Arousal elevation drives the development of oscillatory vocal output.

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Abstract

Adult behaviors, such as vocal production, often exhibit temporal regularity. In contrast, their immature forms are more irregular. We ask whether the coupling of motor behaviors with arousal changes give rise to temporal regularity and drive the transition from variable to regular motor output over the course of development. We used marmoset monkey vocal production to explore this putative influence of arousal on the nonlinear changes in their developing vocal output patterns. Based on a detailed analysis of vocal and arousal dynamics in marmosets, we put forth a general model incorporating arousal and auditory-feedback loops for spontaneous vocal production. Using this model, we show that a stable oscillation can emerge as the baseline arousal increases, predicting the transition from stochastic to periodic oscillations occurring in marmoset vocal development. We further provide a solution for how this model can explain vocal development as the joint consequence of energetic growth and social feedback. Together, our model offers a plausible mechanism for the development of arousal-mediated adaptive behavior.

Keywords

Vocal development, Arousal, Feedback loop, Primate

Introduction

The transformation from variable exploratory motor outputs into structured ones is ubiquitous across different behaviors and species during development. For many altricial animals, vocal production is crucial from the first day of life as it is the only means to solicit attention from caregivers [1, 2]. Early vocalizations are driven by internal needs and act as scaffolding for later, more adaptive communicative functions [1]. In the infants of humans [3, 4], songbirds [5, 6], bats [7], and marmoset monkeys [8-10], vocalizations transition through different states on their way to becoming adult-like. While it is natural to presume that a behavioral state change must be the result of a neural change, the fact of the matter is that the development of vocal behavior involves multiple components, including the biomechanics of the vocal apparatus, the metabolic energy, neural systems, and the social environment [11]. These moving parts interact, form multiple feedback loops, and self-organize into identifiable structures and functions.

Vocalizations, not only during infancy but also through adulthood, are used to coordinate activities among individuals that are crucial for survival. They are tightly correlated with internal states such as arousal [12-16]. Physiological arousal is associated with global changes in brain and body states that modulate sensory processing and allocate metabolic energy to prepare the body for actions [17]. Arousal levels are correlated with a number of acoustic features in marmoset vocalizations, with higher arousal associated with longer, louder, and less noisy sounds [18]. Arousal is also related to the sequential organization of vocalizations on a longer timescale [8, 12, 18]. Spontaneous vocal sequences of infant and adult marmosets exhibit a high coherence with heart rate fluctuations around 0.1 Hz. Male songbirds in a high arousal state induced by the presence of females produce more stereotyped songs than in the undirected context where the songs are more variable [19]. How arousal interacts with the vocal production system to affect the regularity of vocal output is unknown.

In this study, we link these two phenomena of 1) the transitions from stochastic to structured behaviors during development and 2) the arousal correlated vocal sequence organization to understand the pattern of shifts in marmoset vocal development. A useful approach for investigating transitions between behavioral states occurring across development is the application of dynamical systems theory [20]. Stable motor dynamics at timescales of seconds can be understood as consequences of control parameters that change over days and months [21]. The control parameters are components that drive the vocal output into new modes as they cross critical points to enter new states [11, 22-24]. Based on behavioral and physiological data, we describe here a parsimonious model incorporating the vocal production system with arousal dynamics. Using this model, we predicted nonlinear changes in vocal sequences during the development of marmosets and tested it with empirical data. In the end, we discuss the model's insights and offer interpretations of the control parameters.

Results

Vocal output and arousal fluctuation are dynamically coupled

Infant marmosets start vocalizing from the first postnatal day. In a context in which the infants were briefly isolated, they vocalized spontaneously and continuously (data published previously [10]). The initial vocalizations were babbling-like, composed of a variety of acoustic patterns, including both mature and immature-sounding calls (e.g., postnatal day 3 (P3) vocalizations; Fig. 1A). Over the first two months, vocalizations became sparser and more regular, dominated by a tonal contact call, which is the call used for long-distance communication in adults (e.g., P60 vocalizations; Fig. 1A). To quantify the temporal structure of vocal output, we calculated the sound amplitude envelope (Fig. 1B). Previous studies identified high

coherence around 0.1 Hz between spontaneous vocal output and momentary arousal levels (using heart rate as a proxy; see Fig. 1B) in both infants and adults [8, 18], suggesting a robust relationship between these two signals regardless of age. This timescale is much slower than the respiratory rhythm (which is at ~1 Hz), and thus the structural changes occurring during development at low frequencies were attributed to arousal but not respiratory development. We next investigated the temporal coordination between vocalization and arousal in detail.

