Using ephaptic coupling to estimate the synaptic cleft resistivity of the calyx of Held synapse.

Short title: **Electronic model of the synaptic cleft.**

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Abstract

At synapses, the pre- and postsynaptic cell get so close that currents entering the cleft do not flow

exclusively along its conductance, g_{cl}. A prominent example is found in the calyx of Held synapse in the

medial nucleus of the trapezoid body (MNTB), where the presynaptic action potential can be recorded in

the postsynaptic cell in the form of a prespike. Here, we developed a theoretical framework for ephaptic

coupling via the synaptic cleft, and we tested its predictions using the MNTB prespike recorded in

voltage-clamp. The shape of the prespike is predicted to resemble either the first or the second derivative

of the inverted, presynaptic action potential if cleft currents dissipate either mostly capacitively or

resistively, respectively. We found that the resistive dissipation scenario provided a better description of

the prespike shape. Its size is predicted to scale with the fourth power of the radius of the synapse,

explaining why intracellularly recorded prespikes are uncommon in the central nervous system. We show

that presynaptic calcium currents also contribute to the prespike shape. This calcium prespike resembled

the first derivative of the inverted calcium current, again as predicted by the resistive dissipation scenario.

Using this calcium prespike, we obtained an estimate for g_{cl} of ~1 µS. We demonstrate that, for a circular

synapse geometry, such as in conventional boutons or the immature calyx of Held, g_{cl} is scale-invariant

and only defined by extracellular resistivity, which was ~75 Ωcm, and by cleft height. During

development the calyx of Held develops fenestrations. We show that these fenestrations effectively

minimize the cleft potentials generated by the adult action potential, which might otherwise interfere with

calcium channel opening. We thus provide a quantitative account of the dissipation of currents by the

synaptic cleft, which can be readily extrapolated to conventional, bouton-like synapses.

Author Summary

At chemical synapses two neurons are separated by a cleft, which is very narrow. As a result, the concentration of released neurotransmitter can rapidly peak, which is essential for fast synaptic transmission. At the same time, the currents that flow across the membranes that face the synaptic cleft are also substantial, and a fraction of these currents will enter the other cell. We made an electronic model of the synaptic cleft, which indicated that if the shape and fraction of these currents are known, the electrical resistance of the synaptic cleft can be inferred. We tested several of the predictions of our model by comparing the membrane currents during a presynaptic action potential in a giant terminal, the calyx of Held, with the much smaller prespike evoked by these currents in the postsynaptic cell. We estimate the cleft resistance to be about 1 MΩ, which means that the changes in the cleft potential due to the membrane currents can become large enough to have an impact on voltage-dependent calcium channels controlling neurotransmitter release.

Introduction

When sir Sherrington introduced the term 'synapsis' [1], it was still unknown whether or not nerve impulses reach their target through direct electrical coupling [2, 3]. The visualization by electron microscopy of the synaptic cleft, the narrow gap between an axon and its target, put this debate to rest in favor of transmitter-mediated neurotransmission [4-6]. At about the same time, the electrical circuitry of the synapse was outlined, which included a cleft leak conductance (g_{cl}) linking synaptic currents with the interstitial space [7]. Since then, the purported role of the synaptic cleft for synaptic transmission has gradually become more important. Its small volume not only allows a rapid buildup of released neurotransmitter, but protons [8, 9], and potassium ions [10-12] may accumulate as well, whereas calcium ions may become partially depleted [13, 14]. Interestingly, during transmission a cleft potential may arise that may alter the dwelling time of mobile, charged particles within the cleft [15, 16], the distribution of ligand-gated ion channels [17], or the kinetics of voltage-gated channels in the cleft-facing membrane [18]. These features all stem from the small size of the synaptic cleft, which thus forms a barrier for particle movements [19].

Except for some specialized synapses, the contribution of g_{cl} to synaptic transmission has largely been ignored or deemed irrelevant, but its importance will of course depend on its magnitude and on the size of the currents that flow within the cleft. Eccles and Jaeger [7] already indicated that the maximal cleft current will be limited by this conductance. g_{cl} decreases as the path length l_{path} to the extracellular space increases ($g_{cl} \sim A/(R_{ex} \, l_{path})$), where R_{ex} is extracellular resistivity in Ω cm and A the area through which the current flows (see also ref. [20]). Due to their much longer path length, it has been argued that giant synapses have a lower g_{cl} than conventional synapses [21]. Remarkably, estimates of g_{cl} of very different synapses vary over four orders of magnitude (range 1.4 nS to 40 μS) without a clear relation to synapse size [7, 11, 18, 20, 22-30]. Despite the remarkable progress in our understanding of the synaptic cleft, we need a better theoretical understanding of its conductive properties, which is ideally confirmed by experimental observations.

One of the largest synapses in the mammalian brain is the calyx of Held synapse, which spans about 25

μm [31, 32]. Its large size matches its strength, with the maximal glutamatergic conductance exceeding

100 nS, and its fast presynaptic AP allowing firing rates >300 Hz [33]. Large currents, related to the

excitatory postsynaptic current (EPSC) and capacitive currents, need to flow within the cleft through g_{cl},

and under these conditions g_{cl} may limit synaptic transmission. A characteristic feature of this synapse is

that the presynaptic AP can be detected in the postsynaptic recording (Figure 1). This hallmark of the

calyx of Held synapse has been called the prespike [34]. It is probably due to ephaptic coupling of the

presynaptic AP via the synaptic cleft [30, 35], but the precise relation between the presynaptic AP and the

prespike has not been studied analytically. Here, we use the prespike to infer some key electrical

properties of this giant synapse. We provide a theoretical description of ephaptic coupling by defining the

electrical properties of the synaptic cleft, which is validated through dual patch-clamp recordings of the

calyx of Held and its postsynaptic target.

Results

The prespike is the ephaptically coupled calyx of Held AP recorded in its target neuron, a principal neuron

in the medial nucleus of the trapezoid body (MNTB) [34]. Some current clamp (CC) and voltage clamp

(VC) prespike examples are shown in Figure 1A-B. Figure 1C shows that the amplitudes of the CC and

the VC prespikes recorded in the same cell are correlated. Figure 1D shows examples of the CC prespikes

recorded during *in vivo* whole-cell recordings in young rats, whose shape resemble those recorded in the

simultaneous pre- and postsynaptic slice recordings (Figure 1A-B, 2A).

Figure 1 – Examples of prespikes. (A-B) Two examples of prespikes in simultaneous pre- and postsynaptic

patch-clamp recordings. From top to bottom: calyx of Held AP, which was elicited by afferent stimulation, CC

prespike, associated VC prespike. Recordings were made at room temperature in brainstem slices from 7-10

days old Wistar rats [36]. Horizontal scale bar: 0.5 ms. Vertical (from top to bottom): 50 mV, 2 mV, 1 nA.

Arrows indicate the amplitudes that were quantified in (C). (C) Relation between amplitude of the positive

peak of the prespike measured in CC and the amplitude of the first negative peak of the prespike measured in

VC in the same cell. Slope of the fitted line is -4.6 mV/nA. (D). Six examples of CC prespikes from *in vivo*

recordings of a target neuron of the calyx of Held in 3-5 days old Wistar rats [37]. During this period the calyx

of Held expands over the postsynaptic cell. Note the small upward deflection immediately preceding the EPSP

in the bottom three prespikes (arrow heads). Scale bar indicates 0.5 mV and 0.5 ms.

Figure 2. An electrical circuit for the prespike. (A) Drawing of the calyx of Held synapse. The calyx of Held

terminal ('pre') covers a large part of the soma of its target neuron ('post'). In between is the synaptic cleft (not

drawn to scale), a small space where membrane currents will flow during a presynaptic AP. (B) Electrical

circuit to reproduce the prespike. Presynaptic membrane currents can leave the synaptic cleft via the cleft

conductance (g_{cl}) or through the postsynaptic cell. It is assumed that the conductance from the edge of the cleft

to ground is relatively large compared to the cleft conductance itself [38, 39].

The prespike: a dissipation of cleft currents.

The prespike is generated by currents that flow from the calyx to the principal neuron in the

MNTB. While the current flow through the cleft depends on the synapse geometry, which we will

consider later in this article, we can already get a good understanding by analyzing a prespike-

generating, electrical circuit, similar to ref. [30], with the following components (Figure 2B): (1)

the terminal's release face, consisting of voltage-gated conductances, a capacitor (c_{pre}), and a

conductance (g_{pre}); (2) the synaptic cleft, which is separated from the reference by a leak

conductance (g_{cl}), and (3) a passive neuron, which partly faces the synaptic cleft. This part consists

of a capacitor (c_{syn}) and a leak conductance (g_{syn}). The remainder of the postsynaptic cell consists of

a capacitor ($c_{nsyn}=c_{post}-c_{syn}$) and a leak conductance ($g_{nsyn}=g_{post}-g_{syn}$). As the prespike is small, we

ignore the postsynaptic voltage-dependent conductances. We also did not include a direct resistive

connection between calyx and target since there is no physiological evidence for gap junctions at

the calyx of Held synapse [40], and capacitive currents are likely to dominate the electrical

coupling: the susceptance at 1 kHz across an area of 1000 μm^2 is 63 nS ($2\pi f C_\mathrm{m} A$, where f is the

frequency, C_m is the specific capacitance, and A the membrane area), which is much larger than the

postsynaptic resting conductance of ~4 nS.

