Single-molecule imaging of cytoplasmic dynein *in cellulo* reveals the mechanism of motor activation and cargo capture

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Cytoplasmic dynein 1 (dynein) is the primary minus end-directed motor protein in most eukaryotic cells. Dynein remains in an inactive conformation until the formation of a tripartite complex comprising dynein, its regulator dynactin and a cargo adaptor. How this process of dynein activation occurs is unclear, since it entails the formation of a three-protein complex inside the crowded environs of a cell. Here, we employed live-cell, single-molecule imaging to visualise and track fluorescently tagged dynein. First, we observed that only ~30% of dynein molecules that bound to the microtubule (MT) engaged in minus end-directed movement, and that too for a short duration of ~0.6 s. Next, using high-resolution imaging in live and fixed cells, and using correlative light and electron microscopy, we discovered that dynactin remained persistently attached to MTs, and endosomal cargo remained in proximity to the MTs and dynactin. Finally, we employed two-colour imaging to visualise cargo movement effected by single motor binding. We then used these discoveries as the basis for a stochastic model incorporating dynamic motors binding to cargo located along MTs, and also developed a coarse-grained 3-state run-and-tumble particle (RTP) model for the cargo that quantitatively recapitulates the emergent statistics of cargo movement. Taken together, we discovered a search mechanism that is facilitated by dynein's frequent MT binding-unbinding kinetics: (1) in a futile event when dynein does not encounter cargo anchored in proximity to the MT, dynein dissociates and diffuses into the cytoplasm, (2) when dynein encounters cargo and dynactin upon MT-binding, it moves cargo in a short run. Several of these short runs are undertaken in succession for long-range directed movement. In conclusion, we demonstrate that dynein activation and cargo capture are coupled in a step that relies on the reduction of dimensionality to enable minus end-directed transport *in cellulo*, and that complex cargo behaviour emerges from stochastic motor-cargo interactions.

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Introduction

The dynein family of motor proteins comprises axonemal and cytoplasmic dyneins. While cytoplasmic dynein 2 plays a critical role in intraflagellar transport (1), cytoplasmic dynein 1 (dynein henceforth) is responsible for force production and minus-end directed movement of a variety of cargo in cells containing MTs (2). Dynein is a large complex of homodimers comprising 500kDa heavy chains, and other accessory proteins including light chains, light intermediate chains and intermediate chains, which mediate dimerisation of dynein and thereby its processivity.

The activity of motor proteins is typically regulated: several kinesins assume an autoinhibited conformation until attachment to cargo (3); dynein was first found to be regulated for its processivity by the multisubunit complex, dynactin (4). More recent studies have additionally implicated cargo adaptors - which link dynein to a multitude of cargo - in the activation of dynein (5–8). Dynactin is a large multi subunit complex (9) that was first identified as an activator of minus-end directed motility of vesicles (10). Further research indicated that an intact dynactin complex was necessary for dynein's function and that dynactin could interact with MTs via its p150 subunit (11, 12). The N-terminal CAP-Gly domain on p150 was also found to be able to interact with growing MT plus ends via EB1/CLIP-170 pathway (13, 14) and influence intracellular transport (15). However, the interaction between dynein and dynactin, as probed from co-immunoprecipitation assays, was observed to be weak, and over expressing the N-terminal (cytoplasmic) fragment of the cargo adaptor BiCD2 was found to significantly increase dynein-dynactin interaction (16).

Cargo adaptors are proteins with transmembrane and cytoplasmic domains that link membranous cargo to the motor (17), with specific cargo adaptors being unique to different types of cargo (18, 19). Recent single molecule *in vitro* research has established that formation of the dynein-dynactin-cargo adaptor (DDC) complex is essential for processive motion (20, 21). Cryo-EM studies later revealed that the formation of the DDC complex relieved the autoinhibition of dynein, and reoriented the motor for processive movement (22–24). The DDC complex ameliorates the force produced by single dynein motors from 1pN to about 6pN (25). Thus our current understanding suggests that formation of DDC complex is essential in dynein driven transport.

While it is clear that the tripartite complex formation is an essential first step in the activation of the dynein motor,

Fig. 1. Visualisation of single molecules of dynein in living cells. a, Schematic of the protocol followed for the visualisation of single molecules of dynein. **b**, Montage of HILO images showing representative binding and unbinding events of a single fluorescent mDHC molecule ('(binding)' and '(unbinding)'). The single molecule is indicated with yellow arrowheads for the duration of the time it remains bound in the field of view. Time is indicated at the top right of each image in the montage. **c**, Intensity histogram of single molecules of dynein with the Gaussian fits (grey line). The mean ± s.d. of the Gaussian distributions is indicated above the fits. **d**, HILO image (left) and kymograph (right) of a cell expressing mDHC-GFP. Representative stationary, minus end-directed and plus end-directed events are indicated with the white, teal and magenta arrowheads respectively in the kymograph. **e**, Plot of position vs. time for the single-molecule events tracked, showing stationary events (grey), minus end-directed events (teal) and plus end-directed events (magenta). **f**, Histogram with the residence time of dynein on the MT on the x axis and $P(\tau)$=1-cumulative frequency on the y axis. The exponential fit (grey line) gave a mean residence time $\tau_{\mathrm{res}} = 1/r_{\mathrm{off}}^{D} \sim 0.55$ s. In **b** and **d**, 'N' marks the location/direction of the nucleus.

it is unknown how this process occurs in a living cell, amid its crowded environs and the independent dynamics of each component of the tripartite complex. Here, we employ several strategies including single-molecule imaging, correlative light and electron microscopy (CLEM), high-resolution fluorescence microscopy and stochastic modelling to reveal the mechanism of formation of the tripartite complex and therefore, activation of dynein.

Results and Discussion

Dynein interacts transiently with the MT. To visualise the behaviour of dynein in HeLa cells expressing mouse DYNH1C1-GFP (mDHC-GFP, (26)), we adapted and optimised highly inclined and laminated optical sheet (HILO) microscopy ((27), Fig. S1a). When cells expressing low levels of mDHC-GFP were observed under a spinning disk confocal (SD) microscope, the fluorescence signal appeared cytosolic, with no discernible dynein punctae (Fig. S1b). However, when the same cells were observed using our modified HILO microscopy, distinct fluorescent spots were visible.

We adapted our microscopy protocol to obscure dynein diffusing in the cytoplasm and to only observe dynein that was bound to and resided on the MT ((28, 29), Fig. 1a, see Methods). We observed that dynein spots appeared afresh

and remained in the field of imaging for a short duration (Fig. 1b, Video S1). We intuited that these corresponded to events where single dynein molecules previously diffusing in the cytoplasm bound to MTs (Fig. S1c-e).

To confirm that the appearance of fluorescent signal on the MT corresponded to binding of a single molecule of dynein to the MT, we analysed the intensity of these fluorescent spots (Fig.1c). For single dynein molecules, we would expect the intensity histogram to fit to a sum of two Gaussian distributions, one corresponding to a GFP fluorescing from one DHC and the other corresponding to two GFPs fluorescing from both DHCs in the genetic background of these cells. The former primarily arose due to photobleaching of GFP during the course of imaging. Accordingly, the intensity histogram of these fluorescent spots revealed that we indeed observed single dynein molecules since the intensity histogram fit best to a sum of two Gaussians, with the mean of the first Gaussian profile being half that of the second (Fig. 1c).

