

Microscale concert hall acoustics to produce uniform ultrasound stimulation for targeted sonogenetics in hsTRPA1-transfected cells

Aditya Vasan Florian Allein Marc Duque Uri Magaram Nicholas Boechler Sreekanth H. Chalasani James Friend

Aditya Vasan, James Friend

Medically Advanced Devices Laboratory, Department of Mechanical and Aerospace Engineering, Jacobs School of Engineering and Department of Surgery, School of Medicine, University of California San Diego, La Jolla CA 92093 USA

Uri Magaram, Marc Duque, Sreekanth H Chalasani

Molecular Neurobiology Laboratory, The Salk Institute for Biological Studies, La Jolla, CA 92037

Florian Allein, Nicholas Boechler

Department of Mechanical and Aerospace Engineering, University of California San Diego, La Jolla CA 92093 USA

Email address: jfriend@ucsd.edu

Webpage: <http://friend.ucsd.edu>

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1 The field of ultrasound neuromodulation has rapidly developed over the past decade, a consequence of the discovery of strain-sensitive struc-
2 tures in the membrane and organelles of cells extending into the brain, heart, and other organs. Notably, clinical trials are underway for treating
3 epilepsy using focused ultrasound to elicit an organized local electrical response. A key limitation to this approach is the formation of stand-
4 ing waves within the skull. In standing acoustic waves, the maximum ultrasound intensity spatially varies from near zero to double the mean in
5 one half a wavelength, and can lead to localized tissue damage and disruption of normal brain function while attempting to evoke a broader re-
6 sponse. This phenomenon also produces a large spatial variation in the actual ultrasound exposure in tissue, leading to heterogeneous results
7 and challenges with interpreting these effects. One approach to overcome this limitation is presented herein: transducer-mounted diffusers that
8 result in spatiotemporally incoherent ultrasound. The signal is numerically and experimentally quantified in an enclosed domain with and with-
9 out the diffuser. Specifically, we show that adding the diffuser leads to a two-fold increase in ultrasound responsiveness of hsTRPA1 transfected
10 HEK cells. Furthermore, we demonstrate the diffuser allow us to produce an uniform spatial distribution of pressure in the rodent skull. Collec-
11 tively, we propose that our approach leads to a means to deliver uniform ultrasound into irregular cavities for sonogenetics.

12 1 Introduction

13 SOUND diffusers have been applied to concert hall acoustics since the 1800s, when ornamentation along
14 the walls or concave ceilings were used to introduce greater binaural dissimilarity^[1]. Evolution of concert hall architecture eventually led to the development of the Schröder diffuser, arguably the first practical
15 technique to disperse sound in a predictable manner. In the 1970s, Schröder proposed the phase grating dif-
16 fuser^[2], a method to artificially create diffuse reflection. Composed of regular wells of different depths, these
17 structures have periodicity in two dimensions as governed by a pseudostochastic sequence. In the typical
18 configuration, waves incident on this structure undergo phase shifts corresponding to the depth of the wells
19 through which they travel. The structure then scatters sound rather than reflecting it, depending on the mag-
20 nitude of these phase shifts. This method has been widely adopted in architectural acoustics, where sound
21 absorption—the only feasible alternative—is undesirable. This method has also been applied to ultrasound
22 imaging^[3] and microparticle separation^[4] where sound absorption is likewise difficult. More recently, the prin-
23 ciple of applying phase shifts to a coherent ultrasound field has led to development of acoustic holography^[5,6].
24 This novel approach has enabled the generation of customized amplitude profiles based on the location and
25 shape of the target region but does not enable the creation of spatiotemporally incoherent fields within an
26 enclosed cavity.
27

28 Ultrasound transducers have been used for imaging tissue^[7], disrupting blood-brain barriers^[8], invasive^[9] and
29 non-invasive neuromodulation^[10], and thrombolysis^[11]. In these cases, ultrasound is typically focused at a

30 certain depth defined either by the radius of curvature of the transducer, or an acoustic lens¹², or phased ar-
31 rays¹³. The fundamental limitation of these approaches is the formation of standing waves due to resonant
32 reflections within the skull cavity formed by the relatively high impedance of the skull's cortical bone com-
33 pared to the tissue of the brain, and thus regions of either extremely high intensity or zero intensity at every
34 one-half an acoustic wavelength¹⁴. The presence of these local maxima may lead to unintended bioeffects
35 in tissues when applied to neuromodulation, including heating or even tissue damage. Such adverse effects
36 in tissue have been reported during ultrasound-driven thrombolysis and blood brain barrier disruption^{15,16}.
37 Additionally, commonly used transducer materials such as lead zirconate titanate (PZT) also have limitations
38 in high power applications at frequencies above ~ 1 MHz, producing losses, hysteresis, and internal (ohmic)
39 heating as current passes through elemental lead present at the morphological grain boundary¹⁷. Another
40 important consideration for sub-MHz frequencies, which have been used for neuromodulation studies in
41 the past¹⁸, is the lowering of the cavitation threshold to a level that may elicit tissue damage at clinically rel-
42 evant ultrasound amplitudes¹⁹. One approach to overcome these limitations is to build resonant devices us-
43 ing loss-free, single-crystal piezoelectric material operating in the 1–10 MHz range that are capable of deliv-
44 ering a spatiotemporally diffuse ultrasound field for various applications, including sonogenetics.

45 Sonogenetics relies on genetically engineering cells to be more sensitive to mechanical stimuli using mem-
46 brane bound proteins^{20,21}. This technique eliminates the need for focused ultrasound by ensuring that tar-
47 geted neural circuits are the only ones that will respond to an ultrasound stimulus. Recent work has revealed
48 that one protein in particular, human transient receptor potential A1 (*hsTRPA1*), produces ultrasound-evoked
49 responses in several cell types²¹. One limitation of sonogenetics is that existing transducers producing plan-
50 nar or focused ultrasound are unsuitable. Furthermore, in many applications, the transducer must be small
51 to avoid affecting animal behavior. Unfortunately, no small broadband transducers exist^{22,23} that might fa-
52 cilitate the generation of spatiotemporally random ultrasound noise from a similarly random input signal at
53 sufficient power for sonogenetics. Moreover, commonly used animal models like rodents have small heads
54 with a typical mass of 3–4g²⁴, less than half the mass of all commercially available or research-based²⁵ power
55 ultrasound transducers known to the authors. The effective implementation of sonogenetics requires a very
56 different transducer design. It must reduce interference between the radiated and reflected ultrasound, pro-
57 duce diffuse and uniform ultrasound throughout the region, transduce sufficient power to produce over 0.4 MPa
58 acoustic pressure in tissue while remaining sufficiently small and light enough to attach to the head of a live
59 mouse. In addition, these devices also have to avoid generating electromagnetic signals and localized tem-
60 perature changes. If left to appear, these phenomena may conflate with the effects of ultrasound on the cells
61 in sonogenetics experiments, reducing one's confidence in ultrasound's contribution to the observations.
62 Transducers that can be attached to freely moving mice enable the study of neural circuits in their native
63 state, without the confounding effects of anaesthesia as reported in past studies^{26,27}.

