Reshaping the dynamics of follicle-stimulant hormone receptor models in polyunsaturated lipid bilayers. Calculation of conformational free energy landscapes of α-helical domains from all-atom MD simulations.

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Abstract

G-coupled-protein receptors (GPCR) are conspicuous target molecules for novel

therapeutic drugs due to their role as mediators of cellular responses. Structural biology

of GPCR revealed that intracellular signaling stimulated by extracellular ligands involves

subtle conformational changes of the receptor during activation. Nonetheless, transitions

among intermediates evolve in an intricate and rough free energy landscape of the

conformational space. Experimental evidence suggests that the membrane environment

is an active modulator of the receptor dynamics; therefore, the lipid composition may

facilitate conformational transitions towards productive signaling states. In this study, we

setup molecular dynamics simulations to examine the conformational dynamics of the

transmembrane domains, in the context of a membrane of polyunsaturated phospholipid

molecules, for a homology model of the human follicle-stimulating-hormone receptor

(FSHR) and the crystal structure of the Lumi intermediate of the squid rhodopsin (LSRh).

The conformational dynamics of the $\alpha-$helical domains of LSRh was consistent with

interactions stabilizing the crystal structure, which remained well preserved in the

membrane environment. In contrast, conformations in the FSHR model evolved towards

stable states in the membrane environment. To assess the relevance of the

conformational dynamics in the FSHR model, dihedral restraints were imposed for the

helical domains on top of the force field. This strategy was implemented to reoptimize the

interhelical interactions probably overlooked in the modeling process. The conformational

dynamics in the helical domains was evaluated by the TM-score, contact maps, principal

components analysis of $C\alpha$ atoms at the helical domains, and projections of the

conformational free energy on principal components. The roughness of the

conformational landscape in the FSHR model without dihedral restraints, suggested that

alternative interhelical conformational states were populated, whereas imposing

restraints led to a dominant conformational state. Template-based models of GPCR, with

reoptimized interhelical interactions using dihedral restraints, may enhance the

identification of binding sites for potential therapeutic drugs.

Keywords: GPCR, Molecular Dynamics, Lumi Squid-Rhodopsin, Glycoprotein family, Conformational Free Energy Landscape.

Introduction

G-coupled protein receptors (GPCR) encompass the largest family of membrane proteins that share a similar topology: seven-helix transmembrane (TM) domains, intra- and extracellular loops linking the TM helices, and the amino and carboxyl termini. These receptors represent highly attractive target molecules for drug discovery because of their role as mediators of an array of cellular responses triggered by a number of structurally diverse ligands including ions, photons, odorants, lipids, hormones, and neurotransmitters that vary considerably in size from small biogenic amines to peptides to large proteins [1-4]. From the outstanding structural data available in the Protein Data Bank [5], progress in the structural biology of GPCR revealed features on the conformational changes following agonist binding, intermediary stages upon receptor activation, and changes provoked by activated heterotrimeric G-proteins and other membrane-associated intracellular proteins (e.g. β-arrestins) [6]. Active conformations in crystal structures suggest similar intracellular motions in several GPCRs, including the β2-adrenergic receptor (β_2AR), the muscarinic acetylcholine receptor (M2R), and the μ-opioid receptor (μOR), where the transmembrane (TM) helix 6 moves up to 14 Å away from the helical bundle, and TM helices 5 and 7 show a slight rotation together with an inward motion [7]. Crystal structures also revealed tight interactions involving hydrophobic pockets, H-bonds, salt bridges, and location of structural water molecules or other molecules such as detergents required for crystallization [8]. Nonetheless, finding the proper set of crystallization conditions may require up to thousands of trials, and years of intensive research [9]. Computational methods provide practical alternatives for the development of working models useful as templates for finding new agonists and

86   antagonists, which may potentially serve as commercial drugs in pharmaceutical industry

87   [3, 10, 11].