Next, we analysed the two-dimensional position (x and y) versus time (t) of single molecules of dynein that bound afresh from the cytoplasm to the MT (Fig. 1d). Based on automated thresholding (see Methods), we classified the tracks as stationary, minus end-directed and plus end-directed. The plus ends of the MT were predominantly at the periphery in

Fig. 2. The dynactin complex binds along the entire length of the MT. a, Immunofluorescence image of MT (magenta) and p150 (green) obtained using SD microscopy + SRRF. **b**, Enlarged view of the area marked with the white rectangle in **a** and the line profile of p150 intensity along the length of a representative MT from the plus end ('P') to ~6 μm from the plus end of the MT ('0'). **c**, Immunofluorescence image of p62 (green) and p150 (magenta) obtained using SD microscopy + SRRF. The white arrowheads indicate representative p150 spots that also contain p62. **d**, Immunofluorescence image of MT (magenta) and p62 (green) obtained using SD microscopy + SRRF. The white arrowheads indicate representative p62 spots that occur on the MT. **e**, Histogram of the probability of cooccurrence of p62 with p150, indicating a high likelihood of presence of the entire complex at a p150 spot. **f**, Histogram of the probability of cooccurrence of p62 on the MT, which points to a high likelihood for the presence of the entire dynactin complex on the MT. In **a**, **c** and **d**, 'N' marks the location/direction of the nucleus.

these elongated cells (Fig. S1f), and hence we annotated movement towards the cell center as minus end-directed, and movement away as plus end-directed. We observed that ~50% of all the dynein molecules tracked (n=177/329, N=3 independent experiments from >50 cells) remained stationary upon MT binding, while ~30% (n=95/329) moved towards the minus end (Fig. 1e). The remaining ~20% moved towards the plus end and these arose likely due to attachment of dynein to cargo being moved to the plus end by kinesins (Fig. 1e). The velocity measured for minus-end directed movement of single dyneins was 1.2±0.7 μm/s (mean±s.d.), similar to values reported for mammalian dynein previously (30, 31). We also confirmed that the underlying MT was stable, did not undergo sliding, and therefore did not contribute to the dynein behaviour we observed (Fig. S1g).

We then measured the mean residence time (τ_{res}) of dynein on the MTs to be 0.55 s (95% confidence interval (CI): 0.51-0.59 s) (Fig. 1f). The unbinding rate of dynein from MT $r_{\text{off}}^{D} = 1/\tau_{\text{res}}$ was 1.7 s^{-1}, which is similar to the previously reported unbinding rate of single dyneins (32). We verified that this short residence time of dynein on MTs was a true representation of the duration of time that dynein remained attached to the MT and not convolved by GFP's photobleaching time (Fig. S1h and i). Further, by knocking down endogenous HeLa DYNC1H1 (hDHC), we verified that our observations were not an artefact of expression of mDHC-GFP in this background (Fig. S2a-h). To the best of our knowledge, these are the first observations of single molecules of dynein

in mammalian cells and indicate that dynein likely exists in an inactive state inside the cell, similar to reports from *in vitro* studies (16, 20, 21, 23).

The short residence time of dynein that we observed is in contrast with the findings from previous research (20, 21) which report a run length of upto 8 μm for dynein. However, other *in vitro* results have reported dynein run lengths that are comparable with our results (4, 33). It is possible that longer run lengths were observed when dynein formed the DDC complex *in vitro* due to the use of the N-terminal fragment of BiCD2 (BiCD2-N), as opposed to the autoinhibited, full length BiCD2. The former binds more strongly to dynein-dynactin. Moreover, the molar excess of adaptors used, the ionic strength of the buffers, and the source and modifications on the MTs all play a role in motor run length and, therefore, are not directly comparable to *in vivo* results.

Dynactin remains persistently associated with MTs.
Next, we aimed to visualise the dynamics of dynactin, the second player in the tripartite complex. Dynactin was first identified as a complex that was required for dynein-driven motility of vesicles *in vitro* (10). Several recent pieces of research have identified dynactin as an essential part of the active dynein complex (4, 24, 34). Dynactin is a multi-subunit complex which binds to MTs independently of dynein via its N-terminal p150 subunit (9). However, dynactin is unlikely to interact with dynein in the absence of the cargo adaptor (16, 20, 21, 23).

Fig. 3. Endosomes remain close to MTs. a, SD + SRRF image of MTs (green) and dextran vesicles (magenta) in live cells. The white arrowheads indicate representative vesicles on the MT. **b**, SD + SRRF image of p62 (top left), dextran vesicles (top right) in live cells and their merge (bottom, p62 in green and dextran in magenta). The white arrowheads in the merged image indicate representative vesicles that colocalise with p62, and are shown as green and magenta arrowheads in the p62 and dextran images respectively. **c**, Histogram of the probability of cooccurrence of p62 with Rab5, indicating a high likelihood of dynactin being present in a complex with endosomal cargo. **d**, Overlay of confocal images of MT (green) and dextran vesicles (magenta), and EM images of the same cell (grey). **e**, EM image of the region indicated with the white square in **d**, showing a representative MT and dextran endosome. **f**, Confocal fluorescence image of MTs (green) and dextran (magenta) of the region in **e**. **g**, Overlay of confocal images of MT (green) and EGF vesicles (magenta), and EM images of the same cell (grey). **h**, EM image of the region indicated with the white square in **g**, showing a representative MT and EGF endosome. **i**, Confocal fluorescence image of MTs (green) and EGF (magenta) of the region in **h**. In **a** and **b**, 'N' marks the location/direction of the nucleus.

The +TIP protein EB1 has been found to recruit another +TIP protein, CLIP-170, which in turn binds and clusters dynactin via its p150 subunit at growing MT plus ends (14). MT plus ends thus decorated with dynactin also accumulated dynein at these sites, and evidence suggested that cargo transport was initiated when these MT plus ends contacted intracellular cargo (15, 35). However, MT plus end-mediated initiation of dynein-driven transport appears to vary with cell type and context (14, 36, 37). Therefore, using SD microscopy in combination with super-resolution radial fluctuations (SRRF, (38)) we first quantified the localisation of p150 (Fig. 2a). Our high-resolution images revealed that only ~17% of the MT plus ends were enriched with p150, and p150 appeared bound along the entire length of the MT lattice (Fig. 2b). Further, by quantifying the intensities of mDHC-GFP expressed in our cells, we concluded that the significant MT plus end localisation of dynein reported in earlier studies (16, 39) may represent an artefact of dynein overexpression (Fig. S3a, b).

The dynactin complex, however, has been observed to show no MT-binding in the absence of dynein in *in vitro* assays (20). To ascertain that the p150 spots we observed in these cells represented the entire dynactin complex, we used SRRF to visualise p150 in concert with another dynactin subunit, p62 (Fig. 2c). The p62 subunit of dynactin is located in the pointed end complex of dynactin (9), and colocalisation of p62 with p150 would indicate the presence of the complete dynactin complex. We observed that 49±12 % (mean±s.d.) of the p150 spots colocalised with p62 (n=21,934/44,306 spots from N=2 independent experiments with 59 cells, Fig. 2d). Additionally, we used SRRF to visualise the localisation of p62 on MTs. The presence of p62 on MTs would indicate association with the MT of a subunit which does not normally do so unless it is part of the entire dynactin complex. Therefore, occurrence of p62 on the MT would imply localisation of the entire complex on the MT via p150. We observed that 74±18% (mean±s.d.) of the p62 spots (n=59,715/79,639 spots from 1 experiment with 25 cells) were present on MTs (Fig. 2e, f). Therefore, the complete dynactin complex is likely present along the entire length of the MT lattice. Additionally, fluorescently tagged p150 had residence times on the MT that far exceeded that of dynein, with most spots of

Fig. 4. Dextran and EGF vesicles display distinct movement kinetics. a, HILO image from a time-lapse of dextran vesicles in HeLa cells (left) and the corresponding kymograph (right). **b**, HILO image from a time-lapse of EGF vesicles in HeLa cells (left) and the corresponding kymograph (right). **c**, Plot of position versus time of dextran vesicles (grey) and EGF (brown) in HeLa cells. **d**, Probability distribution $P_-(\tau)$ of the minus end-directed run times for dextran (grey) and EGF (brown) vesicles. **e**, Probability distribution $P_0(\tau)$ of the pause times for dextran (grey) and EGF (brown) vesicles. In **a** and **b**, 'N' marks the location/direction of the nucleus. In **d** and **e**, the average run/pause time ($\langle\tau\rangle$) for dextran and EGF vesicles are indicated.

p150 remaining bound to the MT for the entire duration of the time-lapse video ($\sim$ 30 s) (Fig. S3c, Video S2). Further, we observed that depletion of p150 through siRNA mediated silencing reduced the levels of both p150 and p62 along the MT lattice (Fig. S3d-i). To test that the perturbation of dynactin localisation along the MT lattice resulted in reduced dynein activity, we depleted p150 using siRNA mediated silencing. As expected, we observed a reduction in the proportion of dynein molecules that moved towards the MT minus ends with a concomitant increase in the proportion that moved towards the plus ends (Fig. S4a-d).

