Brain strain rate response: addressing computational ambiguity and experimental data for model validation

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Highlights

- A theoretical clarification of two strain rate computational schemes was presented, in which the strain rate was respectively calculated as the time derivative of strain or derived from the rate of deformation tensor.
- Discordant cross-scheme responses were substantiated by experimental principal strain rate and computational principal and tract-oriented strain rates.
- Cross-scheme comparison of strain rate responses is inappropriate.
- The newly calculated experimental strain rate curves can be used for strain rate validation of FE head model

Abstract

Traumatic brain injury (TBI) is an alarming global public health issue with high morbidity and mortality rates. Although the causal link between external insults and consequent brain injury remains largely elusive, both strain and strain rate are generally recognized as crucial factors for brain injury development. With respect to the flourishment of strain-based exploration, concerning ambiguity and inconsistency are noted in the scheme for strain rate calculation within the TBI research community. Furthermore, there is no experimental data that can be used to validate the strain rate responses of finite element (FE) models of the human brain. Thus, the current work presented a theoretical clarification of two commonly used strain rate computational schemes: the strain rate was either calculated as the time derivative of strain or derived from the rate of deformation tensor. To further substantiate this theoretical disparity, these two schemes were respectively implemented to estimate the strain rate responses from a previous-published cadaveric experiment and an FE brain model secondary to a concussive impact. The results clearly showed scheme-dependent responses, both in the experimentally determined principal strain rate and FE-derived principal and tract-oriented strain rates. The results suggest that cross-scheme comparison of strain rate responses is inappropriate, and the utilized strain rate computational scheme needs to be reported in future studies. The newly calculated experimental strain rate curves can be used for strain rate validation of FE head models.

Keywords

Traumatic brain injury; time derivative of strain; rate of deformation tensor; strain rate validation

Graphic abstract

1 Introduction

Traumatic brain injury (TBI), defined as an alternation in brain function or other evidence of brain pathology caused by an external force (Menon et al., 2010), is a major global public health threat. Recent epidemiological studies estimated that in the United States, around 61000 people died from TBI-induced injuries in 2019 (CDC, 2020), and in the European Union, there were approximately 82000 TBI-associated deaths in 2012 (Majdan et al., 2016). Despite collective efforts from the scientific community, clinical services, and policymakers to prevent the occurrence and mitigate the detriment of TBI, no clear improvement has been noted over the past two decades (Roozenbeek et al., 2013). Efforts to better decipher brain injury pathogenesis and develop more effective head protective strategies are vital to resolve this concerning epidemic.

Although the mechanobiological pathway from external insult, localized tissue responses, and resultant brain injury remains to be better explored, both strain and strain rate have been generally recognized as crucial factors for brain injury development (Bayly et al., 2012). For example, through relating controlled mechanical inputs to ensuing neuromorphological changes in a three-dimensional (3D) *in vitro* compression model, Bar-Kochba et al. (2016) found that the localization of neurite blebbing correlated with shear strain along the axial direction of neurons, while neurite swelling was highly loading-rate dependent. Montanino et al. (2019) utilized a composite finite element (FE) model of a single axon to predict the responses of axonal subcomponents secondary to tensile elongation, revealing a dependency of axolemma deformation on both the magnitude and the rate of imposed loadings. Recently, Donat et al. (2021) reported that strain and strain rate at the brain tissue level estimated by a rat FE head model correlated well with the neuroimaging-delineated alteration and histostaining-evidenced pathology from experimental rats secondary to controlled cortical impacts. Hajiaghamemar and Margulies (2021) found FE-derived strain and strain rate along the fiber tracts spatially aligned with the pattern of histology-identified axonal injury in the pig

brains after rapid head rotations. These experimental and computational findings collectively underscored the importance of monitoring both the strain and strain rate during head impacts.

In light of this, various strain-based metrics have been proposed over the past decades, such as maximum principal strain (MPS) in the brain (Bian and Mao, 2020; Kleiven, 2007a), cumulative strain damage measure (CSDM) quantifying the volume fraction of brain experiencing strain levels over a pre-defined threshold (Takhounts et al., 2008; Weaver et al., 2012), maximum tract-oriented strain (MTOS) measuring the stretch along axonal fiber tract (Sullivan et al., 2015; Zhou et al., 2021b). Many studies employed these strain-based metrics as measures to describe brain responses during head impacts. To date, strain-based measurements have been experimentally quantified by Hardy et al. (2007) (recently reanalysed by Zhou et al. (2019)) based on MPS (0.07-0.22) in cadaveric experiments with loading conditions close to traumatic levels, and by Knutsen et al. (2020) based on MPS (0.02-0.05) and MTOS (0.005-0.03) in voluntary loading scenarios. These strain-based metrics were also used to establish kinematic-based brain injury risk functions with the aid of FE head models. For example, the Brain Injury Criteria (BrIC) (Takhounts et al., 2013) was proposed based on the dependency of MPS and CSDM on the angular velocity observed from two FE models of human brain (Mao et al., 2013; Takhounts et al., 2008). Discussions are ongoing to potentially adopt this newly proposed BrIC into the New Car Assessment Program in the vehicle safety standards (Laituri et al., 2016).