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Molecular dynamics (MD) simulations have contributed to an atomistic description

of conformational changes in membrane proteins such as the GPCRs. Analysis on

trajectories at the $\sim10^3$ ns time scale disclosed some features on the activation

mechanisms of the β_2AR [12], and the serotonin-2A receptor [13], including the conserved

TRP switch, the opening of the ionic lock at TM helices 3 and 6, and the dynamics of the

NPxxY motif present in TM helix 7. To overcome the computer time limitations to sample

the conformational space, metadynamics techniques were proposed for describing the

activation mechanism of GPCR [14, 15]. The emerging consensus, therefore, suggests

that conformational transitions of GPCR reach 10^3-10^6 ns time scales [16, 17], which

might be inaccessible in some cases given the system size in explicit membrane

environments. At the $\sim10^2$ ns time scale, MD studies on GPCR still provide valuable

descriptions on fluctuations observed in the secondary structure dynamics, H-bond

dynamics involving internal water molecules, side chains contacts, loop dynamics, and

lipid-protein contacts [18].

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Our research group has been interested in understanding the physiopathogenesis

of dysfunctional gonadotropin receptors due to point mutations in the receptor protein,

from the biochemical and biophysical perspectives [19-21]. Studies on the glycoprotein

hormone receptors subfamily of GPCR have been addressed for the follicle-stimulating

hormone receptor (FSHR), which is the prototype member of these receptors [22, 23].

Our studies include gene sequencing to identify point mutations, receptor expression in

cell lines, binding to ligands and effectors, internalization and recycling of the receptor-

ligand complex, and intracellular signaling [24-27]. In addition, we implemented a

computational approach to determine the impact of point mutations in the interhelical

region of the FSHR [28]. Interestingly, the conformational dynamics of the wild-type

receptor (FSHR-WT), was sensitive to point mutations such as D408Y and D408A,

according to an analysis of autocorrelation matrices of the helical domains [28].

Apparently, when replacing D408 by tyrosine or alanine, fluctuations in protein backbone

atoms showed different patterns in correlated communities in a dynamical network

analysis [28]. More recently, new findings in the naturally occurring I423T missense

mutation located at the TM helix 2 of the FSHR (bearing the frequent polymorphic variant

S680), demonstrated modestly-reduced expression at the cell surface plasma membrane

and, more importantly, a significant impact of the replacement on cAMP/protein kinase A

and β-arrestin-dependent ERK1/2 phosphorylation signaling [29]. According to our

computational analysis, the helical integrity was well preserved in all phenotypes;

nonetheless, deviations from the initial receptor coordinates were detected in some TM

domains [28]. To examine in more detail the significance of conformational fluctuations in

the helical domains of the FSHR-WT model, in this study we extended the simulation

trajectories of the receptor to the $\sim 10^2$ ns time scale. In addition, we setup the crystal

structure of the lumi intermediate of squid rhodopsin (LSRh) using the same membrane

environment and computational protocols set for the FSHR-WT. The LSRh belongs to the

rhodopsin family of GPCR, and exhibits ~40% homology with the consensus sequence

of the TM domains of the glycoprotein hormone receptors (Figure 1), corresponding to

the previous intermediate of the meta I state [30].

Alphafold2 showed accurate predictions for native protein folding [31] in the 14th

Critical Assessment of protein Structure Prediction (CASP14). Nonetheless, structure-

function implications in missense mutations could be beyond the applicability of the

Alphafold2 workflow [32]. Thus, among other strategies to model and detect structure-

function implications of point mutations, MD simulations at all-atom force-field

approximation provide a physicochemical perspective to describe the significance of

disruptive interactions that may impact on the native receptor structure. In this study, we

setup template-based models of the FSHR-WT and the I423T mutant [29], in a lipid bilayer

of polyunsaturated 1-stearoyl-2-docosahexaenoyl-sn-glycero-3-phosphocholine (SDPC)

phospholipid molecules. For comparing the conformational dynamics, trajectories for the

LSRh crystal structure were generated under similar conditions. We assumed that in the

fluid membrane environment, interhelical interactions remain stable in the LSRh structure;

in the FSHR model, interhelical interactions evolved through a rough conformational

space. Previews analysis of trajectories showed that stable fluctuations could populate

states compatible with the roughness of the conformational landscape of the helical

domains [28]. To test the relevance of conformational dynamics in template-based

models, we imposed dihedral constraints in the FSHR models on backbone atoms of the

helical domains. Analysis of structural fluctuations and principal components (PC)

calculation of the interhelical domains, revealed the larger extent of the conformational

dynamics in the unrestrained *vs* the restrained state of the receptor. By using dihedral

constrains, the FSHR model preserved the predicted helical regions in the modeling

process, which was evaluated by the TM-score [33]. We conclude that our computational

implementation provides a strategy to explore the structural-functional implications of

missense mutations, and that it may be useful to develop new therapeutic drugs to

recover receptor function in dysfunctional GPCR.