Using Kirchhoff's law of current conservation ($i_\mathrm{leave} = i_\mathrm{enter}$), we define the equation for v_cl and

subsequently the equations for the prespike (see Appendix A1 for the derivations). The equation for

v_cl is $v_{cl}(g_{cl} + \sum g_{ion}) + (c_{pre} + c_{syn})dv_{cl}/dt = c_{pre}dv_{pre}/dt + \sum g_{ion}(v_{pre} - v_{ion})$. The right-

hand side shows the currents entering the cleft (i_enter), which are the presynaptic capacitive current

and the currents through the voltage-dependent ion channels. The left-hand side of the equation

shows the currents leaving the cleft (i_leave), which are generated by the cleft potential. This response

will lie between two extremes: (1) the capacitive dissipation scenario, where the cleft poses a high

resistance, resistive currents can be neglected, and the currents that leave the cleft are capacitive-

only, and (2) the resistive dissipation scenario where the cleft is highly conductive, and the currents

that leave the cleft are entirely resistive. For the capacitive dissipation scenario, we derive that

$v_{cl} = v_{constant} + (c_{pre} + c_{syn})^{-1} \int i_{enter} dt$; and v_cl then resembles scaled, integrated i_enter, and

therefore a scaled, presynaptic AP and the integrated voltage-gated currents. For the resistive

dissipation scenario, we derive in Appendix A that $v_{cl} = i_{enter}(g_{cl})^{-1}$; v_cl will then resemble a

scaled i_enter, which therefore should resemble the first derivative of the presynaptic AP.

The VC prespike represents the capacitive currents over the cleft-facing postsynaptic capacitance.

In CC, the prespike currents will change the postsynaptic membrane potential. The equations for

the currents entering and leaving the postsynaptic cell are $-i_{rec} = c_{syn}dv_{cl}/dt$ and

$c_{post}dv_{post}/dt + g_{post}v_{post} = -i_{rec}$, respectively. The shape of the VC prespike is entirely

determined by the shape of dv_cl/dt. We can derive how the VC prespike relates to i_enter for the two

extreme scenarios for the v_cl, For the capacitive dissipation scenario, we obtain $-i_{rec} =$

$c_{syn}\left(c_{pre} + c_{syn}\right)^{-1} i_{enter}$, and therefore the VC prespike resembles the cleft currents scaled by the

postsynaptic cleft-facing capacitance relative to the entire cleft-facing capacitance (Appendix A).

Due to the giant size of the calyx of Held synapse, the presynaptic and postsynaptic capacitance

facing the cleft are close to equal ($c_{syn}=c_{pre}=c_{cl}$) [41, 42]. This means that: $-i_{rec} = 0.5\, i_{enter}$, and the

other 'half' of i_{enter} returns through the presynaptic membrane. (In other words, v_{cl} reduces the

capacitive currents from the presynaptic AP (c_{pre} d(v_{pre}-v_{cl})/dt) by half, which dissipate capacitively

to the postsynaptic neuron.) For the resistive dissipation scenario, we obtain $-i_{rec} =$

$c_{syn}(g_{cl})^{-1} di_{enter}/dt$. Under these conditions, the VC prespike resembles the (inverted) second

derivative of the presynaptic AP scaled by the constant $c_{syn}c_{pre}/g_{cl}$, which for the calyx of Held

becomes c_{cl}^{2}/g_{cl}. We therefore conclude that if the VC prespike resembles the (inverted) first

derivative of the presynaptic AP, it indicates that cleft resistance is high, whereas if it resembles the

(inverted) second derivative of the presynaptic AP it indicates a relatively low cleft resistance.

Does the VC prespike resemble the first or second derivative of the presynaptic AP? We analyzed

simultaneous whole-cell presynaptic current clamp recordings with postsynaptic voltage clamp

recordings [36, 43]. In these recordings, the afferent axon of the calyx of Held was stimulated with

an electrode to trigger a train of presynaptic APs, and the postsynaptic responses were recorded

simultaneously (Figure 3A). The prespike generally matched the scaled, inverted second derivative

of the AP (Figure 3B), taking into account that it may have been filtered by the series resistance.

Figure 3C shows that the time course of the second derivative of the AP indeed provides a better

match for the prespike than its first derivative ($n = 13$ paired recordings). This is in agreement with

the idea that the cleft currents mainly dissipate resistively. During the train, the presynaptic AP

becomes smaller and slower [36, 44], resulting in a substantial reduction in the presynaptic

capacitive current amplitudes. To investigate the impact of these changes on the prespike, we

plotted the first two peaks of the prespike against the first two peaks of the first or the second

derivative (Figure 3D). In each case, the second derivative of the presynaptic AP gave a better fit

(Figure 3D-E, AP' versus AP'': Pearson's r = -0.967 ± 0.018 vs. -0.979 ± 0.018, paired T-test, t_{12} =

5.4, p = 0.00017). Furthermore, this fit should cross through the origin, since either $-i_{post} =$

$g_{cl}\tau_{cl}^2 d^2 v_{pre}/dt$ or $i_{post} = 0.5\, c_{pre}\, dv_{pre}/dt$ for the resistive dissipation or the capacitive

dissipation scenarios, respectively. The absolute deviation for the second derivative at 0 V/s^2 was

significantly less for every recording than for the first derivative at 0 V/s (Figure 3F, AP' versus

AP'': deviation = 122 ± 145 pA versus 20 ± 37 pA, paired T-test, t_{12} = 3.3, p = 0.0060). These fits

also yielded an estimate for $g_{cl}\tau_{cl}^2$, namely the slope from the fit of the second derivative and the

prespike amplitudes, which was 0.11 ± 0.07 fA s^2 V^{-1} (Figure 3G). These observations are thus in

agreement with the prespike resembling the inverted second derivative of the presynaptic AP,

suggesting that the cleft currents mostly dissipate through the cleft conductance.

Figure 3 – Relation between the calyceal AP and the prespike. (A) An example of a train of presynaptic APs

(v_{pre}) elicited by afferent stimulation and the postsynaptic voltage-clamp recording (i_{post}). A large EPSC is

elicited by the first AP, which rapidly depresses during the train. Stimulation artefacts were blanked. The first

AP (blue dot) and the last AP (green dot) are shown on an expanded scale in B. Horizontal scale bar: 50 ms.

Vertical scale bar: 90 mV, 1.5 nA. (B) The left and right column show the first and last AP of the train,

respectively. *From top to bottom*: the AP recorded in CC from the calyx of Held, its inverted first (-AP') and

second derivative (-AP''), and postsynaptic voltage clamp recording showing the accompanying prespike (i_{post}).

Vertical scale bars (from top to bottom panels): 40 mV, 0.3 kV/s, 2.5 MV/s^2, 0.2 nA. Horizontal: 0.5 ms. (C)

Relation between the delay between the first negative and positive peak of the prespike and the first two peaks

in the first (blue, circle) and second (black, triangle) derivative of the presynaptic AP of the first AP-prespike

pair in the train. Dashed line is the identity line. (D) Relation between the amplitude between the first negative

and positive peak of the prespike and the first two peaks in the first (blue, circle) and second (black, triangle)

derivative of the presynaptic AP of all AP-prespike pairs from an example train. The red lines indicate the

linear regression lines. Note how the regression line crosses closer to the origin for the second derivative than

for the first derivative. (E) Comparison for each train of Pearson's *r* for the regression lines. (F) Comparison of

the absolute deviation at 0 V/s (AP') versus 0 V/s^2 (AP''). (G) Cumulative distribution of the regression slope.

Paired T-tests: ** p<0.01, *** p<0.001.

The calcium prespike

Voltage-gated currents inside the cleft also contribute to v_{cl}. We explored how the presynaptic

calcium current would affect the prespike (Appendix 1.2 for the derivations). As the prespike

depends on the current entering the cleft i_{enter} irrespective of the origin of the current, a voltage-

dependent calcium current in the cleft should also contribute to the prespike. We call the part of the

VC prespike that is determined by the cleft calcium current, the calcium prespike. For the

capacitive dissipation scenario, we expect that the calcium prespike is $-i_{rec} = i_{ca}$. (Here, i_{ca} is

defined as the recorded current at the presynaptic terminal which is the net current through the

calcium channels and the cleft-facing capacitance, and by definition, this current dissipates across

the postsynaptic membrane.) For the resistive dissipation scenario, the calcium prespike is $-i_{rec} =$

$c_{syn}(g_{cl})^{-1}di_{ca}/dt$. Note that c_{cl}/g_{cl} equals a time constant, which we call the cleft time constant

$\tau_{cl}=c_{cl}/g_{cl}$. We can write $-i_{rec} = \tau_{cl} \, di_{ca}/dt$. The calcium prespike will either resemble an inverted

calcium current, or an inverted, scaled first derivative of the calcium current.

We turned to simultaneous pre- and postsynaptic whole-cell voltage clamp recordings of the calyx

of Held synapse that were originally reported in ref. [44], which were used to evaluate the presence

or absence of a calcium prespike. In these experiments the presynaptic compartment was voltage

clamped with an action potential waveform (APW) command while the presynaptic calcium

currents were pharmacologically isolated (Figure 4A). To isolate the calcium prespike from the

recorded prespike, we averaged the postsynaptic responses to the P/5 AP waveforms and subtracted

the upscaled, passive prespike from the 'active' prespike. The remaining currents indeed resembled

the first derivative of the calcium current (Appendix A4; Figure 4A-B, Figure S1), confirming not

only the existence of a calcium prespike, but also that its shape is in agreement with the cleft

currents dissipating through the different, resistive routes.

Figure 4 – Relation between the calcium current and the prespike. (A) *from top to bottom*: The presynaptic

voltage-clamp waveform (v_{pre}), the inverted second derivative of the waveform ($-v_{pre}''$), the presynaptic calcium

current (i_{Ca}; after P/5 subtraction), its inverted first derivative ($-i_{Ca}'$), the active (black, thick) and passive

prespike (i_{post}; magenta, P/5-scaled), and the calcium prespike which is the difference between the active and

the passive prespike (Δi_{post}). The passive prespike is generated by injecting the presynaptic AP at 1/5th of its

amplitude. While this generates passive currents across the presynaptic cleft-facing membrane, it does not elicit

the presynaptic calcium current. Voltage-gated sodium and potassium currents were blocked (see Methods).

Vertical scale bars: 110 mV, 7 MV/s^2, 2 nA, 15 nA/ms, 0.15 nA, 0.1 nA. Horizontal: 0.5 ms. (B) Relation

between the inverted first derivative of the presynaptic calcium current and the calcium prespike (black dots).

Maximum correlation was obtained by introducing a delay of 60 µs for the first derivative. The slope of the

regression line (magenta) is shown in the bottom right corner.

