (G) Mean displacement within a gait cycle for all gaits - the pronk (blue line), trot (red line), and gallop (black line). Shaded areas correspond to the standard deviation. (H) Mean displacement computed in terms of body lengths per cycle, for each gait.

216 speeds.

217 3.2. A hinged flagellum breaks time-reversal symmetry

218 We first confirmed that our system resides in a low-Reynolds number regime by attaching
 219 3D-printed rigid (unhinged) flagella to the body (Supplementary Video 3). As expected,
 220 reciprocal strokes produced negligible net swimming. The net displacement in the
 221 direction of movement after one complete cycle was 0.07 ± 0.05 cm (0.02 ± 0.01 BL) using
 222 the trot gait (Figure 5B).

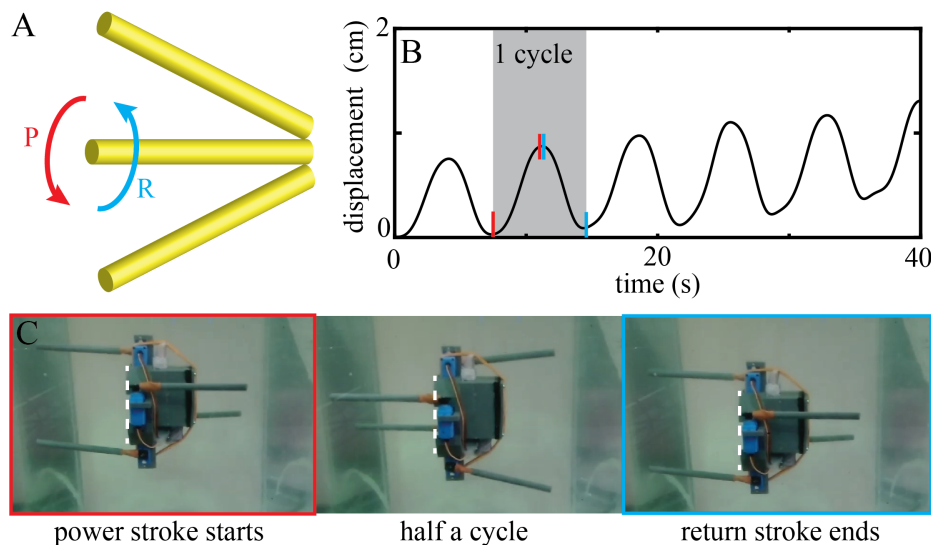


Figure 5: **Kinematic reversibility confirms low-Reynolds regime.** (A) One beat cycle of a single rigid flagellum moving back and forth. P: power stroke, R: recovery stroke. (B) Forward displacement traveled over time in mineral oil. Negligible net displacement per cycle with symmetric stroke pattern. (C) Snapshots of the robot during one gait cycle using rigid flagella. Left panel (outlined in red) shows the robot initiating a power stroke. Middle panel shows the robot during half a cycle. Right panel (outlined in blue) shows the end of the recovery stroke. [Note snapshots shown correspond to our alternative robot swimming in glycerin.]

223 With hinged flagella (Figure 2B), the robot became capable of net forward
224 propulsion. Each gait cycle can be characterized by a power stroke during which the
225 robot gains distance, and a recovery stroke during which it loses distance. We first set
226 out to test the effect of flagella ‘waveform’ on swimming performance, this is expected to
227 scale approximately with stroke amplitude [43, 44].

