

1 Towards the Myoelectric Digital Twin: Ultra Fast
2 and Realistic Modelling for Deep Learning

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12 **Abstract**

13 Muscle electrophysiology has emerged as a powerful tool to drive human
14 machine interfaces, with many new recent applications outside the
15 traditional clinical domains. It is currently a crucial component of control
16 systems in robotics and virtual reality. However, more sophisticated,
17 functional, and robust decoding algorithms are required to meet the fine
18 control requirements of these new applications. Deep learning approaches
19 have shown the highest potential in this regard. To be effective, deep
20 learning requires a large amount of high-quality annotated data for training;
21 the only option today is the use of experimental electromyography
22 data. Yet the acquisition and labelling of training data is time-consuming
23 and expensive. Moreover, the high-quality annotation of this data is often
24 not possible because the ground truth labels are hidden. Data augmentation
25 using simulations, a strategy applied in other deep learning
26 applications, has never been attempted in electromyography due to the
27 absence of computationally efficient and realistic models. Here, we present
28 a new highly realistic and ultra-fast computational model tailored for the
29 training of deep learning algorithms. For the first time, we are able to
30 simulate arbitrary large datasets of realistic electromyography signals with
31 high internal variability and leverage it to train deep learning algorithms.
32 Because the computational model provides access to all the hidden parameters
33 of the simulation, it also allows us to use some annotation strategies
34 that are impossible with experimental data. We believe that this concept
35 of Myoelectric Digital Twin allows new unprecedented approaches to
36 muscular signals decoding and will accelerate the development of human-
37 machine interfaces.

38 1 Introduction

39 Biosignals have been classically used for studying the underlying physiology, for
40 clinical diagnostics, and for monitoring. More recently, they have also been used
41 for interfacing humans with external devices. For example, signals measured at
42 the surface of the skin from skeletal muscle electrical activity, i.e. surface elec-
43 tromyography (sEMG), are used for the control of bionic limbs [1]. In this
44 application, the recorded electrical signals are converted into motion commands
45 using machine learning [2, 3, 4]. In recent years, with the development of deep-
46 learning based methods as well as wearable and cost-effective recording devices,
47 there has been increased interest in using muscular signals as a basis for human-
48 machine interfaces [5, 6]. The potential applications go well beyond the tradi-
49 tional clinical domains of prostheses and orthoses and range from robotic control
50 to gaming and virtual reality [7]. A core challenge of deep-learning methods ap-
51 plied to biosignals is the acquisition of personalized and annotated training data
52 in sufficient quantity and quality. Training data needs to be recorded for differ-
53 ent subjects, at different times, with high variability in electrode configurations
54 and experimental paradigms. In addition, it is challenging and in some cases
55 impossible to properly describe the underlying physiological or neural paramet-
56 ers (e.g. individual muscle forces, fiber physiological parameters, motor neuron
57 impulse timings), which are crucial for the correct annotation of data samples.
58 As a result, acquiring experimental EMG data in sufficient quantity and quality
59 is not only expensive and time-consuming, but in many cases not possible.

60 Data augmentation via simulation is an alternative approach to lengthy data
61 acquisitions, and indeed augmentation techniques have been recently introduced
62 for electrophysiological signals [8, 9, 10, 11]. However, most of these augmenta-
63 tion methods use “black-box” models, which aim to capture essential features
64 of the signal without relating them to the underlying physiology [12]. Thus, the
65 ground truth for most of the crucial parameters is still unknown, greatly limiting
66 the potential use cases of such approaches. More sophisticated biophysical mod-
67 elling methods are based on solving so-called forward equations (e.g., Poisson
68 equation in the electrostatics case). However, this type of biophysical mod-
69 elling has not been considered in the context of data augmentation for machine
70 learning approaches. Indeed, state-of-the-art models are either not sufficiently
71 realistic or not computationally efficient to produce suitable training data. For
72 example, in the case of describing the generation of EMG signals, analytical
73 models based on simple geometries of the tissues [13, 14, 15, 16, 17] provide
74 simulations which reflect the broad characteristics of the signals, but cannot be
75 used to reproduce specific experimental conditions due to the overly simplified
76 anatomy. The more realistic models of EMG generation based on numerical
77 solutions of the Poisson equation with generic volume conductor shapes [18, 19]
78 are currently limited by their prohibitive computational time.

79 Here, we describe an EMG simulation method, based on the numerical so-
80 lution of the forward equations suitable for deep learning data augmentation.
81 It produces highly realistic EMG recordings, provides access to all underlying
82 physiological parameters, and is extremely computationally efficient. Our re-

83 sults show that it is possible to simulate EMG signals for anatomically accurate
84 conductor geometries and multiple muscles with tens of thousands of muscle
85 fibers in a few seconds. As an application scenario, we also demonstrate the
86 use of this model for data augmentation by pre-training neural networks that
87 decompose EMG into the underlying neural activity sent from the spinal cord
88 to muscles [20].

89 Our model is the only realistic and computationally efficient simulator tar-
90 geted to AI training and approaching the concept of a Myoelectric Digital Twin.
91 It allows generating arbitrary large datasets of realistic and personalized EMG
92 signals, with high data variability and with a perfect annotation of diverse hid-
93 den parameters. As a result, our model may allow breakthrough approaches in
94 AI-based EMG signal processing and decoding.

95 2 Results

96 Biophysics

97 To allow the efficient simulation of a large quantity of highly realistic EMG
98 recordings, we have developed a novel approach to solve the forward problem
99 of the volume conductor in electrostatic conditions. Our approach is based on
100 a hierarchical and flexible decomposition of the EMG simulation pipeline which
101 allows the reuse and optimization of individual steps.

102 First, a realistic anatomy, described by bone, muscle, skin, and electrode
103 surfaces, is discretized into a tetrahedral volume mesh. A conductivity tensor,
104 anisotropic for muscles and isotropic elsewhere, is associated with each tetra-
105 hedral of the volume. Unlike the state-of-the-art approaches, which solve the
106 quasi-static Maxwell's equations for each fiber source and for each time instant,
107 we solve them for a set of unit point sources located at each vertex of the mesh
108 associated with the muscle tetrahedrals, which are referred to as basis sources.
109 This computation does not depend on the time variable nor on the fibers and
110 motor unit geometry and their physiological properties. *Therefore, changing*
111 *these parameters does not require recomputing the forward solutions.*

112 Moreover, due to a rewriting of the equations involved using the so-called
113 adjoint method, the solution is obtained by solving as many systems of equa-
114 tions as there are electrodes, rather than basis sources. Because the number
115 of electrodes ($\approx 10^2$) is typically much lower than the number of basis sources
116 ($\approx 10^5$), computational performance is substantially improved.

117 Second, using the same muscle surfaces used to describe the volume conduc-
118 tor, individual fiber geometries can be automatically generated, if this data is
119 not available from other sources (e.g. from diffusion magnetic resonance imag-
120 ing). Moreover, the fibers are grouped into motor units (MUs) following the
121 state-of-the-art models for MU physiology. This step does not depend on the
122 forward computations, and thus altering the related parameters and producing
123 new simulation is highly efficient.

124 Third, the current source density propagating along the fibers is generated

125 using a realistic intracellular action potential model. The contribution of indi-
126 vidual fibers to the EMG recordings is obtained by discretizing each fiber into
127 a set of points, integrating the current source density along its length, and pro-
128 jecting onto the sensor locations using the basis points computed in the first
129 step. This approach effectively decouples the number of fibers and their dis-
130 cretization from the conductor model, allowing the simulator to handle tens of
131 thousands of fibers per muscle. Again, changing the fiber parameters (end-plate
132 location, action potential propagation velocity, tendons length, etc.) does not
133 require recomputing the other blocks of the simulation.