Compared with the many strain-based studies, there are relatively fewer investigations on brain strain rate. This is evidenced by the large body of literature that merely examined the strain-based outputs (e.g., the studies by Zhou et al. (2020) and Bian and Mao (2020)). For those investigations with a focus on strain rate, it is ubiquitous that the strain rate values are directly reported without any clarification on how they are calculated (Beckwith et al., 2018; Kleiven, 2007a; Liu et al., 2021; Mao et al., 2006; McAllister et al., 2012; Post et al., 2015; Viano et al., 2005). Even when limiting to the handful of studies with the strain rate computational schemes elaborated, alarming inconsistency existed. In two experimental

paradigms, Rashid et al. (2013) determined the strain rate from the symmetric part of the

velocity gradient in the simple shear experiments of brain tissue, while LaPlaca and Thibault

(1997) quantified strain rate as the time derivative of strain in an *in vitro* TBI system. Similar

disparity is also noted in the computational works. For example, Zhou et al. (2022) derived the

strain rate from the rate of deformation tensor, while King et al. (2003) instead calculated the

strain rate by differentiating the FE-derived MPS vs. time curves. Concerningly, little or limited

knowledge exists on how this disparity in computational schemes affects the resultant strain

rate values.

When scrutinizing the existing strain rate-relevant investigations employing FE models

of human head, it is common that the model is validated against the brain-skull relative motion

and/or brain strain and then directly used for strain rate prediction. As highlighted by Zhou et

al. (2018) and Zhao and Ji (2020), if the model is evaluated against an experimental parameter

that is not the parameter of interest itself, it is suspected that numerical error might be

introduced and propagated through the chain of calculation from the evaluated parameter to

the parameter of interest (e.g., the relationship between brain motion/strain and brain strain

rate). Thus, direct evaluation of brain strain rate response of the FE model is preferred.

Unfortunately, there is no existing experimental brain strain rate dataset with loading conditions

close to traumatic levels. The reliability of the FE-derived strain rate responses remains to be

appropriately evaluated.

The current study endeavoured to address the ambiguity relevant to the calculation of

brain strain rate and provide experimental data for strain rate evaluation of FE models of the

human brain. Thus, a theoretical clarification of two currently used strain rate computational

schemes was presented, in which the strain rates were computed as the time derivatives of

strains of interest, and the first eigenvalues or components of interest of the rate of deformation

tensor, respectively. These two schemes were then implemented to calculate the experimental

principal strain rate and computational principal and tract-oriented strain rates, through which

the distinction between these two schemes was substantiated. The newly calculated

experimental strain rate curves can be served as a reference to evaluate the strain rate
responses of FE models.

2 Methods

2.1 Theoretical clarification

This section presents a continuum mechanical description of the two strain rate
computational schemes (referred to as scheme 1 and scheme 2 hereafter), in which the
calculation of two specific strain rates are elaborated (i.e., the first principal strain rate and the
tract-oriented strain rate). Two strain rates in combination with two computation schemes thus
yields 4 types of strain rates, all of which are used as response variables in TBI studies. The
mathematical symbols used throughout the text are clarified below: $\dot{\text{MPSR}}$ and $\hat{\text{MPSR}}$
represent the maximum first principal strain rates computed by scheme 1 and scheme 2,
respectively; $\dot{\text{MTOSR}}$ and $\hat{\text{MTOSR}}$ namely denote the maximum tract-oriented strain rates
following scheme 1 and scheme 2.

As is the common context of these two schemes, the motion $\boldsymbol{\varphi}(\mathbf{X},t)$ maps the position
vector of every point in the brain in the undeformed configuration ($\mathbf{X}$) (i.e., $\mathbf{X} = \{X_1 \quad X_2 \quad X_3\}$)
to its deformed configuration ($\mathbf{x}$) at time t (i.e., $\mathbf{x}(t) = \{x_1(t) \quad x_2(t) \quad x_3(t)\}$). The temporal
derivative of $\boldsymbol{\varphi}(\mathbf{X},t)$ at fixed position $\mathbf{X}$ represents the velocity, writing as
$\mathbf{v}(\mathbf{X},t) = \{v_1(t) \quad v_2(t) \quad v_3(t)\}$.

2.1.1 Scheme 1: Time derivative of strain

123   Scheme 1 calculates the strain rates as the time derivatives of the first eigenvalue or
124 the component of interest of the strain tensor $\mathbf{E}(t)$. Given that $\boldsymbol{\varphi}(\mathbf{X},t)$ consists of both
125 deformation and rigid-body motion, the deformation gradient $\mathbf{F}(t)$ is introduced, isolating the

126 deformation through mapping line elements ($d\mathbf{X}$) in the reference configuration into its

127 counterpart ($d\mathbf{x}$) in the current configuration.

128
$$\mathbf{F}(t) = \frac{\partial \boldsymbol{\varphi}(\mathbf{X},t)}{\partial \mathbf{X}} = \begin{bmatrix} \dfrac{\partial x_1(t)}{\partial X_1} & \dfrac{\partial x_1(t)}{\partial X_2} & \dfrac{\partial x_1(t)}{\partial X_3} \\ \dfrac{\partial x_2(t)}{\partial X_1} & \dfrac{\partial x_2(t)}{\partial X_2} & \dfrac{\partial x_2(t)}{\partial X_3} \\ \dfrac{\partial x_3(t)}{\partial X_1} & \dfrac{\partial x_3(t)}{\partial X_2} & \dfrac{\partial x_3(t)}{\partial X_3} \end{bmatrix} \tag{1}$$

129 The Green-Lagrange strain tensor is then calculated as

130
$$\mathbf{E}(t) = \frac{1}{2}(\mathbf{F}^{\mathrm{T}}(t)\mathbf{F}(t) - \mathbf{I}) \tag{2}$$

131 where the subscript T represents the transpose operation and $\mathbf{I}$ is the identify tensor.