Cross-correlation shows that these two signals were significantly correlated in infants with a time difference of around 5 s (n=7 subjects, shuffle test; Fig. 1C). When aligning the signals by the onsets of contact calls (Fig. 1D), we observed a general trend that heart rate was significantly high (>50 percentile of session heart rate) before call onsets and dropped significantly upon call production (Fig. 1E; n=2972 contact calls; shuffle test). Considering all types of vocalizations, we found that louder sounds led to a greater reduction of heart rate (Fig. 1F; n=6763 calls, p=1.98e-12, multiple linear regression (MLR)). These findings have three implications: 1) arousal is discharged upon vocalization (by Fig. 1F); 2) arousal automatically builds up (by Fig. 1E); 3) arousal levels gate vocal production (by Fig. 1E). Based on these observations, we offer a model to describe these processes.

Figure 1. Temporal coordination between vocal output and arousal fluctuations. (A) Two examples of infant vocalizations on two postnatal days. The top two panels are the spectrogram and power of sound amplitude envelope on P3, and the bottom two are on P60. (B) Two examples of infant vocalization amplitude envelope with real-time heart rate. (C) Cross-correlation between sound and heart rate. The highlighted segment is outside the 95% confidence interval (CI) of shuffled data. Shaded area is 95% CI of data. (D) Two examples of heart rate change around the onset of phee call. (E) Heart rate change around phee onset at the population level. Heart rate data were converted to percentiles of the session. Highlighted segments in magenta are outside 95% CI of shuffled data. (F) Linear fitting of largest heart rate drop from call onset vs. acoustic power of the call standardized per session.

A dynamic model for spontaneous vocal production

Infants often produce exploratory movements that increase in structure over time [25]. What refines those movements into structured behaviors is often reinforcement signals in the form of sensory feedback [26-29].

This is true for the specific case of vocal development, as well [30, 31]. The positive sensorimotor feedback loop is crucial for sustained motor activities [32]. In our model, we assumed that the loudness of self-generated vocalization is reinforced and provides delayed feedback to drive the vocal system (Fig. 2A). On the other hand, positive feedback loops are intrinsically unstable. An antagonistic process is that once the action (vocal production) is fulfilled, arousal is discharged – similar to the idea of Lorenz's psycho-hydraulic model [33](Fig. 2A). Thus, three dynamical processes constitute our model: 1) the neural activity generating vocalizations; 2) the sound signal providing positive auditory feedback; and 3) the arousal level modulating the gain of the drive and that is reduced by vocal production.

The neural model describes a population level firing rate with two inputs: the drive signal I and noise (Eq. 1). The drive I consists of the delayed auditory feedback $s(t - \delta_s)$ and a constant bias h multiplied by the gain of a delayed arousal level $a(t - \delta_a)$ (Eq. 2). This I is further transformed by an s-shaped activation function f when added to the neural model. On top of the deterministic dynamics, there is a stochastic component simulating the intrinsic fluctuation of the neural population. This includes Gaussian noise with an activity-dependent amplitude. The sound amplitude is modeled as low-pass filtered muscle activation. Muscle activation is proportional to the neural activity with a time delay δ_v [34]. The arousal level is depleted by muscle activation and slowly replenished to a baseline level (a_B). The rate of arousal reduction is proportional to muscle activation, compatible with the observation that louder sounds reduce arousal more than quieter ones (Fig. 1F). We set the arousal feedback delay δ_a the time difference between heart rate and vocalization, which is 4-5 s (Fig. 1D). The interplay among neural activity, vocal output, and arousal levels together determine the vocal pattern. The complete schematic of this system is shown in Fig. 2A, with equations listed

$$\begin{cases} \tau_v \dfrac{dv}{dt} = -v + f(I) + \sigma\sqrt{v}\,\eta \\[2mm] \tau_s \dfrac{ds}{dt} = -s + Av(t - \delta_v) \\[2mm] \tau_a \dfrac{da}{dt} = [a_B - a] - Bv(t - \delta_v)a \end{cases} ,$$

$$(1)$$

where

$$I = a(t - \delta_a)[h + s(t - \delta_s)],$$

$$(2)$$

f is an s-shaped function taking the form $f(x) = \dfrac{x^2}{1+x^2}$, v, s, a are the time-dependent variables of neural activity, sound, and arousal level respectively, τ_v, τ_s, τ_a are the corresponding time constants, A and B are positive parameters, $\eta(t)$ is Gaussian noise (zero mean unit variance) scaled by an activity-dependent amplitude $\sigma\sqrt{v}$. Parameters are given in Table 1.

We show that this model can generate sound amplitude signals that qualitatively resemble the amplitude envelope of marmoset vocalizations. By varying the parameters, we simulated the temporal patterns corresponding to early (Fig. 2B) and later (Fig. 2C) developmental stages, comparable to the P3 and P60 examples in Fig. 1A.

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Figure 2. Model setup and simulation. (A) Diagram of model setup. The drive integrates arousal level with excitatory input from
auditory feedback. The motor area receives input from the drive and generates the action via the effector (muscles and
biomechanics). Activation of the effector inhibits arousal, which can be restored to a baseline level. The vocal signal is sent back
to the drive via the auditory system. $\delta_{Arousal}$, $\delta_{Auditory}$, and δ_{Motor} are the latencies of the signals. (B) Simulated data for each
compartment of the model. From top to bottom are simulations for $v(t)$, $a(t)$, and $s(t)$. This is a simulated example of early
vocal output. (C) A simulation of later vocal output.