134 Fourth, given a muscle activation profile, we use the size principle to recruit
135 MUs and their associated fibers. This allows a simple and easily interpretable
136 input to the simulation which can be used to simulate EMG recordings associ-
137 ated to specific muscle contractions and their movements.

138 The architecture described above, and detailed in Methods, has several ad-
139 vantages. First, each step of the procedure can be optimized individually, im-
140 proving the performance of the system and the quality of the simulated EMG.
141 In particular, due to the algebraic properties of the computations and their in-
142 dependence, a large part of them can potentially be performed in parallel (on
143 CPU and GPU). Second, simulating data over a range of parameters does not
144 require a full recomputation of the model. This allows the generation of massive
145 EMG datasets covering a range of parameters and using personalized anatomy.
146 In addition, the datasets are perfectly annotated, from overall muscle activation
147 down to individual fiber action potential velocity.

148 As a result, our model is the first that allows the generation of ultra realistic
149 and arbitrarily large (because of its computational performance) datasets of
150 simulated EMG signals that can be used for AI training.

151 The details and all mathematical equations related to the model development
152 are described in the Methods.

153 **The simulator reproduces analytical solutions**

154 To produce realistic EMG data, the simulator leverages a flexible representation
155 of the underlying anatomy and physiology. This flexibility does not only allow
156 the use of realistic and personalized models, but also permits reproducing sim-
157 ple conductor geometry used in analytical solutions. A first validation of our
158 numerical solution is performed by comparing it with its analytical counterpart
159 for a cylindrical volume conductor geometry [21]. The normalized mean square
160 error between the two solutions depended on the depth of the fiber and varied
161 between 3% (1mm depth from the muscle surface) and 5% (11mm depth). Fig. 1
162 illustrates the analytical and numerical solutions for a fiber depth of 1 mm from
163 the muscle surface. Because of the low error, the two waveforms are almost in-
164 distinguishable. It is important to note that the two volume conductor models
165 in this validation are not identical. The theoretical/analytical solution is com-
166 puted for an infinitely long cylinder (repeated periodically when discretized),
167 while the numerical solution uses a cylinder of a large (sufficiently longer than
168 the fiber and the electrode array), yet finite length. Increasing the length of the

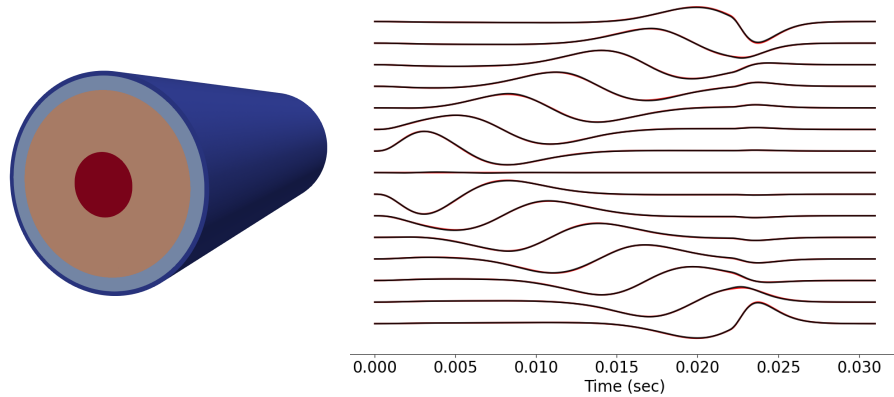


Figure 1: Comparison of the numerical and analytical [21] solutions (on the right) for a four layer cylindrical volume conductor model (on the left): analytical (red) and numerical (black) EMG signals for a differential array electrode montage. The depth of the source fiber in this example is 1 mm from the muscle surface.

169 cylinder did not significantly alter the error.

170 The simulator produces realistic EMG data

171 To evaluate the performance of the simulator at multiple scales, we started by
172 simulating EMG signals associated to a single fiber activation inside the brachio-
173 radialis muscle. The signal recorded by an array of 16 rectangular electrodes (15
174 differential channels) when a single fiber was active is shown in Fig. 2A. The vol-
175 ume conductor model is based on an anatomically accurate forearm geometry,
176 which includes all the muscles, bones, fat, and skin tissues.

177 Different distinctive features are present in the simulated signal that are also
178 observed in experimental EMG signals [22]. In particular, electrodes of channel
179 4 are located on different sides of the neuromuscular junction (NMJ) and thus
180 the respective signals cancel each other out. Channels 7-11 present propagating
181 EMG components resulting from the fiber AP propagating from the NMJ to
182 the tendons. Channels 2-6, as well as channels 12-15, contain non-propagating
183 sEMG components, which are due to the AP generation at the NMJ and its
184 extinction at the tendon (end-of-fiber effect), respectively.

185 A further example is a simulation of an excitation of a single muscle, illus-
186 trated in Fig. 2B. A simple excitation drive for the Brachioradialis muscle is
187 simulated as gradually increasing from 0% to 100% of the maximum voluntary
188 contraction and smoothly decreasing back to 0%. As described in Section 4.5,
189 50000 muscle fibers were realistically distributed into 200 motor units over the
190 muscle volume and recruited according to the size principle [23]. The signal was
191 simulated for 8 circular bipolar electrodes located around the forearm. In this
192 example, the volume conductor effect becomes particularly visible with elec-

193 trodes nearer to the active muscle having higher signal amplitudes. Notice that
194 the electrodes record different signal waveforms as the muscle units are located
195 at varying distances from the electrodes, weighting their contribution to the ob-
196 served EMG signals. We also observe an increase of the signal amplitude with
197 muscle excitation, an important feature of experimental EMG signals, which is
198 a consequence of progressive motor unit recruitment and of an increase in the
199 discharge rates of the active motor units.

200 Finally, we simulated sEMG signals from multiple muscle excitations, corre-
201 sponding to the active wrist flexion and extension and passive wrist abduction
202 against gravity. We used a simple muscle excitation model for three groups of
203 muscles (flexors, extensors and abductors). More details about the experimental
204 design are presented in Section Details of realistic simulation examples. Fig. 2C
205 and Fig. 2D clearly show the qualitative similarities in signal characteristics
206 between experimental and simulated data. Our model was able to reproduce
207 the different signal patterns during both flexion and extension. Beside the dif-
208 ferent activation across the electrodes during flexion and extension, the effect
209 of wrist abduction is also visible in both data sets. Thus, channels 2, 3 and 7,
210 8 present a small signal activity during the whole duration of the simulation,
211 and not only during flexion/extension peaks. Similar activity can also be seen
212 in experimental data, with channels 2 and 7 being the most active.

213 In addition to the analysis in the time domain, simulated data were com-
214 pared against the experimental data in the frequency domain. Fig. 3 illustrates
215 an example of the measured and simulated single channel sEMG. It has to be
216 noted that the spectral characteristics of a signal strongly depends on multiple
217 simulation parameters. In this example, we ran several hundreds simulations
218 by varying the simulation parameters in a realistic range and selected the set of
219 parameters leading to the minimal spectral difference. This approach, which is
220 a simple version of inverse modelling, was possible because of the high compu-
221 tational speed of the simulations.

222 **The simulator is ultra fast**

223 The computational performance of an EMG signal simulation depends on the
224 model properties and the particular experimental setup. Consequently, there is
225 no benchmark to evaluate and compare the performance of different simulation
226 methods. The computational time magnitude of the state-of-the-art methods
227 is, in the best cases, in *the order of hours* for a single simulation (with a fixed
228 set of model parameter values, ≈ 50000 fibers, 5 electrodes) [24, 19].