From the analytical derivations for the resistive dissipation scenario, we know that the presynaptic

calcium current and the calcium prespike are related by the cleft time constant τ_{cl}. To estimate τ_{cl},

we first analyzed the presynaptic calcium current that was evoked by the APW. It had a peak

amplitude of -1.9 ± 0.6 nA and a full-width at half-maximum (FWHM) of 0.51 ± 0.05 ms ($n = 5$

paired recordings). Next, we obtained the values for the calcium prespike. Its peak-to-peak

amplitude was 80 ± 60 pA and its peak-to-peak delay was 0.54 ± 0.11 ms. This peak-to-peak delay

should be about twice τ_{Ca}, which was thus slightly higher than the 0.21 ms reported in ref. [45]. Its

time course closely matched our prediction for the resistive dissipation scenario (Figure 4,

Appendix A2, Figure S1). The amplitude of the calcium current was linearly related to the calcium

prespike peak ($r = 0.98$, p = 0.003); the calcium prespike peak was 1.9 ± 0.7% of the calcium

current amplitude. In Appendix A2 we derived that this percentage should be $\tau_{cl} \tau_{Ca}^{-1} e^{-0.5}$, yielding

an estimated τ_{cl} of 9 ± 5 µs. As an alternative approach, we plot the first derivative of the

presynaptic calcium current against the calcium prespike. Its maximal slope (allowing for a slight

time shift) provided a τ_{cl} of 5.7 ± 1.9 µs (Figure 4B and 9B). The two methods thus gave

comparable values for the cleft time constant.

The ratio of cleft capacitance c_{cl} and cleft leak conductance g_{cl} constitutes τ_{cl}. If we assume that c_{cl}

is in the order of 10 pF, the estimated g_{cl} becomes about 0.5-1 µS. From the circuit analysis we

know that the capacitive VC coming from both the presynaptic AP and the calcium prespike from

the presynaptic calcium current do not scale the same way to c_{cl} and g_{cl} (*cf.* Appendix A1 with A2).

This difference allows us to further define these two parameters. From the calcium prespike we

obtain τ_{cl}. As the equation for the capacitive prespike can be written as $-i_{post} = g_{cl}\tau_{cl}^2 d^2 v_{pre}/dt$,

we find that we can calculate g_{cl}. The first peak of the VC prespike had an amplitude of -120 ± 110

pA, which corresponds to the capacitive currents. We can then calculate a g_{cl} of 2.6 ± 0.8 μS and a

c_{cl} of 24 ± 16 pF.

For τ_{cl} we assumed that the calcium current entirely originated from the membrane facing the cleft.

These estimates should be considered as a lower bound as some of the current may not arise from

the cleft-facing membrane. Immunolabeling for calcium channels identified a fraction of the N-type

and R-type calcium channels to be in the non-cleft-facing membrane [46]. If we assume that only

two-third of the calcium current originated from the cleft, our estimate of τ_{cl} increases to 13.5 ± 7.8

μs. This estimate could be considered as an upper bound, as less than 50% of the calcium current

was mediated by N-type and R-type channels, and the other, P/Q-type, calcium channels were

mostly associated with the release face [46]. When we assume that 2/3 of the total presynaptic

calcium current originated from the cleft, we calculate a g_{cl} of 1.2 ± 0.4 μS and a c_{cl} of 16 ± 11 pF.

With these values $g_{cl}\tau_{cl}^2$ was 0.07, 0.08, 0.11, 0.28 and 0.82 fA s^2 V^{-1} (0.27 ± 0.32 fA s^2 V^{-1}), which

is in agreement with Figure 3G. Based on our limited data set, we determine that g_{cl} ranges between

1-2 μS.

No evidence for a potassium or a sodium prespike

Besides the capacitive currents and the voltage-gated calcium channels, other voltage-gated

channels facing the cleft may also contribute to the cleft currents. We therefore explored the

potential contribution of Na$^+$ and K$^+$ currents to the prespike, the sodium and potassium prespike,

respectively. We estimated their density based on literature values, and assumed a homogeneous

distribution within the terminal. As voltage-gated channels in the cleft-facing membrane sense v_{pre}-

v_{cl} instead of v_{pre}, their gating depends on v_{cl}. This feedback effect of v_{cl} could not be incorporated

in our calculations. Instead, we therefore simulated the impact of voltage-gated currents on the

prespike by numerically solving ordinary differential equations. Figure 5 illustrates that the

presence of these channels at the release face may have a large impact on the shape and size of the

prespike. Due to the differences in their gating properties, the timing of each current during the

presynaptic AP is quite different (Figure 5A), and the same holds for their impact on v_{cl} and on the

prespike (Figure 5B). The prespike examples shown in Figure 1 did not resemble the prespike of

the homogeneous model, which included Na^+, K^+ and Ca^{2+} currents at the cleft-facing membrane,

suggesting that some voltage-gated ion channels are excluded (at least functionally) from the

release face. In particular, the voltage-gated sodium conductance generated a current that partially

counteracted the first prespike peak. Such rapid changes in the prespike during the rising phase of

the presynaptic AP would point at a contribution by sodium channels. However, consistent with

previous observations that the terminal is mostly devoid of sodium channels [47-49], these fast

changes were not observed in the prespike. We therefore conclude that there is no substantial

contribution of sodium channels at the release face to the prespike.

Figure 5 – Impact on the prespike of voltage-gated ion channels located at the release face. (A) Timing of

different voltage-gated currents at the release face during the AP. Note that the amplitudes of each of the

currents depend on cleft potential, thus creating a mutual dependency. Vertical scale bar: 60 mV, 2.4 nA, 1 nA,

1 nA, 3 nA. Horizontal: 0.5 ms. (B) The cleft potential (top), the VC prespike (middle), and the CC prespike

(bottom) in the presence or absence of different voltage-dependent ion channels at the cleft. All present:

homogeneous density for each channel; no VGSCs: no sodium conductance at the release face; only VGCCs:

both calcium and capacitive currents at the release face; no channels: the capacitive-only model. The

conductances that are present in the model are indicated by their currents at the top. Arrow head indicates the

small deflection in the CC prespike that is generated by the calcium prespike. Scale bars: 4 mV, 0.2 nA, 2 mV.;

0.5 ms, except for the homogeneous model where the vertical scale bar corresponds to 1.6 nA.

Potassium channels do not start to open until around the peak of the presynaptic AP. They therefore do not

impact the first part of the VC/CC prespike. The potassium current would contribute to a more positive

synaptic cleft potential, which is opposite to the effect of the calcium current. If the potassium current at

the cleft were significantly larger than the other currents, we would expect a fast inward current in the VC

prespike after the first positive peak, and the CC prespike would become a single positive deflection.

However, CC prespikes are typically biphasic (Figure 1), therefore the contribution of the potassium

current is unlikely to outweigh the contribution of the other currents.

Lastly, we revisit the impact of calcium channels on the prespike. The observation of a calcium prespike

already demonstrated the presence of these channels at the release face. The calcium current occurs during

the repolarization phase of the presynaptic AP [45], and therefore contributes little to the first two peaks of

the VC or to the first peak of the CC prespike. Its contribution is easily obscured by noise or by the EPSC,

which starts rapidly after the calcium-induced glutamate release. Nevertheless, it may generate a small,

upward inflection in the CC prespike (Figure 5B) and such an inflection can sometimes be observed

(Figure 1D).

These simulations show that the prespike waveform may contain information about the presence of

voltage-dependent ion channels at the presynaptic membrane facing the cleft. This means that an

experimental study of the impact of specific channel blockers on the prespike waveform may be

informative about their presence at the release face.

Voltage-dependent kinetics of cleft-facing calcium channels

During development the presynaptic AP narrows from a half width of 0.5 to 0.2 ms [50]. The briefer AP necessitates a much larger capacitive current, which, all else being equal, would generate a larger change in v_{cl}, which may impact the calcium channel kinetics. We again simulated the impact of v_{cl} on the calcium current by solving ordinary differential equations for different values for g_{cl} (Figure 6). For a slow AP, a g_{cl} of 0.67 µS and a c_{cl} of 10 pF, v_{cl} first rises to a maximum of +3.9 mV, and then decreases to a minimum of -4.2 mV. While v_{cl} has little effect on the current amplitude (Figure 6B), it generates a delay of 45 µs in the onset of the calcium current (Figure 6C). In contrast, under the same conditions we observe for a fast AP that v_{cl} rises to a maximum of +16.4 mV, followed by a decrease to a minimum of -8.2 mV. As a consequence, the peak calcium current was delayed by 21 µs and reduced by 19% (-2.0 nA vs -1.7 nA). These effects became even more pronounced at lower g_{cl}. The observed delay corresponded to the delay in reaching the transmembrane $v_{pre} - v_{cleft}$ maximum compared to the v_{pre} maximum; the reduction in calcium current was due to a smaller effective AP, which reduced the maximal open probability of the calcium channels. We conclude that developmental acceleration of the calyceal AP will increase the capacitive current to a level that potentially could delay and impede the presynaptic calcium current substantially if g_{cl} were <1 µS.

Figure 6 – Potential impact of the cleft potential on presynaptic calcium currents. (A) Simulation of the

impact of a slow and a fast AP (top; FWHM: left 1.0, right 0.2 ms) on the cleft potential (second row), the

transmembrane potential sensed by the calcium channels in the cleft (third row) and the presynaptic calcium

current (bottom) at three leak resistances ($1/g_{cl}$: 0 MΩ, 1 MΩ and 5 MΩ). Scale bar: 200 mV, 40 mV, 50 mV, 1

nA. (B) Relation between peak calcium currents and leak resistance for the two APs. Driving force slightly

increases (not shown). (C) The change in delay between the calcium peak current and the AP for the two APs

(solid line). The dashed line shows the delay after correcting for the decrease in the calcium current shown in

B. Calcium conductance density was adjusted to 0.2 nS μm^{-2} and 2.8 nS μm^{-2} for the slow and fast AP,

respectively, in order to elicit a 2 nA-current at 0 MΩ cleft leak resistance. The apposition area was 1000 μm^2.

Other conductances were set to 0 nS.