The persistent association of dynactin to MTs that we obseved is in contrast to *in vitro* observations, where recombinant human dynactin was not found to decorate pig brain-derived MTs (20). Similarly recombinant budding yeast dynactin was observed to have a weak interaction with axonemal MTs (40). But when dynein was added to the mix, movement of dynactin was observed implying that dynactin could not independently bind to the MTs. However, in other instances (41), all p150 fragments containing the CAP-Gly domain bound to the MT. Dynactin that was over expressed in cells was found to bind strongly to and bundle MTs (11, 15). In fact, dynactin-MT interactions have recently been shown to be important for increasing the on-rate of dynein onto MTs (42). Taken together, dynactin's MT-binding function is likely essential for dynein's activity *in cellulo*.

Endosomes remain close to MTs and move in short bursts. We next sought to understand how dynein interacted with the third component of the active complex - the cargo adaptors. The endosomal cargo adaptors, Hook proteins, have been observed to remain persistently bound to their respective cargo (43, 44). Therefore, we used endosomal cargo as a proxy for the cargo adaptor, which confers the additional benefit of visualising the movement of the entire cargo adaptor and cargo complex. To avoid artefacts from overexpressing fluorescently tagged Rab5 to visualise early endosomes (45), we employed cells that had taken up 10kDa dextran or EGF conjugated with Alexa 647 (Fig. S5a).

Small molecular weight dextran enters the cell through all active endocytic mechanisms via fluid phase endocytosis, and uptake is therefore relatively independent of receptor-mediated endocytosis (46, 47). On the other hand, EGF is a ligand that binds to EGF receptor on the cell membrane, which at the concentrations used in this study, is taken up by clathrin-mediated endocytosis (48). While both these cargoes eventually end up at lysosomes near the minus ends of MTs, the timescales of movement of these cargoes is vastly different, with dextran requiring hours (49) to reach lysosomes and EGF doing so in tens of minutes (50). We sought to understand how these two different kinds of cargoes could potentially employ dynein's short run times to reach the lysosomes in different time scales.

Following a short pulse and chase, we observed dextran in Rab5-positive compartments (Fig. S5b, Video S3). EGF has also been shown to be in Rab5-positive compartments within the timescales of our chase observations (51), indicating that both dextran and EGF were in early endosomes in our experiments. We probed the localisation of dextran and EGF vesicles with respect to the MT. First, we imaged

Fig. 5. Cargo capture and dynein activation are coupled a, Image from the first frame of a time-lapse video (left) of dynein (green) and dextran (magenta), and the corresponding kymograph (right). The white arrow heads point to dynein and dextran vesicle moving together towards the minus end. **b**, Plot of position vs time of the dextran vesicle (magenta) indicated in **a**, alongside the intensity of dynein on that vesicle (green), showing a short minus-end directed run of the vesicle upon dynein-binding. **c**, Image from the first frame of a time-lapse video (left) of dynein (green) and EGF (magenta), and the corresponding kymograph (right). The white arrow heads point to dynein and EGF vesicle moving together towards the minus end. **d**, Plot of position vs time of the EGF vesicle (magenta) indicated in **c**, alongside the intensity of dynein on that vesicle (green), showing a short minus-end directed run of the vesicle upon dynein-binding. **e**, Images from time-lapse video of EGF entry and subsequent movement within cells. Time is indicated above the images in mm:ss. **f**, Kymograph of the region indicated by the magenta rectangle in **e**, with the timelapse images sampled every 5 s showing entry of EGF and the net movement of endosomes towards the nucleus. **g**, Kymograph of the trajectory marked by the white dashed rectangle in **f**, where the timelapse images were sampled every 0.5 s. In **a**, **c**, and **e**, 'N' marks the location/direction of the nucleus. **h**, Schematic of dynein's cargo search mechanism: stochastic binding of dynein to the MT at a location proximal to the cargo-adaptor complex leads to a short minus end directed run, which terminates upon the unbinding of dynein.

dextran vesicles and MTs in live cells and observed that the vesicles were in proximity to the MTs (Fig. 3a). We also observed that the dextran vesicles remained close to MTs even while they had no apparent tether to the MTs via motor proteins, and were therefore stationary (Video S4) in our minute-long live cell time-lapse images. We then used SRRF to visualise Rab5 in concert with p62 in living cells, and observed that ~70% of the Rab5 spots colocalised with p62 (n=8,795/12,848 vesicles from N=2 independent experiments with >20 cells each), indicating that cargo and dynactin are in proximity to each other on the MT (Fig. 3b, c). Finally, we employed CLEM to visualise the location of dextran and EGF endosomes with respect to MTs. We observed that in both instances, the endosomes were along MTs, indicating that endosomal cargo remained within $\sim$ 20 nm of MTs (dextran: 22±16 nm, from 3 endosomes, N=2 cells; EGF: 10±2 nm, from 3 endosomes, N=2 cells), likely along with dynactin (Fig. 3d-i).

While the interaction of dynactin with the cargo adaptor independently of dynein was discounted in earlier studies using fractionation (52), here we visualised colocalisation of dynactin with a large proportion of Rab5-labeled endosomes in close proximity to MTs (Fig. 3b).

To understand if endosomal vesicles could diffuse away from their original locations on the MT upon motor unbinding, we depolymerised MTs using nocodazole and visualised the subsequent movement of the vesicles (Fig. S5c, Video S5). Confirming previous findings on Rab5-positive compartments (30, 31), we measured a lower diffusion coefficient for the vesicles in the absence of MTs (Fig. S5d), likely im-

plicating the role of high intracellular crowding in constraining vesicle diffusion (31).

We next visualised dextran and EGF vesicles in 100s-timelapses and observed that the directed runs were sparse (Fig. 4a, b, Videos S6, S7), with only ~33% of dextran vesicles (65/196) and ~43% EGF vesicles (92/214) moving >1 μm during this time (Fig. 3c). Strikingly, both the dextran and EGF vesicles displayed uninterrupted minus end-directed runs that lasted only $\sim$ 0.6 s on an average (Fig. 3d) (0.6±0.2 s (mean±s.d.), n=196 and 214 for dextran and EGF vesicles respectively from N=2 independent experiments). Interestingly, the time between two consecutive runs, the pause time, reflected the time scales of minus end-directed movement of EGF and dextran, with EGF vesicles having an average pause time that was 30% shorter than that of dextran vesicles (0.4 s vs. 0.6 s, Fig. 4e). These results are comparable to the run-and-pause behaviour observed for endosomal cargo in previous studies (30, 31), and will be explored further in our stochastic modelling.

Dynein molecules are activated upon attachment to dynactin and cargo on the MT. Thus far, we have discovered that dynein transiently interacts with the MTs. However, dynactin and cargo are localised together and remained close to MTs. We therefore sought to understand how dynein interacted with the dynactin-cargo complexes on MTs. First, we performed dual-colour imaging of dynein and endosomal vesicles, and observed that vesicles that had dynein signal were, as expected, more likely to move towards the minus end of MTs (Fig. S5e-g). Next, by comparing the intensity of dynein on endosomal vesicles to single molecule bind-

Fig. 6. MCD model for dynein's cargo capture and activation a, Schematic of the MCD model depicting the MT (grey), with lattice points separated by a distance of Δx. Cargo complexes (blue) are positioned randomly along the MT and do not diffuse away. Dynein (green) and kinesin (magenta) motors bind and unbind stochastically to the MT lattice points with rates r_{on}^D, r_{on}^K and r_{off}^D, r_{off}^K respectively. If a cargo-complex is present at the site of motor binding (i), the motor could hop to the adjacent site $i+1$ or $i-1$, with a rate v^K or v^D for kinesin and dynein respectively. Note that we if two motors are already present at a lattice point (e.g.: 'Kinesin pair'), no more motors can bind to the same site. We used periodic boundary conditions to simulate this model on an MT of length $N\Delta x$. **b**, An example position vs. time trajectory of a dextran vesicle obtained from experiment (grey line), showing representative pauses (blue rectangles), minus end-directed runs (solid magenta lines), plus end-directed runs (solid green lines), and transitions from minus-to-plus and plus-to-minus end-directed runs (black arrowheads). **c**, Position vs. time trajectories of dextran vesicles from experimental data (green, top) and from the MCD simulation (magenta, bottom). **d**, Position vs. time trajectories of EGF vesicles from experimental data (green, top) and from the MCD simulation (magenta, bottom). Note that the experimental data in **c** and **d** are identical to those in Fig. 4c and have been reproduced here for comparison.

ing events, we estimated that, on an average, there were 1-2
dynein molecules bound to a vesicle (Fig. S5h). Finally, to
verify if single molecules of dynein could be activated and
perform minus end-directed movement when they stochas-
tically bound to MTs and encountered dynactin-cargo com-
plexes, we performed fast dual-colour HILO imaging of sin-
gle dynein molecules, along with dextran and EGF vesicles.
In these videos (Videos S8, S9), we observed instances where
previously stationary dextran and EGF vesicles started mov-
ing together with mDHC-GFP towards the MT minus ends
upon dynein binding (Fig. 5a-d, n=8/9 and n=3/3 events from
N=2 independent experiments with at least 15 cells).