229 By exploiting the mathematical properties of the forward equations and
230 source model, we were able to achieve a computational performance of *the or-
231 der of minutes* per simulation. Moreover, in our model, changing most of
232 the simulation parameters does not require recomputing the whole model and
233 reduces the computational time of new simulations to *the order of seconds*,
234 if the volume conductor remains constant. As a result, it becomes practically
235 possible to simulate arbitrary large datasets of highly realistic EMG signals with
236 high variability in the simulation parameters. Details on the computational time

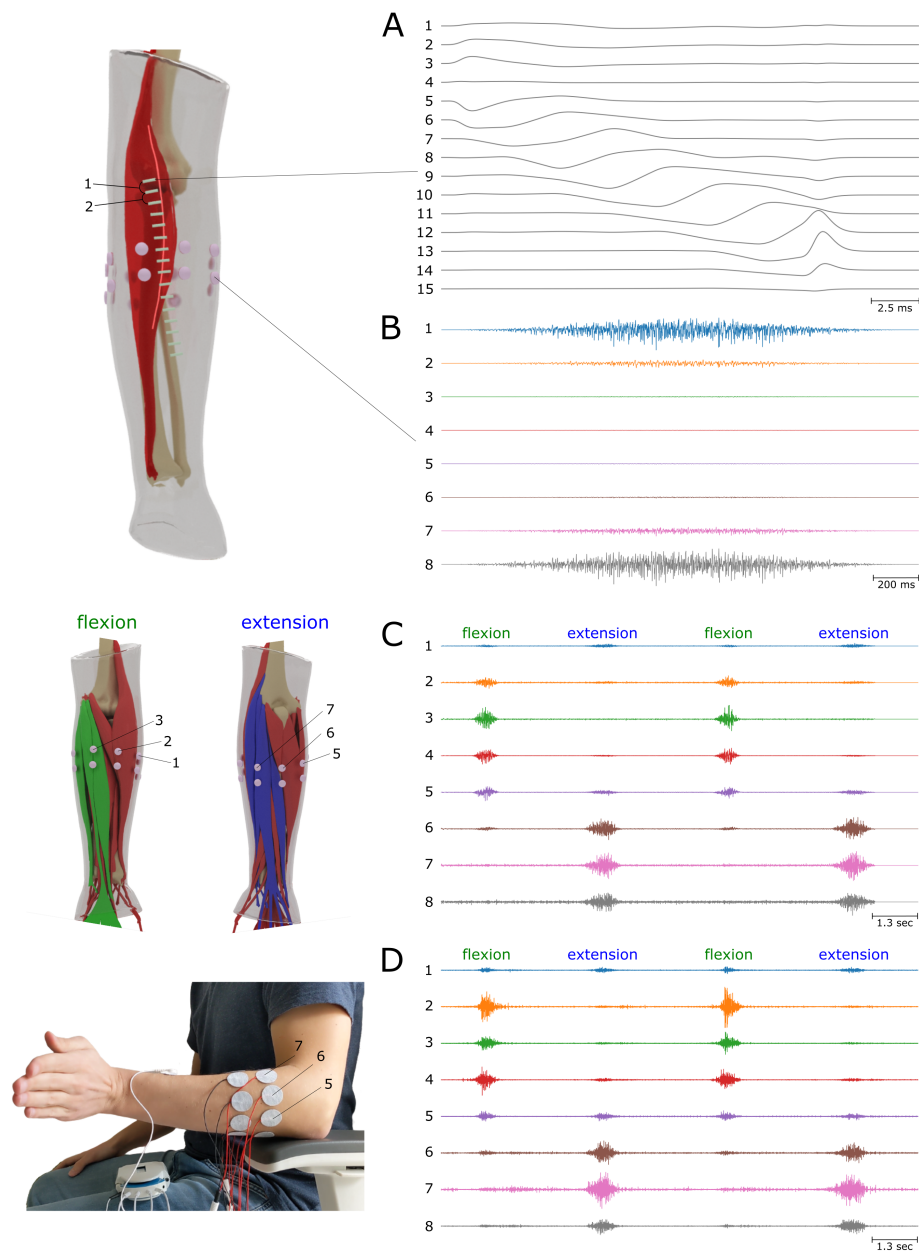


Figure 2: Simulation examples at multiple activation scales. (A) Single fiber activation in the Brachioradialis muscle measured by an electrode array with 15 differential channels. (B) 2-second long activation of the Brachioradialis muscle, reaching 100% of maximum voluntary contraction (MVC). 8 bipolar electrodes located around the forearm are simulated. (C) Simulation of wrist flexion and extension by activating the corresponding flexor and extensor muscles. (D) The experimental EMG signals of wrist flexion/extension.

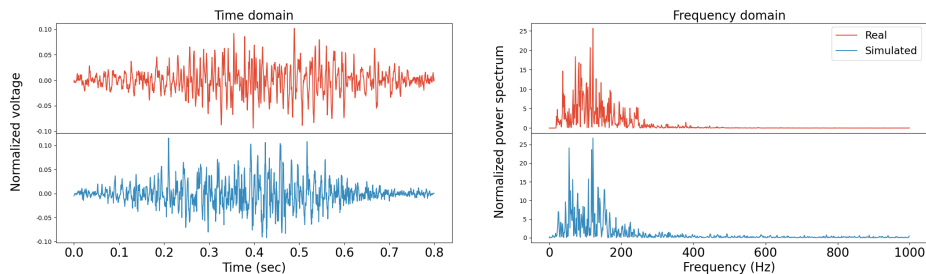


Figure 3: Comparison of experimental (red) and simulated (blue) single channel normalized sEMG signals in the time and frequency domains. The simulator parameters have been optimized to match the experimental signals.

237 in several conditions are provided in Methods (section Computational perfor-
238 mance).

239 The proposed model is also highly scalable for multiprocessing, and the cur-
240 rent computational time can be further reduced by several orders of magnitude
241 by implementing parallel computation on CPU and GPU.

242 **Realistic and fast EMG simulations open unique perspec-** 243 **tives for deep learning**

244 Here, we show a potential use of high volumes of simulated surface EMG data
245 for deep learning, utilising the proposed model to generate MUAP templates
246 which can be used to pre-train neural networks. This methodology is used in
247 other deep learning domains, such as the use of the ImageNet image database
248 to pre-train object classifiers prior to adaptation to specific applications [25].

249 The myoelectric digital twin simulations (Fig. 4B) were used *to pre-train*
250 *a neural network* that could extract motor unit activations from unprocessed
251 HD-sEMG signal [26]. This pre-trained network was then trained to decom-
252 pose experimentally measured HD-sEMG signals collected at the dominant wrist
253 from nine participants (Fig. 4A). This procedure was then repeated, but with a
254 randomly initialised version of the network instead of the pre-trained weights.
255 See section 4.8 for details.

256 The simulation pre-trained network outperformed random initialisation in
257 decomposition accuracy when compared to the original decomposition as mea-
258 sured by the rate of agreement (RoA) metric [27] (Fig. 4C). The median (IQR)
259 RoA of the pre-trained network was 93.8% (84.8 to 100.0), compared to 82.4%
260 (71.6 to 100.0) in the random initialisation network, a significant difference ac-
261 cording to the Wilcoxon signed-rank test ($p < 0.001$). Of the 39 decoded motor
262 units, 22 had improved RoAs with pre-training and one had a worse RoA, with
263 the remaining 16 showing no change, generally because the initial RoA was al-
264 ready 100% without pre-training. The pre-trained network had a much lower
265 variance in the accuracy of predictions on the test sets than random initialisa-
266 tion, quickly optimising to a model effective for generalisation to new signals.