Table 1. Parameters

A	B	τ_v	τ_s	τ_a	δ_v	δ_s	δ_a	a_B	σ	h
1	30	0.1	0.1	1	0.02	0.01	4.2	0.5-5	1	0.1

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152 **Changes in a parameter correlate with shifts in vocal dynamics**
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To see how parameters change during development, we fitted s and a in our model to the real vocalization
and heart rate data. We used an unscented Kalman filter for simultaneous state and parameter estimation of
$v(t)$, $s(t)$, $a(t)$, and a_B while fixing other parameters [35](Fig. S1). Because the state variables s and a are
unit-less, we normalized the sound amplitude and scaled the heart rate for each session with an optimized
scaling factor such that the tracking error of vocalization and heart rate was minimized (Methods). The data
fitting shows that the estimated a_B increased significantly with postnatal days (p=1.6e-8, MLR; Fig. 3A).
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Although we scaled the data for each session, it is possible that the scaling factor was constant, and
changes in a_B simply reflect an elevation of baseline heart rate during development. To test this, we estimated
the allometric scaling of a_B with body mass and compared it to the scaling exponent of heart rate. The
allometric scaling is the power-law relationship between a biological parameter y with body mass m: $y = am^\beta$. Infants tripled their body masses during the first two months [10]. If a_B is proportional to mean heart
αm^β. Infants tripled their body masses during the first two months [10]. If a_B is proportional to mean heart
rate, we should expect identical scaling exponents. The scaling exponent for heart rate was $\beta_{HR} = 0.44 \pm$
0.02 ($R^2 = 0.78$, p=1.6e-31), whereas the scaling exponent for a_B was $\beta_{aB} = 0.77 \pm 0.11$ ($R^2 = 0.37$,
p=3.1e-10), significantly greater than β_{HR} (p=3.5e-3; Fig. 3B). Thus, a_B cannot be simply interpreted by
baseline heart rate. We will discuss two potential factors that may contribute to the increase of a_B in the last
section.

Figure 3. Parameter estimation for a_B. (A) Fitted a_B using unscented Kalman filter as a function of postnatal days (PND). Each dot is a session. Red line is cubic spline fitting of the estimated a_B. Shaded area is 95% CI of the curve fitting. (B) Linear fitting in the logarithmic scale to estimate allometric scaling of a_B and heart rate. β_{aB} and β_{HR} are slopes. a_B is scaled to the same starting point as heart rate. Shaded areas are 95% CI.

The model predicts a Hopf bifurcation in vocal development

One prediction of this model is that increasing a_B can drive the system across the boundary of a Hopf bifurcation (see Supporting information). The deterministic behavior of the model is shown in the phasegram of Fig. 4A, which is constructed by intersection points between phase portraits under different a_B and a Poincare section [36]. The system is attracted to stable fixed points when a_B is low and to limit cycles as a_B crosses a boundary. It can be proved that the oscillation period of the limit cycles is approximately twice the longest delay along the feedback loops (Supporting information). Thus, for time delay ~5 s, the oscillation period is ~10 s, consistent with the 0.1 Hz oscillation of marmoset vocal output. When adding the amplitude-dependent Gaussian noise to the model, the system is mainly driven by noise at low a_B (as in Fig. 2B). As a_B is increased above the boundary, the system is dominated by a periodic oscillation (as in Fig. 2C).

Consistent with our model prediction, the phasegram of the vocalization data as a function of postnatal days shows the emergence of periodic oscillations after about P30 on average (Fig. 4B). Similar transitions occurred in individual subjects (Fig. S2). When including the stochastic component, the signal-to-noise ratio (SNR) increases with a_B (Fig. 4C). Here we treated the power spectral density (PSD) near 0.1 Hz as signal and PSD at other frequencies as noise. In agreement with this prediction, the PSD of the data became more concentrated at 0.1 Hz over time (Fig. 4D) and the SNR increased with postnatal days (p=1.5e-8, MLR; Fig. 4E). Thus, the decreases in the stochasticity and the increases in the regularity of infant vocalization during development can be explained by the crossing of Hopf bifurcation in the feedback loop system.

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200 Figure 4. Hopf bifurcation of the temporal vocal structure during development. (A) Phasegram of the deterministic model as a
201 function of a_B. (B) Phasegram of data as a function of PND. Color represents probability density normalized per day. Insets are
202 the phase portrait of two examples before and after the bifurcation. (C) Simulated SNR of the stochastic model as a function of
203 a_B. Error bars are standard deviations. (D) PSD of vocal amplitude envelope averaged by postnatal weeks. The peak around 0.1
204 Hz became sharper over time. (E) SNR of data as a function of PND. The curve is fitted by a cubic spline. Shaded area is 95% CI.