Materials and Methods

Preparation of simulation boxes

Simulation boxes containing the GPCR protein, SDPC lipid molecules forming the

membrane, water molecules for hydration of the lipid heads, and Na^+ or Cl^- ions for system

neutrality were prepared similarly. Modeling of the FSHR-WT was performed as

described previously [27, 28]. Briefly, the human FSHR sequence D317 to N695 was

uploaded to the GPCR-I-TASSER server [34], selecting the highest scored model for

subsequent refinement by MD simulation. The GPCR-I-TASSER modeling procedure

includes multiple sequence alignments and experimental information stored in the GPCR

database such as contact maps, as suggested by mutagenesis experiments [35]. In

addition, the GPCR-I-TASSER performs *ab-initio* helix modeling, which improves the TM-

score for helical domains in GPCR [34]. The initial coordinates corresponded to the

highest TM-scored (0.4 ± 0.14) model calculated for the helical domains. Then, the FSHR-

WT (S680) model was inserted in a pre-equilibrated lipid bilayer of SDPC molecules

obtained from previous simulations [36]. Importantly, interhelical water molecules stable

in the squid rhodopsin setup were preserved [37]. Disulfide bonds were defined between

cysteines C338-C356 and C442-C517. Palmitoylated tails were bonded to cysteines

C644 and C646 by forming tioester bonds [27]. For the mutant I423T, we used the initial

receptor coordinates, and isoleucine side chain at 423 was replaced by the threonine side

chain, using the mutate command in the PSFGEN script [38]. Initial box dimensions were

90 Å x 105 Å x 114 Å, respectively, for the x-, y- and z-axis. A total of 128979 atoms were

included: 29231 water molecules, 248 SDPC lipid molecules, a sodium ion for charge

neutrality, and the receptor FSHR-WT or mutant I423T, with 379 residues encompassing

D317 to N695 [28].

The initial coordinates of the LSRh were taken from the crystal structure

PDB:4WW3 chain A [39]. A pre-equilibrated membrane lipid bilayer of SDPC lipids was

taken from a previews simulation of the dark state [36, 37]. The interhelical water

molecules resolved in the crystal structure were preserved for the setup. Disulfide bonds

between cysteines C96-C184, and C108-C186 were preserved as exhibited by the

corresponding bonds in other members of the rhodopsin family of GPCR. A palmitoylated

cysteine was defined at C337 and retinal was bound to K305. Initial box dimensions were

88 Å x 102 Å x 120 Å, respectively, for the x-, y- and z-axis. A total of 112118 atoms were

included in the LSRh simulation box, with 23705 water molecules, 249 SDPC

phospholipids, and two chloride ions for charge neutrality.

Dihedral constraints were imposed in the FSHR-WT or mutant I423T, for the helical

domains, using the ssrestraints plug in of VMD [40], with options h-bonds, a dihedral force

constant of 200 kcal/mol rad^2, and including helix 1: 362 to 392; helix 2 : 397 to 432; helix

3 : 437 to 472; helix 4 : 487 to 507; helix 5: 532 to 557; helix 6: 567 to 597; and helix 7:

607 to 626 as the selected helical regions for the starting FSHR model. Helical restraints

were also included for the I423T setup; therefore, restrained and unrestrained runs were

executed for FSHR-WT and the I423T mutant (Table 1).

Force Field parameters

207 CHARMM36 all-atom force field parameters were used for the lipid molecules [41] and

208 the protein atoms [42, 43], including CMAP correction [44, 45]. The retinal parameters

209 used in the LSRh setup corresponded to those reoptimized in reference [46], which

210 include atom charges from reference [47]. Water molecules were modeled using the