228 3.3. Flagellar undulation pattern affects swimming performance

229 We implemented two distinct flagellar undulation patterns - as defined by the maximal
230 sweep range of the segments. For simplicity and to prevent axial rotation, we then
231 reduced our quadriflagellate robot to a biflagellate robot, by removing one pair of
232 flagella (Supplementary Video 4). The remaining pair of flagella was programmed to
233 follow a breaststroke pattern (Figure 6A). We prescribed and compared the swimming
234 performance for two different sets of motor angles for the proximal segment: i) $[0^\circ,$
235 $180^\circ]$, and ii) $[45^\circ, 135^\circ]$ (Figure 6A inset). The motion of the distal segment always
236 follows passively, with the hinge breaking time-reversal symmetry. We tracked the
237 flagella ‘waveform’ in the two cases and calculated the angles generated by each flagellum
238 segment over time (from motor to joint and from joint to tip, Figure 6A). The two sweep
239 amplitudes produced two distinct gaits θ_1 - θ_2 shape space Figure 6B. A reduced sweep
240 range results in a higher beat frequency ($\omega = 0.14$ Hz for motor angles of $[0^\circ, 180^\circ]$, and
241 $\omega = 0.41$ Hz for motor angles of $[45^\circ, 135^\circ]$). The rescaled displacement shows swimming
242 performance increases with amplitude (Figure 6C). As predicted, the larger-amplitude
243 breaststroke achieves a greater displacement after each gait cycle, consistent with the
244 notion that non-inertial locomotion is dictated by simple geometric mechanics. Here,

245 movement is kinematic, and net displacement is determined largely by the gait and its
246 associated low-dimensional properties [45].

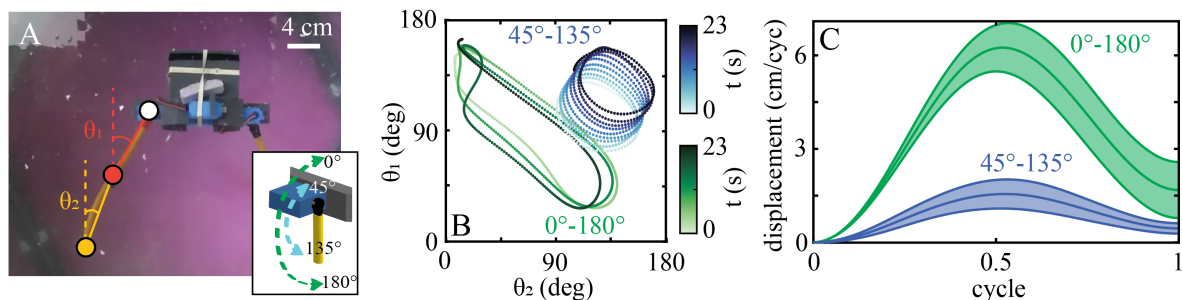


Figure 6: **Swimming performance increases with stroke amplitude.** (A) Quadriflagellate robot modified as a biflagellate robot, performing a breaststroke pattern with one pair of flagella. Angles θ_1 and θ_2 correspond to the angles generated by the flagella segment from the motor (white circle) to the joint (dark orange circle) and the segment from the joint (dark orange circle) to the tip (light orange circle). Inset shows variation of prescribed angles from 0° to 180° (green) and from 45° to 135° (blue). (B) θ_1 as a function of θ_2 , colored by time. Green dots corresponds to angles from 0° to 180° . Blue dots corresponds to 45° to 135° . (C) Mean displacement as a function of a gait cycle. Green line corresponds to angles from 0° to 180° . Blue line corresponds to 45° to 135° . Shaded areas correspond to the standard deviation. These experiments were conducted in glycerin with the alternative robot, to ensure the flagella beat pattern can be tracked.

247 3.4. Roboflagellate swimming performance depends on gait and appendage placement

248 To test if swimming performance is dominated by gait or by other factors such as flagellar
249 stiffness or shape dynamics, we prescribed the gaits exhibited by each algae species to our
250 roboflagellates. We explored the effect of varying appendage phase coordination (gait)
251 for two different configurations of four flagella, in which motors are positioned either in
252 a parallel or a perpendicular orientation with respect to an identical body.

253 These configurations are inspired by naturally-occurring arrangements of basal
254 bodies and flagella found in extant algal flagellates (Figure 7). All three species of
255 algae studied here correspond to configuration A, in which the approximate plane of
256 flagellar beating is perpendicular to the surface of the robot body. The main difference
257 is that when viewed from the anterior of the cell, the four flagella are inserted with a
258 clockwise twist or offset for *Carteria*, but an anticlockwise offset for *Pyramimonas* [37].
259 Algal species reported to exhibit configuration B [37] were not available in culture and
260 were not represented in the present study. Appendage coordination was prescribed in the
261 robot by specifying the phase differences between flagella, to produce each of the three
262 gaits: pronk, trot, or gallop, as previously described Figure 3.