To test that the repetitive binding of dynein to an endo-
some, followed by its short run towards the minus end was
sufficient to effect large-scale movements towards the nu-
cleus, we visualised the uptake and subsequent movement of
EGF endosomes over a $\sim$ 17-min period using HILO mi-
croscopy. We observed robust movement of the EGF en-
dosomes in the net minus end direction during the duration
of imaging (Fig. 5e, f, Videos S10). Strikingly, movements
that appeared as long-range at a low sampling rate (Fig. 5f)
showed intrinsic short runs when sampled at a higher rate
(Fig. 5g), similar to those seen in Fig. 4. Therefore, the
higher temporal resolution of our experiments enabled us to
discern the stop-and-go behaviour of cargo, that is ultimately
sufficient to drive movement to its intracellular destination in
time frames consistent with those previously observed.

Conceptual model of dynein's cargo attachment and subsequent movement within cells. Taken together, we observed that: (i) single molecules of dynein bind and unbind stochastically with MTs (Fig. 1b), (ii) if a dynactin-cargo complex is found close to this attachment location of dynein to the MT, a minus end-directed run of the dynactin-cargo-motor tripartite complex is effected (Fig. 5a-d), (iii) the detachment of dynein from the cargo concludes this run (Fig. 5h), (iv) the detached dynein motor is free to diffuse back into the cytoplasm, whereas dynactin and the cargo remain paused and remain close to the MT (Fig. 3), (v) the long-range movement of endosomal cargo requires the repeated binding-unbinding of dynein to the dynactin-cargo complex, (vi) the resulting motion of the cargo consists of short minus (and plus) end-directed runs punctuated by long pauses (Fig. 4c).

    Several remarks are in order. First, the steps outlined above constitute an effective intermittent mechanism to ensure cargo transport by stochastic motor dynamics. Such strategies have been described in the past for transcription factors seeking specific sequences along the DNA (53, 54). Second, since cytoplasmic dynein is the only minus end-directed motor that participates in membrane trafficking in many cell types, how dynein interacts with and transports different types of cargo is an interesting question. Recent research suggests that cargo specific adaptors like BicD2 (for Rab6-positive cargo), and Hook1/3 (for Rab5-positive cargo) have differential interaction with dynein (20). These cargo

Fig. 7. The 3-state RTP model and its fits for the movement of dextran vesicles. **a**, Schematic of the RTP model depicting the 3 states: stationary, minus end-directed and plus end-directed, and the transition rates between them, i.e. $\gamma_{0\mathrm{m}}$, $\gamma_{\mathrm{m}0}$– rates of transition from the stationary to minus end-directed state and vice versa; $\gamma_{0\mathrm{p}}$, $\gamma_{\mathrm{p}0}$– rate of transition from the stationary to the plus end-directed state and vice versa; γ_{mp}, γ_{pm}– rates of transition from the minus end-directed to the plus end-directed state and vice versa. In **b-g**, experimental data (green squares) are compared with the output from MCD simulations (magenta circles) and the predictions of the RTP model (grey lines) for the probability distributions of the cargo waiting time (τ) while transitioning from the stationary to the minus end-directed state (**b**), minus end-directed to stationary (**c**), plus to minus end-directed (**d**), stationary to plus end-directed (**e**), plus end-directed to stationary (**f**), and minus end-directed to plus end-directed (**g**). Further, the probability distributions of the pause-times (**h**), the run times in the plus end-directed state (**i**), and the run times in the minus end-directed state (**j**) are shown in **h-j**, while **k** and **l** show the distributions of the the run lengths (l) in the plus end-directed and minus end-directed states respectively. The solid grey lines represent fits to the experimental data from the RTP model. The MCD model results allowed pairs of motors to transport cargo-complexes, i.e., M=2. Note that he plots (**b-l**) are in log-linear scale.

adaptors could modulate the interaction time of dynein with the cargo thereby leading to different kinetics. Such stochastic binding and unbinding allow the same dynein molecule to sample and interact with a wide range of cargo, akin to round robin scheduling. Third, differences in the internalisation route for cargo bound by the same adaptor, such as EGF and dextran as seen in this study, could also alter the number of cargo adaptors on the cargo, which in turn changes the motor on-rates to the cargo. Movement to the plus or minus ends of MTs could therefore also be tuned by slightly increasing the bias in one direction, for example by increasing the number of adaptors for dynein on a cargo destined towards the minus end. Fourth, cellular morphology is likely to play a significant role in the kinetics of motor attachment and detachment from the MT. For instance, in highly polar, narrow axons of neuronal cells, reattachment of dynein to cargo localised close to the MT is likely to occur frequently due to the confinement, and therefore reduced space in which dynein could diffuse. This would result in higher on rates for the motor to the cargo, and therefore more processive cargo trajectories. Finally, it is very likely that a similar mechanism is at work in transporting plus end directed cargo by kinesin motors. More generally, both dyneins and kinesins can simultaneously affect cargo movement, even for those cargo that are transported towards the minus end on the average such as dextran and EGF.

Our experimental techniques - HILO, SD and SRRF, and CLEM - provided high spatiotemporal data for the dynamics of motors and dynactin-cargo complexes. In fact, this not only allowed us to quantify the stochastic kinetics of individual dynein motors to extract its detachment rates (Fig. 1f), but also allowed us to explore the long-time dynamics of cargo with high precision (Fig. 4c). As such, we sought to quantify the statistical dynamics of endosomal cargo by measuring the distributions of the cargo residence time of both dextran

and EGF in the moving and paused states. Before discussing these results in detail (Fig. 7-9), we develop stochastic models at two levels of descriptions that can quantitatively account for these distributions of cargo dynamics.

A stochastic model captures the kinetics of motor–driven cargo movements. To study the stochastic dynamics of motors and cargo units capable of activating the motor (which include dynactin and adaptor for dynein, henceforth referred to as 'cargo complexes'), we modelled the MT as a discrete one-dimensional lattice of N binding-sites each of size Δx (=0.05 μm) (Fig. 6a). Our model included both dyneins and kinesins interacting with the MT and cargo complexes. At each site i of the MT, exchange kinetics with a large cytoplasmic reservoir caused dynein (kinesin) motors to bind at a Poisson rate r_{on}^D (r_{on}^K) to the MT and unbind from the MT at a Poisson rate r_{off}^D (r_{off}^K) to the cytoplasm. Note that this binding-unbinding process was stochastic and occurred uniformly throughout the length of the MT. If a cargo complex was found at the binding location of a motor protein, then the motor-cargo complex translocated along the MT with a speed v^D towards the minus end in the case of dyneins and with a speed v^K towards the plus end in the case of kinesins. Dynein (kinesin) motors detached from the MT after a mean residence time $\sim 1/r_{\mathrm{off}}^D$ ($\sim 1/r_{\mathrm{off}}^K$), which also ended the motion of the cargo complex. From this point on, the cargo complex remained stationary, but still in proximity to the MT. In other words, we completely neglected the 3D diffusion of the cargo-adaptor complexes since our data suggested that cargo diffusion is negligible (Fig. S5d). Furthermore, we only considered a single cargo complex at a time and disregarded any steric interactions between them. The chemical master equation describing this Markovian stochastic process is described in detail in the Supplementary Text. When we restricted the number of motors that could be bound

to each site of the MT (M) to one, this stochastic motor-cargo complex dynamics model (referred to as the 'MCD model' henceforth) generated cargo trajectories that underwent pauses, minus end-directed and plus end-directed runs (Fig. S6). However, in this instance, the only transitions that occurred were between paused and moving (plus or minus) states of the cargo complex, which were not truly reflective of our experimental observations.