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206 **Interpretations of the control parameter**
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Arousal is essential to allocate energy to relevant behaviors. Contexts inducing higher arousal should have
more energy allocated and vice versa. Thus, one natural interpretation for a_B is the baseline of context-
dependent energy allocation. Physical development is accompanied by increases in general metabolic energy,
and thus more energy can be allocated for vocalization. During marmoset development, not only did heart
rates become higher, but the volume of the heart also expanded. By Fick's equation, the rate of oxygen
consumption is $\dot{V}_{O_2} = V_h f_h (C_a O_2 - C_v O_2)$, where V_h is the stroke volume of the heart, f_h is the heart rate,
and $C_a O_2 - C_v O_2$ is the oxygen extraction [37, 38]. Thus, if oxygen extraction in this context is constant
across development, then $\dot{V}_{O_2} \propto V_h f_v \propto m^{\beta_v + \beta_{HR}}$, $\beta_v > 0$. In other words, one should expect the growth of
metabolic rate to be faster than heart rate (Fig. 3B). This is consistent with the allometric scaling of a_B. The
ability to produce regular vocal output thus can be a consequence of general metabolic development.
218
Previous studies also show that vocal development can be accelerated by increased contingent
parental vocal feedback, independent of body growth [10, 39]. In a study of an opposite extreme (lack of
parental input), hand-raised marmoset triplets continued producing immature calls and vocal sequences even
at one year [40, 41], suggesting that body growth is not a sufficient condition for vocal development. If we
decompose the arousal state $a(t)$ into an energetic level $E(t)$ and a socially tuned factor γ that changes much
more slowly: $a(t) = \gamma E(t)$. We can rewrite the last equation of Eq. (1) into

$$225 \qquad \tau_a \frac{dE}{dt} = [E_B - E] - Bv(t - \delta_v)E,$$

226 $\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad$ (3)

227 where $E_B = a_B/\gamma$. E_B thus can be understood as the baseline metabolic energy allocated for vocal
228 production. The drive I can also be rewritten as

$$230 \qquad I = E(t - \delta_a)[h' + \gamma s(t - \delta_s)],$$

$\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad$ (4)

where $h' = \gamma h$. The socially tuned factor γ is hence simply the strength of the auditory feedback to the vocal
production system. In other words, it is the reinforcement strength of the self-produced vocalizations.

An intriguing corollary of this formulation is that within the same range of E_B increment, a higher γ,
or a steeper increase of γ can advance the timing of bifurcation crossing (Fig. 5 red lines). This simulates the
acceleration of vocal maturation due to more contingent parental feedback [10, 39]. Compatible also with
empirical findings in hand-raised marmosets, our model shows that, even with a high E_B, if the weight of
auditory feedback is too weak, the system operates below the bifurcation and thus would still produce
immature vocal behavior (Fig. S3). Therefore, the development of energetics may provide the necessary
condition for the animal to produce adult-like vocal behavior, while social feedback provides the correct
tuning of it. Both factors may act as control parameters for the development of vocal dynamics.

244
245 Figure 5. Bifurcation diagram of the deterministic model in the space of E_B and γ. Blue region is the stable fixed point solution,
246 and yellow region is the limit cycle solution. The red line with a steeper slope represents the developmental trajectory with
247 more contingent parental feedback. The arrow shows the earlier crossing of the bifurcation boundary with more contingent
248 feedback.

249
Discussion

In this study, we propose a parsimonious model that integrates sensorimotor processes with the dynamics of
internal states based on simultaneous measurements of vocalizations and arousal in infant marmosets. We
demonstrated that increases in energy allocation--the baseline arousal level--during development could cause a
Hopf bifurcation in vocal sequence production, leading to a stable oscillatory pattern at 0.1 Hz. We further
provided insights into how increases in metabolic rate and social feedback can contribute to increasing the
value of this parameter. Together, these investigations show that motor development can be affected by the
development of arousal, which is a quantity for socially guided energy allocation.

Organisms continuously regulate physiological parameters to maintain them within tenable ranges [42]. This is key to survival from the very beginning of life and universal for all species. Arousal is a grand parameter that coordinates physiology at the systems level to accommodate behavioral states to internal demands [43]. Arousal fluctuations correlate with dynamical changes of brain states and behaviors [44-46]. Spontaneously produced motor behaviors are tightly coupled with internal physiological fluctuations, and they are actively coordinated via functionally and structurally interconnected large-scale brain networks [47]. For dedicated brain functions, arousal modulates the gain of visual [48-50] and auditory systems [51], enabling adaptive behaviors [18]. Arousal-modulated sensory processing has been suggested to be regulated via cholinergic input from the ascending arousal system [50, 52]. In the song system of birds, cholinergic modulation of premotor nucleus HVC increases firing rate prior to syllable onsets and enhances song stereotypy, resembling the effect of females' presence [53]. These changes are consistent with our model's increases in baseline arousal level. Another system implementing feedback loops for motor vigor is the dopaminergic signaling through the basal ganglia, thalamus, and back to the neocortex, which contributes to reward-based motor learning [54, 55]. In birds, the dopamine-secreting neurons in the periaqueductal grey (PAG), which are believed to encode information about social context and arousal, enable vocal learning by releasing dopamine in HVC [56]. This mechanism can facilitate the selection of socially relevant behaviors during development. The neuromodulation for primate and other forms of mammalian vocal development is still unclear and requires future investigation.