132 The first principal strain ($\mathrm{PS}(t)$) is computed as the maximum eigenvalue of $\mathbf{E}(t)$ and

133 the tract-oriented strain ($\mathrm{TOS}(t)$) is obtained by projecting $\mathbf{E}(t)$ along the real-time fiber

134 orientation ($\mathbf{a}(t)$) (equation 3) (Zhou et al., 2021a):

135
$$\mathrm{TOS}(t) = \mathbf{a}^{\mathrm{T}}(t) \cdot \mathbf{E}(t) \cdot \mathbf{a}(t) \tag{3}$$

136 Resultantly, the first principal strain rate ($\dot{\mathrm{PSR}}(t)$) and tract-oriented strain rate ($\dot{\mathrm{TOSR}}(t)$)

137 were computed as the derivative of $\mathrm{PS}(t)$ and $\mathrm{TOS}(t)$ with respect to time (equation 4 and 5),

138 respectively. Of note, a five-point stencil neural network derivative function was used for all the

139 derivative operations in the current study following previous studies (Hajiaghamemar et al.,

140 2020b; Hernandez et al., 2019; Miller et al., 2021).

141
$$\dot{\mathrm{PSR}}(t) = \frac{\partial \mathrm{PS}(t)}{\partial t} \tag{4}$$

142
$$\dot{\mathrm{TOSR}}(t) = \frac{\partial \mathrm{TOS}(t)}{\partial t} \tag{5}$$

The maximum value of $\dot{\mathrm{PSR}}(t)$ and $\dot{\mathrm{TOSR}}(t)$ across the simulated time domain was identified

(i.e., $\dot{\mathrm{MPSR}}$ and $\dot{\mathrm{MTOSR}}$, respectively). Scheme 1 has been used in previous TBI studies to

extract the maximum principal strain rate (Hajiaghamemar and Margulies, 2021;

Hajiaghamemar et al., 2020a; Hajiaghamemar et al., 2020b; Hernandez and Camarillo, 2019;

Hernandez et al., 2019; King et al., 2003; Miller et al., 2021; Sahoo et al., 2016; Sullivan et al.,

2015; Vogel et al., 2020; Wu et al., 2021a; Zhang et al., 2003) and maximum tract-oriented

strain rate (Hajiaghamemar and Margulies, 2021; Hajiaghamemar et al., 2020a;

Hajiaghamemar et al., 2020b; Hernandez and Camarillo, 2019; Hernandez et al., 2019; Sahoo

et al., 2016; Sullivan et al., 2015; Wu et al., 2021a), respectively.

2.1.2 Scheme 2: Rate of deformation tensor

153 Scheme 2 calculates strain rate as the first eigenvalue or the component of interest of

154 the rate of deformation tensor $\mathbf{d}(t)$. Of note, the rate of deformation tensor is synonymous with

155 rate of strain tensor (Holzapfel, 2000). In contrast with the deformation gradient $\mathbf{F}(t)$ based

156 on displacement field, velocity gradient $\mathbf{l}(t)$ is derived from velocity field $\mathbf{v}(\mathbf{X},t)$ by taking its

157 spatial gradient according to equation 6:

158
$$\mathbf{l}(t) = \frac{\partial \mathbf{v}(\mathbf{X},t)}{\partial \mathbf{x}} = \begin{bmatrix} \dfrac{\partial v_1(t)}{\partial x_1} & \dfrac{\partial v_1(t)}{\partial x_2} & \dfrac{\partial v_1(t)}{\partial x_3} \\ \dfrac{\partial v_2(t)}{\partial x_1} & \dfrac{\partial v_2(t)}{\partial x_2} & \dfrac{\partial v_2(t)}{\partial x_3} \\ \dfrac{\partial v_3(t)}{\partial x_1} & \dfrac{\partial v_3(t)}{\partial x_2} & \dfrac{\partial v_3(t)}{\partial x_3} \end{bmatrix} \tag{6}$$

159 The symmetrical part of $\mathbf{l}(t)$ is defined as the rate of deformation tensor $\mathbf{d}(t)$:

160
$$\mathbf{d}(t) = \frac{1}{2}(\mathbf{l}(t) + \mathbf{l}^{\mathrm{T}}(t)) \tag{7}$$

161 The first principal strain rate ($\hat{\text{PSR}}(t)$) and tract-oriented strain rate ($\hat{\text{TOSR}}(t)$) were

162 respectively computed as the maximum eigenvalue of $\mathbf{d}(t)$ and the projection of $\mathbf{d}(t)$ along

163 the real-time fiber orientation ($\mathbf{a}(t)$) (equation 8).

$$\hat{\text{TOSR}}(t) = \mathbf{a}^{\text{T}}(t) \cdot \mathbf{d}(t) \cdot \mathbf{a}(t) \tag{8}$$

The time-accumulated peaks of $\hat{\text{PSR}}(t)$ and $\hat{\text{TOSR}}(t)$ were determined and labelled as

$\hat{\text{MPSR}}$ and $\hat{\text{MTOSR}}$, respectively. Scheme 2 has been previously used to extract the strain

rate in the brain (Reese et al., 2002; Zhou et al., 2022) and other biological tissues (e.g., leg

muscles (Malis et al., 2020; Mazzoli et al., 2018; Sinha et al., 2020; Sinha et al., 2018),

myocardium (Dou et al., 2003; Robson and Constable, 1996; Wedeen, 1992)).