263 For the perpendicular configuration, example trajectories as well as the cumulative
264 forward displacement over time for each gait as shown in Figures 8(B)-(C), (E)-(J). We
265 also analyzed the detailed within-cycle dynamics for each gait (Supplementary Video
266 5). The pronk and both the clockwise (CW) gallop and counterclockwise (CCW) gaits
267 produce significant forward displacement during the power stroke (up to 5.7 cm for
268 the pronk, 4 cm and 2 cm for the CW and CCW gallops respectively after 1/2 a
269 gait cycle), but also produce a significant backward displacement during the recovery

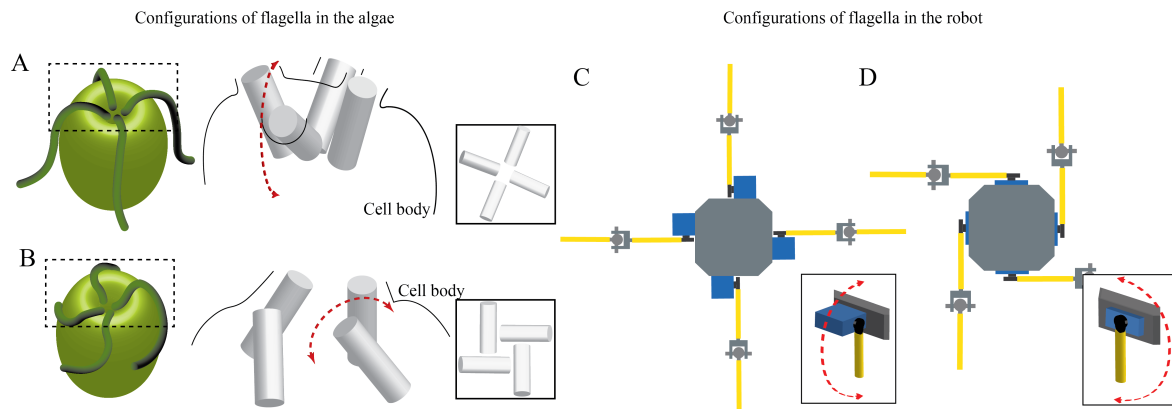


Figure 7: **Modelling appendage placement.** (A) Illustration of two configurations of flagella and basal bodies that are found in quadriflagellates [37]. The flagella emerge from basal bodies (cylinders) that are oriented largely perpendicular (A), or parallel (B) to the cell body. Insets show anterior views (A: cruciate arrangement, B: turbine or windmill-like). Double-arrow indicates the approximate beat plane of the individual flagella. Similarly, two roboflagellate designs are presented. Motors and attached ‘flagella’ are oriented perpendicular (C) or parallel (D) to the central body. Again, double-arrow indicates oscillation plane.

270 stroke, generating overall small displacement from cycle to cycle (0.33 ± 0.04 BL/cyc,
271 0.16 ± 0.05 BL/cyc, and 0.15 ± 0.08 BL/cyc for the pronk, the CW gallop, and CCW gallop
272 respectively). On the other hand, while the trot does not achieve a greater displacement
273 (only 2.3 cm after 1/2 a gait cycle) than the pronk or gallop during the power stroke, it
274 loses a much smaller distance during the recovery stroke. This is because while one pair
275 flagella is moving towards the body and consequently producing backward motion, the
276 other pair of flagella moves away from the body so as to resist this motion. This can also
277 be observed in the trajectories, where the pronk and gallop gait shows backward motion,
278 unlike the trot gait. Due to this, of the three gaits investigated the robot achieves the
279 greatest hydrodynamic performance (0.6 ± 0.08 BL/cyc) using the trot gait (Figure 8K),
280 just as in the algae.