Typical experimental trajectories of the cargo showed discrete paused and moving states (Fig. 6b for dextran). In addition to the transitions between pauses and runs, we observed direct transitions between minus end-directed runs and plus end-directed runs. This, however, is not possible unless the cargo complex was moved by at least a pair of motors consisting of one dynein and one kinesin. In other words, our experimental data (Fig. 6b) indicated that there were antagonistic pairs of motor proteins bound to the cargo complex (see Fig. S6 for the output of the MCD model with M=1). In fact, our experiments with dual imaging data of single molecules of dynein and cargo indicated that there were one to two dyneins bound to the cargo (Fig. S5h). It was thus reasonable to assume that we would also have pairs of kinesins, and dynein-kinesin pairs bound to a cargo complex. Note that cargo translocation could be effected by both single motors as well as pairs of motors.

We therefore extended our MCD model to the case where a maximum of M=2 motors could be bound to each site of the MT. Note that even with M=2, we assumed that the motors hopped at the same speeds, *i.e.*, at rates $v^D/\Delta x$ and $v^K/\Delta x$, across the lattice sites of the MT. In particular, we also assumed that the antagonistic pair could move in either direction at the appropriate rates. Further, we neglected any cooperative binding between multiple motors. We discuss in detail the chemical master equation associated with this model in the Supplementary Text, and also its stochastic simulation via the Gillespie algorithm (55, 56). With values for the parameters r_{on}^D, r_{on}^K, r_{off}^D, r_{off}^K, v^D and v^K that compared reasonably well with their estimates from experiments (see table in Fig. 9c), the MCD model generated cargo trajectories that were comparable to empirical results (Fig. 6c and d). However, taking advantage of our long-term and high temporal resolution imaging data, we next asked whether the model could quantitatively account for the residence time distributions of the cargo in the various states, and if there was a simplified description for the dynamics of the cargo complexes that themselves behaved as "super-motors", displaying motion which was markedly different from simple Brownian motion on the MT.

The emergent dynamics of cargo is captured by a 3-state run-and-tumble model. Typical trajectories, such as the one shown in Fig. 6b, revealed short periods of directed motion of the cargo complexes interspersed with frequent pause events. This kind of stochastic motion was *not* characteristic of the simple, and well-known, thermal Brownian motion. Rather, this movement was reminiscent of the classic run-and-tumble motion wherein bacteria run in almost straight line paths for a certain run-duration after which they tumble while being stationary for a tumble-duration, and then run again in a randomly chosen direction (57). Statistical physics models of such run-and-tumble particles (RTP) have revealed very interesting features of these persistent and active random walks (58–60). However, in these well-studied RTP models, the tumble is, typically considered to be an instantaneous event. Several studies have explored models, with three distinct states– two moving states (with movements in opposite directions), and a paused state– for cargo that arise from underlying motor movement (61, 62).

Our MCD model discussed in the previous section led to an emergent picture for the dynamics of the cargo complexes as RTPs moving on a one-dimensional MT with three distinct states: *minus end-directed*, *paused* and *plus end-directed*. In more detail, these states resulted from the following situations: (i) a minus-end directed run of the cargo complex occurred when there were motor configurations dominated by dyneins, (ii) cargo complexes were paused when there were no motors (of either kind) within a proximal region on the MT, and (iii) a plus-end directed run occurred when kinesins were the primary movers of the cargo complexes. As such, the effective movement of the cargo was captured by three discrete states with distinct transition rates between them.

We captured the dynamics of the cargo complexes discussed above in a 3-state RTP model. Specifically, as depicted in Figure 7a, stationary cargo could transition to the plus (minus) end-directed state at a rate $\gamma_{0\mathrm{p}}$ ($\gamma_{0\mathrm{m}}$), while moving cargo could transition to the stationary state at rates $\gamma_{\mathrm{m}0}$ (from the minus end-directed state) and $\gamma_{\mathrm{p}0}$ (from the plus end-directed state). Further, we considered direct transitions between the plus and minus end-directed states at rates γ_{pm} and γ_{mp} respectively, since these transitions were apparent in our experimental data (Fig. 6b). In the plus end-directed state, the cargo translocated with a velocity $u = +u^{\mathrm{p}}$ and in the minus end-directed state, the translocation velocity was $u = -u^{\mathrm{m}}$. In the stationary state, by definition, the cargo was immobile, i.e., $u = 0$. Therefore, denoting $(\mathcal{P}_{\mathrm{m}}, \mathcal{P}_0, \mathcal{P}_{\mathrm{p}})$ as the probabilities of finding the cargo in the minus end-directed, stationary and plus end-directed states respectively, the master equations governing their Markovian time-evolution were:

$$\frac{d}{dt}\begin{pmatrix} \mathcal{P}_{\mathrm{m}} \\ \mathcal{P}_0 \\ \mathcal{P}_{\mathrm{p}} \end{pmatrix} = \Gamma \begin{pmatrix} \mathcal{P}_{\mathrm{m}} \\ \mathcal{P}_0 \\ \mathcal{P}_{\mathrm{p}} \end{pmatrix} \qquad \textbf{(1)}$$

where the transition matrix Γ was:

$$\Gamma = \begin{pmatrix} -\gamma_{\mathrm{m}0} - \gamma_{\mathrm{mp}} & \gamma_{0\mathrm{m}} & \gamma_{\mathrm{pm}} \\ \gamma_{\mathrm{m}0} & -\gamma_{0\mathrm{m}} - \gamma_{0\mathrm{p}} & \gamma_{\mathrm{p}0} \\ \gamma_{\mathrm{mp}} & \gamma_{0\mathrm{p}} & -\gamma_{\mathrm{p}0} - \gamma_{\mathrm{pm}} \end{pmatrix} . \qquad \textbf{(2)}$$

Note that in these emergent discrete states, the cargo complexes moved with the appropriate velocity. In other words, the position x of the cargo complex changed with time according to $dx/dt = u$, where the velocity u could take either of the three possibilities $\{-u^{\mathrm{m}}, 0, u^{\mathrm{p}}\}$.

The stochastic dynamics captured by equation Eq. (1) predicts that the distribution $P_{\mathrm{ab}}(\tau)$ of the residence time

Fig. 8. Fits of the 3-state RTP model for EGF vesicle trajectories. In **a-f**, experimental data (green squares) are compared with the output from MCD simulations (magenta circles) and the predictions of the RTP model (grey lines) for the probability distributions of the cargo waiting time (τ) while transitioning from the stationary to the minus end-directed state (**a**), minus end-directed to stationary (**b**), plus to minus end-directed (**c**), stationary to plus end-directed (**d**), plus end-directed to stationary (**e**), and minus end-directed to plus end-directed (**f**). Further, the probability distributions of the pause-times (**g**), the run times in the plus end-directed state (**h**), and the run times in the minus end-directed state (**i**) are shown in **g-i**, while **j** and **k** show the distributions of the the run lengths (l) in the plus end-directed and minus end-directed states respectively. The solid grey lines represent fits to the experimental data from the RTP model. The MCD model results allowed pairs of motors to transport cargo-complexes, i.e., M=2. Note that the plots **a-k** are in log-linear scale.

(equivalently waiting time) in the a-state before jumping to the b-state is an exponential distribution with the rate constant γ_{ab} (here a, b could take either of the values m, 0 or p). For instance, the distribution of the residence times of the cargo complex in the paused state before transitioning to the minus end-directed state would be $P_{0m}(\tau) = A_{0m} e^{-\gamma_{0m}\tau}$, where A_{0m} is an amplitude factor. Similar considerations applied to the waiting time distribution for the other transitions.