While collectively considering these two schemes, it is worth mentioning that the

material time derivative of the Green-Lagrange strain tensor (i.e., $\mathbf{E}$ in scheme 1) is the pull-

back of the rate of deformation tensor (i.e., $\mathbf{d}$ in scheme 2) (equation 9) (de Rooij and Kuhl,

2016; Holzapfel, 2000). $\dot{\mathbf{E}}$ in equation 9 is termed as material strain rate tensor in the literature

(Holzapfel, 2000) and, to the best of the authors' knowledge, has also been used to inform the

computation of brain strain rate in two successive studies from one research group (Donat et

al., 2021; Ghajari et al., 2017). Given its limited usage within the TBI community, this scheme

is thus not incorporated in the current study.

$$\dot{\mathbf{E}} = \mathbf{F}^{\text{T}} \cdot \mathbf{d} \cdot \mathbf{F} \tag{9}$$

2.2 Experimental strain rate

To substantiate the disparity between scheme 1 and scheme 2 from an experimental

perspective and provide experimental data for strain rate evaluation of FE brain model, the

principal strain rates in the cadaveric experiments were respectively computed using both

schemes. In Hardy et al. (2007), thirty-five impacts were conducted on eight inverted and

perfused head-neck complexes of post-mortem human subjects, in which the motion of the

radiopaque neutral density targets (NDTs) implanted within the cadaveric brain was measured

using a high-speed, bi-planar X-ray system with a temporal resolution of 1 ms. To expedite the

estimation of local responses, Hardy et al. (2007) deployed 7 NDTs in a cluster configuration

with the center target being 10 mm from the resting six counterparts, roughly occupying 1.5 ml

of tissue. As represented by C1 and C2 in Figure 1A, two NDT clusters were implanted within

each cadaveric brain.

Following the selection criteria proposed by Zhou et al. (2018), 15 NDT clusters from

14 impacts with both impact kinematics and motions of all 7 associated NDTs recorded over

40 ms were chosen to calculate the principal strain rates. As illustrated by a representative

cluster in Figure 1B, representative FE nodes were generated based on the NDT starting

positions. Based on the NDT deploying characteristics in Hardy et al. (2007), 8 tetrahedron

elements were developed for each cluster by connecting each NDT to its neighbouring

counterparts (Zhou et al., 2019). The tracked motions were used to excite the tetrahedron

model, which was solved by the LS-DYNA software (Livermore Software Technology

Corporation, LSTC). From the completed simulation, the first principal strain rates were

respectively extracted as the time derivatives of the first principal strain (i.e., $\dot{PSR}(t)$ for the

scheme 1) and maximum eigenvalue of the rate of deformation tensor (i.e., $\hat{PSR}(t)$ for scheme

2). The strain rates of all available tetrahedron elements were averaged to provide a general

estimation of the strain rate within the volume occupied by the NDT cluster. The maximum

values of $\dot{PSR}(t)$ and $\hat{PSR}(t)$ were identified (i.e., $M\dot{PSR}$ and $M\hat{PSR}$). Due to the

unavailability of fiber orientation information of these cadavers, it is infeasible to compute tract-

oriented strain rate from this experimental data.

207
208 Figure 1. (A) Representative sketch of NDT cluster implanting schemes. For a given head, two NDT
209 clusters were implanted as cluster 1 (C1) and cluster 2 (C2). (B) Strain-rate estimation with the NDTs
210 as black spheres while tetrahedron elements in grey and mesh line in black. To visualize the center
211 NDT, all tetrahedron elements are shown in translucency.

2.3 Computational strain rates

2.3.1 The ADAPT model

We next implemented scheme 1 and scheme 2 to respectively derive the principal

strain rate and tract-oriented strain rate from an anatomically detailed head model (i.e., the

ADAPT model in Figure 2). The ADAPT model was developed previously at KTH Royal

Institute of Technology in Stockholm using the LS-DYNA software. Details about the model

discretization and material properties are available in Li et al. (2021). Briefly, the model

includes the brain, skull, meninges (i.e., dura mater, pia mater, falx, and tentorium),

subarachnoid cerebrospinal fluid (CSF), ventricle, and sinuses (i.e., superior sagittal sinus and

transverse sinus). As is the structure of interest in the current study, the brain was further

grouped with explicit representations of various components, including four lobes of the

cerebral gray matter (GM) (Figure 2A), cerebellar GM, cerebral white matter (WM), cerebellar

WM, and brainstem (Figure 2B). A second-order Ogden hyperelastic constitutive model

combined with six-order Prony series was employed to describe the nonlinearity and

viscoelasticity of brain material (Kleiven, 2007b). The biofidelity of the model has been

evaluated against experimental data of intracranial pressure (Nahum et al., 1977), brain-skull

relative displacement (Hardy et al., 2007), and brain strain (Zhou et al., 2019). The strain rate

responses of the model were additionally evaluated by the newly calculated experimental

strain rate data in Appendix.