Vocal communication is a flexible motor behavior that can be adapted to various contexts through modifying acoustic properties. Our model suggests that drastically changed vocal behavior can be effectively tuned by a slowly varying parameter [57]. This offers an efficient solution for adaptive vocal control. In marmoset monkeys, physical distance consistently induces different vocalizations with shorter, softer, and noisier calls produced in shorter distances [18, 58]. Adult marmosets that are alone (or in an "infinitely far" context) induces the highest arousal levels compared to other distances between conspecifics [18]. Consistent with our model, animals in this high arousal condition also produce more regular vocalizations [18]. Infant marmosets in the high arousal condition are more likely to reach the limit of their motor and energetic capacity due to a generally high-level motor output. This might explain why we observed a drastic vocal shift in this context during the first two months when the infants were rapidly growing. In comparison, the association between social context and arousal can be a learned trait through experience [10, 39-41]. Without adequate social feedback, the reinforcement of self-produced vocalizations may be downregulated, making the vocal output less vigor. Social context and arousal are, maybe in this way, naturally linked via the strength of sensory feedback.

How the system "decides" on the acoustic output could be via the medial brain pathway regulating arousal and motor aspects of vocal production [59-61]. Along the medial wall, several subcortical and cortical areas controlling cardiac activity are related to socially relevant vocal perception and production [47, 62]. The anterior cingulate cortex (ACC), for example, is a robust node along the primate vocal pathway [61, 63]. It reacts strongly to conspecific calls [62] and regulates cardiovascular activity [64-66]. In the simple scenario depicted in our model, the drive is a conjunctive consequence of sensory input and arousal level. Social context may affect vocal output by shifting the arousal baseline. The integration of social context, arousal, and sensory feedback may occur within ACC. The drive signal can be passed to the downstream motor areas such as PAG, which has been found to initiate vocal production in rodents and primates [67-69]. The PAG dynamics drive the brainstem central pattern generators (CPG) that oscillate on a faster respiratory timescale [57]. The model in this study focuses on the slower dynamics that generate call bouts instead of individual utterances whose boundaries are set by respiration. The acoustic structure is further unpacked via interactions between the CPGs and the vocal apparatus, which changes during development and affects the vocal outcome as well [22].

Energy is often an ignored factor when it comes to motor control. Energy impacts motor output at various timescales. Sound production, for instance, is a prerequisite for self-preservation in some species; its importance is reflected in the generally high energetic cost attained by a variety of physiological and

biomechanical mechanisms [70]. In some arthropods and frogs, the cost of calling can be up to 20 times the resting metabolic rate [71]. At short timescales, this dependence on energetics affects acoustic features such as power, frequency, duration, and rate of call emission [72]. At longer timescales, the internal goals and needs of the animal change, and energy allocated for sound production varies according to the time of the day [73], nutrition [74], seasons [75], state of learning [76], and age [77]. In our study here, we offered an energetic interpretation for our model to raise the attention to this critical factor. There are several mechanisms by which metabolic energy affects neural activity. At the cellular and tissue levels, energy availability can influence neural excitability and information processing [78, 79]. At the whole-body level, diet, exercise, and blood sugar levels affect neural signaling [80]. The effect of metabolism on behavioral phenotypes has been linked to biochemical pathways operating in the brain. For example, in insects, inhibition of oxidative phosphorylation (a process producing maximal ATP) causes increased aggression [81]. Moreover, the effect of oxidative phosphorylation inhibition is modulated by the social environment [81], suggesting the context-dependent energy demand for producing motor behavior. Future studies bridging energetics and behavior at the levels of genetics, biochemistry, neurons, and physiology will help fill in the details of our theory.

Material and methods

Subjects

The subjects used in the original study were 17 infants up to two months of age, captive common marmosets (*Callithrix jacchus*) housed at Princeton University. Vocalization recordings from ten infants have been published previously [10]. Another seven infants were used for heart rate measurements throughout their development. All experiments were performed in compliance with the guidelines of the Princeton University Institutional Animal Care and Use Committee.

Vocalization recordings

Ten infant marmosets participated in this experiment. On each day starting from the first postnatal day, an infant was temporarily separated from its parents and taken to the experiment room, where it was placed in a test box. Undirected vocalizations were recorded for 5 min per day. The types of vocalizations that were observed were identical to those produced when the infant was naturally separated from its parents. A speaker was placed in the room broadcasting pink noise (amplitude decaying inversely proportional to frequency) of 45 dB (at 0.88m from the speaker to the microphone) in order to mask occasional noises produced outside of the testing room. The sound pressure level (SPL) was measured at the microphone location and used as a reference to calibrate the SPL of marmoset vocalizations. The digital recorder (ZOOM H4n Handy Recorder) was placed directly in front of the transfer cage at a distance of 0.76m. Audio signals were acquired at a sampling frequency of 96kHz.