We next asked whether our high-temporal resolution, long-term data for the positions of the cargo complexes was consistent with the waiting time distributions obtained from this coarse-grained 3-state RTP model. To this end, we partitioned the trajectory of the cargo complexes into discrete states (Fig. 6b, see Methods), measured the waiting time for each transition depicted in Fig. 7a, and fitted exponential functions to the probability density functions of the residence times τ. We found a good match between our empirical data and the distributions $P_{ab}(\tau)$ as shown in Fig. 7(b-g) for the case of dextran vesicles. We employed this procedure to extract the effective rate constants γ_{ab} for a transition of the cargo complex from the a−state to the b−state (see the table in Fig. 9c). Notice that the cargo could exit the stationary/paused state via two independent pathways: transitioning into the minus end-directed state or transitioning into the plus end-directed state. As such, the distribution of the pause times was a mixed distribution $P_0(\tau) = w_0 P_{0m}(\tau) + (1-w_0) P_{0p}(\tau)$, where $0 \le w_0 \le 1$ is a weight factor. To check whether this prediction worked, we fixed the rates γ_{0p} and γ_{0m} to the values obtained by the fits in Fig. 7b and Fig. 7e respectively, and fit the form of $P_0(\tau)$

to the empirical data. The excellent fit seen in Fig. 7h confirms the validity of our 3-state RTP model. Similar considerations applied for the distribution $P_m(\tau) = w_m P_{m0}(\tau) + (1-w_m) P_{mp}(\tau)$ of minus end-directed run times (Fig. 7i), and the distribution $P_p(\tau) = w_p P_{p0}(\tau) + (1-w_p) P_{pm}(\tau)$ of plus end-directed run times (Fig. 7j). Further, in the minus (plus) end-directed states, we measured the distribution of the run lengths $\mathbb{P}_m(l)$ and ($\mathbb{P}_m(l)$). Since the cargo complexes only moved when they were in the $m-$ or $p-$states, the run length distributions were expected to be related to the run time distributions. Specifically, we expected: $\mathbb{P}_m(l) = [\omega_m P_{m0}(l/u^m) + (1-\omega_m) P_{mp}(l/u^m)]/u^m$, and $\mathbb{P}_p(l) = [\omega_p P_{p0}(l/u^p) + (1-\omega_p) P_{pm}(l/u^p)]/u^p$. The excellent fits of the empirical data to these distributions seen in (Fig. 7j-k) allowed us to extract the emergent speeds u^m and u^p of the cargo complexes.

We obtained more compelling evidence to this emergent picture when we were able to match the results of the MCD simulation, depicted in Fig. 6a, for the various residence time and run length distributions of the empirical data and the RTP model. We repeated this procedure for the case of EGF vesicles and obtained similar agreement between the empirical data, the MCD simulation and the RTP model as shown in Fig. 8. The parameter values both in our MCD simulations and those obtained in the RTP model are available in Fig. 9c. The predictive power of our approach was revealed when we examined the dependence of the average run length $\langle l \rangle$ as a function of the run time τ in the case of both dextran and EGF. Since all parameters in the RTP model were known, we compared the empirical data and the MCD simulation data with the predictions $\langle l \rangle = u^p \tau$ and $\langle l \rangle = -u^m \tau$ for plus end-

Fig. 9. Run lengths of cargo scale with run time, Plot of average run length $\langle l \rangle$ as a function of run time τ for dextran vesicles (**a**) and EGF vesicles (**b**), with data from experiments (green squares) and output from the MCD model with $M = 2$ (magenta circles). The black solid line represents the prediction of the RTP model using the parameter listed in the table below. **c** Table containing the parameters that were measured for single dynein and kinesin molecules in this study and other experiments ('Experiments'), those that were employed in the MCD model ('Stochastic model'), and those for the cargo movement that obtained from fits to the experimental data from the RTP model ('RTP model') for both dextran and EGF.

Experiment			MCD model			RTP model		
Parameter	Value	Source	Parameter	Value		Parameter	Value	
				Dextran	EGF		Dextran	EGF
r_{on}^{D}	5 /s*	Ref. (31)	r_{on}^{D}	1.2 /s	1.6 /s	γ_{0m}	2.5±0.7 /s	5.5±1.0 /s
r_{off}^{D}	1.7±0.2 /s‡	This study	r_{off}^{D}	1.0 /s	1.0 /s	γ_{m0}	4.1±0.8 /s	4.8±0.9 /s
r_{on}^{K}	5 /s*	Ref. (31)	r_{on}^{K}	1.1 /s	1.5 /s	γ_{0p}	3.6±0.7 /s	4.7±1.0 /s
r_{off}^{K}	0.8 - 2.5 /s	Ref. (31)	r_{off}^{K}	1.0 /s	1.0 /s	γ_{p0}	4.3±0.6 /s	5.0±1.0 /s
v^{D}	1.2±0.7 µm/s†	This study	v^{D}	0.2 µm/s	0.2 µm/s	γ_{mp}	4.5±0.6 /s	4.5±0.9 /s
v^{K}	1.0±0.7 µm/s†	This study	v^{K}	0.2 µm/s	0.2 µm/s	γ_{pm}	4.0±0.8 /s	5.0±0.7 /s
						u^{p}	0.49±0.04 µm/s	0.49±0.03 µm/s
						u^{m}	0.48±0.05 µm/s	0.52±0.09 µm/s

Estimated from comparing experimental data to a stochastic model, ‡ 95% confidence interval, † Standard deviation

directed and minus end-directed runs respectively. The remarkable agreement seen in Fig. 9a and b confirmed that our simulation and modelling approaches provide an accurate description of the experimental data. So too, the scaling of emergent cargo velocities with run time points to the participation of at least 2 motors in the cargo trafficking process. Indeed, the output from the MCD model with $M = 1$ did not match the experimental data (Fig. S6) or fit the RTP model (Figs. S7, S8a).

It is thus apparent that our MCD model, which explicitly included the dynamics of motors, led to a coarse-grained emergent view wherein the dynamics of the cargo complexes were governed by a 3-state RTP model. However, an important point to note is that it is a non-trivial task to find analytical relations between the emergent parameters at the RTP level (γ_{p0}, γ_{m0}, γ_{0m}, γ_{0p}, γ_{pm}, γ_{mp}, u^{m} and u^{p}) and the parameters in the model (r_{on}^{D}, r_{off}^{D}, r_{on}^{K}, r_{off}^{K}, v^{D} and v^{K}). Nevertheless, the parameters used in our simulations revealed that the on-rates for both dynein and kinesin are higher by $\sim 30\%$ for EGF as compared to dextran. Remarkably, our experimental data showed that EGF vesicles had an average pause time which was $\sim 30\%$ shorter than that of dextran vesicles (Fig. 4e). This difference in on-rates and thereby the pause times, could have arisen from an increased number of adaptors for dynein (and kinesin) on EGF-containing vesicles compared to dextran-containing vesicles, and will be the subject of future studies. Together, we developed a framework for considering cargo movement in light of a 3-state RTP model which could be used to infer not only underlying motor kinetics, but also predict long-term cargo movement and fates, such as in differential cargo sorting.

Conclusion

In this work, we established a technique to visualise single molecules of dynein in living cells, and observed their interaction with the MT, and other components of the tripartite complex required for its activation. We discovered distinct localisation and kinetics of dynein, dynactin and the cargo complex: While dynein interacted stochastically and transiently with the MT, dynactin and cargo remained persistently bound and in proximity to the MT. When dynein interacted with a dynactin-cargo complex upon MT-binding, the entire complex moved towards the minus end of MTs in a short run lasting a little over half a second. Repeated rounds of such stochastic interactions of the motor with the MT and cargo complex are required for long-range movement of cargoes.

We measured single-molecule parameters pertaining to dynein, and used these and other parameters available in literature to develop a MCD model for cargo trajectories based on assumptions from experiments. The experimental data

and output from the MCD model for two different kinds of cargo – dextran and EGF vesicles – were in good agreement. Importantly, the experimental data and MCD model output quantitatively agree with the predictions of a 3-state RTP model. Taken together, we demonstrate that stochastic interactions of motors with MTs and cargo complexes are sufficient to elicit complex cargo trafficking behaviour in living cells.

Methods

Cell culture. HeLa cells were cultured in Dulbecco's Modified Eagle Media (DMEM) containing 100 units/mL Penicillin, 0.1 mg/mL Streptomycin, 2 mM L-Glutamine, 400 µg/mL Geneticin (HiMedia Laboratories) and supplemented with 10% fetal bovine serum (Sigma Aldrich). Cells were grown in an incubator at 37°C under 5% CO_2. During live-cell imaging, cells were maintained in live-cell imaging solution (140 mM NaCl, 2.5 mM KCl, 1.8 mM CaCl2, 1 mM MgCl2, 4 mg/mL D-Glucose, 20 mM HEPES) at 37°C using a stage-top incubator (Oko Labs).