Relation between synapse and prespike size

So far, we modeled the cleft as a single, equipotential compartment with capacitive and conductive

properties. However, it is expected that the voltage profile within the cleft will depend on synapse

geometry [20], and that larger synapses have larger prespikes [30].We therefore need an analytical

description of the voltage profile within the cleft. Previous work considered the impact of cleft geometry

on the EPSC, a current that runs radially through the cleft and enters at a centered region containing the

glutamate receptors [7, 20, 30]. Our situation is different as we consider a homogeneously distributed

current: the capacitive currents arise at every membrane patch facing the cleft. For an analytical

description we start by assuming a radially-symmetric synapse with a radius r and a cleft height h. Within

this synapse we only consider capacitive currents from the cleft-facing membrane (Figure 7A). As only a

small fraction of these currents enters the postsynaptic cell (the resistive dissipation scenario), we simply

assume that all currents run radially to the synaptic membranes. The electric field at location x relative to

the cleft center will be $dv_{cl}/dx = xR_{ex}C_m(2h)^{-1}dv_{pre}/dt$ (Appendix B1.3) where R_{ex} is the

extracellular cleft resistivity and C_m is the specific membrane capacitance. The electric field grows

linearly with the distance from the center. The cleft potential is equal to $\int_x^r dv_{cl}\, dx$, which gives $v_{cl} =$

$(r^2 - x^2)R_{ex}C_m(4h)^{-1}dv_{pre}/dt$ (Appendix B1.4). The size of the cleft potential thus depends on r^2,

indicating that for larger terminals the cleft potential may become increasingly important (Figure 7B). The

voltage profile also shows that not all voltage-gated channels in the cleft will be affected, but only the

channels located towards the cleft center.

The temporal changes in the cleft potential will generate a capacitive current at the postsynaptic side. To

obtain the VC prespike, we derive $i_{post} = r^4\pi R_{ex}C_m^2(8h)^{-1}d^2v_{pre}/dt^2$ (Appendix B2.2). The VC

prespike therefore grows with the fourth power of the radius, which explains why prespikes are only

observed in giant synapses such as the calyx of Held synapse. With our geometrical relation for the VC

prespike we can relate g_{cl} to extracellular resistivity: $g_{cl} = 8\pi h(R_{ex})^{-1}$ (Appendix B3). This relation

leads to an estimate for the intra-cleft resistivity of 75 Ω cm for an h of 30 nm [41] and a g_{cl} of 1 µS.

Considering that there may be extended extracellular spaces within the synaptic cleft [51], and taking into

account the variation in g_{cl} (1-2 µS), we estimate an R_{ex} in the range of 50-100 Ω cm.

Figure 7. Synapse geometry determines the profile of the cleft potential. (A) Drawing of a radially-symmetric synapse with cleft height h and radius r. The capacitive currents will generate a cleft potential v(x) at location x. (B) Cleft voltage profile for synapses with different radii (color-coded) at the first peak in the capacitive current. The horizontal axis depicts the location from the center (0) relative to the edge (1.0). In (B), R_{ex} is 100 Ω cm, C_m is 1 μF cm^{-2}, and h is 30 nm.

Fenestration of the calyx of Held reduces cleft potentials.

Giant synapses typically fenestrate during development [52-55]. The relation between fenestration and the coupling between pre- and postsynaptic membranes has been studied in numerical simulations [30]. We initially assumed that fenestration minimizes v_{cl} by reducing g_{cl} through reducing the path length to the extracellular fluid (r). However, we analytically derived that for a radially-symmetric synapse, g_{cl} does not depend on its radius. The relation between the conductance and the resistivity is $g_{cl} \sim A\,(r\,R_{ex})^{-1}$, where A is the area through which the current escapes. As A of a radially-symmetric cleft is $2\pi rh$, when r is doubled, A is also doubled, leaving the conductance independent of synapse size. The cup-shape of the young calyx makes the radially-symmetric synapse a reasonable approximation. This also means that our estimate of the resistivity of the synaptic cleft is relevant for smaller synapses with a comparable radially-symmetric shape.

Following a period of fenestration, the calyx of Held becomes, as its name implies, a synapse with extended fingers [31, 56]. However, for finger-like synapses with a sheet-like or cylindrical cleft, the geometrical relations are slightly different. For these synapses we assume that the currents only run

transversally, *i.e.* perpendicular to the longitudinal axis. The cleft potential then becomes $v(x) =$

$(r^2 - x^2)(2h)^{-1}R_{ex}C_m dv_{pre}/dt$ (Appendix B4.4), where $2r$ equals the width of the sheet (or r the length

of the cylinder). As the escape route for the current is more restricted than for a radially-symmetric

synapse, its voltage gradient is steeper. However, the diameter of a calyceal finger will be much less than

the diameter of the juvenile calyx, effectively reducing the cleft potential. For a calyceal finger of 4 μm

diameter [31, 57] the cleft potential at the center of the finger-like synapse would be only ~2.5% of the

cleft potential in the cup-shape calyceal synapse (equation B5.2). As a consequence, the VC prespike

would become ~3.3% assuming that the total contact surface is conserved (equation B5.4). The cleft

potential and VC prespike with more complex synapse morphology may be investigated using finite-

element modeling, but this is beyond the scope of this study.

Fenestration will effectively reduce the amplitude of the cleft potentials, and therefore we expect that the

prespike in the adult calyx of Held synapse should be significantly smaller than the juvenile prespike. To

study the adult prespike, we revisited *in vivo* whole-cell current-clamp recordings of MNTB principal

neurons from wild-type C57BL/6J mice previously reported in refs. [58, 59]. CC prespikes consisted of a

small positive deflection, preceding the start of the rapid rising phase of the EPSP by about 0.3 ms. In 11

of 12 recordings, the amplitude of the positive deflection was <0.2 mV. In contrast to the juvenile

recordings (Figure 1), the positive peak was not followed by a negative peak. An example is shown in

Figure 8. The good correlation between the size of CC and VC prespikes (Figure 1C) in combination with

the absence of a large decrease in the input resistance of adult principal neurons [58, 59] suggests that the

VC prespike can be expected to also be relatively small. Our anecdotal evidence from *in vivo* whole-cell

recordings in adult mice therefore confirms that the prespike is smaller than the juvenile prespike and that

it possibly has a different shape (*cf.* Figure 1 and 8).

Figure 8 – Adult prespike. (A) Black traces: Spontaneous, sub- and suprathreshold events recorded during a

whole-cell *in vivo* recording of a principal neuron in the MNTB from a P52 C57BL/6J mouse [58]. Events were

aligned on the start of their EPSP and (<1 mV) differences in the resting membrane potential between events

were subtracted for the display. Minimum interval between events was 20 ms. Red trace: average of >1000

events. Resting potential was -64 mV. Arrow points at the small prespike. Vertical scale bar: 10 mV. Horizontal:

1 ms. (B) As (A), showing close-up of the prespike (arrow). Vertical scale bar: 0.2 mV. Horizontal: 0.2 ms.

Discussion

Here, we studied how electrical currents that are needed to depolarize a presynaptic terminal or trigger transmitter release are dissipated in the synaptic cleft. We derived equations for the ephaptic coupling between the presynaptic terminal and the postsynaptic cell. We found that the capacitive component of the prespike is closely approximated by the second time derivative of the presynaptic action potential. Its size scales with the fourth power of the radius of the synapse, explaining why intracellularly recorded prespikes are uncommon in the CNS. We showed that presynaptic calcium currents can contribute to the prespike and that their contribution is closely approximated by the scaled first derivative of these currents. We used their contribution to infer an estimate of the cleft conductance, which was ~1 µS for the juvenile calyx of Held synapse. We showed that the cleft conductance of a radially-symmetric synapse does not depend on its radius, but is entirely defined by its height and the extracellular resistivity. This led to an estimate of the cleft resistivity that is at most 50% larger than of the interstitial fluid, which means that the cleft conductance is large enough to allow for the large EPSCs of the calyx of Held, but small enough to have a potential impact on the opening of presynaptic voltage-dependent calcium channels. Finally, we showed that developmental fenestration of the calyx of Held constitutes a very effective adaptation to minimize synaptic cleft potentials.

Electrical properties of the calyx of Held synapse

Based on our electronic model, we discriminated two extreme scenarios, a scenario in which currents entering the cleft dissipate via the membrane capacitance, and one in which they exclusively dissipate via the cleft conductances. We compared the shape of both the passive and the calcium component of the prespike with both scenarios, and found that the resistive dissipation scenario provided a much better prediction for the shape of the prespike, since the passive prespike resembled the inverted, second derivative of the presynaptic AP much better than the inverted, first derivative, whereas the calcium prespike resembled the inverted, first derivative of the presynaptic calcium currents much better than the

inverted calcium currents themselves. There are several other compelling arguments for this scenario. For

the calyx of Held synapse to be able to function as a relay synapse in the auditory system, it needs to be

both fast and precise, which means a fast presynaptic AP and a large size with many release sites and large

EPSCs. We derived for the capacitive dissipation scenario that the size of the prespike would be about half

of the size of the presynaptic currents entering the cleft during the presynaptic AP. Because of the large

size of the calyx of Held and the fast time course of the presynaptic AP, these currents are in the nA range.

The observed size of the VC prespike is therefore clearly not in line with the predictions of the capacitive

dissipation scenario. Secondly, the calyceal EPSC is very fast and can easily be 10 nA [36, 50, 60-62].

This EPSC should flow not only through the postsynaptic glutamate receptor channels, but also through g_{cl}

[7]. For the synapse to be able to conduct an EPSC of 10 nA at a driving force of 100 mV, the cleft leak

conductance must be at least >0.1 μS. This is consistent with the observed range for g_{cl}. These arguments

are therefore in line with our conclusion that the resistive dissipation scenario provides a much better

description than the capacitive dissipation scenario. Less obviously, the model allowed us to use the

calcium component of the prespike to obtain an estimate of the cleft conductance of about 1 μS. This

value is subject to some uncertainty. For the calculation we assumed that most, but not all, presynaptic

calcium channels are facing the synaptic cleft. The juvenile calyx contains N- and R-channels, which are

not located near the release sites based on immunostaining and their lower efficacy at triggering release,

which is in contrast to the P/Q-type calcium channels [46]. Our assumption that only two-third of the

calcium current originated from the cleft is therefore reasonable, considering that the N- and R-channels

make up about half of the calcium currents of the juvenile calyx.