Electrocardiogram (ECG) recording and heart rate analysis

A different group of infant marmosets (n=7) participated in this experiment. Infant marmosets were placed in a test box and recorded as they were spontaneously vocalizing for 10 mins. To record the electrocardiogram (ECG) signal, we used Ag-AgCl surface electrodes (Grass Technology) placed on the chest, close to the heart. The electrodes were sewn onto an elastic band that was secured around the marmoset's thorax. The ECG signals were acquired by the OmniPlex data acquisition system (Plexon) at a sample rate of 40 kHz. The ECG signals were downsampled and bandpass-filtered between 5 and 39 Hz. Heartbeats were detected using template matching by cross-correlation within a 50 ms sliding window. The heartbeats were binarized and convolved with a Gaussian window with $\sigma = 0.25$ s. The mean heart rate of the session was estimated as the mean of this convolved time series.

Acoustic feature processing

To calculate the acoustic power of marmoset vocalizations, we removed non-vocal signals by identifying the boundaries of individual utterances. Vocal utterances were defined as continuous sound separated by silence. The boundaries of utterances were determined by threshold crossings of the sound amplitude envelope and checked manually. Sound pressure was calculated as the root-mean-square (RMS) envelope of the sound signal and calibrated with the 45-dB pink noise background. It was then converted to power in milliwatts using a standard equation that accounts for measurement distance (0.76m) and assumes hemispherical

spreading $P = IA = \frac{4\pi r^2 p^2}{\rho v}$, where p the RMS sound pressure in Pa, ρ=1.2 kg m^{-3} the air density, and v=344 m s^{-1} the speed of sound [82].

Spectral analysis

The acoustic power of the sound amplitude envelope was downsampled to 1 Hz. The PSD of was calculated using MATLAB implementation of the multi-taper method *pmtm*. Mean PSDs were calculated for every postnatal week up to the 8th week across all ten marmosets. The SNR was estimated as the ratio between the sum of PSD within the range of 0.05 - 0.2 Hz and the sum of the rest within 0.02 – 0.5 Hz [83].

Heart rate analysis

To analyze the heart rate dynamics around vocal production, we first converted detrended momentary heart rates into percentiles. This was done by binning 100 percentiles of the session heart rate and assigning the heart rate of each time point into a percentile bin.

To correlate heart rate drop with acoustic power, we calculated the heart rate change from call onset to the lowest heart rate during call production. The acoustic power was standardized by session, and the mean power of a call was calculated for the correlation.

Model setup

We modeled the response function of the arousal gained auditory feedback to the neural population as a sigmoidal curve. We added noise $\sqrt{v(t)}\eta(t)$ to the deterministic system, which is a delta-correlated Gaussian white noise with zero mean and an activity-dependent amplitude. The square-root scaling is due to the central limit theorem, as we assume that the activity of each neuron in this population is an independent stochastic process [84]. We only introduced noise in the modeling of neural activity for simplicity; in principle, there is also noise in arousal and vocalization.

The stochastic differential equation with multiplicative noise was solved in two steps. The stochastic part was integrated first to generate a sample of neural activity v^* with guaranteed positivity, using a trick developed in ref. [85]. In the next step, v^* was used as a starting point for the integration of the deterministic part. The time step was set at dt=0.01.

Parameter estimation

To estimate a_B of the model, we performed an unscented Kalman filter for noisy data to track observations of vocalization and heart rate [35]. We also included covariance inflation due to the noisiness of the data [86]. The remaining parameters were fixed according to Table 1. For each session, we normalized the sound amplitude by the maximum. We detrended the heart rate signal, band-passed it between 0.02 and 1 Hz, and converted it to percentiles. We then scaled the signal from 1 to 10 and passed each scaled version through the Kalman filter tracking. To find the best scaling factor, we calculated the tracking error of sound and heart rate as the sum of the L2 norm of these two errors and selected the scaling factor yielding the lowest error. We initialized the covariance matrix for real-time tracking of state variables using the sample covariance matrix of sound and heart rate signals. The process noise for parameter tracking was set to 0.015 and was reset each