RNAi experiments. For all RNAi experiments, cells were first grown in glass-bottom imaging dishes (Cat. no. 81518, Ibidi) for 48 h, transfected with the appropriate concentration of siRNA using Jet Prime transfection reagent (Cat. no. 114, Polyplus) and imaged after 48 h. For hDHC RNAi the following siRNA sequence was used: 5'-ACA-UCA-ACA-UAG-ACA-UUC-A-3'10. For p150 RNAi the following siRNA sequences were used: 5'-GUA-CUU-CAC-UUG-UGA-UGA-A-3', 5'-GAU-CGA-GAG-ACA-GUU-AUU-A-3' [14]. siRNA oligonucleotides were procured from Eurogentec, Belgium. To quantify the knockdown of proteins, the cells were lysed 48 hours after siRNA transfection to obtain protein solution and SDS-PAGE, western blotting was performed. The following primary antibodies were used to visualise proteins of interest and controls: Rabbit DYNC1H1 Polyclonal Antibody (PA5-49451, Invitrogen, 0.5 µg/mL), Rabbit Dynactin 1 Polyclonal Antibody (PA5-37360, Invitrogen, 0.5 µg/mL), Mouse GAPDH Loading Control Antibody (MA5-15738, Invitrogen, 0.5 µg/mL), Rabbit Actin Polyclonal Antibody (ab8227, AbCam, 0.5 µg/mL). The following secondary antibodies were used: Anti-Rabbit HRP (31460, Invitrogen, 0.16 µg/mL) and Anti-Mouse HRP (62-6520, Invitrogen, 0.3 µg/mL).

Immunofluorescence. To immunostain p150 and tubulin, cells were first grown for 48 h in glass-bottom imaging dishes (Cat. no. 81518, Ibidi) and fixed in ice-cold methanol at -20°C for 3 min. Cells were washed thrice in phosphate buffered saline (PBS) for 5 min each and incubated in blocking buffer (5% BSA in PBS) for 60 min. Subsequently, the cells were incubated with p150 and α-tubulin antibodies in antibody dilution buffer (1% BSA in PBS) for 60 min. Cells were then washed thrice in PBS for 5 mins and incubated with secondary antibodies in antibody dilution buffer for 45 min. At the end of incubation with secondary antibodies, the cells were washed with PBS and imaged immediately. The following primary antibodies were used. Rabbit Dynactin 1 Polyclonal Antibody (PA5-21289, Invitrogen, 2 µg/mL), Mouse α Tubulin Monoclonal Antibody (32-2500, Invitrogen, 2 µg/mL). The secondary antibodies used were Donkey Anti-Rabbit A555 Antibody (A31572, Invitrogen, 0.4 µg/mL) and Goat Anti-Mouse A647 Antibody (A28181, Invitrogen. 2 µg/mL).

To immunostain p150 and p62, p62 and α-Tubulin and EB1 and α-Tubulin the following method was used. Cells were first grown for 48 h in glass bottomed imaging dishes and fixed in ice-cold methanol at -20°C for 3 min. Subsequently, the cells were washed thrice in PBS for 5 min each and incubated in antibody dilution buffer (2% BSA, 0.1% Triton-X100, in PBS) for 10 min. Next, the cells were incubated in antibody dilution buffer containing the appropriate antibodies for 60 min. Cells were washed thrice in PBS for 5 min each and incubated with appropriate secondary antibodies in antibody dilution buffer for 30 min. At the end of incubation with secondary antibodies, the cells were washed well in PBS and imaged immediately. The following primary antibodies were used. Rabbit Dynactin 1 Polyclonal Antibody (PA5-21289, Invitrogen, 2.5 µg/mL), Rabbit Alpha Tubulin Monoclonal Antibody (PA5-22060, Invitrogen, 2 µg/mL), Mouse Dynactin 4 Monoclonal Antibody (MA5-17065, Invitrogen, 2 µg/mL) and Mouse EB1 Monoclonal Antibody (412100, Invitrogen, 2 µg/mL). The secondary antibodies used were Donkey Anti-Rabbit A647 Antibody (A32795, Invitrogen, 2 µg/mL) and Goat Anti-Mouse A555 Antibody (A21422, Invitrogen. 2 µg/mL).

To immunostain Rab5 and p62, cells were first fixed in 4% Paraformaldehyde (PFA) for 15 mins. The cells were washed well in PBS and incubated for 45 min in antibody staining solution (0.2% Saponin, 0.1% BSA and 0.02% Sodium Azide in PBS) containing primary antibodies against Rab5 and p62. Subsequently, cells were washed well in PBS and incubated for 45 min in antibody staining solution containing appropriate secondary antibodies. Cells were washed well in PBS and imaged immediately. The following primary antibodies were used. Rabbit Rab5 Monoclonal Antibody (3547, Cell Signaling Technology, 14 µg/mL), Mouse Dynactin 4 Monoclonal Antibody (MA5-17065, Invitrogen, 2 µg/mL). The secondary antibodies used were Donkey Anti-Rabbit A647 Antibody (A32795, Invitrogen, 2 µg/mL) and Goat Anti-Mouse A555 Antibody (A21422, Invitrogen. 2 µg/mL).

Dextran and EGF uptake. To visualise fluorescent dextran vesicles, HeLa cells that were grown in glass-bottom imaging dishes (Cat. no. 81518, Ibidi) for 48 h were transferred to serum-free DMEM for 4 h in a 37°C CO_2 incubator. Subsequently, the cells were pulsed for 10 min in complete DMEM containing 200 µg/mL A647-Dextran (D-22914, Invitrogen). Cells were then washed well with live-cell imaging solution before proceeding for microscopy. In all experiments, imaging was completed within 45 min of dextran uptake.

To visualise fluorescent EGF vesicles, HeLa cells that were grown in glass-bottom imaging dishes (Cat.

no. P35G-1.5-14-C, Mattek) for 24 h were transferred to serum-free DMEM for 1 h in a 37°C CO$_2$ incubator. Cells were washed in serum-free DMEM, then incubated with 1 nM EGF (Cat. No. E9644, Sigma-Aldrich) labelled with Alexa647 (Cat. No. ab269823, Abcam) in phenol-free, serum-free DMEM for a minimum of 8 minutes prior to imaging. In all experiments, cells were imaged between 8-20 minutes of EGF uptake.

Transfection. To visualise fluorescently tagged proteins, cells were transfected with the appropriate plasmids using Jet Prime Transfection Reagent (Cat. no. 114, Polyplus). 3 h after transfection, the cells were washed well in PBS, grown in complete DMEM at 37°C and imaged ~20 h later. The following plasmids were used in this study. (i) mCherry-tubulin was a gift from Mariola Chacon, TUD, Dresden. (ii) Gal4T-mCherry was a gift from Thomas Pucadyil, IISER Pune. (iii) mCherry-DCTN1 was a gift from Prof. Kozo Tanaka, Tohoku university, Japan. (iv) mCherry-Rab5 was a gift from Gia Voeltz (Addgene plasmid - 49201).

Microscopy. Imaging was performed on a Nikon Ti2E inverted microscope equipped with a Toptica MLE laser combiner, Nikon H-TIRF Module, Yokogawa CSU-X1 spinning disk module, an Andor iXon 897 EMCCD camera and an Oko Lab stage top incubator. The microscope was controlled using Nikon NIS Elements software or Micromanager (63). Alternately, a Zeiss ELYRA controlled using Zen Black software was used.

HILO microscopy and particle tracking. For HILO microscopy, a Nikon 100X 1.49NA TIRF objective was used. We optimised the HILO microscopy setting for each cell: First, to avoid overexpression artefacts, we visualised cells expressing low levels of mDHC-GFP. The diameter of the illuminated area was kept constant at 30 μm and the long axis of cell was aligned perpendicular to the excitation laser. Finally, we adjusted the angle of incidence of the excitation laser, such that the fluorescent spots in a plane ~0.5 μm from the cover slip appeared bright and distinct. The incidence angle and orientation of excitation laser beam were adjusted using the H-TIRF module and the laser power (50 mW at fiber) was kept at 40%. To visualise single molecules of dynein, time-lapse images were acquired at 50 fps with 20 ms exposure per frame.