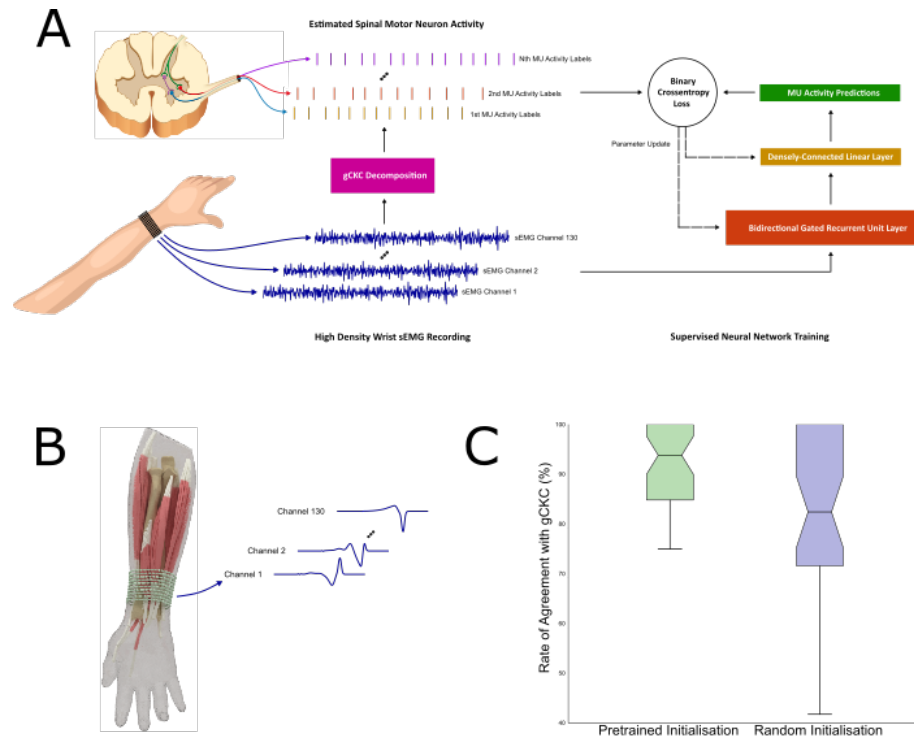


Figure 4: (A) Decomposition of experimental high density EMG recordings into underlying spinal motor neuron activities. The results obtained with the neural network (NN) were tested against the decomposition by a reference blind source separation method and manual editing by an expert operator. (B) Myoelectric digital twins were used to generate MUAP templates for different muscles and different model parameters (tissue conductivities, fiber properties, tendon sizes, etc). 64 sets, each containing 5 simulated MUAPs for 130 electrodes, were used for pre-training. (C) Rate of agreement (%) between the neural network MU activity predictions and the decomposition algorithm on one second of wrist flexor HD-sEMG signal. Median and interquartile range plotted over 39 motor units from nine participants. Both outputs were converted to timestamps using a two class K-means clustering. The neural network using a gated recurrent unit (GRU) network that was pre-trained using simulated EMG signal significantly outperformed a GRU with random initialisation ($p < 0.001$).

267 3 Discussion

268 We have proposed an efficient computational approach to highly realistic surface
269 EMG modeling. The method provides the solution to the generation of EMG
270 signals from anatomically accurate volume conductor properties and number of
271 muscle fibers, within limited computational time compatible with real-time sig-
272 nal generation. The proposed model is the only available EMG simulator with
273 realistic description of the volume conductor and optimized for such computa-
274 tional efficiency. The main value of the model is that it eliminates bottlenecks
275 of the state-of-the-art methods and opens unprecedented perspectives for using
276 simulated sEMG for data augmentation in the deep learning framework and
277 therefore for building a myoelectric digital twin.

278 The computational efficiency in the volume conductor solution has been recog-
279 nized as an important component of EMG modeling, and some attempts to
280 decrease the computational time in EMG simulations have been described. For
281 example, the approaches developed by Dimitrov & Dimitrova [28] and Farina et
282 al. [29, 21] substantially decreased the computational time in analytical EMG
283 modeling for simple volume conductor geometries. These models provide simu-
284 lations which reflect the broad characteristics of EMG signals, but can not be
285 anatomically accurate because of the restrictions on the volume conductor and
286 fiber source geometry. Realistic models using numerical solutions have also been
287 recently proposed. The previous most complete and efficient model has been
288 proposed by Pereira Botelho et al. [19]. These authors have used an anatom-
289 ically accurate model to simulate EMG signals generated during index finger
290 flexion and abduction. They gained computational speed by using the principle
291 of reciprocity. In fact, one part of our calculations also includes the adjoint
292 method, which is an algebraic representation of this principle. By using reci-
293 procity, Pereira Botelho et al. [19] reported a computational time of 1 hour for
294 simulating the activation of nearly 15500 fibers for 5 electrodes. This time, how-
295 ever, remains impractical for simulating arbitrary large data sets for a variety of
296 parameter values. The model we proposed in this paper substantially surpasses
297 the computational efficiency reported in [19]. We achieved it by efficiently
298 exploiting mathematical properties of the forward equations, in particular by
299 introducing the concept of basis points and by separating model parameters
300 and variables into independent computational blocks. The approach does not
301 only reduce the computational time for a full simulation, but also allows us to
302 scale the solution, so that new solutions for the same volume conductor can be
303 obtained without re-computing the volume conductor transformation. In this
304 way, the generation of EMG signals within the same volume conductor, but
305 varying all other simulation parameters, can be performed in extremely short
306 time. Complex EMG signals from tens of thousands of muscle fibers located
307 in multiple muscles, can be generated (and regenerated with different param-
308 eter values) in a computational time of the order of seconds. In contrast to
309 previous models, our proposed simulator does not compromise accuracy and
310 computational speed.

311 Some limitations remain in the current state of the presented model. It does

312 not include some sources of variability that are present in experimental EMG
313 signals and strongly impact their processing and analysis. For example, the
314 model does not include advanced noise and artifacts descriptions, biomechanical
315 modeling of the musculoskeletal system, and non-stationary volume conductor
316 properties and fiber geometry. While these aspects are beyond the scope of this
317 paper, they are relevant features to include in future developments.

318 The advances presented in this work, together with the proposed future de-
319 velopments, naturally lead to the concept of a myoelectric digital twin - a hyper-
320 realistic, personalized, computationally-efficient model which generates EMG
321 data in a quality and quantity sufficient not only to augment but to replace real
322 data, with utility for AI training in the various real world applications. Here we
323 have illustrated the potential of this approach by augmenting training data for
324 deep neural networks, with the aim of identifying the discharge times of spinal
325 motor neurons from surface EMG signal. By using the simulator to augment
326 training (through a pre-training procedure), we showed a substantial increase
327 in the performance of the decomposition network when applied to experimental
328 data, demonstrating a highly relevant use of the proposed approach for de-
329 creasing the need for experimental training data in human-machine interfacing
330 applications.