To extract tract-related responses, it is imperative to monitor the real-time direction of

fiber tracts during the impact (Zhou et al., 2021a). Thus, the primary orientation of the WM fiber

tracts was extracted from the ICBM DTI-81 atlas (Mori et al., 2008) and then embedded within

the WM elements in the form of truss elements (Zhou et al., 2021a). Through this, potential

variation in fiber orientation during the impacts was reflected by the temporal direction of the

truss element, which was updated at each solution cycle of the time-marching simulation.

These temporally dynamic fiber orientation informed the computation of tract-oriented strain

rate.

239
240　Figure 2. Finite element model of the human head with embedded truss elements. (A) Finite element
241　model of the human head with the skull open to expose the subarachnoid cerebrospinal fluid (CSF) and
242　brain. A skull-fixed coordinate system and corresponding axes are illustrated with the origin at the head's
243　centre of gravity. (B) Brain components of cerebral white matter (WM), cerebellar WM, brainstem, pia
244　mater, falx, tentorium, superior sagittal sinus, and transverse sinus. (C) Brain with fiber tracts.

2.3.2 Impact simulation

The ADAPT model was employed to simulate a concussive impact, through which both

the maximum principal strain rates and maximum tract-oriented strain rates were respectively

computed using scheme 1 and scheme 2. The impact loading was measured by instrumented

mouthguard from a collegiate American football player with self-reported post-concussive

symptoms (Hernandez et al., 2015). Both translational and rotational acceleration profiles

(Figure 3) were applied to a node located at the head's centre of gravity and constrained to the

rigidly modelled skull. The model responses were output at every 0.1 msec, the same as the

output interval in Hajiaghamemar et al. (2020b).

254
Figure 3. Head model loading conditions for one concussive impact with the kinematics measured from
instrumented mouthguard. Note the X, Y, and Z axes are the same as those in the skull-fixed coordinate
system in Figure 1A.

Once the simulation was complete, the strain tensor and rate of deformation tensor of

all brain elements and the temporal orientation of the truss elements embedded within the WM

elements were obtained. For each brain element, the first principal strain rates were calculated

as the time derivatives of the first principal strain (i.e., scheme 1) and the first eigenvalue of

the rate of deformation tensor (i.e., scheme 2) with their maximum values across the whole

impact reported (i.e., $\dot{\mathrm{MPSR}}$ and $\hat{\mathrm{MPSR}}$). The tract-oriented strain rates were additionally

computed for all WM elements. In scheme 1, the strain tensor of each WM element was

iteratively resolved along the temporal direction of its embedded truss element (i.e., equation

3) with the outcome being discretely differentiated between time points to obtain the tract-

oriented strain rate (i.e., equation 5). In scheme 2, the rate of deformation tensor of each WM

element was projected along the real-time direction of its embedded truss element (i.e.,

equation 8). For both schemes, the maximum values of the tract-oriented strain rates across

the whole impact were reported, i.e., $\dot{\mathrm{MTOSR}}$ and $\hat{\mathrm{MTOSR}}$.

2.3.3 Data analysis

To quantify the cross-scheme disparity, relative differences in strain rate peaks

computed by scheme 1 and scheme 2 were computed for all brain elements based on the

principal measures and for all WM elements based on the tract-oriented measures,

respectively. To test whether cross-scheme strain rate results were readily translatable, linear

regression analyses were conducted between $\dot{\text{MPSR}}$ and $\hat{\dot{\text{MPSR}}}$, and between $\dot{\text{MTOSR}}$

and $\hat{\dot{\text{MTOSR}}}$, respectively. In addition to the correlation coefficient (r), the root-mean-square

error normalized by the mean value of strain rate peaks calculated by scheme 1 (i.e., NRMSE)

was reported. The threshold for significance was $p<0.05$.

3 Results

3.1 Experimental principal strain rates

Experimental principal strain rates computed by both schemes for all the 15 clusters

are depicted in Figure 4. Using scheme 1, the time history curves of the principal strain rate

were generally biphasic with both positive and negative phases, while the curves were

monophasic with purely positive values when using scheme 2. The strain rate responses of all

15 clusters attained the peak values within the pulse durations except for those in C241-T5 C1

and C393-T1 C1 solved by scheme 2. A cluster-wise comparison of maximum experimental

principal strain rates for the resting 13 clusters is shown in Figure 5. The maximum values

ranged from 9.3 s^{-1} (C241-T6 C2) to 25.8 s^{-1} (C380-T4 C1) in scheme 1 and spanned from

15.6 s^{-1} (C380-T2 C1) to 35.6 s^{-1} (C064-T1 C1) in scheme 2.

Figure 4. Comparison of experimental principal strain rates computed as time derivate of first principal strain ($\overset{\bullet}{\mathrm{PSR}}(t)$ in scheme 1) and first eigenvalue of the rate of deformation tensor ($\overset{\wedge}{\mathrm{PSR}}(t)$ in scheme 2).

Figure 5. Cluster-wise comparison of maximum experimental principal strain rates from 13 NDT cluster, each of which the strain rates was computed by schemes 1 ($\dot{\mathrm{MPSR}}$) and scheme 2 ($\hat{\mathrm{MPSR}}$), respectively.