281 For the parallel configuration (Supplementary Video 6), example trajectories as well
282 as the forward displacement over time for each gait can be seen in Figure 9(B)-(C), (E)-
283 (J). Similar to the perpendicular configuration, the pronk gait allows the robot to gain a
284 significant amount of distance during the power stroke (up to 5 cm after 1/2 a gait cycle)
285 but also lose a significant amount of distance during the recovery stroke, generating little
286 net displacement from cycle to cycle. The gallop gait in the counterclockwise displays a
287 similar oscillatory pattern, however there is a discrepancy between the counterclockwise
288 and clockwise gallops (5.3 cm after 1/2 a gait cycle for the CW gallop, but only 1 cm
289 for the CCW gallop). This is likely due to rotation-translation coupling in the second
290 configuration (in which the flagella are inserted in the CCW sense), generating significant
291 motion laterally and causing axial rotation of the robot. Similar to the perpendicular
292 robot, the trot gait gains less distance during the power stroke (only 1.5 cm after 1/2 a gait
293 cycle) and loses more distance during the recovery stroke, relative to the perpendicular
294 configuration. The phasing between appendages in the trot gait again aids the robot
295 in traversing a greater distance from cycle to cycle than the pronk (0.15 ± 0.4 BL/cyc),

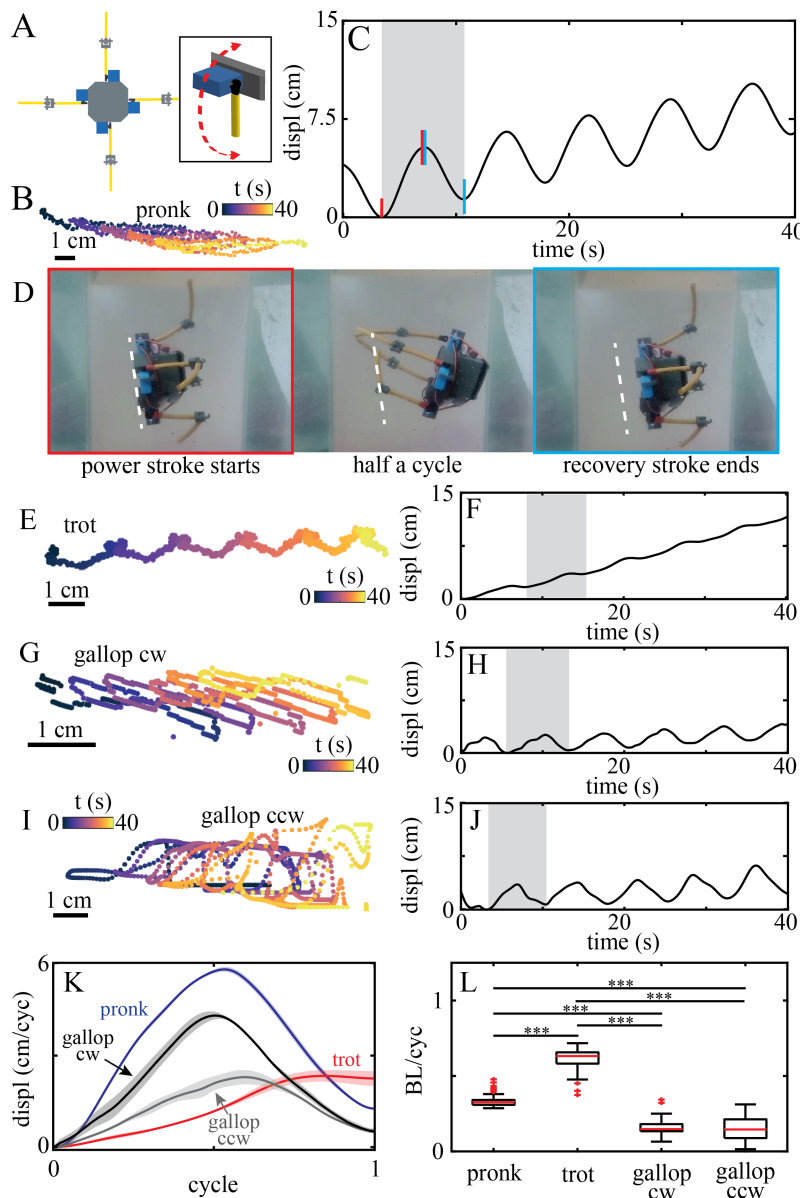