331 4 Methods

332 4.1 Forward problem

333 The fiber extracellular potentials that are measured by EMG electrodes are
334 generated by transmembrane currents. The properties of bioelectric currents
335 and potential fields can be determined from solutions of the Maxwell's equations,
336 taking into account the electrical properties of biological tissues. Because of the
337 relatively low frequencies of signal sources of biological origin, the quasi-static
338 assumption can be applied [30, 31], so that the electric potential and the primary
339 current sources are related by the following Poisson equation [30, 32, 33] with
340 Neumann boundary conditions:

$$\begin{cases} \nabla \cdot (\sigma \nabla \phi) = -I & \text{in } \Omega \\ \sigma \frac{\partial \phi}{\partial \mathbf{n}} = \sigma \nabla \phi \cdot \mathbf{n} = 0 & \text{on } \partial \Omega \end{cases} \quad (1)$$

341 where $\Omega \subset \mathbb{R}^3$ is a volume conductor domain of interest, $\partial \Omega$ its boundary
342 with outward pointing normal unit vector \mathbf{n} , $\phi(\mathbf{r})$ [V] is the electric potential,
343 $I(\mathbf{r})$ [A/m³] is the current source density (CSD), $\sigma(\mathbf{r})$ [S/m] is a conductivity
344 tensor. The second line of the equation (boundary condition) reflects the as-
345 sumption that no current flows out of the domain of interest. In the context of
346 EMG modeling, this implies that there is no current flow between the skin and
347 air. The current source density $I(\mathbf{r})$ is interpreted as the volume density of cur-
348 rent entering or leaving the extracellular medium at position $\mathbf{r} \in \Omega$. A negative
349 CSD corresponds to current leaving the extracellular medium (due to the fiber

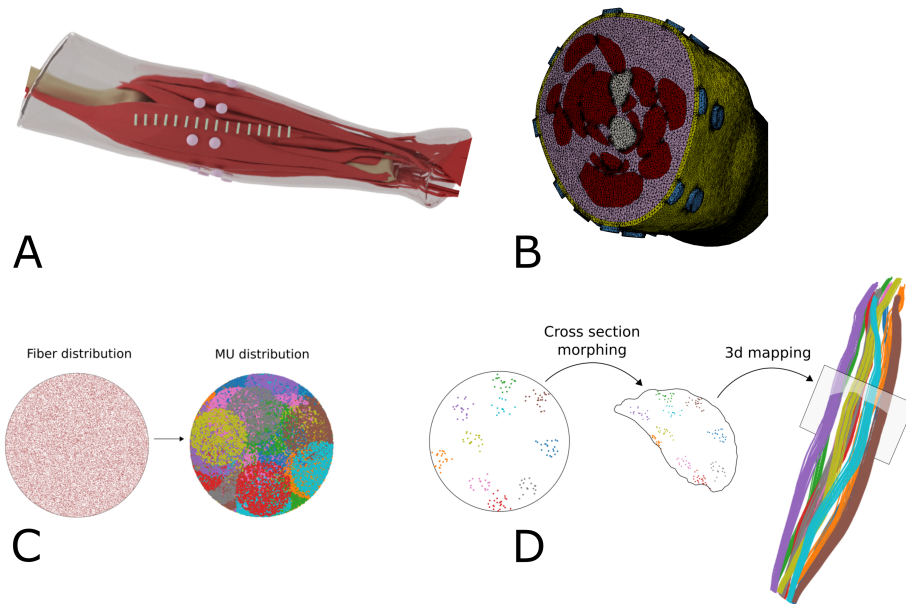


Figure 5: (A) Surface geometry of muscles, bones, subcutaneous tissue, skin and electrodes used for arm modeling (taken from BodyParts3D, The Database Center for Life Science (<http://lifesciencedb.jp/bp3d/>)). (B) Cross-section of the volume mesh generated from the arm surfaces. (C) Uniformly distributed fibers inside a unit circle are grouped into motor units of different sizes, locations and territories. (D) Example of mapping of 10 small motor units from the circle into an arbitrary muscle by morphing the unit circle into the muscle cross section.

350 transmembrane currents) and is thus conventionally called a sink. Likewise,
 351 current entering the extracellular medium is called a source [34, 35].

352 Equation (1) cannot be solved analytically for general volume conductor
 353 geometries, but several numerical methods can be used to approximate its so-
 354 lution. Here, we use the finite element method (FEM) [36], which discretizes
 355 the volume conductor Ω as a tetrahedral mesh Ω_t . Given this mesh, we use the
 356 Galerkin method to project the potential ϕ onto the space of piecewise affine
 357 functions defined on Ω_t . Fig. 5A and Fig. 5B illustrate an example of a realistic
 358 forearm model and corresponding discretized volume mesh respectively.

359 This discretization process converts the continuous operator problem of
 360 Eq. (1) to a finite system of linear equations:

$$A\mathbf{v} = \mathbf{b} \quad (2)$$

361 where A is a symmetric and sparse $n_v \times n_v$ matrix, n_v is the number of mesh
 362 vertices, $\mathbf{v} \in \mathbb{R}^{n_v}$ is a vector of potential values at mesh nodes, and $\mathbf{b} \in \mathbb{R}^{n_v}$ is a
 363 vector containing source information. Because the electric potential is defined

364 up to a constant, the matrix A always has a one dimensional null space. To
 365 obtain a unique solution to the system of Eq. (2), we constrain potentials \mathbf{v} to
 366 have a zero sum.

367 In the context of EMG, we are not interested in finding electric potentials
 368 everywhere in the conductor, but only at the electrode locations. Let S be a
 369 selection matrix with a shape $n_e \times n_v$ which only selects the values at EMG
 370 electrode locations (n_e is the number of electrodes). Each row of S can be
 371 designed to select a single point location or to integrate over an area (e.g. the
 372 electrode-skin interface) depending on the location and number of its non-zero
 373 elements. Also, let $\mathbf{b}(\mathbf{r})$ correspond to a point source at location \mathbf{r} . The resulting
 374 EMG signal is thus given by:

$$\mathbf{v}_{point}(\mathbf{r}) = SA^{-1}\mathbf{b}(\mathbf{r}). \quad (3)$$

375 Let us analyze in more detail the structure of A and \mathbf{b} from Eq. (2). Let
 376 $\{w^i(\mathbf{r}), i = 1 \dots n_v\}$ be a set of n_v P^1 (piecewise linear) basis functions over the
 377 tetrahedral mesh Ω_t . Note, that w^i is 1 at the i -th vertex of the mesh, is 0 at
 378 all other vertices and is linear at all tetrahedra adjacent to the i -th vertex. In
 379 this case, A and \mathbf{b} have the following structure:

$$A_{ij} = \int_{\Omega_t} \sigma(\mathbf{r}) \nabla w^i(\mathbf{r}) \nabla w^j(\mathbf{r}) d\mathbf{r}$$

$$b_i = \int_{\Omega_t} I(\mathbf{r}) w^i(\mathbf{r}) d\mathbf{r}.$$

380 First, let us notice that A is symmetric and, in general, a very large matrix
 381 which can be stored only because it is sparse. Indeed, the functions w^i have a
 382 compact support and their pairwise scalar product is non-zero only for “neigh-
 383 bor” functions. Since the pseudo-inverse (or the inverse) of a sparse matrix
 384 is usually not a sparse matrix, it is impractical to compute it because of the
 385 amount of memory needed to store it. Thus, iterative methods are typically
 386 used to solve the system of Eq. (2) for every given \mathbf{b} .