64 We have overcome the limitations of existing transducers by incorporating a machined diffuser *on the trans-*
65 *ducer face* in order to produce spatiotemporally incoherent ultrasound. Diffusers are typically used in reduc-
66 ing coherent reflected sound—echoes—and their use *on the sound generator itself* has not been reported to
67 the knowledge of the authors. A diffuser is ideally suited for sonogenetics as it nearly losslessly eliminates
68 the presence of regions of either high or low intensity within an enclosed cavity, in both *in vitro* assays and
69 within the rodent skull for longer-term applications. First, we discuss the design of the diffuser and validate
70 its performance using numerical simulations. We then address the challenge of fabrication of a complex three-
71 dimensional structure at sub-millimeter scales, as conventional photolithography, three-dimensional print-
72 ing, and classic machining techniques are unsuitable for this task. Next, we characterize a device which has
73 been coupled to a lithium niobate transducer operating in the thickness mode. We use lithium niobate, a sin-
74 gle crystal material with low losses and no hysteresis over the MHz operating frequency range required for
75 ultrasound neuromodulation²⁸.

76 Forming uniform acoustic pressure fields in enclosed cavities is a key challenge in a variety of biomedical ap-
77 plications. Coherent propagation and reflection of ultrasound naturally forms antinodes and nodes in a cav-
78 ity, corresponding to regions of strong and weak ultrasound. Uniform ultrasound reduces the risk of over or
79 under-exposing the tissue regions that contain ultrasound-sensitive proteins. We present an application of

80 the Schröder diffuser to screen for ultrasound-sensitive ion channels in human embryonic kidney (HEK293)
 81 cells *in vitro* for the purposes of identifying and isolating targets for sonogenetics in freely moving mice. We
 82 also verify the presence of nearly equivalent acoustic pressures across two deep brain regions in an *ex vivo*
 83 model.

84 2 Results

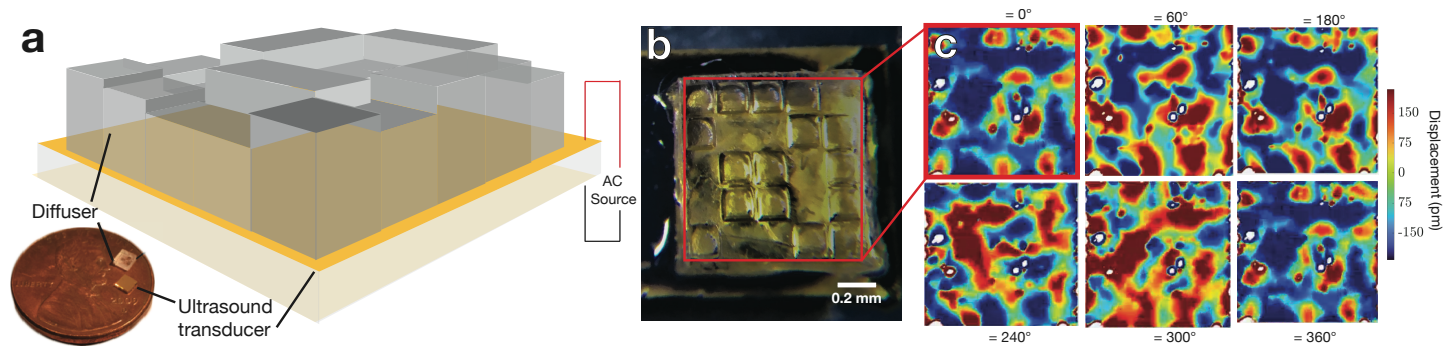


Figure 1: (a) A diffuser design based on Schröder’s method of quadratic-residue sequences to determine well depth. The wells were machined in glass using a KrF excimer laser system with a custom metal mask to restrict beam width. The machined depth of the pillars can be up to 309 μm . (b) The glass diffuser block was then (c) bonded to a transducer operating in the thickness mode at 7 MHz using an ultraviolet light-curable epoxy. (c) A scanning laser Doppler vibrometer image of the diffuser face in the time domain shows phase differences corresponding to pillar heights (normalized autocorrelation > 0.73).

85 The design of the Schröder diffuser is based on quadratic-residue sequences defined by $s_n = n^2$, where n^2 is
 86 the least non-negative remainder mod N , with N always an odd prime. One of the properties of this number
 87 sequence relevant to the design of an optimum diffuser is that both the Fourier transform of the exponential
 88 sequence $r_n = \exp(i2\pi s_n/N)$, and by extension the scattered wave produced by it, have a constant magni-
 89 tude^[129] $|R_m|$ such that

$$|R_m|^2 = \left| \frac{1}{N} \sum_{n=1}^N r_n e^{-\frac{2\pi i m n}{N}} \right|^2 = \frac{1}{N}, \quad (1)$$

90 where $\iota = \sqrt{-1}$.

91 We may then use this to define the wells’ depths, $d(x_n, y_n)$, corresponding to the number sequence. In one
 92 dimension, the depth of the n^{th} well is given by^[30]

$$L_n = \frac{\pi c n^2 (\text{mod } N)}{N \omega_r}, \quad (2)$$

93 where ω_r is the design frequency, N is a prime number, and c is the speed of sound in the medium. Extend-
 94 ing the concept of a diffuser defined per the above numerical sequence to two dimensions involves replacing
 95 n^2 in the above formula with $n^2 + m^2$, where m represents the number of wells in the second dimension. A
 96 representative image of a diffuser fabricated using a two-dimensional sequence is shown in Fig. 1.