Figure 8: Swimming gait kinematics and performance for robot with flagella in perpendicular orientation. (A) Diagram of robot with motors oriented perpendicular to the body. Inset illustrates beating plane. For the pronk gait, (B) shows a sample trajectory of the robot, colored by time (5 cycles), and (C) the forward displacement traveled over time. For one gait cycle, red vertical lines highlight power stroke, and blue vertical lines highlight return stroke. (D) Snapshots of the robot during one cycle of the pronk gait. Left panel (outlined in red) shows the robot initiating a power stroke. Middle panel shows the robot during half a cycle. Right panel (outlined in blue) shows the robot completing the recovery stroke. (Arrow: swimming direction.) Trajectory of the robot during the trot gait, colored by time (5 cycles) (E), and forward displacement traveled over time of the robot during the trot gait (F). Trajectory of the robot during the clockwise gallop gait, colored by time (5 cycles) (G), and forward displacement traveled over time of the robot during the trot gait (H). Trajectory of the robot during the counter-clockwise gallop gait, colored by time (5 cycles) (I), and forward displacement traveled over time of the robot during the trot gait (J). (K) Mean displacement over a gait cycle for all gaits - the pronk (blue line), trot (red line), clockwise gallop (black line), and counter-clockwise gallop (grey line). Shaded areas correspond to the standard deviation. (L) Body length per cycle as a function of swimming gait.