387 Consider the case of $I(\mathbf{r}) = \delta(\mathbf{r} - \bar{\mathbf{r}})$ which corresponds to a unit point
 388 current source at location $\bar{\mathbf{r}}$. Without loss of generality, we assume that this
 389 source is inside a tetrahedron formed by the vertices i_1, \dots, i_4 of the mesh. In
 390 this case, we obtain:

$$b_i = \begin{cases} \lambda_j, & \text{if } i \in \{i_1, \dots, i_4\}, \\ 0, & \text{otherwise} \end{cases}$$

391 where $\{\lambda_j, j = 1, \dots, 4\}$ are the barycentric coordinates of the point $\bar{\mathbf{r}}$ inside the
 392 tetrahedron $\{i_1, \dots, i_4\}$. Applying this expression to Eq. (3), we obtain:

$$\mathbf{v}_{point}(\bar{\mathbf{r}}) = SA^{-1}\mathbf{b}(\bar{\mathbf{r}}) = SA^{-1}\bar{B}\boldsymbol{\lambda}.$$

393 where \bar{B} is a $n_v \times 4$ matrix with $\bar{B}_{i_j, j} = 1$ for $j = 1, \dots, 4$, and 0 otherwise. This
 394 implies that the solution of the system of Eq. (2) for any unit point source can be

395 computed as a barycentric sum of solutions on the vertices of the corresponding
 396 tetrahedron. Therefore, it is sufficient to compute solutions of Eq. (2) for “basis”
 397 sources located on mesh vertices, to be able to evaluate a solution for any point
 398 inside this mesh efficiently. Let n_s be the number of such basis sources. For the
 399 most general case, when the source can be located anywhere inside the mesh
 400 and $n_s = n_v$, let B be a $n_v \times n_s$ identity matrix. The objective is to compute
 401 “basis” solutions:

$$V_{basis} = SA^{-1}B \quad (4)$$

402 where V_{basis} is a $n_e \times n_s$ matrix, whose columns contain the solutions of Eq. (2)
 403 for a unit point source located at the corresponding mesh vertex. Hence, the
 404 potentials for any source location \mathbf{r} is given by:

$$\mathbf{v}_{point}(\mathbf{r}) = V_{basis}\boldsymbol{\lambda}(\mathbf{r}) \quad (5)$$

405 where $\boldsymbol{\lambda}(\mathbf{r}) \in \mathbb{R}^{n_s}$ is a vector, whose four non-zero elements contain the barycen-
 406 tric coordinates of point \mathbf{r} inside a corresponding tetrahedron. Note, that one
 407 may restrict potential sources to be located inside specific subdomains of the
 408 whole mesh (which is the case for EMG). In this case, n_s corresponds to the
 409 number of vertices of these subdomains, and the matrix B is a submatrix of the
 410 identity matrix.

411 The most straightforward way to compute V_{basis} from Eq. (4) is to solve a
 412 problem of the form $A\mathbf{x} = \mathbf{b}_i$ for each column of the matrix B . It would thus
 413 require solving n_s systems of linear equations. For realistic conductor geome-
 414 tries, which have a large number of vertices, solving a single system may take up
 415 to a few minutes and solving n_s systems quickly becomes impractical. There-
 416 fore, we propose the use of the adjoint method [37], which requires solving n_e
 417 systems only. In the context of EMG, the number of electrodes is usually signifi-
 418 cantly smaller than the number of vertices in the muscle subdomain meshes, i.e.
 419 $n_e \ll n_s$. Let us define $K = SA^{-1}$, which is a matrix of size $n_e \times n_v$. Because
 420 A is symmetric, and the inverse of a symmetric matrix is also symmetric, we
 421 can write $K^T = A^{-1}S^T$. Then, K can be found by solving the system:

$$AK^T = S^T. \quad (6)$$

422 The matrix S^T has n_e columns and, thus, only n_e linear systems need to be
 423 solved to find K . The basis solutions can then be found as:

$$V_{basis} = KB. \quad (7)$$

424 4.2 EMG signal of a single fiber activation

425 The action potential generated by the flow of ionic currents across the muscle
 426 fiber membrane is the source of excitation. For a given intracellular action
 427 potential (IAP) model $V_m(z)$, the transmembrane current source per unit length
 428 is proportional to the second derivative of $V_m(z)$, where z is a fiber arc length

429 measured in mm . A general description of the current source density traveling
 430 at velocity v along the fiber with the origin at the neuromuscular junction at
 431 location z_0 is [27, 29, 38]:

$$I(z, t) = \sigma_{in} \pi r^2 \cdot \frac{\partial}{\partial z} \left[\psi(z - z_0 - vt) w_{L_1}(z - z_0 - \frac{L_1}{2}) - \psi(-z + z_0 - vt) w_{L_2}(z - z_0 + \frac{L_2}{2}) \right] \quad (8)$$

432 where $z \in [0, L]$ is a location along the fiber of length L , $\psi(z) = \frac{d}{dz} V_m(-z)$, L_1
 433 and L_2 are the semi-lengths of the fiber from the end-plate to the right and to
 434 the left tendon, respectively, σ_{in} is the intracellular conductivity, and r is the
 435 fiber radius. We have chosen w_L to be a Tukey window, as proposed in [24].
 436 The IAP V_m [$\frac{mV}{mm}$] can be mathematically described in the space domain as
 437 proposed in [39]:

$$V_m(z) = 96z^3 e^{-z} - 90.$$

438 Let $\mathbf{r}(z)$ be a fiber geometry parametrized with respect to the fiber arc length
 439 z . Combining the transfer function of a point source in Eq. (3) with the fiber's
 440 current density in Eq. (8), we obtain the equation for the EMG signal resulting
 441 from a single fiber activation:

$$\mathbf{v}_{fiber}(t) = \int \mathbf{v}_{point}(\mathbf{r}(z)) I(z, t) dz. \quad (9)$$

442 This integral can be efficiently approximated by discretizing the fiber geometry
 443 into sufficiently dense spatial samples $\{\mathbf{r}(z_i)\}_i$ and assuming that $\mathbf{v}_{point}(\mathbf{r}(z))$
 444 is piecewise constant around these points. If we also rewrite Eq. (8) in a shorter
 445 form as $I(z, t) = \sigma_{in} \pi r^2 \cdot \frac{\partial}{\partial z} F(z, t)$, Eq. (9) becomes:

$$\begin{aligned} \mathbf{v}_{fiber}(t) &\approx \sum_i \mathbf{v}_{point}(\mathbf{r}(z_i)) \int_{z_i - \Delta_i}^{z_i + \Delta_i} I(z, t) dz = \\ &\sum_i \mathbf{v}_{point}(\mathbf{r}(z_i)) \int_{z_i - \Delta_i}^{z_i + \Delta_i} \sigma_{in} \pi r^2 \cdot \frac{\partial}{\partial z} F(z, t) dz = \\ &\sigma_{in} \pi r^2 \sum_i \mathbf{v}_{point}(\mathbf{r}(z_i)) \left(F(z_i + \Delta_i, t) - F(z_i - \Delta_i, t) \right). \quad (10) \end{aligned}$$

446 Note, that $\mathbf{v}_{point}(\mathbf{r}(z_i))$ can be efficiently computed from Eq. (5). Moreover,
 447 once $\mathbf{v}_{point}(\mathbf{r}(z_i))$ are computed for all given fibers, we can change the param-
 448 eters of the current source density (action potential waveform shape, propagation
 449 velocity, location of neuromuscular junction), and compute the corresponding
 450 EMG signal with Eq. (10) by only matrix multiplication complexity.

4.3 Geometrical and physiological modeling of motor units

The motor unit action potential (MUAP) is the summation of the single fiber action potentials (APs) of the muscle fibers in the MU. Different types of MUs can be modeled [40, 41]. Our approach consists in generating fiber and motor unit distributions inside a unit circle, and then projecting it into arbitrary 3D muscle geometry (Fig. 5D), using methods similar to those described in [42]. This provides a high level of control for the fiber and MU distribution parameters independently of a particular muscle geometry. A common way to simulate fibers and MUs is to start by defining MU positions, sizes and territories, and then simulate fibers inside these MUs according to their parameters [43, 44]. We, however, propose another approach. First, we simulate uniformly distributed fibers inside a unit circle. Then, MU centers and their circular territories are generated and, finally, we associate each fiber to an MU. A fiber is associated to one of the MUs that contains it inside its territory with a probability proportional to the MU density (Fig. 5C). This approach has two main advantages. First, it guaranties (by construction) the uniform fiber distribution inside a circular muscle cross-section. Second, once fibers are generated and projected into a muscle geometry, different MU distributions can be generated very quickly, without regenerating fibers and recomputing transfer functions $v_{point}(\mathbf{r}(z_i))$ for their nodes.