97 While a one-dimensional diffuser creates a uniform two-dimensional pressure field, a two-dimensional dif-
 98 fuser with varying well depths creates a uniform three-dimensional pressure field. Ultrasound neuromodu-
 99 lation typically relies on frequencies in the 1–10 MHz range^[31] and this requires sub-millimeter well depths
 100 as defined by eqn. (2). Although structures based on the quadratic-residue sequence have been achieved
 101 at the macro-scale in two dimensions and at the microscale in one dimension^[4], it has not been achieved in
 102 two-dimensional structures on the micron to sub-mm scale due to the lack of established fabrication tech-
 103 niques for these dimensions^[32]. Conventional photolithography is good for creating patterns that have the
 104 same depth or, at most, a few different depths. It becomes challenging when features of varying depths are

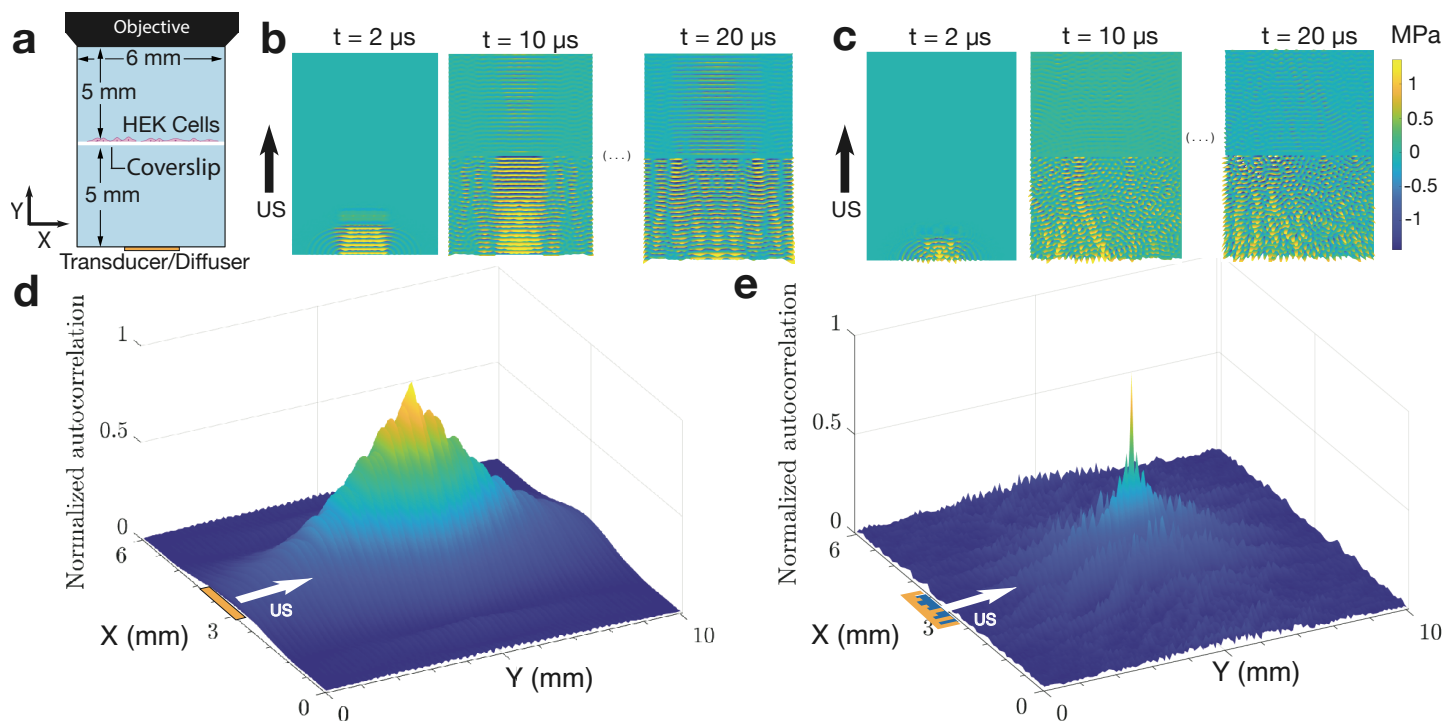


Figure 2: Two-dimensional pressure profile for the (a) domain (b) with and (c) with the diffuser. Human embryonic kidney (HEK) cells were placed in the middle of the (light blue) fluid domain with an objective lens for an inverted microscope at top. The pressure field was generated by defining a sinusoidal pressure displacement to the transducer face, located at the bottom of the domain. Pressure maps were extracted from the results in $1\mu\text{s}$ time steps over grid points specified within the domain. A two-dimensional autocorrelation was performed on this grid over time; each X,Y point plotted from the (d,e) results of the autocorrelation corresponds to a single point in the (a) domain. Spatial and temporal periodicity are observed through the existence of a large value of autocorrelation over the domain (d) without the diffuser compared to the generally low autocorrelation with (e) the diffuser introduced onto the transducer.

105 desired because multiple photolithography and etching steps are required. Alternate approaches, including
 106 three-dimensional or two-photon printing methods are unable to produce acoustically low-loss structures
 107 with sufficient dimensional accuracy at these scales. We sought to address these limitations by using an ex-
 108 cimer laser to machine sub-millimeter pillars of varying heights in glass in two dimensions.

109 A time-domain scan using laser Doppler vibrometry (LDV; *see* Methods) in Fig. 1 shows significant phase cor-
 110 relation (normalized autocorrelation > 0.73) with the machined geometry. The transducer was driven at its
 111 resonance frequency with a sinusoidal input power range of $0.5 - 2\text{ W}$ and a peak pressure output of 0.6 MPa
 112 as measured with a fiber optic hydrophone³³.

113 Finite element analysis (COMSOL 5.5, Comsol Inc., Los Angeles, CA USA) was used to validate the design of
 114 the diffuser. The domain was chosen to mimic an experimental setup used for identifying ultrasound sensi-
 115 tive ion channels in an *in vitro* setup. This consists of an inverted fluorescence microscope with a custom
 116 perfusion chamber to house a cover slip and transducer. The simulation domain is illustrated in Fig. 2 and
 117 specific dimensions of the domain and simulation parameters are included in the *Methods*. The transducer
 118 and the diffuser assembly was fixed at the bottom of the domain. A custom perfusion chamber that contains
 119 a slot for a coverslip was mounted over the ultrasound source. The transducer was coupled to the coverslip
 120 through water and there was a layer of media above the cover slip. The walls were defined to be hard bound-
 121 aries with the acoustic impedance $Z_i = \infty$ such so that the normal derivative of the total acoustic pressure,
 122 $\frac{\partial p_t}{\partial n} = 0$.