3.2 Computational principal and tract-oriented strain rates

Figure 6A-F visually compares the maximum principal and tract-oriented strain rates at a representative horizontal cut-section. For the maximum principal strain rate (Figure 6A-C), although both schemes predicted high strain rate peaks (i.e., peaks over or approaching 40 s^{-1}) at the cortical region, substantial differences (i.e., differences over 10 s^{-1} or less than -10 s^{-1}) between $\dot{\mathrm{MPSR}}$ and $\hat{\mathrm{MPSR}}$ were noted in the left frontal lobe, right occipital lobe, and paraventricular region. For example, $\dot{\mathrm{MPSR}}$ was larger than $\hat{\mathrm{MPSR}}$ in the left temporal lobe (solid arrow in Figure 6C), while a reverse trend was noted in the left frontal cortex (dash arrow in Figure 6C). When switching to maximum tract-oriented strain rate (Figure 6D-F), $\dot{\mathrm{MTOSR}}$ and $\hat{\mathrm{MTOSR}}$ differed in magnitude and spatial distribution. In particular, scheme 1 revealed substantially larger tract-oriented strain rate peaks than scheme 2 across the whole WM in the frontal region (solid arrows in Figure 6F), while $\dot{\mathrm{MTOSR}}$ was relatively smaller than $\hat{\mathrm{MTOSR}}$ in the external capsules (dash arrows in Figure 6F).

This visible difference was further quantified by linear regression analysis between two schemes (Figure 6G-H). Although a significant correlation ($p<0.001$) was attained for both types of strain rates, the correlation coefficients (r) were 0.82 for maximum principal strain rate

315 and 0.77 for maximum tract-oriented strain rate. Across the whole brain elements, a NRMSE

316 value of 16.69% was computed between $\overset{\bullet}{\text{MPSR}}$ and $\hat{\text{MPSR}}$, while across the whole WM

317 elements, a NRMSE value of 27.35% was calculated between $\overset{\bullet}{\text{MTOSR}}$ and $\hat{\text{MTOSR}}$.

318
319 Figure 6. Comparison of time-accumulated maximum principal strain rate and tract-oriented strain rate

320 by both schemes. Upper row: Horizontal cross-sections of $\overset{\bullet}{\text{MPSR}}$ (A), $\hat{\text{MPSR}}$ (B), and their relative

321 difference (C). Middle row: Horizontal cross-sections of $\overset{\bullet}{\text{MTOSR}}$ (D), $\hat{\text{MTOSR}}$ (E), and their relative
322 difference (F). Lower row: Linear regression plots for element-wise principal strain rate in the whole
323 brain (G) and WM region (H), respectively.

4 Discussion

The current study addressed the ambiguity of brain strain rate computation by presenting a theoretical clarification of two alternative schemes, in which the strain rate was either computed as the time derivative of strain or derived from the rate of deformation tensor. This theoretical disparity was further substantiated by discordant cross-scheme responses in experimental principal strain rate, and computational principal and tract-oriented strain rates. It was thus verified that strain rates were dependent on the specific computational schemes, underscoring that cross-scheme comparison of strain rate results was inappropriate. The newly calculated experimental strain rate curves, especially those with the peak values attained within the pulse durations, can be used to evaluate the strain rate response of FE models of human head, as exempfied for the ADAPT model in Appendix.

Many *in vivo*, *ex vivo*, and *in vitro* models of TBI have reported a strain rate dependency of morphological or functional alterations of experimental tissues secondary to mechanical insults with precisely controlled loading parameters (Montanino, 2020; Morrison III et al., 2011; Wu et al., 2021b), while several computational studies have also reported strain rate as an appropriate predictor for brain injury (Hajiaghamemar and Margulies, 2021; Hajiaghamemar et al., 2020b; King et al., 2003; McAllister et al., 2012; Sullivan et al., 2015). However, most TBI studies didn't report details of strain rate calculation, making an appropriate interpretation of strain rate results impossible. For the handful of investigations with the strain rate calculation procedures elaborated, disparate inconsistency was noted in the computational scheme, posing inevitable challenges while comparing strain rate-related finding across studies. Our work demonstrated cross-scheme disparities in the resultant strain rate responses from the theoretical perspective as well as experimental and computational substantiation. These disparities highlighted that cross-scheme comparison of strain rate results is inappropriate. For example, the strain rate thresholds established by Patton et al. (2015) using scheme 2 were not comparable with the strain rate results derived by scheme 1 in the study by Hernandez and Camarillo (2019), even though both studies used the exact same FE model (Kleiven, 2007b).

Similarly, interpreting the FE-derived strain rate outputs with reference to the experimentally determined strain rate thresholds computed by another scheme was also inappropriate. We recommended that, if applicable, clarifying the strain rate computational details be regarded as standard content in future studies.