To improve the signal-to-noise ratio, a five-frame sliding average of the time-lapse images was used for particle tracking. In these videos, the appearance and subsequent disappearance of intensity at a particular location was classified as a binding event. Such binding events were visually identified and tracked from the start to end frame using Low Light Tracking tool (64) in Fiji/ImageJ (65). The fluorescent spots were classified as stationary, minus-end moving or plus-end moving by marking a region close to the nucleus as the minus end and calculating the displacement of fluorescent spots with respect to the minus end. Particles moving away from the nucleus for >0.32 μm were classified as plus-end moving, and particles with displacement >0.32 μm towards the nucleus were minus-end directed, and particles with net displacement <0.32 μm were classified as stationary.

For dual channel imaging of dynein and cargo, videos were acquired at 20 fps with 25 ms exposure per channel for dynein and dextran (in cells depleted of hDHC using RNAi), and at 6 fps with 33 ms exposure per channel for dynein and EGF. To observe dynein-dextran interaction with a better signal-to-noise ratio, a two-frame sliding average of time-lapse images was used.

For 100 s dextran and EGF imaging experiments, an alpha Plan-Apochromat 100x/1.46 Oil Elyra TIRF objective was used. Consistent HILO setting were used across all long-term dextran and EGF uptake experiments. Fluorophores were excited using a 642 nm laser (150mW at source) at a TIRF angle of 63°. Exposure times of 50 ms with a frame interval of 100 ms were used (total frame time, 125 ms with frame transfer), and imaging undertaken for 1000 frames (125 s in total). The movement of dextran vesicles was classified in the following way: the movement between two consecutive frames was classified as zero if it was less than 20 nm (the tracking accuracy of low light tracking tool under the imaging conditions used (64).). Then, displacement for four consecutive frames towards or away from the nucleus was classified as a minus or plus end-directed run, and a zero displacement for 2 or more consecutive frames was classified as a pause. The same criteria for runs and pauses was used for analysing the trajectories obtained from the MCD model.

Spinning disk (SD) confocal microscopy and image analysis. To quantify the correlation between expression level of dynein and its clustering, z-stack images of live HeLa cells expressing mDHC-GFP were acquired using a 60x, 1.4 NA objective with 50 ms exposure. The intensity comparison was done for the lowest plane at which clusters were distinctly visible. The dynamics of mCherry-p150 in cells were visualised by acquiring 60-s long movies of the cell with a 1.4 NA 60x objective, 100 ms exposure per frame and 1 s interval between frames. The dynamics of dextran vesicles along the MTs was visualised by acquiring 30-s long movies of the cell with a 60x, 1.4 NA objective, 100 ms sequential exposure/channel and 1 s interval between frames.

The SD+SRRF images were obtained by taking the mean of the radiality map of 100 images of a single field acquired under a 100x, 1.49NA objective with the SD confocal microscopy setup. The radiality magnification in SRRF was set to 4 for all experiments. To visualise the association between MTs and dextran vesicles in live cells, the SRRF acquisition settings used were 20ms exposure, ring radius of 0.5 for MTs and 3 for dextran. To visualise p62 with tubulin and EB1 with tubulin, 20 ms exposure and ring radius of 0.5 was used. To visualise p150 with tubulin, and p150 with p62 the SRRF acquisition settings used were 100 ms exposure, ring radius of 1. Finally, to visualise p62 with Rab5, 20ms exposure was used in both channels. A ring radius of 0.5 and 3 were used for p62 and Rab5, respectively.

The cooccurrence of p150 with p62 was quantified by first thresholding SRRF images of p150 for the top 5% of intensity. The number of p150 spots was counted N_{p150}, and a mask was created. Then, the SRRF image of p62 in the same cells was combined with the p150 mask using AND operation in Fiji. The resultant image was thresholded and the number of spots counted (N_{p62}). $N_{p62}/N_{p150}*100$ gave the percentage of p150 spots with p62. A similar procedure was used to quantify the percentage of p62 spots that occurred on the MT and also the percentage of Rab5 spots that were associated with p62. To quantify the effect of p150 knock-down on the levels of p150, p62 along the MT, ROIs were drawn along distinct MT segments that were randomly chosen in SRRF images of MT, and the mean intensities of p150 or p62 in these ROIs were measured.

Correlative Light and Electron Microscopy. HeLa cells were grown for 24 hours on gridded glass-bottom imaging dishes (Cat. No. P35G-1.5-14-CGRD, Mattek), following which cells were transfected with EGFP-tubulin (Cat. No. 56450, AddGene) overnight. Cells were transferred into serum-free, phenol free DMEM for 1 h, following which A647-dextran or EGF-647 were added to cells for 10 min. Cells were extensively washed in PBS and fixed in 4% PFA containing 0.2% glutaraldehyde (Cat. No. 354400, Sigma-Aldrich) for 1 h at room temperature. Following fixation, cells were dyed with MitoTracker Orange CMTMRos (Cat. No. M7510) to fluorescently label mitochondria through the cell volume, then imaged on a Zeiss LSM 800 with Airyscan detector, with a pixel size of 42.5 nm and z-slice spacing of 150 nm. The MitoTracker-stained mitochondria served as markers for correlation and registration of the relevant confocal z-slice with that of the corresponding TEM z-slice based on identification of mitochondrial morphology and distribution at each slice (Fig. S8). After imaging, cells were further fixed in 2.5% glutaraldehyde in 0.1 M sodium cacodylate for 1 h at room temperature. All the subsequent processing steps were carried out in a BioWave microwave (Pelco). Post fixation, samples were washed with cacodylate buffer, additionally fixed in reduced osmium (1% osmium tetroxide-1.5% potassium ferricyanide) solution followed by buffer wash, resuspended in 2% aqueous OsO_4 (osmium tetroxide) (ProScitech). Washed samples were then stained with 2% w/v Phosphotungstic acid (PTA) in 30% ethanol at 60° for 1 h, followed with a 30% ethanol wash at room temperature. Samples were then stained with 2% aqueous uranyl acetate prepared in 30% ethanol, incubated at 4°C for 1 h. Serial dehydration was then continued at this point in increasing percentages of ethanol, following which cells were serially infiltrated with Durcupan ACM (Cat. No. 44610, Sigma). Fresh 100% resin was then added and polymerised at 60°C overnight. Ultrathin sections were cut on an ultramicrotome (UC6: Leica) and imaged at 100 kV on a JEOL1400 transmission electron microscope fitted with a Phurona EMSIS CMOS camera using a 5x5 stitching matrix. Light and electron microscopy images were subsequently overlayed and correlated using the Correlia plugin for Fiji/ImageJ (66) and is briefly described here. Using the signal from the orange fluorescent mitochondrial dye, a TEM section of the cell 65-nm thick was correlated with the best fit confocal z-slice within the realms of the confocal slice thickness of 150 nm. An overlay of the TEM image and corresponding confocal image showing the mitochondrial signal only at the relevant z-slice was created for cells that had taken up both A647-dextran and A647-EGF. Once the overlay confirmed the right cell volume, then the green signal from EGFP-tubulin and the red signal from A647-dextran or EGF-647 were overlayed onto the same TEM image and this enabled visualisation of the proximity of dextran and EGF endosomes with respect to MTs.

Curve fitting and data visualisation. For the intensity fits in Fig. 1c, the intensity values of tracked dynein molecules was plotted as a histogram. This distribution was fit by a sum of two Gaussians with the same standard deviation (σ). The smaller mean was constrained between 15-30 and the larger mean was constrained between 35-50. For the exponential fit in Fig. 1f, the probability vs residence time (τ) histogram of dynein on the MTs was obtained from the lengths of all single molecule binding events. From this $P(\tau) = 1$-cumulative frequency of the residence time was obtained, and was fit to $\lambda e^{-\lambda t}$ to obtain λ, the off-rate from MTs. The average residence time was given by $1/\lambda$. The fits in Figs. 7, 8 and 9 were similarly obtained using in-built Python functions. All data was plotted using MATLAB (MathWorks) or Python. Figures panels were prepared using Adobe Illustrator.

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