211 TIP3P potential [48].

Molecular Dynamics details

All simulations were performed with the NAMD 2.12 software [49]. Simulation trajectories

were generated in the isothermal-isobaric ensemble (NPT) with Langevin dynamics to

maintain a constant temperature, and Nosé-Hoover Langevin piston to maintain a

constant pressure of 1 bar. Anisotropic cell fluctuations in the x-, y- and z-axis were

allowed. Non-bonding interactions were calculated with a cutoff of 11.5 Å, and a shifting

function starting at 10.0 Å. A multiple time step integration for solving the motion equations

was used with one step for bonding interaction, two steps for short range nonbonding

interactions, and four steps for electrostatic forces. All hydrogen atoms were fixed using

the SHAKE algorithm [50]. Electrostatic interactions were evaluated using PME [51], with

a 4th order interpolation on a grid of ~1 Å in the x-, y- and z-directions, and a tolerance of

10^{-6} for the direct evaluation of the real part of the Ewald sum. The unrestrained

trajectories were generated for the FSHR-WT at 300K as well as at the physiological

temperature of 310 K (Table 1). Trajectory of the FSHR-WT at 300 K correspond to *in-*

vitro experimental conditions [29]. Trajectories using dihedral constrains were generated

for FSHR-WT at 300K, and I423T at 310 K. Trajectory of the I423T mutant at 310 K using

dihedral constrains was expected to recover interhelical interactions at physiological

temperature upon disruption of side chain contacts. The LSRh trajectory was generated

without helical restraints at 300K; LSRh system was considered as a comparable

reference system given that it shows ~40% of sequence homology with the glycoprotein

hormone subfamily of GPCR (Fig 1). The set of simulations were prepared to explore the

conformational space of the seven-helix domain in the FSHR-WT, and the novel mutant

I423T [29].

Table 1. Trajectory lengths for the simulation boxes of ~120 thousand atoms
containing the receptor, solvent water molecules, and SDPC lipids.
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System	Trajectory length	Trajectory length helical-restraints	
	ns	Ns	Temperature
LSRh	145.9		
FSHR	158.3	160.5	300 K
FSHR	144.3		
I423T	180.8	140.0	310 K

Trajectory analysis

Conformational changes were calculated by the root mean square deviation (RMSD) for

the Cα atoms of the helical domains in FSHR-WT and I423T mutant: Y362 to Y392 (helix

1), P397 to Y432 (helix 2), Q437 to T472 (helix 3), A487 to G507 (helix 4), M532 to R557

(helix 5), D567 to L597 (helix 6), and A607 to Y626 (helix 7). In LSRh: D31 to T62 (helix

1), P68 to L98 (helix 2), A106 to G138 (helix 3), H149 to P170 (helix 4), D194 to R240

(helix 5), L246 to P288 (helix 6), and P294 to V317 (helix 7). For each TM helix, contact

maps were calculated for side chain heavy atoms using a cutoff of 3.5 Å. Hence, TM

interactions were detected as function of simulation time whenever side chain heavy

atoms were within the cutoff distance.

To assess deviations from the initial helical structure of the transmembrane

domains, we calculated the TM-score defined as:

$$TM - score = Max \left(\frac{1}{L_N} \Sigma_{i=1}^{L_T} \frac{1}{1 + \left(\frac{d_i}{d_0} \right)^2} \right)$$

where L_N is the length of the reference structure, L_T is the length of the aligned residues,

d_i is the distance between the *i-th* pair of aligned residues, and d_0 is a scale to normalize

the match difference. 'Max' denotes the maximum value after optimal spatial

superposition. The TM-score takes values within 0 and 1, with TM-score > 0.5 for well-

preserved helical domains [33].

The conformational landscapes of LSRh, FSHR-WT, and I423T were obtained by

a principal component analysis of the Cα carbons (N) using the GROMACS 2019.4 suite

of programs [52, 53]. By diagonalization of covariance matrices, the resultant matrix

transformation contained eigenvectors in 3 x N columns as well as the eigenvalues

corresponding to the mean square fluctuation in direction of the principal modes, and

ordered from the largest to the lowest contribution to the total fluctuation [54]. As a

measure of the total fluctuation, the root mean square fluctuation (RMSF) was calculated

from the square root of the sum of eigenvalues divided by the number of Cα atoms (N).

The simulations trajectories were divided in blocks of 2^9, 2^{10}, … ,2^{16} ps in order to

calculate averages and standard deviations as a function of the block size [55].