Even though we calculated the AP-related capacitive currents that run through the cleft-facing membrane,

we did not record the AP itself at the cleft-facing membrane. The AP passively invades the terminal [47],

and therefore may be slowed down and have a reduced amplitude at the cleft-facing membrane. We

consider it unlikely that this filtering is significant for two reasons. First, passive filtering from one end to

the other end of the calyx reduces the AP peak by only 6 mV with a delay of <0.2 ms [63]. Second, the

axonal AP is able to invade the fenestrated endings of the mature calyx [64], which should have a

significantly higher impedance than the cup-shaped calyx. Our calculation of the presynaptic capacitive

currents as the rate of change of the recorded calyceal AP relative to the cleft potential is therefore likely

to be accurate.

Comparison of cleft electrical properties across synapses

Interestingly, with similar values for cleft height and extracellular resistivity other studies came to a wide

range of estimated values for the cleft conductance, both at least an order of magnitude smaller than 1 μS

[11, 26, 28-30, 65], but in some cases also considerably larger than 1 μS [7, 66]. In some of the synapses,

an unusually large and narrow path to the interstitial fluid was responsible for the low estimates. Examples

are mossy fiber boutons penetrated by a thin (<1 μm) spinule [28], the unusually long calyx-type

vestibular synapse [11], or the cone-to-horizontal cell synapse, for which a large part of the horizontal cell

is located within an invagination of the cone photoreceptors [67]. While cleft resistivity should be

independent of morphology, the cleft conductance for membrane currents at the deep end of these

synapses would be significantly less than for the homogeneously distributed current addressed here. In

these morphologically complex synapses, slow removal of ions or neurotransmitter may serve a function,

for example in the case of vestibular hair cells, where potassium ions accumulate [10, 11], or in the cone

photoreceptor synapse, where prolonged effects of glutamate are important, and where postsynaptic

currents are typically relatively small [68]. At the other end of the range lies the high estimate for the cleft

conductance of the neuromuscular junction, to which the postjunctional folds were suggested to contribute

[7]. The unique morphological properties of these synapses have therefore made a large contribution to the

large range of estimated values for the cleft conductance.

Our calculations indicated that for homogeneously-distributed currents the associated cleft conductance is

fully determined by the extracellular resistivity, the cleft height and the morphological type of the synapse,

but not by the size of the synapse. The cleft height is constrained within narrow boundaries because of the

opposing requirements of low volume for high neurotransmitter concentrations and high conductance to

allow large synaptic currents [16]. The cleft height typically ranges across synapses between 10-30 nm,

but there is some uncertainty regarding the impact of chemical fixation. Assuming a cleft height of 30 nm,

we estimated a resistivity of 75 Ω cm based on the prespike, only about 30% larger than of cerebrospinal

fluid or Ringer solution [16, 69]. While we are not aware of other estimates of the extracellular resistivity

within the synaptic cleft, our estimate is similar to the intracellular resistivity of neurons [70-72] and the

periaxonal resistivity of myelinated axons [72]. The extracellular resistivity can also be estimated based on

the diffusional mobility of small molecules. Our estimate agrees with those made with an imaging method

in hippocampal mossy fiber synapses, which also estimated that small molecules move on average about

50% slower than in the Ringer solution [73]. Estimates of neurotransmitter diffusion constants within the

synaptic cleft range from 1.5 to 4 times smaller than in free solution [74-76], but for neurotransmitters

diffusion within the synaptic cleft may be slowed down by neurotransmitter-binding proteins. We

therefore suggest that a cleft resistivity of ~75 Ω cm may apply to other synapses as well. Given the range

of cleft heights and the estimated cleft resistivity, the cleft conductance can be inferred for a distributed

current using our analytical relations (B3.2). It should be not so different from the calyx of Held

considering that cleft height is conserved across synapses, while the geometry of the juvenile calyx

translates well to bouton synapses by simple downscaling. Cleft conductances for circular synapses would

then range between 0.3-1.0 μS.

The prespike: what is it good for?

The extracellular AP of the calyx of Held can be observed in extracellular recordings at the target neuron

[77], and the characteristic shape of the complex waveform consisting of the presynaptic AP, the EPSP

and the postsynaptic AP has facilitated identification of *in vivo* single unit recordings at the MNTB [59].

Importantly, it was shown that the calyceal AP can be picked up in postsynaptic whole-cell recordings as

well [34]. This extracellular field potential outside the cleft cannot generate a 0.1 nA-prespike amplitude

intracellularly without assuming an obviously non-physiological postsynaptic admittance (~0.2 μS).

Incidentally, this also indicates that other compartments, such as the presynaptic heminode, contribute

very little to the generation of the prespike. In postsynaptic voltage clamp recordings the prespike showed

a resemblance to a high-pass filtered version of the presynaptic AP [36], but we showed here that it is

closer to the second derivative of the AP. At the related endbulb of Held synapse, the presynaptic AP is

also readily picked up in extracellular recordings [78]. However, no prespikes were reported during

postsynaptic whole-cell recordings [79] possibly due to the smaller size of the endbulb compared to the

calyx of Held. In the visual system, prespikes can be seen in a subset of lateral geniculate neurons [80].

However, in many synapses in which the presynaptic AP is readily observed in extracellular recordings, it

is not reported in intracellular postsynaptic recordings [18, 81-85], attesting to the efficiency of the

synaptic cleft in dissipating these currents. Interestingly, especially in the hippocampus, neurons can

display so-called spikelets with an amplitude of mV in intracellular recordings. They have a biphasic

shape resembling the calyceal prespike recorded in current-clamp. These spikelets are thought to result

from ephaptic, capacitive coupling with one or more nearby neurons [86]. The equations we derived here

should be applicable to these spikelets as well: we predict them to resemble the second derivative of the

AP in the nearby neurons [87], and to critically depend on contact area, but also on proximity and on the

resistivity of the extracellular fluid.

We believe that the major reason for the lack of prespikes at most synapses is that the prespike amplitude

is proportional to the 4th power of the synapse radius (Appendix B2). During the postnatal development of

the MNTB, it can therefore be argued that the absence of postsynaptic recordings with multiple prespikes

cannot be used as firm evidence against the presence of multiple calyces on these cells. We previously

found anecdotal evidence for the presence of multiple large inputs during early postnatal development

[88]. In a small fraction of the recordings multiple large inputs were observed, one with and one without a

prespike. It might be that the one without the prespike had a smaller surface area or a higher cleft height,

going below the radar. This highly nonlinear dependence on apposition area may thus provide an

explanation for the discrepancy between anatomical evidence favoring multiple large inputs during

development [89] and the lack of recordings with multiple prespikes.

Divide and conquer

Virtually all synapses have in common that neurotransmission is initiated by a presynaptic action potential

that results in capacitive cleft currents. The size of these currents depends on the amplitude and speed of

the presynaptic AP. The properties of the presynaptic APs have mostly been studied at giant synapses,

where they tend to be large and fast [50, 90-95]. These giant synapses are typically relay synapses that are

specialized for rapid signaling, but recent measurements at more conventional bouton synapses in the

neocortex show that their presynaptic action potentials are large and fast as well [96]. These capacitive

currents, together with the presynaptic ion channel currents, will lead to a change in the cleft potential

whose time course will closely follow the time course of the presynaptic membrane currents, and whose

amplitude will critically depend on the cleft resistivity. Because of this cleft potential, the presynaptic ion

channels facing the cleft will sense a reduced action potential. We showed that this should result in a less

effective and delayed opening of the presynaptic calcium channels facing the synaptic cleft. Since rapid

and large changes in the voltage sensed by the presynaptic ion channels serve to rapidly turn transmitter

release on and off again, large cleft potentials seem undesirable, even though they may serve other, more

beneficial functions such as influencing the neurotransmitter mobility within the cleft [15, 16], or rapidly

inhibiting postsynaptic action potential generation [18, 97].

Above, we argued that many radial-like synapses should have a cleft conductance similar to what was

measured here. An important difference between small and large synapses, however, is that the total

current will be much larger in the large synapse, as the capacitive current scales with the presynaptic

apposition area. A large synapse thus has to dissipate this much larger current over a cleft conductance

that is similar to the cleft conductance of a small synapse. There are three obvious different neural

strategies to mitigate significant cleft potentials. Firstly, a strategy may be widenings of the synaptic cleft

in between synaptic junctions. These intrasynaptic space enlargements [4], also called channels of

extracellular spaces [98, 99] or subsynaptic cavities [100], have been found at many synapses, including

the calyx of Held synapse, and may increase cleft conductance. Interestingly, a knockout of brevican, a

component of the extracellular matrix of perineuronal nets, leads to a reduction of the size of the

subsynaptic cavities at the calyx of Held synapse and an increased synaptic delay, even though the relation

with a change in cleft conductance remains to be established [100]. The second strategy is to divide the

cleft membrane currents over many postsynaptic junctions. For example, in the case of the squid giant

synapse, the postsynaptic axon has numerous processes contacting the presynaptic axon, each with a width

that is only ~1 μm [101]. At each of these processes the cleft potential will be small, thus effectively

mitigating its impact. The last strategy is to fenestrate presynaptically. Most large synapses fenestrate

presynaptically during development, leading to extended extracellular spaces [99]. Well-known examples

include the neuromuscular junction, where the terminal arbor assumes a pretzel shape [102], the calyx of

Held, which develops from a cup shape to the calyx shape [52, 53], the endbulb of Held synapse [54] or

the avian ciliary ganglion [55]. In Appendix B5 we quantified the impact this will have on the size of the

cleft potential change and on the prespike. Previously, other roles for these adaptations have been

emphasized, such as speeding up of transmitter clearance or minimization of crosstalk between

postsynaptic densities, both for the calyx of Held synapse [52, 103, 104] and at other synapses [105, 106].

The digitated adult morphology of the calyx of Held is accompanied by an increase in the number of

puncta adhaerentia [104] and the formation of swellings connected via necks to the fingers [42, 51]. In the

adult calyx of Held, a large fraction of the release takes place at these swellings, which should facilitate

the dissipation of capacitive currents by the cleft, and reduce the impact of the cleft conductance on the

postsynaptic currents.