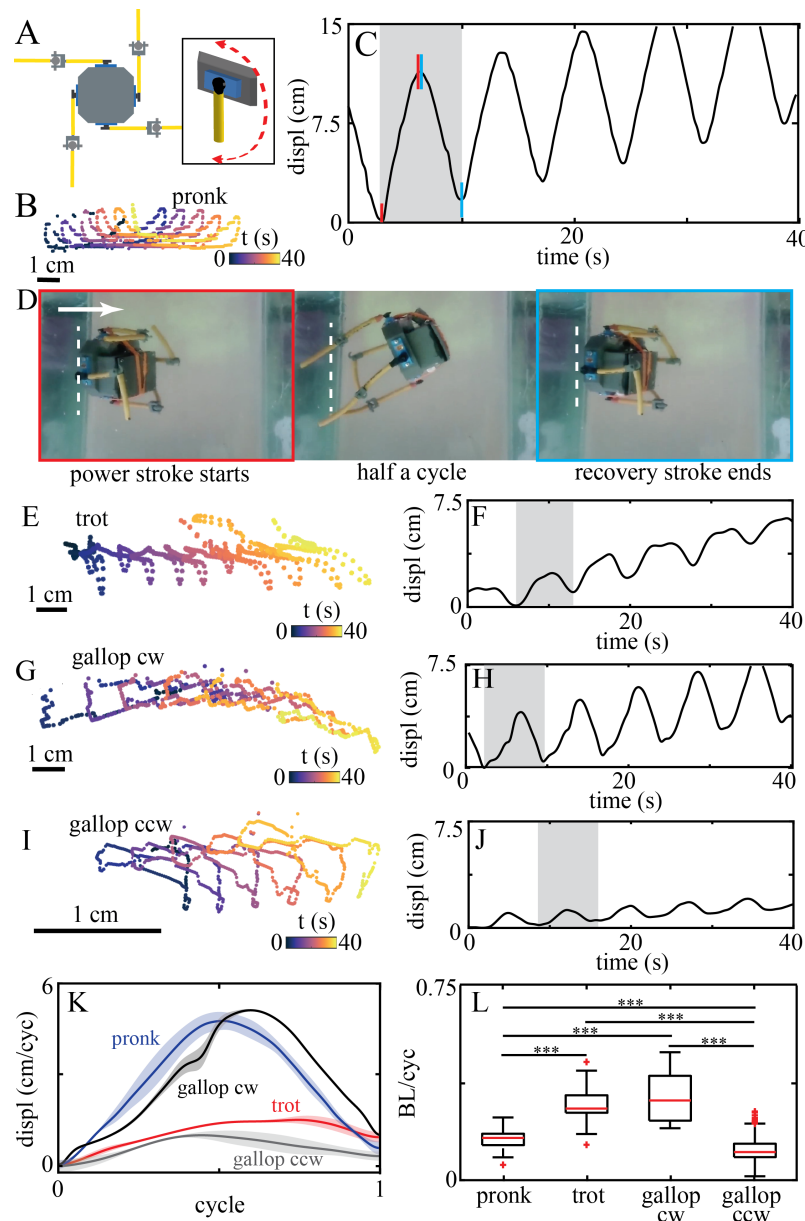


Figure 9: **Swimming gait kinematics and performance for robot with flagella in parallel orientation.** (A) Diagram of robot with motors oriented perpendicular to the body. Inset illustrates beating plane. For the pronk gait, (B) shows a sample trajectory of the robot, colored by time (5 cycles), and (C) the forward displacement traveled over time. For one gait cycle, red vertical lines highlight power stroke, and blue vertical lines highlight return stroke. (D) Snapshots of the robot during one cycle of the pronk gait. Left panel (outlined in red) shows the robot initiating a power stroke. Middle panel shows the robot during half a cycle. Right panel (outlined in blue) shows the robot completing the recovery stroke. (Arrow: swimming direction.) Trajectory of the robot during the trot gait, colored by time (5 cycles) (E), and forward displacement traveled over time of the robot during the trot gait (F). Trajectory of the robot during the clockwise gallop gait, colored by time (5 cycles) (G), and forward displacement traveled over time of the robot during the trot gait (H). Trajectory of the robot during the counter-clockwise gallop gait, colored by time (5 cycles) (I), and forward displacement traveled over time of the robot during the trot gait (J). (K) Mean displacement over a gait cycle for all gaits - the pronk (blue line), trot (red line), clockwise gallop (black line), and counter-clockwise gallop (grey line). Shaded areas correspond to the standard deviation. (L) Body length per cycle as a function of swimming gait.

296 and also greater than the average of the CW and CCW gallop gaits (0.15 ± 0.9 BL/cyc).
 297 (We assume that by symmetry, this average between the two chiralities should cancel
 298 any rotational effects.) Thus, the trot remains a hydrodynamically effective gait for the
 299 parallel robot (0.26 ± 0.08 BL/cyc) (Figure 9K).

300 We conclude that the swimming performance of the roboflagellate is highly sensitive
 301 to both gait and flagellar orientation (which defines the principal beat plane) of the
 302 flagella. It is possible that the organisms can access different regimes by controlling the
 303 3D beat plane of their flagella, and that divergent flagellar placement evolved in different
 304 species as a result of different environmental selection pressures. In flagellates such as
 305 *Volvox*, nearby basal bodies (from which the flagella emerge) have rotated 90 degrees
 306 compared to the ancestral configuration found in the unicellular *Chlamydomonas*, likely
 307 to facilitate coordinated flagellar beating as an intact colony [5, 46].