Visualization of the trajectory projected over the first six PC was generated using the VMD

software [40]. Free energy landscapes were calculated from projections of the principal

modes, PC1 to PC6. Our computational approach, including the system set-up,

generation of MD trajectories, and analysis of fluctuations allowed us to apply a strategy

to evaluate both the conformational dynamics of a membrane receptor relaxed from a

crystal structure, such as the LSRh, or a homology model for the FSHR-WT, and the

impact of the I423T mutation located at the TM helix 2 of the helical bundle

Results

By comparing the conformational dynamics of LSRh and a homology model of the FSHR-WT and I423T mutant, we analyzed the dynamics of the Cα atoms at the helical domains (Fig 1). Such strategy was implemented to assess the conformational fluctuations as a means to account for the biological relevance of helical motion in hormone receptor models *vs* a photon-activated x-ray structure (Fig 2). From a previous study, the integrity of the helical domains of the FSHR-WT was evaluated according to calculations of the secondary structure, number of residues forming each TM helix and root mean squared deviations (RMSD) of each TM helix, RMSF, and other parameters. We found that helical domains were well preserved [28]. Nonetheless, using the TM-score we detected values below the 0.5 helicity threshold for both, the FSHR-WT and I423T mutant, suggesting a disruption of the TM folding of these FSHRs (Fig 3). Interestingly, the TM-score in LSRh showed values close to one for TM helices 1, 2, 5, and 6, whereas TM helices 3, 4, and 7 exhibited significant variability, yet above the 0.5 threshold. TM-score values for TM helices in LSRh suggest that the secondary structure was well preserved. As the initial LSRh structure was obtained from x-ray diffraction at 2.8 Å resolution [39], helical interactions were compatible with the tight packing of the crystal structure, and remained stable in the fluid environment of the membrane. Importantly, TM-score calculations for LSRh were performed without any restrained force on top of the force field.

Fig 1. A) Sequence alignment of the squid-rhodopsin (LSRh) and the glycoprotein family: thyroid stimulant hormone receptor (TSHR), follicle stimulant hormone receptor (FSHR), and luteinizing hormone receptor (LHR). The Ballesteros and Weinstein numbering system was used, which assigns 50 to the most conserved position on TM helices 1-7 [56]. B) Schematic representation of the seven TM domains [57]. Conserved position 50 of on each TM helix is highlighted (red circle).

Fig 2. Membrane receptors in the lipid bilayer environment at 300 K. FSHR-WT and
LSRh are depicted, respectively, as green and magenta ribbons. SDPC lipid molecules
are depicted as spheres in the background, with carbon atoms in gray, and phosphate
atoms as ice-blue. The solvent water medium is depicted as a continuum in ice blue.
The structural alignment of the receptors shows consistency in the location of the helical
domains in the hydrophobic region of the membrane. The structural alignment between
these receptors show some differences, for example, the extension of intracellular
domains of TM5 and TM6, the tilt angle of TM1, and kink angle of TM7.

Table 2. TM-score values for each TM helix in restrained and unrestrained runs,
and at 300 K and 310 K. Values larger than 0.5 represent well preserved helicity.

System	TM-score							
	TM1	TM2	TM3	TM4	TM5	TM6	TM7	Temp
LSRh	0.87±0.02	0.76±0.03	0.80±0.03	0.60±0.05	0.79±0.04	0.86±0.02	0.71±0.04	300 K
FSHR	0.46±0.03	0.54±0.03	0.64±0.04	0.46±0.04	0.73±0.05	0.71±0.05	0.31±0.02	
FSHR+ restraints	0.86±0.04	0.93±0.02	0.91±0.03	0.72±0.07	0.89±0.03	0.91±.0.03	0.82±0.04	
FSHR	0.51±0.03	0.59±0.05	0.70±0.05	0.41±0.04	0.77±0.04	0.73±0.05	0.53±0.03	310 K
I423T	0.49±0.04	0.67±0.03	0.81±0.06	0.40±0.04	0.73±0.04	0.69±0.05	0.32±0.03	
I423T+ restraints	0.87±0.04	0.93±0.01	0.90±0.03	0.66±0.08	0.88±0.03	0.89±0.04	0.79±0.05	

In FSHR-WT and I423T, the TM-score values displayed significant differences

depending on the simulation conditions. On the one hand, TM-score values for TM helices

1, 4, and 7 were lower than 0.5 in the unrestrained runs (Table 2); for TM helices 2, 3, 5,

and 6 TM-score the corresponding values fluctuated above the 0.5 helicity threshold