123 The diffuser consists of seventeen elements, the heights of which were calculated from eqn. (2). The cover-
 124 slip in serves as a solid boundary and allows the evaluation of the acoustic field in the closed domain below
 125 and the open domain above it, corresponding to the different boundary conditions assigned to the model.

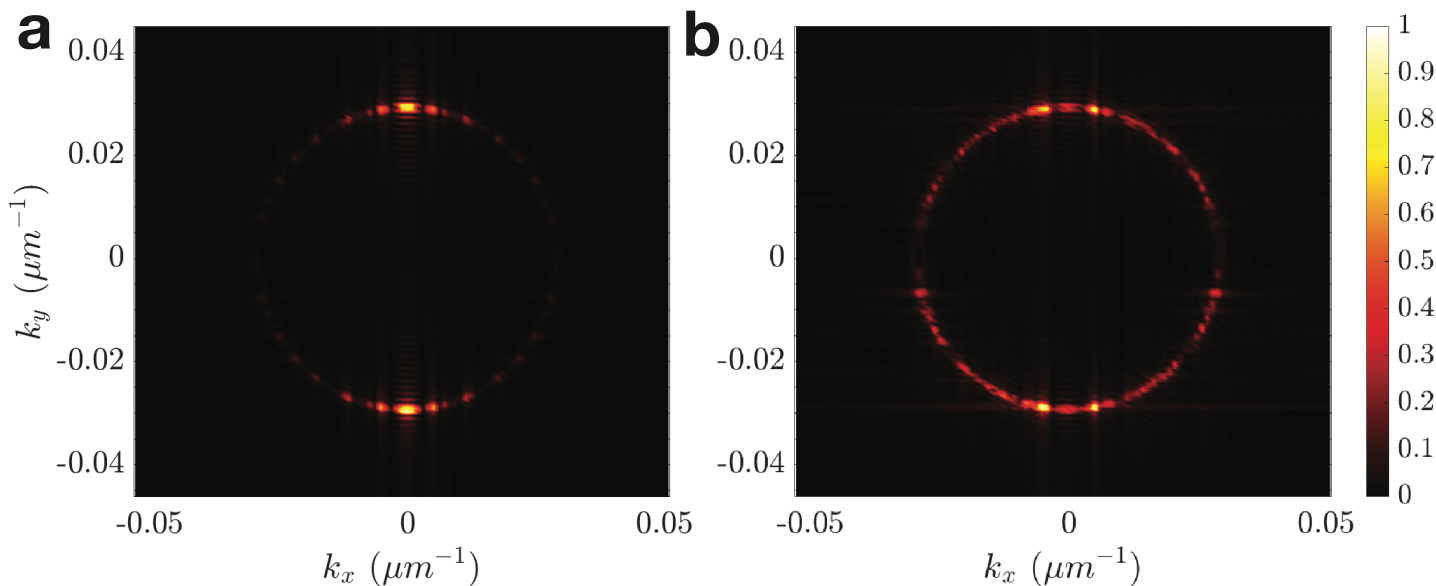


Figure 3: Isofrequency contour at the driving frequency (a) without and (b) with the diffuser. The circular profile traced by both cases corresponds to the wave vector in water at the driving frequency. The transducer (b) with diffuser produced wave vectors spread around this circular profile, indicating a more uniform distribution of the ultrasound. Without the diffuser, most of the wave is isolated to propagation along the Y axis.

126 The time variation of the pressure field with and without the diffuser was evaluated (supplementary video S1
 127 and S2). Several points in the fluid domain were chosen and the time evolution of the pressure field for the
 128 two cases was compared using the techniques described in the *Methods* section. A two-dimensional autocor-
 129 relation was calculated in order to determine if there were any locations within the domain with coherence
 130 (echoes) or localized increases or decreases (constructive and destructive interference) in ultrasound inten-
 131 sity.

132 Spatial and temporal patterns that form over the duration of the stimulus are represented by a two-dimensional
 133 autocorrelation in Fig. 2. It is evident that there is both spatial and temporal periodicity with the transducer
 134 alone (Fig. 2a and supplementary video S1) that is greatly reduced when the diffuser is introduced (Fig. 2b
 135 and supplementary video S2). Videos of the sample autocorrelation in the domain over the stimulus dura-
 136 tion are presented in the supplementary information (videos S3 and S4) and show that there is greater au-
 137 tocorrelation over the duration of the stimulus without the diffuser (video S3). This indicates that the ultra-
 138 sound field with the diffuser is temporally inconsistent. The autocorrelation plot is an instructive technique
 139 to determine periodicity in time and space but does not quantify spatial dispersion in the domain at the driv-
 140 ing frequency.

141 For the purpose of quantifying the dispersion at 7 MHz, an isofrequency contour plot is provided in Fig. 3(a)
 142 without and (b) with the diffuser. Without the diffuser, wave vectors are only present in the vicinity of $k_x = 0$,
 143 along the direction of propagation of the pressure wave in the medium: the Y axis. The angular spread is 20°
 144 on either side of the direction of propagation without the diffuser. Particularly, the majority of the wave can
 145 be seen to be propagating along the Y axis, with significant sidelobes immediately to the left and right and
 146 much smaller sidelobes slightly farther away. Including the diffuser produces wave vectors beyond the main
 147 direction of propagation (Fig. 3b), with significant components oriented along directions from the Y axis
 148 (along k_x) to the X axis (along k_y). The previously significant sidelobes remain significant, but are aug-
 149 mented by wave propagation beyond 45° in the XY plane. This indicates strong dispersion in the domain when in-
 150 cluding the diffuser.