415 iteration. The time-series data were repeated for each session such that enough iterations were given for
416 stabilizing the tracking. The estimated a_B was calculated as the mean of the second half of the tracking.
417
418 <u>Bifurcation analysis of data</u>
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To construct the orbit diagram, we first calculate the envelope of the acoustic power using the RMS with a 1-
s sliding window. In this way, we smooth out the rhythms at the respiratory timescale. We then Hilber-
transformed the signal and estimated the time-dependent phase angle $\phi(t)$. This step mitigated the
observational variability of acoustic power. We finally took $y(t) = \sin(\phi(t))$ to construct a phase portrait
in the space of $y(t)$ and $y(t + \tau)$, where τ is a time delay. τ is typically one-tenth to one-half the oscillation
period [87], and in this analysis, we used $\tau = 1.5$ s. The Poincare section was chosen as the x-axis of the
phase space. We then calculated a Gaussian-filtered 2D histogram using MATLAB function *ndhist* (file
exchange #45325) in the space of postnatal days (PND) and $\sin(\phi)$. The histogram was normalized within
the same postnatal day.
429
430 <u>Statistical analysis</u>
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We used the MATLAB *fitlm* routine to fit linear models: $y = b_0 + b_1 x + b_2 ID + \epsilon$, where y is the response,
x is the independent variable, and ID is the subject, to find the scaling exponent of heart rate vs. body mass
and a_B vs. body mass. We assumed that each individual had its own slope and intercept and thus included the
identities of the subjects as a predictor in the model. A similar linear fitting was applied to test the
dependence of a_B and SNR on PND. For linear fitting of heart rate drop vs. acoustic power, we also
included PND as a predictor in the MLR model. The statistical significance was evaluated based on the p-
values reported in the MATLAB function output.
439
440 The test for significantly high or low heart rate was performed by randomizing the physiological
441 signals and calculating the mean 1000 times to construct the 95% CI of the null distribution. A similar
442 procedure was applied to find the 95% CI of the cross-correlation between sound amplitude and heart rate.
443 We adopted the Type I error $\alpha = 0.05$ for the criterion of statistical significance.
444
445

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447
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Supplementary material
665
666   <u>Derivation of bifurcations</u>
667
668   We analyze the stability of the system around fixed points using linearization. The fixed point solution of the
669   deterministic system is given by the following equations that can be solved numerically
670

671
$$\begin{cases} \bar{v} = f\left(\frac{a_B}{1+B\bar{v}}[h + A\bar{v}]\right) \\ \bar{s} = A\bar{v} \\ \bar{a} = \frac{a_B}{1+B\bar{v}} \end{cases} ,$$
(1)

672
673   where $(\bar{v}, \bar{s}, \bar{E})$ is the fixed point. There is a very narrow region within our parameter range that has three
674   fixed points, and the rest has a unique fixed point. The following analysis does not include the three-fixed-
675   point situation as it is relatively transient.
676   Oscillations occur where in parameter space this fixed point becomes unstable. To investigate the stability of
677   the fixed point, we linearize the system around it by writing
678
679
$$v = \bar{v} + \tilde{v}, s = \bar{s} + \tilde{s}, a = \bar{a} + \tilde{a}.$$
680
681   For simplicity, we neglect the tilde and still use v, s, a in the following but for perturbations around the fixed
682   point. The linearized system is
683

684
$$\begin{cases} \tau_v \frac{dv}{dt} = -v + f_s'(\bar{s})s(t - T_s) + f_a'(\bar{a})a(t - \delta_a) \\ \tau_s \frac{ds}{dt} = -s + Av(t - \delta_v) \\ \tau_a \frac{da}{dt} = -a - B[\bar{v}a + v(t - \delta_v)\bar{a}] \end{cases} ,$$
(2)

685
686   where $f_s' = \frac{\partial f}{\partial s}$ and $f_a' = \frac{\partial f}{\partial a}$.
687
688   Now we look for solutions in the form $v = v_0 e^{\lambda t}$, $s = s_0 e^{\lambda t}$ and $a = a_0 e^{\lambda t}$. Plugging into the linearized
689   equations gives
690

691
$$\begin{pmatrix} -1 - \tau_v\lambda & f_s'(\bar{s})e^{-\lambda\delta_s} & f_a'(\bar{a})e^{-\lambda\delta_a} \\ Ae^{-\lambda\delta_v} & -1 - \tau_s\lambda & 0 \\ -Be^{-\lambda T_v}\bar{a} & 0 & -1 - \tau_a\lambda - B\bar{v} \end{pmatrix} \begin{pmatrix} v_0 \\ s_0 \\ a_0 \end{pmatrix} = 0.$$
(3)

692
693   λ thus must be a root of the transcendental equation
694

695    $$\lambda^3 + A_1\lambda^2 + \left(B_1 + B_2 e^{-\lambda(\delta_v+\delta_s)} + B_3 e^{-\lambda(\delta_v+\delta_a)}\right)\lambda + C_1 + C_2 e^{-\lambda(\delta_v+\delta_s)} + C_3 e^{-\lambda(\delta_v+\delta_a)} = 0,$$

696    $$A_1 = \frac{\tau_v\tau_s + \tau_v\tau_a + (1+B\bar{v})\tau_v\tau_s}{\tau_v\tau_s\tau_a}, B_1 = \frac{(1+B\bar{v})(\tau_v + \tau_s) + \tau_a}{\tau_v\tau_s\tau_a}, B_2 = -\frac{Af_s'(\bar{s})}{\tau_v\tau_s}, B_3 = \frac{B\bar{a}f_a'(\bar{a})}{\tau_v\tau_a}, C_1 = \frac{1+B\bar{v}}{\tau_v\tau_s\tau_a}, C_2 =$$