As both schemes are in use for strain rate computation and no consensus has been reached within the research community, we thus computed the experimental strain rates from the cadaveric tests reported earlier (Hardy et al., 2007) using both schemes. The newly calculated strain rate responses are then compared with values reported in other experimental studies (Table 1). When computing the strain rate as the time derivate of strain (i.e., scheme 1), the current experimental strain rate peaks largely exceeded those measured from the volunteer tests whose loading conditions were far from injury levels (Feng et al., 2010). Sullivan et al. (2015) estimated the strain rate using the scheme 1 in six horizontally dissected piglet brains during rapid rotational events, in which the imposed angular velocities were one order of magnitude larger than those in Hardy et al. (2007). Consequently, the reported strain rates in Sullivan et al. (2015) largely surpassed the current experimental results. When switching to scheme 2, Reese et al. (2002) reconstructed the velocity gradient from phase-contrast magnetic resonance imaging due to the vascular pulsatile and then computed 3D strain-rate tensor of brain parenchyma. This pulsatile-induced strain rate was two orders of magnitude smaller than the current experimental strain rates calculated by scheme 2.

Table 1. Comparison of brain strain rates from existing studies. Note that the strain rate in Feng et al. (2010) was estimated from Figure 10 in the original article using scheme 1, while the computational scheme for strain rate in Reese et al. (2002) was clarified in the caption of Figure 2 in the original article, corresponding to scheme 2 in the current study.

Scheme	Approach	Study	Loading conditions	Strain rate
Scheme 1: Time derivative of strain	Volunteer test	Feng et al. (2010)	Linear acceleration: 1.5 g Angular acceleration: 120-140 rad/s^2	Peak principal strain rate: ~1.2-2.7 s^{-1}
	Animal test	Sullivan et al. (2015)	Angular velocity: 121 ± 3.22 rad/s	Mean principal strain rate: 150-189 s^{-1}
	PMHS test	Current study	Linear acceleration: 38-291 g Angular acceleration: 2370-24206 rad/s^2 Angular velocity: 20.3 ± 5.7 rad/s	Peak principal strain rate: 9.3-25.8 s^{-1}
	Volunteer test	Reese et al. (2002)	Vascular pulsatility	Principal strain rate: 0.1 s^{-1}

Scheme 2: Rate of deformation tensor	PMHS test	Current study	Linear acceleration: 38-291 g Angular acceleration: 2370-24206 rad/s^2 Angular velocity: 20.3 ± 5.7 rad/s	Peak principal strain rate: 15.6-35.6 s^{-1}

FE brain models are valuable instruments to explore the pathway from external

mechanical insult and resultant brain trauma by linking FE-derived variables (e.g., strain rate)

to recorded injury information. The reliability of such exploration can be partially guaranteed

only if the FE-derived variables are validated against relevant experimental data. It has been

long recognized that an FE head model validated against experimental pressure does not

necessarily guarantee an acceptable predictability of the brain displacement or deformation

(Bradshaw and Morfrey, 2001; Kleiven and Hardy, 2002; Zhao et al., 2015). More recently,

several independent research groups advocated that a model to be used for strain prediction

should be evaluated against the experimental brain strain, and not just the brain-skull relative

motion (Zhao and Ji, 2020; Zhou et al., 2019; Zhou et al., 2018). When it comes to strain rate,

there is only one FE model of piglet brain (Sullivan et al., 2015), to the best of the authors'

knowledge, having the strain rate peaks (instead of the whole strain rate curves) evaluated

against the experimental strain rates. As is the typical practice, FE head models validated

against brain-skull relative motion and/or brain deformation are directly used for strain rate

prediction. Although the strain rate can be derived from strain or displacement field via a chain

of calculation (as detailed in section 2.1), a direct evaluation of FE-derived strain rate

responses is preferred if the model is intended to be used for strain rate prediction. This is

partially supported by the incongruent strain and strain rate validation performances of the FE

model of piglet brain (Sullivan et al., 2015), in which the computational strain rate peaks

exhibited superior correlation with their experimentally determined counterparts with respect

to that between the experimental and computational strain peaks. Note that, in the study by

Sullivan et al. (2015), the strain rate was computed as the time derivative of strain (i.e., scheme

1 in the current study). Taken together, to ensure the confidence of using FE head model for

injury prediction based on strain rate, a direct evaluation of FE-derived stain rate response

against its experimental counterpart (e.g., the newly calculated experimental strain rate curves

in the current study) is preferred, as exemplified by the ADAPT model in Appendix.

Several recent studies based on FE models of piglet brain (Hajiaghamemar and

Margulies, 2021; Hajiaghamemar et al., 2020b; Sullivan et al., 2015) have reported that the

maximum tract-oriented strain rate exhibited superior injury predictability over maximum

principal strain rate, while the majority of the existing FE investigations on human brain remain

to solely scrutinize the maximum principal strain rate. We thus examined the relationship

between the prevalently used principal strain rate and the recently advocated tract-oriented

strain rate across all WM elements in the ADAPT model. As quantitatively illustrated in Figure

7, no consistent scaling from maximum principal strain rate to its tract-aligned counterpart was

noted regardless of the computational scheme. This supports the potential incorporation of

tract-oriented strain rate as an independent injury metric in future studies of human brain injury,

in addition to the principal strain rate.

411
412 Figure 7. Element-wise comparisons of maximum principal strain rates and tract-related strain rates
413 computed by scheme 1 (A) and scheme 2 (B) based on results of WM elements.

4.1 Limitation and future work

Several limitations in this study should be noted, which require further investigation.

First, although our results clearly demonstrated cross-scheme disparities, we have no

recommendation on which scheme is superior. This can be alternatively done by evaluating

the injury predictability of these strain rate-based metrics computed by these two schemes.