308 *3.5. Speed of roboflagellate is comparable to that of real algae*

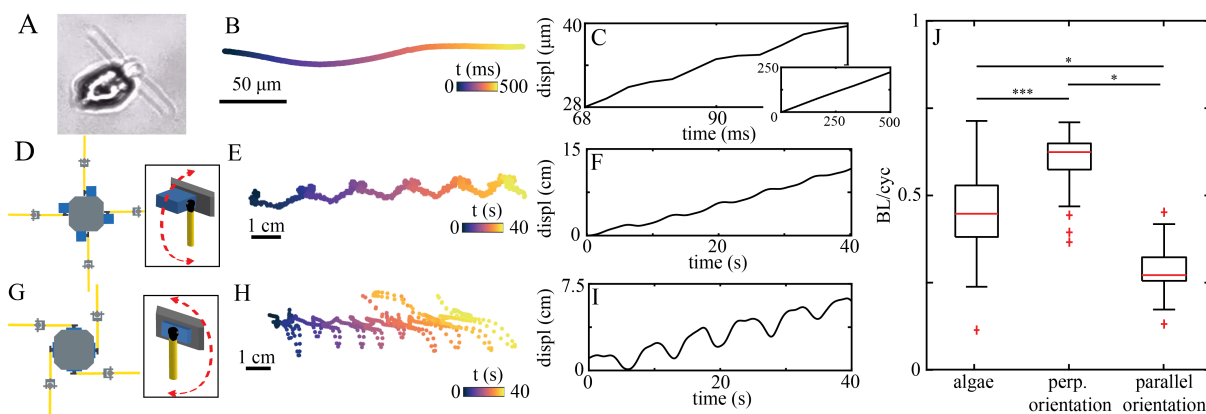


Figure 10: **Comparing the trot gait in the algae and robot.** (A) The alga *Pyramimonas parkeae* swimming using the trot gait. (B) Trajectory of *P. parkeae*, colored by time. (C) Forward displacement traveled over time by *P. parkeae*. (D) Diagram of robot with motors oriented perpendicular to the body. Inset illustrates beating plane. (E) Trajectory of the robot with perpendicular configuration using the trot gait, colored by time. (F) Forward displacement traveled over time of the robot with perpendicular configuration using the trot gait. (G) Diagram of robot with motors oriented parallel to the body. Inset illustrates beating plane. (H) Trajectory of the robot with parallel configuration using the trot gait, colored by time. (I) Forward displacement traveled over time of the robot with parallel configuration using the trot gait. (J) Body length per cycle for the trot gait for the algae, the perpendicular configuration, and the parallel configuration.

309 The above results show that a change in flagellar configuration can significantly
 310 change the performance of a given swimming gait. Focusing only on the trot, we note
 311 that the trot gait yielded the highest hydrodynamic performance for the algae and for
 312 the perpendicular robot, (Figure 10).

313 In both robot configurations, significant axial rotation and lateral movement was
 314 observed in the free-swimming trajectories (Figure 10E,H) showing that our robotic
 315 models do not swim as smoothly as their algal counterparts (Figure 10B). This
 316 suggests that the algal cytoskeleton could play a role in gait stabilization. The
 317 cumulative displacement over time for a trotting cell and our perpendicular robot

318 are comparable (Figure 10C,F). Meanwhile the parallel configuration displays larger
319 amplitude oscillations in which a greater distance gained during each the power stroke is
320 negated during the subsequent recovery stroke (Figure 10I). This is likely due to three-
321 dimensional effects as mentioned above. In all, we find that the performance of the algae
322 and both roboflagellate configurations are comparable in absolute terms, as measured
323 in terms of body lengths per stroke cycle. This agreement is surprising as we did not
324 precisely match the dimensions of our robots to that of the algal cell, and unlike the algal
325 flagella the robot ‘flagella’ were not capable of active bending - being comprised only of
326 rigid tubing and a 3D-printed hinge.

327 4. Conclusion and future work

328 Microscopic organisms have evolved to harness many different ways of swimming at low-
329 Reynolds number. Despite their apparent simplicity and lack of centralized control,
330 many species of single-celled algal multiflagellates can perform robust free-swimming gaits
331 analogous to animal gaits, which require specific temporal ordering of a small network
332 of flagella [18]. Since structural and genetic information about many of the species of
333 interest is lacking, we turned to robophysical modelling to understand how motility and
334 gait are controlled in these unicellular microswimmers.

335 Robotics approaches at the macroscale have long been used successfully to explore
336 and validate fluid dynamical theories of self-propulsion [47, 48]. Scenarios can be tested
337 in robots that may not be possible in the live organism [49, 50]. Dynamically-scaled
338 robotic models of microswimmers operating in viscous media have been used to mimic
339 bacterial swimming [51], to examine flows induced by bundles of rotating flagella [52],
340 and to reveal the role of elasticity [53], or to investigate metachronal actuation of rigid
341 appendages [54].

342 Here, we created for the first time a novel self-powered, untethered robot (no external
343 forces or torques) modelled upon quadriflagellate algae, which could not only self-propel
344 at low-Reynolds number but also recapitulate gait-dependent differences in swimming
345 performance that were observed in different species of microalgae. These results reveal
346 that differences in phase coordination of propulsive appendages alone has a significant
347 impact on hydrodynamic performance. The orientation of moving appendages on the
348 propelling body also influences net propulsive speed. In the perpendicular configuration
349 that best matches the algae, the trot gait is consistently faster than either the pronk
350 or gallop gait, and that the net displacement achieved by the robot in terms of body
351 lengths per cycle is similar in absolute terms to the algae. Thus our dynamically-scaled
352 robophysical model is a good model of the biological microswimmer.