(Table 2). Interestingly, TM-score values of helix 5 were similar to those calculated for the

LSRh run (Fig 3). On the other hand, in the restrained runs all TM-score values were

stable and close to one (Table 2). TM helices 1, 4, 6 and 7 showed similar trends to those

calculated for the LSRh (Fig 3). For TM helix 5, our analysis suggested that it might not

necessary to use helical restraints at all, since the TM-score values fluctuated as those

in LSRh (Table 2 and Fig 3). In TM helix 2, the TM-score was larger in the receptor models

than the values calculated for LSRh (Table 2 and Fig 3). Regarding the magnitude of the

force constant of the helical restraints, TM-score variability was well reproduced as that

observed in LSRh, except for TM helix 2; for example, FAHR-WT helix 1, 4, 6, and 7

showed fluctuations within magnitudes similar to those in LSRh (Table 2 and Fig 3). In

TM helix 2, a weaker force constant might improve the comparison against helix 2 of

LSRh. Thus, TM-score was a parameter sensitive to distortions of helicity in the TM

domains in the FSHR-WT and I423T mutant. Finally, the response of TM-score due to

the increase in temperature from 300 K and 310K in FSHR-WT showed similar values

within error bars, except for TM helix 7 of FSHT-WT, whose helicity was better preserved

at 300K (S1 Figure in S1_info, Table 2, and Fig 3). Similar trends in the TM-score of I423T

at 310K were detected, except for the higher helicity of TM helix 3 relative to that in FSHR-

WT at 300K (Table 2, and Fig 3). The overall results showed that helicity measured by

the TM-score was higher in LSRh relative to those in FSHR-WT and I423T without helical

restraints, which suggested that the intermolecular interactions stabilizing the helical

domains in LSRh were well represented by the force field; hence, helicity disruption in the

receptor models such as the FSHR-WT and I423T were largely related to suboptimal

interhelical interactions.

355
356   **Fig 3**. TM-score of the TM domains of the FSHR-WT and I423T models, respectively at
357   300 and 310 K, and the LSRh at 300K. Unrestrained runs correspond to dash lines in
358   green. Blue dash lines represent the TM-score for the restrained runs. TM-score
359   showed well reserved helicity in the restrained runs.
360
Dihedral restraints had an impact on the conformational changes according to

RMSD calculations of the TM domains. Fig 4 shows time series for the RMSD for all runs.

After 20 ns of trajectory RMSD fluctuated at an average of 1.2 ±0.10 Å in LSRh, and 0.83

± 0.11 Å and 0.86 ±0.12 Å, respectively, for the restrained runs of FSHR-WT and I423T.

For the unrestrained runs, after 50 ns of trajectory, RMSD values reached an average

1.75±0.19 Å and 1.99 ± 0.12, respectively, for FSHR-WT and I423T at 310 K; for FSHR-

WT at 300 K, RMSD reached 2.42±0.15 Å. The larger RMSD averages in unrestrained

runs were expected because restrained forces attenuated displacements of TM domains.

To measure the positional displacements from the starting structure, RMSD calculations

for each TM helix were also performed. In the restrained runs, RMSD averages were

calculated from 0.29 to 0.51 Å for the FSHR-WT helices, and 0.31 to 0.53 for the I423T

mutant (S2 and S3 Figures in S1_info). In LSRh, RMSD averages from 0.6 to 1.2 Å

suggested slight positional displacements of the TM helices in the fluid environment of

the membrane (S4 Fig in S1_info). Without restrained forces, TM helices showed stable

average configurations in both, FSHR-WT and I423T at 310 K: RMSD from 0.6 to 1.2 Å

for helices 3, 4, 5, and 6; and RMSD 1.3 to 2.0 Å for helices 1, 2, and 7. Larger RMSD

values for individual TM helices correlated with decreasing helicity, for example, in helices

1 and 7 (Fig 3).