151 To verify the effects of the diffuser *in vitro*, we used an upright optical imaging setup including an immer-
 152 sion objective, a custom perfusion chamber, and the diffuser assembly. The diffuser assembly and the test
 153 setup are represented in Fig. 4a; we used lithium niobate due to its relatively high coupling coefficient and

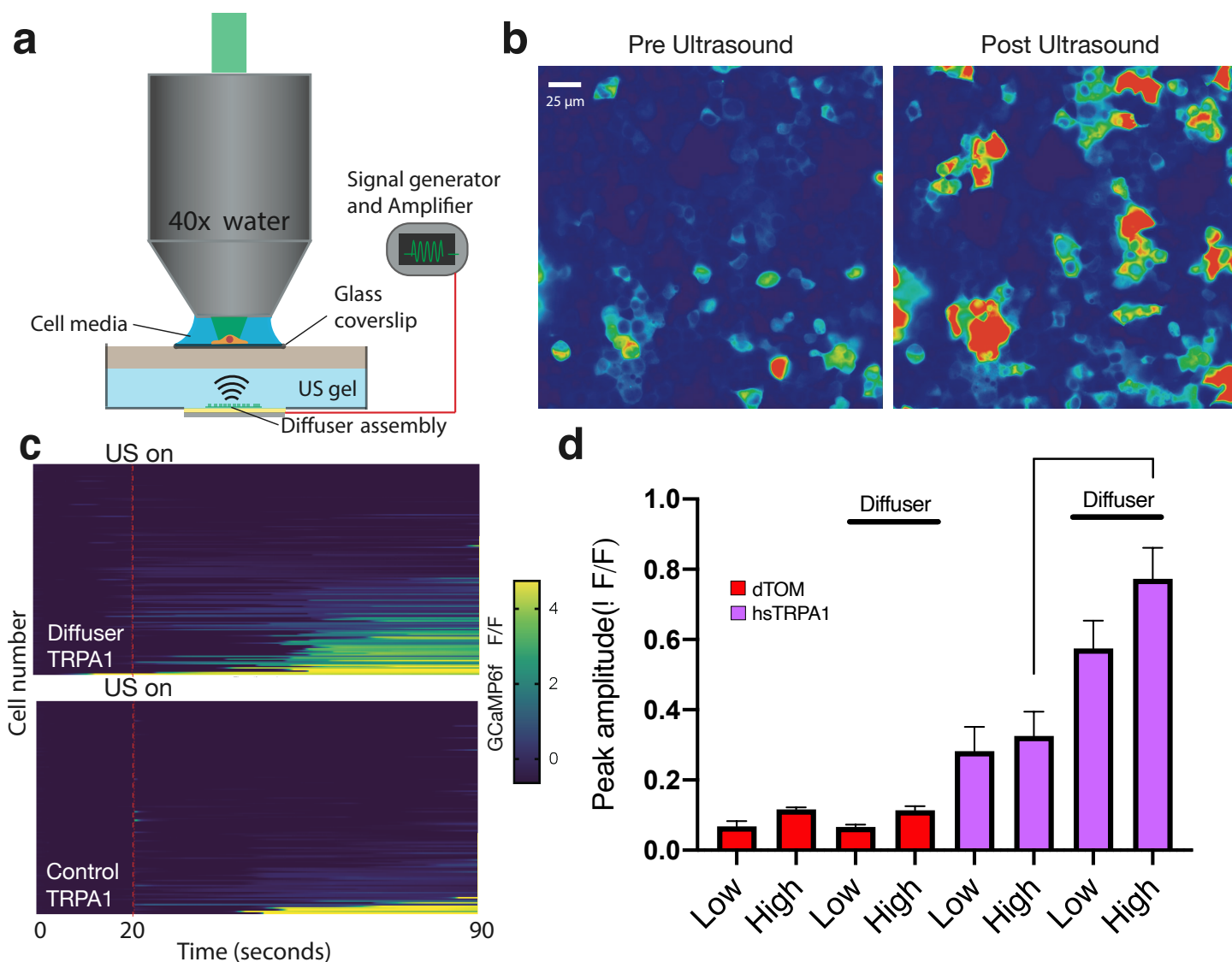


Figure 4: (a) The experimental setup for confirming the utility of the diffuser in an *in vitro* setting consists of an upright epifluorescent microscope, an immersion objective, and a chamber that houses cells on a coverslip and the diffuser assembly. Standing wave components may exist between the transducer and the cover slip and between the cover slip and the immersion objective. The calcium concentration before and after ultrasound stimulation in the same field of view is (b) shown for HEK cells expressing hsTRPA1. Comparison of fluorescence changes as measured using GCaMP6f reporters with respect to time for two cases, (c) without (control) and with the diffuser show an increase in both number and magnitude of cells being activated upon introduction of the diffuser. (d) HEK cells expressing TRPA1 show a greater response to ultrasound stimuli with a diffuser present in comparison to both no diffuser and dTom-based controls. The magnitude of the response when the diffuser is used is significantly greater (over twice as high) than when the diffuser is not used ($n = 76 - 221$, $p < 0.0001$ by a Mann-Whitney test).