353 Our work raises open questions about why the different quadriflagellate species
354 have distinct motility repertoires in the first place. Freely-locomoting organisms at all
355 scales, switch dynamically between multiple gaits, [55, 56], e.g. to escape predation.
356 Even bacteria motility exhibits strong heterogeneity across species [3]. While marine
357 *Pyramimonas* species exhibit sporadic bursts of fast activity with extended quiescent
358 phases [42], freshwater alga *Carteria* (closely related to *Chlamydomonas*) and other
359 *Volvoclean* algae do not show such rest periods [18]. More generally, ciliary strokes

360 that are optimised for swimming may not be optimised for other tasks such as feeding
361 or fluid pumping, and vice versa [57, 58]. Distinct gaits are unlikely to have evolved to
362 achieve ever-faster swimming, but rather reflects a more nuanced relationship between
363 the organism's metabolic requirements and its habitat.

364 We conjecture that differences in gait confers an evolutionary advantage even at
365 the microscale. Of the three algae studied here, two (*P. tetrarhynchus*, *C. carteria*)
366 are freshwater species and one is a marine species (*P. parkeae*). *P. tetrarhynchus* (type
367 species) was originally isolated from a peaty pool and cultured in a biphasic soil medium
368 [39]. *C. crucifera* is also a freshwater species that forms surface associations with leaves
369 and other decaying material. In contrast *P. parkeae* is most abundant in Arctic surface
370 water and in tidal rock pools, where it can access sufficient sunlight for photosynthesis.
371 *P. parkeae* also exhibits a unique diurnal vertical settling behaviour [59]. The latter
372 behaviour, along with phototaxis, accentuates the requirement for vigorous swimming
373 and hence the fast trot gait. Field data has shown that marine *Pyramimonas* routinely
374 blooms in and around sea ice, where the unique polar environment (extreme fluctuations
375 in temperature, light, salinity etc) is associated with a highly heterogeneous distribution
376 of different *Pyramimonas* species even within the same water column [60]. The habitats of
377 these algae may therefore be a key evolutionary driver leading to significant diversification
378 of gait, even across species with apparently convergent morphology and size [61, 62].
379 Further experiments using both lab strains and wild isolates, controlling more precisely for
380 culturing medium, are need to test this hypothesis. In parallel, we will use roboflagellates
381 to explore mix-mode propulsion strategies and unsteady effects, such as nutrient dispersal.

382 We highlight two limitations of the current design. The first concerns boundary
383 and finite-size effects, particularly due to fluid-structure interactions between moving
384 appendages and the bounding tank, and between different parts of the robot. The
385 presence of no-slip boundaries will alter the flow fields around a beating appendage,
386 and change propulsion efficiency [63]. The rigid insertion of the robot flagella around the
387 central body likely introduced an additional (unwanted) rotational movement. Second,
388 the current robot relies on a simple 2-link flagellum facilitated by a rigid 3D printed
389 joint which has a very limited number of degrees of freedom. The rigid joints have
390 limited ability to resist torsion - which may be gait-dependent. The compliance of real
391 (eukaryotic) flagella and cilia, which deform actively by distributed motor elements,
392 also improves propulsive force generation and efficacy. These organelles can actively
393 maintain their shape even when subject to significant hydrodynamic forces. In future
394 work we will resolve these limitations with improved roboflagellate designs, in parallel
395 with hydrodynamic simulations and modelling to understand gait optimisation.

396 In conclusion, we have presented a macroscopic robot capable of self-propulsion
397 at low-Reynolds number, and used this to model aspects of microorganism swimming
398 behaviour. We have applied this physical model to different permutations of gait
399 coordination patterns and flagellar placement and their influence on hydrodynamic
400 swimming performance. This approach has transformative potential for testing new
401 hypotheses relating to low Reynolds number self-propulsion in other small-scale biological
402 systems, such as mechanisms of gait selection and stimulus-dependent steering [16]. These
403 insights could have profound implications for how morphological computation may be

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404 achieved in a neural or early nervous systems. From a technological perspective, these
405 diverse propulsion strategies can provide unique, innovative solutions to the formidable
406 challenge of navigating viscous fluids.

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