Fig 4. Root mean square deviations (RMSD) of the helical domains in the FSHR-WT
(top plot) and I423T mutant (mid plot), with and without helical restraints, at 300 K or
310K. RMSD for the LSRh helical domains at 300K (bottom plot). Using helical
restraints modulate the positional displacements.
384

Interactions stabilizing the 3D structure in the FSHR-WT and I423T models were

evaluated by measuring contact maps for each TM-helix. Importantly, TM helices work as

molecular switches since interhelical interactions may lead to conformational changes

[58]. Fig 5 shows the results for contact maps in the FSHR-WT for the restrained and

unrestrained runs. Consistently detected contacts were exhibited by the TM helical

domains: TM1 made permanent contacts with TM2 and TM7; TM2, with TM3, TM4, and

TM7; TM3, with TM4, TM5, and TM7; TM4, with TM5; TM5, with TM6; and TM6, with

TM7. In TM2 of I423T mutant, contacts with TM7 involved V612, H615, L616, S619, Y626,

T630, and F633 (Fig 6), with a difference in the restrained *vs* the unrestrained run

consisting in fewer partner exchanges and interhelical contacts detected every 3 or 4

residues apart. Contacts between TM7 and side chains of TM6 were absent in the

restrained run, for example with side chains of M576, T580, and K598 (Fig 6). TM3

showed significant contacts with TM6 in the unrestrained run, involving L577, D581,

C584, and M585, and I588 (Fig 6). For the light receptor LSRh, Fig 7 shows contacts

among the TM helices: TM1 was in permanent contact with TM2, and TM7; TM2, with

TM3, and TM7, and the intracellular half of TM4; TM3, with TM4, TM5, and TM6; TM5,

with TM6; and TM6, with TM7. In summary, contacts detected in the LSRh were

consistent with contacts identified in the FSHR-WT and I423T models, although slight

differences were detected in restrained *vs* the unrestrained runs. No disruption of

interhelical contacts due to a temperature increase were detected, as replica of the

FSHR-WT at 310 K showed similar interhelical contacts as trajectory at 300 K (S5 Fig in

S1_info).

Fig 5. Contact maps as function of simulation time for the TM-helices in the FSHR-WT at
300 K. Data points were depicted whenever residues of TM-X were at 3.5 Å of residues
of TM-Y. In the vertical axis the TM helices were identified from 1 to 7, for each plot of
TM1-7. Analysis performed for the last 20 ns of trajectory.
414

416 **Fig 6**. Contact maps as function of simulation time for the TM-helices in the I423K at 310
417 K. Data points were depicted whenever residues of TM-X were at 3.5 Å of residues of
418 TM-Y. In the vertical axis the TM helices were identified from 1 to 7, for each plot of TM1-7.
419 Analysis performed for the last 20 ns of trajectory.
420
Fig 7. Contact maps as function of simulation time for the TM1-TM7 of LSRh. Data points
depicted whenever residues of TM-X were at 3.5 Å of residues of TM-Y. In the vertical
axis the TM helices are identified as from 1 to 7, for each plot of TM1-7. Analysis
performed for the last 20 ns of trajectory. TM1 was in contact mainly with TM2 and TM7;
TM2 made contact with TM3, and TM7; TM3 with TM4-6; TM5 with TM4 and TM6; TM6
with TM7.
427
428 **Fig 8.** Normalized density distributions for vectors PCA1-6 (nm) for the FSHR-WT, and
429 the I423T mutant. Distributions for the restrained and unrestrained runs are shown.
430 Unimodal distributions were found in the restrained runs, except for TM2 of FSHR-WT
431 and TM1 of I423T. PC distributions unveil the population of conformational states
432 sampled in the simulation trajectories.
433
434

Because of the helical restraints imposed on FSHR-WT and I423T mutant,

conformational states would differ from those sampled in the unrestrained runs. Even

though the unrestrained runs preserved the overall structure of the receptor [28], in

present study we examined more deeply the relevance of conformational fluctuations of

receptor models using the Cα atoms at the helical domains only. Fig 8 shows distributions

for PC1-6 of the FSHR-WT, and I423T mutant obtained from PCA calculations; the

normalized histograms for restrained *vs* unrestrained runs were overlapped for

comparisons. Profiles of the restrained runs showed unimodal distributions, except in PC2

of FSRL-WT and PC1 of I423T, which showed distributions with two and three maxima,

respectively. Unrestrained runs showed wider distributions than those obtained for the

restrained trajectories, which suggests that the receptor evolved to states not accessible

when helical restraints were imposed (Fig 8). Fig 9 shows the conformational change due

to PC1, which corresponds to the motion that includes the largest contribution to the total

fluctuation. In LSRh and the restrained FSHR-WT, helical motions showed small

amplitudes in all TM helices, with consistent displacements in TM helices 1-4, as well as

the extracellular half of TM5 and TM6 (Fig 9). The main difference between LSRh and

restrained FSHR-WT was observed in the intracellular side of TM of helices 5 and 6: in