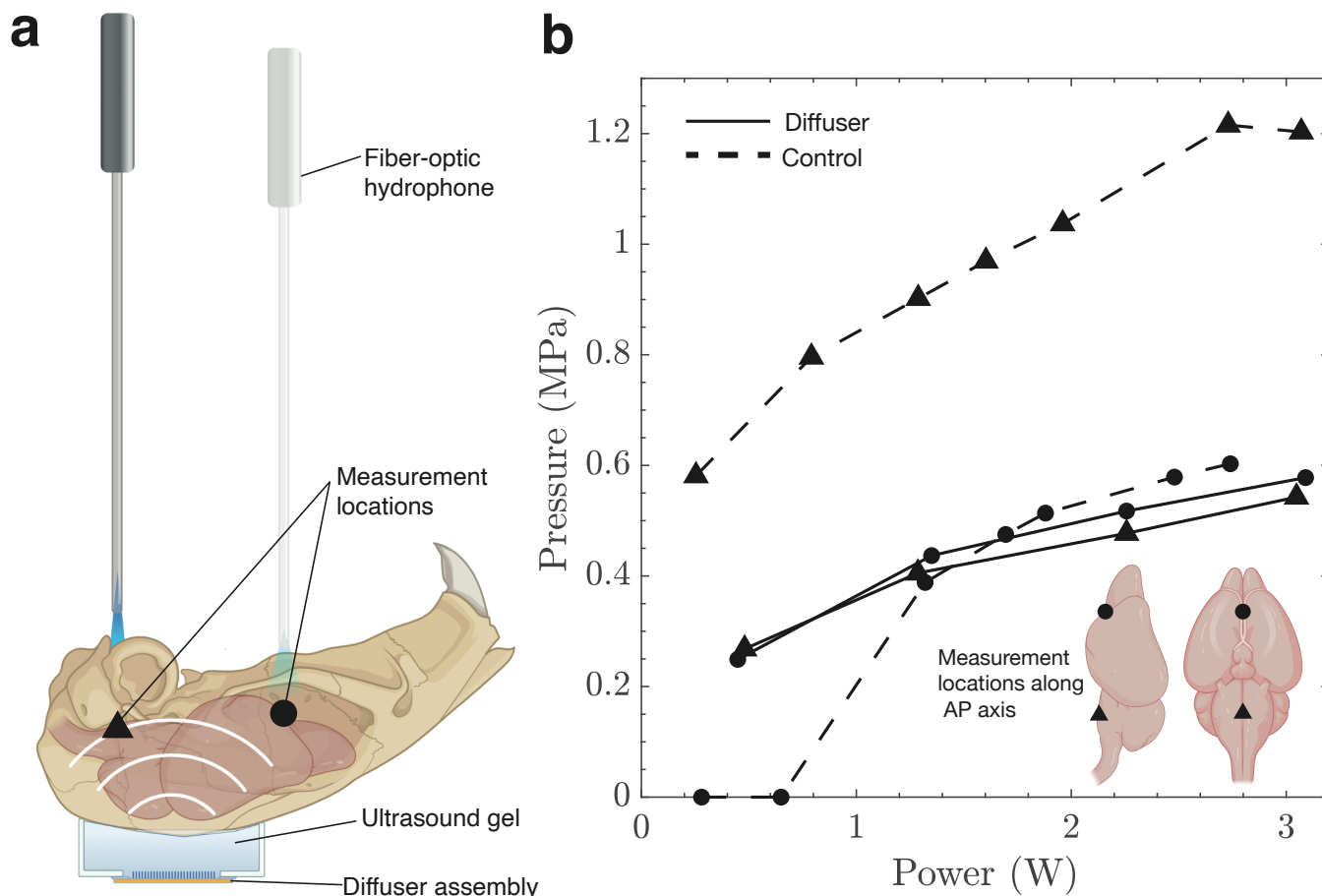


Figure 5: The pressure was measured using (a) a fiber optic hydrophone at two different locations along the anterior-posterior axis: the ventral surface of the pons (triangle) and the ventral surface of the anterior olfactory bulb (circle). The measured pressure is (b) uniform across different brains for different input powers above 0.2 MPa (minimum detectable pressure using our setup), indicating that the diffuser creates a uniform acoustic field within the skull cavity. This eliminates the influence of the cranial structure and ensures that only the brain regions that have been infected with hsTRPA1 will be sensitive to ultrasound stimuli. In comparison, the control case without the diffuser shows a three-fold deviation in pressure values for the same input power for different brain regions along the AP axis.

154 zero hysteresis³⁴ which implies no heating from the piezoelectric material itself. Human embryonic kidney
 155 (HEK293) cells expressing GCaMP6f³⁵ were transfected with hsTRPA1. We compared fluorescence changes
 156 ($\frac{\Delta F}{F}$) for four cases, with and without the channel, without the diffuser assembly (transducer alone), and with
 157 the diffuser assembly. Representative GCaMP6f images of HEK293 cells transfected with hsTRPA1 are shown
 158 in Fig. 4b and heat maps of fluorescence intensity with respect to time are presented in Fig. 4c, with a clear
 159 increase in both the magnitude and number of cells being activated with the diffuser assembly. Cells express-
 160 ing hsTRPA1 and controls were tested at two different pressures, 0.32 MPa and 0.65 MPa. There was a con-
 161 sistent increase in fluorescence intensity with an increase in acoustic pressure for both the control and the
 162 hsTRPA1 condition, whether or not the diffuser was present. Including the diffuser increased the mean fluo-
 163 rescence amplitude by at least a factor of two for cells that had been infected with hsTRPA1 ($p < 0.0001$,
 164 Fig. 4d).

165 We also tested the effects of ultrasound on mouse primary cortical neurons. Neurons were infected with adeno-
 166 associated viral (AAV) vectors to express hsTRPA1 and a genetically encoded calcium indicator, GCaMP6f³⁵,
 167 or a control with only the calcium indicator. We found that ultrasound triggered an increase in calcium up-
 168 take in both cases, with the hsTRPA1 neurons showing a greater number of activated cells in comparison to
 169 the control (Fig. S2). A video of real-time calcium response in hsTRPA1-expressing neurons is presented in
 170 the supplementary information (Video S5).

171 The uniform nature of the ultrasound field created by the diffuser was also verified *ex vivo* in a mouse skull
172 *see Methods*. Pressure measurements were taken at two different locations as indicated in Fig. 5 along the
173 anterior-posterior axis, at the ventral surface of the pons and the ventral surface of the anterior olfactory
174 bulb. With the diffuser, the pressure at both these locations was uniform, with minimal deviation between
175 them and a uniform increase with input power to the transducer (Fig. 5). However, the transducer alone pro-
176 duced diverging values of pressure at these positions, so much so that the pressure at the pons (triangle) ex-
177 ceeded the pressure at the anterior olfactory bulb (circle) by a factor of 3 at an input power of 3 W, yet fell be-
178 low the hydrophone's minimum measurement value, 0.2 MPa, at the anterior olfactory bulb when using less
179 than 1.25 W of power. By contrast, the diffuser had minimal deviation in pressure values at these locations,
180 with pressure values ranging from 0.25 - 0.5 MPa at the ventral surface of both the pons and the anterior ol-
181 factory bulb. These brain regions were chosen not for their function, but as they were as they were the far-
182 thest from each other and the device, while keeping as much of the mouse skull intact during preparation
183 [see methods]. Collectively, these results demonstrate that the diffuser is capable of delivering uniform ultra-
184 sound fields *in vivo* in comparison to the a transducer alone, thus enabling sonogenetic studies across large
185 brain regions.

186 3 Discussion

187 Existing non- and minimally-invasive techniques to stimulate brain regions, such as transcranial magnetic
188 stimulation and transcranial direct current stimulation, offer poor spatial resolution. This is a problem for
189 precisely targeting brain regions that have specific functions. Ultrasound-based stimulation enables target-
190 ing brain regions with sub-millimeter-scale accuracy. This precision can be achieved in different ways, either
191 by using an array to focus ultrasound to a specific region³⁶ or by using sonogenetics to engineer cells to lo-
192 cally be more sensitive to mechanical stimuli.