MU recruitment model

During muscle contraction, the MUs are recruited according to the size principle [23]. This can be simulated by associating a threshold of excitation to each MU, as described for example by Fuglevand et al. [45]. Linear or non-linear rate coding models can be used [45, 46, 47].

The excitation rate as a function of time for each muscle is converted into the firing rates of the active MUs. Inter-discharge intervals are then generated with variability of the discharges around the mean firing interval [48].

4.4 Implementation remarks

The implementation of the main steps presented in the previous section can be summarized as follows. Once the matrices S , A and B are computed, the matrix K is determined using Eq. (6) by solving n_e linear systems. Then, Eq. (7) is used to find the solutions for n_s basis points, which is a fast matrix multiplication operation. For any given point source location \mathbf{r} , we compute its barycentric coordinates in associated tetrahedron and apply Eq. (5) to get values of electrical potentials at electrode locations. Finally, for a given fiber geometry, the single fiber action potential as recorded by the EMG electrodes is computed using Eq. (9).

The results presented in this study are obtained using a Python implementation of the proposed strategy. Assembling the matrix A and solving the system (6) is delegated to the FEniCS computing platform [49, 50]. The

General basis points	Fibers basis points	Fibers EMG response	MUAPs assembling	Raw sEMG assembling
7 min (13 sec/elec)	2 min	30 sec	0.8 sec	2.6 sec

Table 1: Computational performance of each of the main steps of a raw EMG simulation. General basis points computation refers to equation (7); fiber basis points are computed with equation (5); fibers EMG response is computed with equation (9).

492 forearm geometry that is here representatively used as a conductor model is
493 taken from the website of BodyParts3D, The Database Center for Life Science
494 (<http://lifesciencedb.jp/bp3d/>). The volume mesh is generated from the
495 surface meshes of the forearm tissues using the CGAL C++ library [51].

496 4.5 Computational performance

497 In this section, we report the computational time of the proposed model for a
498 specific simulation case. The exact computational time values strongly depend
499 on the implementation, experiment design, model parameters etc. The order of
500 magnitude, however, stays the same. Note, that no multiprocessing tools were
501 used in these computations. Each step, however, is highly scalable and can
502 be efficiently distributed between parallel processes, which would significantly
503 increase the performance. Computations for each muscle and fiber are indepen-
504 dent and can be performed in parallel. Parallel computing would also apply to
505 the electrodes in the general basis points computation.

506 For the purpose of demonstration, we simulated a 1-min-long, 100% maxi-
507 mum voluntary contraction (MVC) excitation of the Brachioradialis muscle with
508 50000 individual fibers and 200 motor units. The mesh of the volume conductor
509 contained 2.1M vertices, which formed 13M tetrahedra. 16 rectangular and 16
510 circular electrodes were included in the model. The sampling frequency of the
511 simulated signals was 2000 Hz. Table 1 shows the computational time for each
512 of the main steps in this simulation.

513 An important property of our model is that each step depends only on the
514 data produced by the previous steps. This property can be exploited to change
515 some simulation parameters without recomputing every step of the simulation.
516 For example, it is not necessary to recompute solutions for the fiber basis points
517 if fibers geometry and conductor model stay the same and only the parameters
518 related to the fiber properties (AP velocity, end-plate location, tendon sizes,
519 etc.), MU distribution or recruitment model are modified. In this example, the
520 **total** simulation time for this new set of parameters will only take approximately
521 $30 + 0.8 + 2.6 = 33.4$ s.

522 A brief description of the main parameters required at each step follows. The
523 full arm and electrode geometry as well as the tissue conductivities define the

524 computation of general basis points. To compute fibers basis points solutions,
525 the 3D geometry of the fibers is required. Computing the fiber EMG responses
526 requires the shape of the intracellular AP waveforms, AP propagation velocity,
527 sizes of tendon and active fiber parts, neuromuscular junction location, fiber di-
528 ameter and intracellular conductivity, and sampling frequency. To compute the
529 MUs action potentials, the MU distribution in the muscle, i.e. the association of
530 fibers to each motor unit, need to be defined. In the proposed model, once the
531 number of MUs, their sizes and territory areas are selected, the MU distribu-
532 tion is randomly generated. Finally, to synthesize the sEMG signal, the muscle
533 excitation drives and recruitment model parameters (motor unit recruitment
534 thresholds and firing rates) are required.

535 **4.6 Comparison with the cylindrical analytical solution**

536 First, we compared our numerical solution with its analytical counterpart for a
537 simple volume conductor geometry [21]. We used a four layer cylindrical model
538 with layers corresponding to bone ($r = 0.7\text{cm}$), muscle ($r = 2\text{cm}$), fat ($r =$
539 2.3cm) and skin ($r = 2.4\text{cm}$) surfaces. 16 point electrodes were simulated on the
540 skin surface directly above a fiber. The fiber was located at varying depths into
541 the muscle tissue, in the range 1 mm to 11 mm. Differential sEMG signals were
542 simulated using the analytical and numerical solutions of the forward problem.

543 **4.7 Details of realistic simulation examples**

544 For the single muscle excitation example, 50k muscle fibers were generated inside
545 the muscle and distributed within 200 motor units. The size of MUs varied
546 exponentially from 11 to 1150 fibers. The areas of MU territories varied from
547 10% to 50% of the muscle cross-sectional area. The muscle excitation drive
548 was decomposed into MU impulse trains according to the size principle. In this
549 example, the firing rate for each MU ranged from 8 Hz to 35 Hz and all MUs
550 were recruited when an excitation level of 75% MVC was reached.

551 For the multiple muscles experiment, the flexor group included the Palmaris
552 longus, Flexor carpi ulnaris (ulnar head), Flexor carpi ulnaris (humeral head),
553 and Flexor carpi radialis muscles. The extensor group included the Extensor
554 digitorum, Extensor carpi ulnaris, Extensor carpi radialis brevis, and Exten-
555 sor carpi radialis longus muscles. During a wrist flexion, the muscles of the
556 flexor group reached an excitation level of 90% MVC. During extension, ex-
557 tensor group was activated with the same excitation level. Moreover, a small
558 but constant excitation of the abduction muscle group was added to simulate
559 the wrist resistance against gravity. The abduction muscle group included the
560 Flexor carpi radialis, Extensor carpi radialis brevis, and Extensor carpi radialis
561 longus muscles. For each muscle, a number of muscle fibers between 32k and
562 78k was simulated, depending on the muscle cross-sectional area. Muscle fibers
563 were distributed within motor units, whose number varied from 150 to 300 per
564 muscle.

565 4.8 Details of deep learning experiment

566 To evaluate the effect of using the simulation-pre-trained network, an experimen-
567 tally collected high-density surface electromyography (HD-sEMG) signal dataset
568 was used, originally created to test wrist-wearable interfaces[52]. The experi-
569 mental protocol was designed in agreement with the Declaration of Helsinki and
570 was approved by Imperial College London ethics committee (JRCO: 18IC4685).
571 Nine participants (4 females, 5 males, ages: 23-31) took part in the study after
572 signing informed consent forms. The participants performed 5-second isomet-
573 ric contractions of their dominant-hand index finger at 15% of maximal force,
574 with sEMG activity measured using two flexible 5x13 electrode grids with 8-mm
575 spacing placed on the circumference of the wrist, immediately proximal to the
576 ulnar head. HD-sEMG signal was sampled at 2048Hz, whilst force profiles were
577 sampled with a custom load cell at 10Hz. The signal was then decomposed into
578 motor neuron activity using convolutive blind source separation [53]. For the
579 purpose of training and testing the supervised decomposition pipeline, motor
580 neuron activity was accepted if it was present for at least 80% of the contraction
581 window. For each participant the HD-sEMG signal and accompanying decom-
582 posed motor neuron activity (as a sparse binary matrix) was then split into a 4
583 second training window and a 1 second testing window.