In conclusion, we show here that the prespike of the calyx of Held synapse can be used to infer critical

properties of its synaptic cleft; apart from the cleft conductance, it may also provide information about the

presence or absence of ion channels at the cleft, or the effectiveness of morphological changes in reducing

the size of cleft potentials during development. We thus provide a quantitative description for how the

synaptic cleft, despite its narrow, limited space, can sustain the large and fast currents that are at the heart

of chemical neurotransmission.

Materials & Methods

The AP template shown in Figure 4A was based on an AP template previously used in voltage-clamp

experiments [43].

For the experimental procedures and recording conditions for the dual recordings of the calcium current

and the postsynaptic currents, we refer the reader to the original paper in which these recordings were

reported [44]. For the current paper, we used trains of 45-50 identical APWs at an interval of 100 Hz.

These APWs were what was called the 'control AP' in ref. [44], i.e. the first AP of a train of presynaptic

APs. To limit the possible impact of EPSCs on the prespike during the trains, the first 20 APWs were

skipped, and the responses to only the next 25 APWs were averaged. The presynaptic passive (capacitive)

component was estimated by presenting the same train five times at 5x smaller amplitude, and summing

the responses (P/5 method). The capacitive component of the prespike was isolated analogously from the

same reduced APW train. *In vivo* whole-cell current-clamp recordings of developing MNTB principal

neurons with afferent stimulation were previously reported in ref. [37] ($n = 6$). *In vivo* whole-cell current-

clamp recordings of MNTB principal neurons from adult C57BL/6J mice were previously reported in ref.

[58] ($n = 7$, only wild-type animals) or in ref. [59] ($n = 5$). We also refer to these papers for the respective

ethics statements.

For numerical simulations with voltage-gated channels the following two differential equations were

defined, following [30]:

621
$$c_{post} \frac{dv_{post}}{dt} = c_{cl} \frac{dv_{cl}}{dt} - g_{post}(v_{post} - v_{rest})$$

622
$$2c_{cl} \frac{dv_{cl}}{dt} = c_{cl} \left(\frac{dv_{pre}}{dt} + \frac{dv_{post}}{dt} \right) - g_{cl}v_{cl} + \sum g_{max}p_{open}(v_{pre} - v_{cl} - v_{channel})$$

Where g_{max} represents the maximal conductance of a voltage-gated channel, p_{open} the open probability of

the voltage-gated channels, which is calculated using a multiple-state model (see below), and $v_{channel}$ the

reversal potential of the voltage-gated channel. The reversal potentials were 50 mV, -90 mV, 40 mV and -

70 mV for the sodium conductance, potassium conductance, calcium conductance, and the resting

membrane potential of the postsynaptic neuron, respectively. v_{pre} was defined by an AP waveform.

Conductance densities in Figure 5 were 0.2 nS μm^{-2} for the sodium conductance, 0.08 nS μm^{-2} for the

high-threshold potassium conductance, 0.24 nS μm^{-2} for the low-threshold potassium conductance, and

0.04 nS μm^{-2} for the calcium conductance with a cleft-facing membrane area of 1000 μm^2.

Voltage-gated conductances were modeled with multiple states following ref. [107] except for the calcium

conductance. In the models C indicates a closed state, O the open state and I the inactivated state. The

sodium channel followed a five states-model, the high-threshold and low-threshold potassium channel

were modeled as a six-states channel, and the calcium channel as a two-state channel for which the open

probability was squared:

Five states-model: $$C_1 \underset{\beta}{\overset{2\alpha}{\rightleftharpoons}} C_2 \underset{2\beta}{\overset{\alpha}{\rightleftharpoons}} C_3 \underset{\delta}{\overset{\gamma}{\rightleftharpoons}} O \underset{\epsilon}{\overset{\theta}{\rightleftharpoons}} I$$

Six states-model: $$C_1 \underset{\beta}{\overset{4\alpha}{\rightleftharpoons}} C_2 \underset{2\beta}{\overset{3\alpha}{\rightleftharpoons}} C_3 \underset{3\beta}{\overset{2\alpha}{\rightleftharpoons}} C_4 \underset{4\beta}{\overset{\alpha}{\rightleftharpoons}} C_5 \underset{\delta}{\overset{\gamma}{\rightleftharpoons}} O$$

Two states-model: $$C \underset{\beta}{\overset{\alpha}{\rightleftharpoons}} O$$

The rate constants indicated by the Greek letters in the state models were calculated as

$$k = a\, e^{b(v_{pre} - v_{cl})}$$

where k represents a particular rate constant, and *a* and *b* are channel-specific parameters listed in Table 1.

To take into account a possible effect of the cleft potential on the channel kinetics, the rate constant was

calculated using the voltage difference between the presynaptic potential and the cleft potential.

The ordinary differential equations were numerically solved in Igor using IntegrateODE with default

settings (fifth-order Runga-Kutta-Fehlberg; step size was adapted to keep $Error_i < 10^{-6}$; $Error_i < 10^{-10}$ did

not lead to substantial differences in the prespike waveforms).

Statistics are reported as mean and standard deviation. Pearson's correlation coefficient was calculated to

quantify correlations. Paired t-tests were performed to compare correlation coefficients from the same data

and to compare the absolute deviation from the origin.

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prespike.

Table 1. Rate constants for the voltage-gated channels		a (ms⁻¹)	b (mV⁻¹)
		a (ms^{-1})	b (mV^{-1})
Sodium conductance	α	48.9	1/128.4
	β	1.46	-1/8.9
	γ	392.8	1/36.0
	δ	0.73	-1/15.5
	ε	7.8	1/35.9
	θ	0.01	-1/18
High-threshold potassium conductance	α	1.097	1/57.404
	β	0.794	-1/79.264
	γ	33.750	0
	δ	74.360	0
Low-threshold potassium conductance	α	1.204	1/37.574
	β	0.360	-1/230
	γ	245.488	0
	δ	132.566	0
Calcium conductance	α	1.78	1/23.3
	β	0.14	-1/15
Parameters were taken from ref. [107], except for the calcium conductance [45].			

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Appendix A1. Analysis of the electronic model

Following Kirchhoff's current law, the sum of the currents that enter the cleft must be equal to the sum of the currents leaving the cleft:

$$i_{leave} = i_{enter} \quad (A1.1).$$

The currents that enter the cleft are either resistive or capacitive, coming either from the presynaptic or the postsynaptic side; we also included the conductances of presynaptic voltage-gated channels (g_{ion}) (Figure 2B):

$$i_{enter} = g_{pre}(v_{pre} - v_{pre,rest}) + c_{pre}\frac{dv_{pre}}{dt} + g_{syn}(v_{post} - v_{post,rest}) + c_{syn}\frac{dv_{post}}{dt}$$

$$+ \sum g_{ion}(v_{pre} - v_{ion}) \quad (A1.2)$$

where v_{pre} and v_{post} are the pre- and the postsynaptic membrane potential relative to the interstitial fluid, respectively, $v_{pre,rest}$ and $v_{post,rest}$ are the pre- and the postsynaptic resting membrane potential relative to the interstitial fluid, respectively, c_{pre} and c_{syn} are the capacitances of the pre- and the postsynaptic membrane facing the cleft, respectively, and g_{pre} and g_{syn} and the resting, leak conductance of the pre- and postsynaptic membrane facing the cleft, respectively (Figure 2). The summation sign indicates the summation of the different types of voltage-gated channels, each with their own time- and voltage-dependent conductance g_{ion} and their own reversal potential v_{ion} relative to the extracellular fluid.

The currents leaving the cleft are resistive through a cleft leak conductance (g_{cl}), and resistive and capacitive through the presynaptic and postsynaptic side:

$$i_{leave} = v_{cl}\left(g_{cl} + g_{pre} + g_{syn} + \sum g_{ion}\right) + (c_{pre} + c_{syn})\frac{dv_{cl}}{dt} \quad (A1.3)$$

This gives the full equation for the cleft (combine Eqs. A1.1-3), similar to ref. [30]:

987
$$v_{cl}\left(g_{cl} + g_{syn} + g_{pre} + \sum g_{ion}\right) + (c_{pre} + c_{syn})\frac{dv_{cl}}{dt}$$

988
$$= g_{pre}(v_{pre} - v_{pre,rest}) + c_{pre}\frac{dv_{pre}}{dt} + g_{syn}(v_{post} - v_{post,rest}) + c_{syn}\frac{dv_{post}}{dt}$$

989
$$+ \sum g_{ion}(v_{pre} - v_{ion}) \quad (A1.4)$$

We simplify the full equation by assuming that the postsynaptic currents entering the cleft (g_{syn} (v_{post} -

$v_{post,rest}$) + $c_{syn}dv_{post}/dt$) are much smaller than the presynaptic currents entering the cleft, as is the case

when the prespike is very small, and therefore can be neglected. For the calyx of Held synapse, the

postsynaptic, cleft-facing, resting conductance is indeed rather small ($g_{syn} < 3$ nS, excluding the

conductance of the EPSC), and the prespike is at most a few mV. The presynaptic capacitive currents

during the presynaptic AP are much larger than the resistive currents, as $g_{pre} < 5$ nS, generating currents

less than 0.6 nA (<120 mV x 5 nS), while the presynaptic capacitive current can easily be 5 nA (500 V/s x

10 pF). We thus derive the following, simplified equation:

998
$$v_{cl}\left(g_{cl} + \sum g_{ion}\right) + (c_{pre} + c_{syn})\frac{dv_{cl}}{dt} = c_{pre}\frac{dv_{pre}}{dt} + \sum g_{ion}(v_{pre} - v_{ion}) \quad (A1.5)$$

1000   For a given synapse, we can distinguish two extreme scenarios depending on the route currents entering

1001   the cleft may take: resistive versus capacitive dissipation. If the cleft currents mainly dissipate via

1002   capacitive routes (membrane, $v_{cl}(g_{cl} + \Sigma g_{ion}) \ll (c_{pre} + c_{syn})dv_{cl}/dt$). Eq. A1.5 can be simplified to:

1003
$$(c_{pre} + c_{syn})\frac{dv_{cl}}{dt} = c_{pre}\frac{dv_{pre}}{dt} + \sum g_{ion}(v_{pre} - v_{ion}) \quad (A1.6)$$

1004
$$\frac{dv_{cl}}{dt} = \frac{c_{pre}}{c_{pre} + c_{syn}}\frac{dv_{pre}}{dt} + \frac{1}{c_{pre} + c_{syn}}\sum g_{ion}(v_{pre} - v_{ion}) \quad (A1.7)$$

1005
$$v_{cl} = \int \left(\frac{c_{pre}}{c_{pre} + c_{syn}} \frac{dv_{pre}}{dt} + \frac{1}{c_{pre} + c_{syn}} \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) dt$$

1006
$$= \frac{c_{pre}}{c_{pre} + c_{syn}} v_{pre} + \frac{1}{c_{pre} + c_{syn}} \int \left(\sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) dt + v_{constant} \qquad (A1.8)$$

1007 The VC prespike becomes (using eq. A1.7):

1008
$$-i_{rec} = c_{syn} \frac{dv_{cl}}{dt} = \frac{c_{syn}}{c_{pre} + c_{syn}} \left(c_{pre} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) \qquad (A1.9)$$

Note that in the capacitive dissipation scenario, cleft currents dissipate as a postsynaptic as well as a

presynaptic capacitive current. The relative fraction that dissipates through the presynaptic and

postsynaptic membrane equals $c_{pre} (c_{pre}+c_{syn})^{-1}$ and $c_{syn} (c_{pre}+c_{syn})^{-1}$, respectively. In other words, the

presynaptic currents are partially reflected back to the presynaptic side in this scenario. For the calyx of

Held synapse, where the presynaptic and postsynaptic capacitance are nearly equal, the relative fraction

equals 0.5.

1015

1016 In the resistive dissipation scenario we consider the case where the cleft currents mainly dissipate via

1017 resistive routes ($v_{cl}(g_{cl} + \Sigma g_{ion}) \gg (c_{pre} + c_{syn})dv_{cl}/dt$). We can neglect the contribution of the capacitive

1018 currents induced by the cleft potential to derive the equation for the cleft potential. Eq. A1.5 can then be

1019 simplified to:

1020
$$v_{cl} \left(g_{cl} + \sum g_{ion} \right) = c_{pre} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \qquad (A1.10)$$

1021
$$v_{cl} = \left(g_{cl} + \sum g_{ion} \right)^{-1} \left(c_{pre} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) \qquad (A1.11)$$

1022 An additional simplification is justified if $g_{cl} \gg \Sigma g_{ion}$:

1023
$$v_{cl} = (g_{cl})^{-1} \left(c_{pre} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) \qquad (A1.12)$$

1024   In this case, v_{cl} becomes a scaled sum of the first derivative of the presynaptic AP and ionic currents

1025   through voltage-gated channels.

1026   The postsynaptic prespike, when recorded in voltage-clamp, still equals the capacitive current over the

1027   postsynaptic membrane, which now is assumed to be much smaller than the currents that are present in the

1028   cleft:

1030
$$-i_{rec} = c_{syn} \frac{dv_{cl}}{dt} = c_{syn} \frac{d}{dt} \left((g_{cl})^{-1} \left(c_{pre} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) \right)$$

1031
$$= \frac{c_{syn} c_{pre}}{g_{cl}} \frac{d^2 v_{pre}}{dt^2} + \frac{c_{syn}}{g_{cl}} \sum \frac{d}{dt} \left(g_{ion} (v_{pre} - v_{ion}) \right) \qquad (A1.13)$$

1029

1032

1033   For the CC prespike, in both situations, the current entering the postsynaptic cell via the cleft membrane

1034   must be equal to the current leaving the cell:

1035
$$c_{post} \frac{dv_{post}}{dt} + g_{post} \left(v_{post} - v_{post,rest} \right) = -i_{rec} = c_{syn} \frac{dv_{cl}}{dt} + g_{syn} v_{cl} \qquad (A1.14)$$

Where c_{post} and g_{post} are the total postsynaptic capacitance and conductance, respectively (Figure 2). The

shape of the CC prespike will depend on the cleft potential, the cleft-facing capacitance and the

conductance of the cleft-facing membrane, allowing a large range of different prespike shapes (reflecting

the integral of the cleft potential, the cleft potential itself, or the first derivative of the cleft potential).

We can simplify the notation of some equations by defining a cleft time constant ($\tau_{cl} = c_{cl}/g_{cl}$), and and c_{cl}

$= c_{syn} = c_{pre}$. The equations then become remarkably simple:

$$1042 \qquad v_{cl} = \frac{1}{g_{cl}} \left(c_{cl} \frac{dv_{pre}}{dt} + \sum g_{ion} \left(v_{pre} - v_{ion} \right) \right) = \frac{i_{cleft}}{g_{cl}} \qquad (A1.15)$$

$$1043 \qquad -i_{rec} = c_{cl} \frac{dv_{cl}}{dt} = \tau_{cl} \frac{di_{cleft}}{dt} \qquad (A1.16)$$

$$1044 \qquad c_{post} \frac{dv_{post}}{dt} + g_{post} v_{post} = -i_{rec} = \tau_{cl} \frac{di_{cleft}}{dt} = c_{cl} \frac{dv_{cl}}{dt} \qquad (A1.17)$$

1045   If we furthermore assume that $g_{post} (v_{post} - v_{post,rest}) \ll c_{post}\, dv_{post}/dt$:

$$1046 \qquad c_{post} \frac{dv_{post}}{dt} = \tau_{cl} \frac{di_{enter}}{dt} = c_{cl} \frac{dv_{cl}}{dt}$$

$$1047 \qquad v_{post} = \frac{c_{cl}}{c_{post} g_{cl}} i_{enter} + v_{post,rest} = \frac{c_{cl}}{c_{post}} v_{cl} + v_{post,rest} \qquad (A1.18)$$

1048   The cleft potential is then defined by Ohm's law: the ratio of the net current that runs through the cleft and

1049   the cleft leak conductance. The VC prespike is defined by the first derivative of the net currents scaled by

1050   the cleft time constant. The CC prespike is defined by the cleft potential and the fraction of the

1051   postsynaptic capacitance facing the cleft.

1052

1053   **Appendix A2. The calcium prespike**

1054   This Appendix deals with the question how the recorded calcium current contributes to the VC prespike.

1055   The current i_{ion} that runs through the presynaptic cleft-facing membrane can leak away either capacitively

1056   or resistively similar to the capacitive currents from the presynaptic AP. The recorded calcium current

1057   typically closely follows a Gaussian function [45, 108]:

$$1058 \qquad i_{ca} = i_{max} e^{t^2/(-2\tau_{Ca}^2)} \qquad (A2.1)$$

1059   For the capacitive dissipation scenario, we rewrite A1.9:

1060
$$-i_{rec} = c_{syn}\frac{dv_{cl}}{dt} = c_{pre}\frac{d(v_{pre} - v_{cl})}{dt} + i_{Ca} \qquad (A2.2)$$

1061 When we subtract the passive prespike ($i_{passive-cap}$) from the prespike recording, we arrive at the calcium

1062 prespike:

1063
$$-i_{rec} + i_{passive-cap} = i_{ca} \qquad (A2.3)$$

1064 For the resistive dissipation scenario, we rewrite A1.19 and we again subtract the passive prespike ($i_{passive-}$

1065 $_{res}$) to arrive at the calcium prespike:

1066
$$-i_{rec} + i_{passive-res} = \tau_{cl}\frac{di_{ca}}{dt} = \tau_{cl}i_{max}\frac{d}{dt}\left(e^{t^2/(-2\tau_{Ca}^2)}\right) = \frac{\tau_{cl}i_{max}}{\tau_{Ca}^2} t\, e^{t^2/(-2\tau_{Ca}^2)} \qquad (A2.4)$$

1067 This equation shows two peaks at $t = \pm\tau_{Ca}$ with an amplitude of

1068
$$-i_{rec}(\tau_{Ca}) + i_{passive-res}(\tau_{Ca}) = \pm\frac{\tau_{cl}i_{max}}{\tau_{Ca}^2}\tau_{Ca}\, e^{(\pm\tau_{Ca})^2/(-2\tau_{Ca}^2)} = \pm\frac{\tau_{cl}}{\tau_{Ca}\sqrt{e}}i_{max} \qquad (A2.5)$$

1069 We thus define a simple scaling factor $\tau_{cl}(\tau_{Ca}\sqrt{e})^{-1}$ that relates the amplitude of the presynaptic calcium

1070 current to the amplitude of the calcium prespike.

1071

A

B

Figure S1 – Relation between the calcium current and the prespike, continued. Our data set was composed

of five double recordings. The presynaptic voltage-clamp waveform is shown in Figure 4A. (A) *from top to*

bottom: The presynaptic calcium current (after P/5 subtraction), its inverted first derivative, the calcium

prespike, and the active (black, thick) and passive prespike (magenta, P/5-scaled). Voltage-gated sodium and

potassium currents were blocked (see methods). The left-most column is also shown in Figure 4A. The right-

most is the only one plotted on scaling of the axes on the right. Notice how even when the presynaptic calcium

current does not smoothly follow a Gaussian function, its first derivative resembles the calcium prespike as

predicted from our analytical relations for the resistive dissipation scenario. (B) Relation between the inverted

first derivative of the presynaptic calcium current and the calcium prespike (black dots). Maximum correlation

was obtained by introducing a delay for the first derivative. This delay and the slope of the regression line are

shown in the top left corner of each panel. Each panel corresponds the example in the same column in A.

1084

Appendix B. Geometry

Question B1: What is the voltage profile in the synaptic cleft for a simplified synapse?