For the time being, choices may be prioritized differently depending on the intended purpose,

but we strongly recommended that the exact strain rate computational schemes be clearly

stated in future studies. Secondly, the cadaveric experiments analysed in this study were not

without errors, as was acknowledged in the original studies (Hardy et al., 2001; Hardy et al., 2007) and exhaustively evaluated in several recent independent studies (Dutrisac et al., 2021; Zhao and Ji, 2020; Zhou et al., 2019; Zhou et al., 2018). Moreover, the principal strain rate responses estimated from the experimental data and the ADAPT model exhibited certain deviations (Appendix), although this is the first FE human brain model performing strain rate evaluation. Even so, we'd like to highlight that the main advocation of the current work, i.e., cross-scheme comparison of the strain rate results was not valid, is theoretically supported by the content presented in the "Theoretical basis" section in the current study. Such a advocation should not be disapproved by these experimental and computational substantiations with certain deficits. Nevertheless, we acknowledge that further demonstration of cross-scheme disparity is deemed necessary by employing other experimental data and better validated FE models. Lastly, the current work focused on the disparity of strain rate as outputs. Concerning inconsistency was also noted in the constitutive modeling of strain rate-dependent material, in which the strain rate served as an input to describe the stress-strain relations of a given material (Destrade et al., 2013; Rashid et al., 2013).

5 Conclusion

Given the ambiguity and inconsistency of strain rate computational schemes in existing brain injury studies, the current work presented a theoretical clarification of two rival schemes: strain rates computed as the time derivative of strain and derived from the rate of deformation tensor. Both schemes were employed to calculate experimental principal strain rate and computational principal and tract-oriented strain rates, exhibiting highly scheme-dependent responses. Our work underscores that cross-scheme comparison of strain rate should be avoided and the exact strain rate computational scheme needs to be clarified in future studies. The newly calculated experimental strain rate data can be used to validate the strain rate response of FE models of human head.

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Author contributions

Zhou Zhou: Conception and study design, Theoretical investigation, Finite element simulation, Data processing, Visualization; Writing – review & editing; **Xiaogai Li:** Conception and study design, Methodology, Review & editing; **Yuzhe Liu:** Theoretical investigation, Review & editing; **Warren Hardy:** Experimental data resources; **Svein Kleiven:** Conception and study design, Methodology, Review & editing.

Conflict of Interest

The authors declare that they have no conflict of interest.

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Appendix: Strain rate evaluation of finite element head model

Strain rate responses of the ADAPT model were validated against the newly calculated

experimental results. Three representative impacts were selected, including C1 in C288T3

(sagittal impact), C1 in C380T2 (horizontal impact), and C1 in C380T4 (coronal impact). To

better reflect the cranial geometry of experimental specimens, the model was scaled to match

the reported anthropometric measurements of the cadaveric heads. The nodes nearest to the

initial position of the experimental NDT markers were identified and further informed the

formulation of the eight tetrahedral elements representing the tissue encased by the

experimental NDT cluster. Following the same approach of computing experimental strain rate,

the 3D motion of the identified nodes with respect to the skull was acquired from the whole

head model simulation and further exercised the tetrahedral elements, through which the first

principal strain rate of each element was respectively extracted as the time derivatives of the

first principal strain (i.e., $\dot{\mathrm{PSR}}(t)$ for the scheme 1) and maximum eigenvalue of the rate of

deformation tensor (i.e., $\hat{\mathrm{PSR}}(t)$ for the scheme 2). For each cluster, the strain rates from all

eight elements were averaged.

667   The correlation between the computationally predicted strain rate responses and

experimentally determined measurements was quantitatively evaluated by the cross-

correlation method in CORA (correlation and analysis) with technical details available in

(Gehre et al., 2009). Briefly, the cross-correlation method in CORA scored the level of

correlation between a pair of time history curves by equally weighting their similarities in shape,

size, and phase. The CORA score ranges from 0 (no correlation) to 1 (perfect match). The

exact setting of relevant parameters for CORA implementation and biofidelity classification

based on CORA scores are respectively presented in Table A1 and Table A2.

675       Table A1. Parameter setting for the cross-correlation in CORA.

Parameter	Value	Parameter	Value
A_THRES	0.030	INT_MIN	0.70
B_THRES	0.075	K_V	3
A_VAL	0.010	K_P	1
B_DELTA_END	0.200	K_G	1
T_MIN / T_MAX (ms)	0/40	G_V	0.333
D_MIN	0.01	G_P	0.333
D_MAX	0.25	G_G	0.333

676 Table A2. Biofidelity rate and corresponding range of CORA score defined by ISO/TR 9790.

Biofidelity rating	CORA score
Excellent	0.86-1.00
Good	0.65-0.86
Fair	0.44-0.65
Margin	0.26-0.44
Unacceptable	0-0.26

677   Figure A1 illustrates the comparison of model-predicted strain rate curves against the

experiment-derived results for both schemes. The CORA score ranged from 0.65 to 0.71 for

scheme 1 and from 0.64 to 0.94 for scheme 2. According to the biofidelity classification based

on CORA scores (Table A2), the strain rate response of the ADAPT model can be rated as

"good" for scheme 1, and "good" to "excellent" for scheme 2.

Figure A1. Comparison of experimental and simulated strain rates computed by scheme 1 (upper row)
and scheme 2 (low row).