584 A gated recurrent unit (GRU) network was used as the deep learning model
585 due to previous studies showing good performance with this data type[26]. After
586 hyperparameter optimisation by grid search, a minimally-parameterised model
587 was found to perform optimally, likely due to the short length of the training
588 data available. Input HD-sEMG signal was first encoded by a single layer GRU
589 with a hidden dimension of 1024 in length[54]. To make a time instant predic-
590 tion a densely-connected linear layer with sigmoid activation function took as
591 an input a moving 20 sample-wide window from the GRU output, centred on
592 the time instant of interest. Predicted activity was converted to spike times-
593 tamps using a two-class K-means clustering algorithm. Binary cross entropy was
594 used as the loss function and Adam with weight decay used as the optimising
595 algorithm[55].

596 To improve model generalisation an early-stopping framework was used,
597 based on 10% of the training data retained as a validation set. Training, val-
598 idation and test data was z-score standardised using the mean and standard
599 deviation calculated from the training set. During training the input signal was
600 augmented with noise of standard normal distribution. To account for the high
601 sparsity of the output matrix, samples containing motor neurons were artifi-
602 cially oversampled, with each each input batch of 512 time instants containing
603 at least 20% motor neuron activation. All machine learning was implemented
604 using the pytorch library in python. Final performance was assessed using the
605 rate of agreement metric (RoA).

606 The optimised architecture of the GRU network was used for pre-training,
607 which was conducted using multi-task learning in a hard parameter sharing
608 paradigm[56]. Four digital twins were created for simulation using different
609 model parameters (tissue conductivities, MU distribution, fiber properties, etc.),

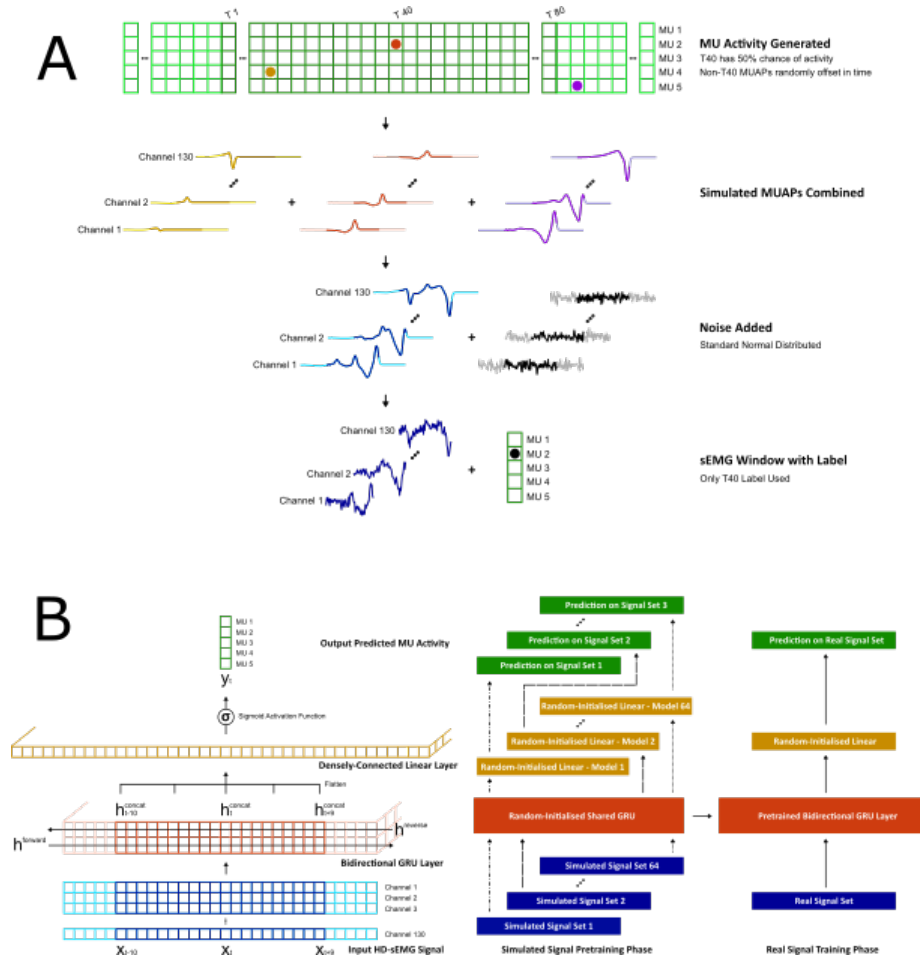


Figure 6: (A) Methodology used to build windows from the simulated MUAP template set for the pre-training phase. Each simulated template was 160 samples wide at a 2048Hz sampling rate and with 130 channels. First either a MUAP template was placed in the centre of the window or it was left empty at a 50% probability. Then MUAP templates from other MU classes were added to the window at a random offset to generate superpositions. Finally standard normal distributed noise was added to the window, with the central 80 samples then paired with the label for supervised learning. (B) The neural network architecture and pre-training methodology used to improve the performance of a deep learning-based HD-sEMG decomposition algorithm. The neural network consists of a single gated recurrent unit layer, with predictions made using a 20-sample wide window of the hidden vector output, which is flattened before being passed to a sigmoid-activated densely-connected linear layer. In the pre-training phase a multi-task learning regimen is used to optimise the parameters of the gated recurrent unit using the simulated sEMG. This pre-trained layer can then be used to improve the optimisation performance on real sEMG data.

610 with the generated motor unit activation (MUAP) templates from flexor dig-
611 itorum profundus and superficialis used to create 64 sets, each containing 5
612 MUAPs. Each set was used to generate windows of signal with a range of
613 MUAP superpositions (Fig. 6A). In signal windows with motor neuron activity
614 a MUAP template was placed in the centre of the window, before being ad-
615 ditively superimposed with a random number of MUAP templates from other
616 motor units at random time offsets. In windows without activity no template
617 was placed in the centre of the window. During multi-task learning training, the
618 same GRU layer (and parameters) were shared between the 64 recordings, but
619 each recording had its own output layer, operating on a 20 sample-wide window
620 as in the experimental recordings (Fig. 6B). In this way the GRU layer was
621 trained to act as a more general feature extractor, whilst the individual linear
622 output layers made class predictions specific to each recording. Training again
623 used noise augmentation, binary cross-entropy and Adam with weight decay.

624 To use the simulation-pre-trained network in the experimental data the GRU
625 parameters from the pre-trained network were used, whilst the linear output
626 layer used a normal random initialisation. This was compared to a normal
627 random initialisation of both the GRU and output layer. In both instances
628 the network was trained using the methodology specified above, with the only
629 difference being whether the GRU layer was simulation-pre-trained or not.

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637 **Author contributions**

638 KM, SDG, and DF conceptualized the study. KM and SDG developed the soft-
639 ware implementation of the simulator. KM, AC, IMG, SDG, and DF performed
640 the experimental measures and conceptualized the data analysis. AC and IMG
641 performed the data analysis. KM, SDG, and DF prepared the first draft of the
642 manuscript. All authors edited the manuscript for important scientific content
643 and all approved the final version.

644 **Competing interests**

645 KM and SDG are founders of the company Neurodec which specializes in EMG
646 simulation and analysis.

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