193 Still, the limitation with focused ultrasound is the alteration in the position and shape of the focal zone due
194 to spatial variations in acoustic impedance³⁷. Sonogenetics is an attractive option because of the potential
195 of having a toolkit of specific proteins that can be engineered to be sensitive to ultrasound stimuli at different
196 frequencies or pressures. Current ultrasound transducers and how ultrasound interacts with the skull cavity
197 are important limitations to translate Sonogenetics. Standing waves in the skull cavity produce nodes and
198 antinodes, each separated by one-half of the acoustic wavelength and responsible for pressure minima and
199 maxima, respectively. This can lead to haemorrhage and heating in tissue¹⁵ in past studies. A further limita-
200 tion is that ultrasound transducers cannot be driven using broadband white noise to produce a spatiotempo-
201 rally random acoustic field at pressures sufficient to elicit cellular responses. The development of sonogenet-
202 ics that started with the *TRP4* channel has expanded to include a library of proteins that are sensitive to ul-
203 trasound stimuli at different ultrasound stimulation parameters. Examples include MSC³⁸, TREK³⁹, Piezo⁴⁰,
204 and other TRP channels⁴¹, all have which have been shown to be sensitive to ultrasound *in vitro*.

205 Designed via computational analysis and fabricated with an excimer laser, a microscale Schröder diffuser
206 was devised to eliminate the spatiotemporally heterogeneous distribution of ultrasound by placing it upon
207 the emitting transducer. The transducer alone was shown to produce standing waves in the absence of the
208 diffuser. With the diffuser in place, autocorrelation of the ultrasound field quantifies the elimination of the
209 standing waves and consequent suppression of antinodes associated with potential tissue damage. This re-
210 sult is confirmed for the resonant frequency of operation of the transducer using an isofrequency contour.
211 We verified the predictions of the simulation *in vitro* using HEK293 cells and neurons that were transfected
212 with a sonogenetic candidate, hsTRPA1.

213 Development of sonogenetics in larger animal models—such as primates—will require ultrasound trans-
214 ducers that are capable of delivering an acoustic field that is spatially and temporally incoherent, as we have
215 shown with the diffuser assembly. This ensures that the pressure in different brain regions is uniform over

216 the stimulus duration, thus eliminating the aberrations in the acoustic field due to the skull cavity. Function-
217 alization of specific brain regions using ultrasound-sensitive proteins can offer sub-millimeter spatial pre-
218 cision. Localization of Sonogenetic proteins in combination with an acoustic field provided by a diffuser as-
219 sembly will ensure that the observed neuromodulatory effects are solely due to ultrasound activation of tar-
220 geted regions of tissue and not due to the confounding effects of reflection or interference from the geometry
221 of the skull.

222 4 Methods

223 Ultrasound transducers

224 Ultrasound transducers used in this study were single crystal lithium niobate transducers operating in the
225 thickness mode with lateral dimensions of 5x5 mm and thickness 500 μm . The 128YX cut of lithium niobate
226 was used and the fabrication process involved cleaning of the wafer with acetone, isopropyl alcohol, and
227 ultra-pure deionized (DI) water followed by sputtering both sides with an adhesion layer of 20 nm titanium
228 followed by 1 μm gold. The deposition parameters were (with a Denton Discovery 635, Denton Vacuum LLC,
229 New Jersey, USA) 5–10 nm of Ti at 1.2–1.6 A/s with the power set to 200 W, with argon as the gas in the cham-
230 ber at 2.3 mT and the stage rotating at 13 RPM to ensure uniform deposition over the sample. The thickness
231 of gold deposited was 1 μm at a rate of 7 – 9 A/s.

232 Diffuser fabrication and characterization

233 Code (MATLAB, Mathworks, Massachusetts, USA) was used to define the well depths based on the medium
234 of choice and was used to define a computer-aided design program that controlled the operation of a laser
235 system. The excimer machining laser used for this application was a 6 mJ, 200 Hz, 248 nm (Lasershot, Optec
236 Inc., San Diego, CA USA and Frameries, Belgium) KrF laser machining system. A grid was defined using the
237 method of quadratic residues described above and the well depth was calculated for each increment along
238 the X and Y directions. The parameters needed for determining well depth are the speed of sound in glass,
239 $c = 4550$ m/s, and the operating frequency of the transducer, $\omega_r = 2\pi f$ where f in this case is 7 MHz, corre-
240 sponding to the fundamental frequency for 500 μm thick 128YX lithium niobate^[34]. The well depth ranged up
241 to 309 μm and required between three and ten laser machining passes. The machined diffuser was bonded to
242 the transducer face using a UV-curable epoxy (NOA81, Norland Products Inc, Cranbury, NJ USA) using pre-
243 viously described techniques^[42]. This fabrication technique enables the miniaturization of devices that are
244 capable of producing diffuse acoustic fields irrespective of the nature of the enclosed volume, as determined
245 using surface and domain measurements as follows. A scanning laser Doppler vibrometer (UHF-120SV, Poly-
246 tec, Waldbronn, Germany) was used to characterize the displacement of the substrate when actuated. Mea-
247 surements of the pressure output from the transducer were performed using a fiber optic hydrophone (FOHS92,
248 Precision Acoustics, Dorchester, UK).

249 Ultrasound field simulation

250 Finite element analysis (COMSOL 5.5, Comsol Inc., Los Angeles, CA USA) was used to simulate the system as
251 a linear media with a time-dependent acoustic pressure field present in two dimensions. The boundaries be-
252 tween the coupling fluid and the cover slip, and the cover slip and the media above it were defined as acoustic-
253 structure boundaries, where there is fluid load on the structure due to pressure waves originating from the
254 ultrasound source and structural acceleration on the fluid domain across the fluid-solid boundary. This re-
255 sults in stress build-up in the cover slip that is translated to the fluid domain above it for the duration of the

256 stimulation. The coverslip was defined to have the elastic properties of silica glass, with an isotropic Young's
257 modulus of 73.1 GPa, a density of 2203 kg/m³ and a Poisson's ratio of 0.17. The distance between the trans-
258 ducer and the upper boundary is 5 mm. The simulations were conducted in the time domain, with a 20 μs
259 burst followed by an 10 μs dwell to observe changes in the pressure field both during and after the stimu-
260 lus. The acoustic field was modeled as a sinusoidal input to the transducer, $d = 5 \sin(\omega t)$ nm, and the fluid
261 domain was defined to have the properties of water ($\rho = 1000$ kg/m³, $c = 1500$ m/s). The maximum mesh
262 size was chosen so that all element sizes were always less than one-eighth of a wavelength, and the data were
263 exported every 0.05 μs so that a frequency range of up to 10 MHz could be analyzed. The cover slip was de-
264 fined to be 500 μm thick and spanned the entire width of the domain. The spatial step chosen for plotting
265 the isofrequency contours was less than $k_{\max} = \omega c^{-1}$.