697    $$-\frac{A(1+B\bar{v})f_s'(\bar{s})}{\tau_v\tau_s\tau_a}, C_3 = \frac{B\bar{a}f_a'(\bar{a})}{\tau_v\tau_s\tau_a}.$$
(4)

698
699   For the system to have limit cycle solutions, the fixed point should lose its stability, i.e., Re $\lambda > 0$, while there
700   is a nonzero solution for the imaginary part. We numerically identify the regions that satisfy these criteria by
701   letting $\lambda = \mu + i\omega$ and solving the real and the imaginary parts of this transcendental equation separately
702

703 $$\mu^3 - 3\mu\omega^2 + A_1\mu^2 - A_1\omega^2 + B_1\mu + B_2\mu e^{-\mu\delta_{vs}}\cos\omega T_{vs} + B_2\omega e^{-\mu\delta_{vs}}\sin\omega\delta_{vs} +$$
704 $$B_3\mu e^{-\mu\delta_{va}}\cos\omega\delta_{va} + B_3\omega e^{-\mu\delta_{va}}\sin\omega\delta_{va} + C_1 + C_2 e^{-\mu\delta_{vs}}\cos\omega\delta_{vs} + C_3 e^{-\mu\delta_{va}}\cos\omega\delta_{va} = 0,$$
705 $$(8)$$

707 $$3\mu^2\omega - \omega^3 + 2A_1\mu\omega + B_1\omega + B_2\omega e^{-\mu\delta_{vs}}\cos\omega\delta_{vs} - B_2\mu e^{-\mu\delta_{vs}}\sin\omega\delta_{vs} + B_3\omega e^{-\mu\delta_{va}}\cos\omega\delta_{va} -$$
708 $$B_3\mu e^{-\mu\delta_{va}}\sin\omega\delta_{va} - C_2 e^{-\mu\delta_{vs}}\sin\omega\delta_{vs} - C_3 e^{-\mu\delta_{va}}\sin\omega\delta_{va} = 0,$$
$$(9)$$

where $\delta_{vs} = \delta_v + \delta_s$ and $\delta_{va} = \delta_v + \delta_a$. Two regions are identified in the parameter space of a_B, one with
a positive real part and a nonzero imaginary part, and one with only negative real parts.

<u>Derivation of period</u>

We estimate the angular frequency ω_c near the bifurcation as $\mu \to 0^+$. To further simplify the problem, as the
delay is dominated by the energy delay δ_a and $\omega\delta_{vs} \ll 1$, equations (8) and (9) can be reduced to

$$A_1\omega_c^2 - B_3\omega_c\sin\omega_c\delta_{va} - (C_1 + C_2 + C_3\cos\omega_c\delta_{va}) \approx 0, \tag{10}$$

$$\omega_c^3 - (B_1 + B_2 + B_3\cos\omega_c\delta_{va})\omega_c + C_3\sin\omega_c\delta_{va} \approx 0. \tag{11}$$

Suppose $\frac{B_3\omega_c}{C_3} = \tau_s\omega_c \ll 1$, the equations can be further simplified into

$$C_3\cos\omega_c\delta_{va} \approx A_1\omega_c^2 - (C_1 + C_2), \tag{12}$$

$$C_3\sin\omega_c\delta_{va} \approx (B_1 + B_2)\omega_c - \omega_c^3. \tag{13}$$

Thus,

$$\cot\omega_c\delta_{va} \approx \frac{A_1\omega_c^2 - (C_1 + C_2)}{(B_1 + B_2)\omega_c - \omega_c^3}. \tag{14}$$

As $\delta_{va} \gg 1$, the smallest solution for ω_c is then obtained as $\omega_c\delta_{va} \to \pi$, which leads to

$$\frac{2\pi}{T} = \omega_c \approx \frac{\pi}{\delta_{va}} - \frac{\pi(B_1 + B_2)}{(C_1 + C_2)\delta_{va}^2} + \mathcal{O}(\frac{1}{\delta_{va}^3}). \tag{15}$$

The oscillation period T is thus approximately twice of δ_{va}.

Supplementary Figures

 <u>Figure S1</u>. Kalman filter tracking of vocalization $(s(t))$ and heart rate $(a(t))$ and estimation of a_B. Red traces
 are state tracking, blue traces are observation. a_B errors are predetermined process noise.

 <u>Figure S2</u>. Phasegram of individuals. These are 8 out of 10 individuals, the other 2 individuals have too few
 time points to plot. The last two subjects already started producing periodic oscillatory vocalizations from P1.

749   <u>Figure S3</u>. Left: an example of high E_B and high γ. Right: an example of high E_B but low γ.