We assume a radially-symmetrical synapse with radius r and cleft height h. The capacitive current that

enters the synaptic cleft equals:

$$i_{cap} = c_m v'_{pre} = C_m \pi r^2 v'_{pre} \qquad \text{(B1.1)}$$

with v'_{pre} indicating the first temporal derivative of the presynaptic AP. We assume that the current is

homogeneously distributed over the cleft-facing membrane and that all currents within the cleft run

radially towards the interstitial fluid (resistive dissipation). The current density $J(x)$ can be defined as a

quasi-static function of the radial location x, where 0 indicates the center and r the synapse edge. $J(x)$ will

be equal to the current that runs through the cleft-facing membrane from the center to radius x divided by

the surface through which it escapes, $A_{esc}(x)$. This surface is the surface of an open cylinder with radius x

and height h:

$$J(x) = \frac{i(x)}{A_{esc}(x)} = \frac{\pi x^2 C_m v'_{pre}}{2\pi x h} = \frac{C_m v'_{pre}}{2h} x \qquad \text{(B1.2)}$$

Based on $J(x)$, the electric field can be defined using a homogeneous extracellular cleft resistivity R_{ex},

where the negative sign reflects the convention that current runs down the voltage gradient:

$$-\frac{dv(x)}{dx} = R_{ex} J(x) = \frac{R_{ex} C_m v'_{pre}}{2h} x \qquad \text{(B1.3)}$$

The voltage profile relative to the cleft edge is the integral of equation B1.3 from x to r:

$$v(x) = \int_x^r \left(-\frac{R_{ex} C_m v'_{pre}}{2h} x\right) dx = -\frac{R_{ex} C_m v'_{pre}}{2h} \int_x^r (x) dx = \frac{R_{ex} C_m v'_{pre}}{4h} (r^2 - x^2) \qquad \text{(B1.4)}$$

We can also rewrite equation B1.4 as a relative distance to the cleft edge:

1104
$$v(x/r) = \frac{r^2 R_{ex} C_m v'_{pre}}{4h}(1 - (x/r)^2) \quad (B1.5)$$

The voltage profile will apply as long as all current runs radially to the extracellular matrix. When this is

not the case -for instance, when currents enter into the postsynaptic cell or return through the presynaptic

cleft-facing membrane- the current density will be reduced, and the voltage gradient will be reduced as

well.

Question B2: How much current will be generated at the postsynaptic membrane?

We assume that the currents that run through the postsynaptic membrane are purely capacitive:

1111
$$i_{post} = \int_{x=0}^{r} i_{cap}(x)dx = \int_{x=0}^{r} (C_m A(x)v'(x))dx \quad (B2.1)$$

1112 Here, the current through an annulus of membrane (at x, with thickness dx) will depend on the radial

1113 location as the surface of the annulus equals $2\pi x$dx and the cleft potential is also defined by the radial

1114 location (equation B1.4). Only v'_{pre} will be affected by the time derivative, which will become the second

1115 time derivative (v''_{pre}) . The equation then becomes:

1116
$$i_{post} = C_m \int_{x=0}^{r} \left(2\pi x \frac{R_{ex} C_m v''_{pre}}{4h} (r^2 - x^2) \right) dx$$

1117
$$= \frac{\pi R_{ex} C_m^2 v''_{pre}}{2h} \int_{x=0}^{r} (x (r^2 - x^2))dx$$

1118
$$= \frac{\pi R_{ex} C_m^2 v''_{pre}}{2h} \left[\frac{1}{2}r^2 x^2 - \frac{1}{4}x^4 \right]_{x=0}^{r}$$

1119
$$= \frac{\pi R_{ex} C_m^2 v''_{pre}}{8h} r^4 \quad (B2.2)$$

1120

1121 **Question B3: How does the cleft resistivity relate to the cleft conductance?**

1122   We noted that B2.2 resembles A1.19 without the voltage-gated currents: :

1123
$$i_{post} = \frac{c_{cl}^2}{g_{cl}} v''_{pre} = \frac{(C_m \pi r^2)^2}{g_{cl}} v''_{pre} \quad \text{(B3.1)}$$

1124   By combining equation B2.2 and B3.1 we find that g_{cl} relates to the cleft resistivity:

1125
$$g_{cl} = \frac{8\pi h}{R_{ex}} \quad \text{(B3.2)}$$

1126   Combining equations B3.1 and B3.2 yields an alternative way to calculate the extracellular resistivity:

1127
$$R_{ex} = \frac{i_{post}}{c_{cl}^2} \frac{8\pi h}{v''_{pre}} \quad \text{(B3.3)}$$

1128   This equation holds as long as the cleft currents dissipate resistively and the contribution of voltage-gated

1129   currents are negligible.

1130

Question B4. What are the relations for a finger-like synapse?

We assume a sheet-like synapse with a cleft height h, a constant width of $2r$, and a length L that is closed

at the longitudinal ends. Any current that enters the synaptic cleft cannot escape along the longitudinal

axis, leaving only the transversal route to the interstitial space. Examples of such synapses are finger-like

synapses, such as the neuromuscular junction or the endbulb of Held synapse for which we ignore the

longitudinal endings, or deeply-invaginating, cylindrical synapses (spinules) with a circular longitudinal

axis (*i.e.* a rolled-up sheet forming a cylinder). Because of the constant width, the voltage profile will be

identical along the longitudinal axis, varying only along the synapse width. The capacitive current will be

equal to $i_{cap}(x) = C_m v'_{pre} dr dl$ (B4.1).

The current density J at each location will be:

1141
$$J(x) = \frac{i(x)}{A_{esc}(x)} = \frac{C_m v'_{pre} x dl}{h dl} = \frac{C_m v'_{pre}}{h} x \quad \text{(B4.2)}$$

1142   The voltage gradient then is:

$$-\frac{dv(x)}{dx} = R_{ex}J(x) = \frac{R_{ex}C_m v'_{pre}}{h}x \qquad \text{(B4.3)}$$

1143

1144   And the voltage at x is:

$$v(x) = -\frac{R_{ex}C_m v'_{pre}}{h}\int_x^r (x)dx = \frac{R_{ex}C_m v'_{pre}}{2h}(r^2 - x^2) \qquad \text{(B4.4)}$$

1145

The changes in the cleft potential will generate a postsynaptic, capacitive current. To calculate the total

current from the finger-like synapse, we sum the current along the two dimensions:

$$i_{fin} = \int_{l=0}^L 2\int_{x=0}^r i_{cap}(x)dx\, dl = \int_{l=0}^L 2\int_{x=0}^r C_m A v'(x)dx\, dl$$

$$= \int_{l=0}^L 2\int_{x=0}^r C_m \frac{R_{ex}C_m v''_{pre}}{2h}(r^2-x^2)dx\, dl = \frac{R_{ex}C_m^2 v''_{pre}}{h}\int_{l=0}^L\int_{x=0}^r (r^2-x^2)dx\, dl$$

$$= \frac{R_{ex}C_m^2 v''_{pre}}{h}\int_{l=0}^L\left[xr^2 - \frac{1}{3}x^3\right]_{x=0}^r dl = \frac{2R_{ex}C_m^2 v''_{pre}r^3}{3h}\int_{l=0}^L dl$$

$$= \frac{2R_{ex}C_m^2 v''_{pre}r^3 L}{3h} \qquad \text{(B4.5)}$$

For the sheet-like synapse, x=0 indicates the cleft center and currents can run in two, transversal

directions. We therefore multiply the capacitive current by 2 at the start of the derivation of B4.5. For a

spinule synapse, no current can escape from the 'deep' end of the cylinder (at x=0) and therefore the

currents run in a single, transversal direction. The right-hand side of B4.5 should therefore be halved for

spinule-like synapses or other deeply-invaginating synapses, neglecting any contribution from the cup-

shaped cleft at the deep end. The characteristic difference between the radial-like or sheet-like synapse is

whether the escape area along the longitudinal axis increases or stays constant, respectively. Following the

logic outlined in B1, B2 and B4, one can calculate the cleft potentials and postsynaptic currents for other

synaptic geometries.

1161

1162 **Question B5. How does fenestration change the cleft potential and the VC prespike?**

1163 To assess the impact of fenestration on the cleft potential and the VC prespike, we compare a radially-

1164 symmetric synapse and a synapse with n equally-sized fingers. The relative cleft potential can be

1165 calculated using equation B1.5 for the radially-symmetric synapse (v_{sym}) and equation B4.4 for the finger-

1166 like synapse (v_{fin}), while assuming that the cleft potential of each finger is independent from each other:

1167
$$\frac{v_{fin}}{v_{sym}} = \frac{r_{fin}^2 R_{ex} C_m v'_{pre}(1 - x^2/r_{fin}^2)}{2h} \frac{4h}{r_{sym}^2 R_{ex} C_m v'_{pre}(1 - x^2/r_{sym}^2)}$$

1168
$$= \frac{2r_{fin}^2(1 - x^2/r_{fin}^2)}{r_{sym}^2(1 - x^2/r_{sym}^2)} \quad \text{(B5.1)}$$

1169 This gives at the center (x=0) a ratio of:

1170
$$\frac{v_{fin}}{v_{sym}} = 2\left(r_{fin}/r_{sym}\right)^2 \quad \text{(B5.2)}$$

1171 To compare the VC prespike of a radially-symmetric synapse and a fenestrated synapse, we assume that

1172 the currents from each finger are synchronized and that they sum linearly. We use equation B2.2 and

1173 equation B4.5:

1174
$$\frac{n\, i_{fin}}{i_{sym}} = \frac{2nR_{ex} C_m^2 v''_{pre} r_{fin}^3 L}{3h} \frac{8h}{\pi R_{ex} C_m^2 v''_{pre} r_{sym}^4} = \frac{16nL r_{fin}^3}{3\pi r_{sym}^4} \quad \text{(B5.3)}$$

1175 We assume that the total contact surface is conserved, so the surface of the fenestrated synapse (n $2r_{fin}$ L)

1176 and the radially-symmetric synapse (πr_{sym}^2) should be equal. Therefore L = πr_{sym}^2(n $2r_{fin}$)$^{-1}$, which we

1177 insert into equation B5.3:

1178
$$\frac{n\, i_{fin}}{i_{sym}} = \frac{16n r_{fin}^3}{3\pi r_{sym}^4} \frac{\pi r_{sym}^2}{n\, 2r_{fin}} = \frac{8}{3}\left(r_{fin}/r_{sym}\right)^2 \quad \text{(B5.4)}$$

1179