266 **Imaging rig for ultrasound stimulation**

267 An upright epi-fluorescent microscope (Imager M2, Carl Zeiss GmbH, Gottingen, Germany) was used for
268 the *in vitro* experiments. For this application we used our transducer assembly placed in a heated stage fix-
269 ture set to 37°C underneath the cell chamber, which ensures homeostatis. Stimulus frequency and duration
270 was controlled by a waveform generator (33600A Series, Keysight, CA USA), and the pressure was controlled
271 through a 300-W amplifier (VTC2057574, Vox Technologies, Richardson, TX USA). Simultaneous calcium
272 imaging was performed using a 40x water dipping objective at 16.6 frames per second with a camera (Orca
273 Flash 4.0, Hamamatsu Photonics K.K., Japan) and a GFP filter.

274 **HEK293 cell culture and transfection**

275 HEK293 cells (ATCC CRL-1573) were cultured using a standard procedure in DMEM supplemented with 10%
276 fetal bovine serum (FBS) and 20 mM glutamine in a 37°C and 5% CO₂ incubator. Cells beyond passage thirty
277 were discarded and a new aliquot was thawed. A stable calcium reporter line was generated with a GCaMP6f
278 lentivirus (Cellomics Technology PLV-10181-50) followed by fluorescence-activated cell sorting (FACS). For
279 diffuser experiments, GCaMP6f-expressing HEK cells were seeded on a twelve-well cell culture plates with
280 18-mm glass coverslips coated with poly-D-lysine (PDL) (10 μg/μl; P6407, Sigma-Aldrich, St. Louis, Missouri,
281 USA) for 1-2h. Coverslips were washed with (Milli-Q) ultrapure water and cells were seeded at a density of
282 250000 cells/well. Cells were transfected with lipofectamine LTX Reagent (15338100, ThermoFisher Scien-
283 tific, Massachusetts, USA) according to the manufacturer's protocol and 24 hours after plating, using 500 ng
284 DNA of the clone of interest for each well. Cells were kept at 37°C for an additional 24 h before imaging on
285 our ultrasound stimulation setup. For imaging, cover slips were mounted on a specialized chamber featur-
286 ing an ultrasound transducer approximately 5 mm below the cover slip and a 10 mL reservoir of media above
287 the cover slip. Once cells were in focus, an ultrasound pulse of 100 ms duration was delivered as described
288 in previous sections while imaging with a 40X immersion objective, and a cell membrane profile was recon-
289 structed and analyzed from these images (ImageJ, National Institutes of Health, Bethesda, Maryland, USA).

290 **Rat primary neuronal culture**

291 Rat primary neuronal cultures were prepared from rat pup tissue at embryonic day (E) 18 containing com-
292 bined cortex, hippocampus and ventricular zones. The tissue was obtained from BrainBits (catalog #: SDE-
293 HCV) in Hibernate-E media and used the same day for dissociation following their protocol. Briefly, tissue
294 was incubated in a solution of papain (BrainBits PAP) at 2 mg/mL for 30 min at 37°C and dissociated in Hi-
295 bernate-E for one minute using one sterile 9" silanized Pasteur pipette with a fire-polished tip. The cell dis-
296 persion solution was centrifuged at 1100 rpm for 1 min, and the pellet was resuspended with 1 mL NbActiv1
297 (BrainBits NbActiv1 500 mL). The cell concentration was determined using a haemocytometer and neurons

298 were plated in 12-well culture plates with 18-mm PDL-coated cover slips (Neuvitro Corporation GG-18-PDL)
299 at a concentration of 900,000 cells/well. Neurons were then incubated at 37°C, 5% CO₂, performing half me-
300 dia changes every 3-4 days with fresh NbActiv1 supplemented with Primocin™ (InvivoGen ant-pm-1). For
301 calcium imaging experiments, cells were infected with AAV9-hSyn-GCaMP6f (Addgene #100837-AAV9) at day
302 3 *in vitro* (DIV3) and a half media change was performed the next day. Neurons infected with GCaMP6f as
303 stated above were infected with AAV9-hSyn-Cre (Addgene #105553-AAV9) and AAV9-hSyn-TRPA1-myc-DIO
304 (Salk GT3 core) at DIV4 and a half media change was performed the next day. Cultures were incubated at
305 37°C, 5% CO₂ until DIV10-12 and then imaged using the same equipment as for HEK cell experiments.

306 Calcium imaging

307 Calcium imaging analysis was performed using custom scripts written as ImageJ macros. Transfected cells
308 were segmented and cell fluorescence over time in the GCaMP6f channel was measured and stored in comma-
309 delimited text (csv) files. Calcium data was analyzed using custom Python scripts. Calcium signals were nor-
310 malized as $\Delta F/F$ using a 6 s baseline for each region of interest (ROI) and a peak detection algorithm with a
311 fixed threshold of 0.25 was used to identify responsive cells after ultrasound stimulation.

312 Ex-vivo hydrophone measurements

313 Hydrophone measurements were performed with a fiber-optic hydrophone (FOHS92, Precision Acoustics,
314 Dorchester, UK) *ex vivo*. C57BL/6 mice (JAX 000664), aged 10-14 weeks, were sacrificed and decapitated.
315 The skin over the skull was removed, followed by removal of the lower mandible, soft palate and hard palate.
316 Once the ventral part of the brain was exposed, the mouse head preparation was placed dorsal side down
317 on the diffuser assembly coupled with ultrasound gel, and the hydrophone tip was lowered into the ven-
318 tral portion of the brain using a micromanipulator (Fig. S3). Animals used in this trial were group housed in
319 an American Association for the Accreditation of Laboratory Animal Care approved vivarium on a 12 hour
320 light/dark cycle, and all protocols were approved by the Institutional Animal Care and Use Committee of the
321 Salk Institute for Biological Studies. Food and water were provided ad libitum, and nesting material was pro-
322 vided as enrichment. Mice were euthanized using CO₂ according to approved protocols before decapitation
323 and dissection.

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336 Competing Interests

337 All authors declare no competing interests.

338 Author contributions

339 J.F. and A.V. designed the research. A.V. built, characterized the devices and conducted experiments. A.V., F.A.,
340 and N.B. analyzed data. M.D. and S.H.C. performed in-vitro studies. A.V. and U.M. performed the ex-vivo
341 study. A.V. and J.F. wrote the paper.

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