LSRh, TM5 moved toward TM6, while TM6 moved inwards whereas in FSHR-WT, both

TM5 and TM6 displaced outwards. The PC1 collective motion in restrained I423T showed

small amplitudes in all helical domains, except in TM6 which moved outwards. In the

unrestrained runs of FSHR-WT at 300 K and 310 K, displacements were consistent in

amplitudes and directions with those observed in the restrained run, although the upper

half of TM6 moved with larger amplitude at 310 K (Fig 9). Despite the slight displacements

in the 3D space, the transformation of the covariance matrix performed in PCA allowed

to identify the axis (eigenvectors) that account for the variability of the $C\alpha$ coordinates. In

the rotated coordinates, the PC disclosed the collective low frequency modes of motion

for the helical domains. Simulation studies on resolved structures of GPCR performed at

the micro- and millisecond time scales provided some clues on the activation mechanism,

such as the dynamics of the conserved TRP switch and the NPxxY motif, and the opening

of the ionic lock at TM3 and TM6 [12, 13]. Such long simulation times scales were required

to overcome configurational free energy barriers [59, 60]. Using a subset of coordinates

such as the $C\alpha$ atoms of the helical domains, at the 10^2 ns time scale, the LSRh trajectory

was expected to relax into a dominant conformational intermediary, which was evident

from the calculation of unimodal distributions in PC1-6 (Fig 10). Therefore, from the

profiles in the distributions of the PC1-6 in the FSHR-WT and I423T models, and the

LSRh, we noticed that the unrestrained runs showed significant differences in the

dynamics of the helical domains, probably related to suboptimal interhelical interactions.

Further, from the calculation of the total fluctuation (RMSF) of Cα atoms in blocks of 2^9 to

2^{16} time frames (1 ps each), we determined the block size for the convergence of the

RMSF (S6 Figure in S1_info). In LSRh, total fluctuation reached constant value of 4 pm

for 2^{16} time frames (65.5 ns). In contrast, for the unrestrained runs, the total RMSF

increased with the block size. When helical restraints were imposed, the total RMSF in

both FSHR-WT and the I423T showed a similar behavior as the LSRh, reaching a 3 pm

plateau for a block size of 2^{16} time frames. In summary, the PCA1-6 of unrestrained runs

showed broad distributions, whereas the restrained runs as well as LSRh showed mostly

unimodal distributions in the PC1-6 principal modes.

Fig 9. Top view of helical domains, TM1-TM6, for receptors LSRh, FSHR-WT and I423T.
For a reference coordinates x, y, and z are shown, respectively, as red, green, and blue
arrows; all views share the same relative position of TM helices as in panel A. A. Helical
domains of the LSRh, helices TM1 to TM6, are depicted as cyan ribbons. Blue arrows
depicted on the Cα atoms represent the direction of the motion associated to the first
principal component whose sizes correspond to the magnitude of the displacement. B.
Helical domains of the FSHR-WT receptor model, at 300K. C. Helical domains of the
FSHR-WT including helical restraints, at 300K. D. Helical domains of the FSHR-WT
receptor model, at 310K. E. Helical domains of the I423T receptor model including helical
restraints, at 310K. F. Helical domains of the I423T receptor model at 310K.
493
494 **Fig 10**. Normalized density distributions for vectors PCA1-6 (nm) for the LSRh. PC
495 distributions unveil the population of conformational states sampled in the simulation
496 trajectories. Unimodal distributions of the LSRh suggests that fluctuations were detected
497 for one main conformational state.
498

By choosing specific atoms of the protein structure we implemented PCA to

determine the relevance of positional fluctuations expected for atoms at the helical region

in a light activated receptor and in hormone receptor models. Including the overall 3D

structure of a membrane protein to detect conformational changes may require long time

scales to sample a complex configurational space. Nonetheless, PCA analysis on

subdomains (e.g. binding pocket residues, backbone atoms, etc.) may provide

characteristics of the dynamics on a local scale, even at the ~10^2 ns time scale [18].

Interesting, projections of the free energy landscape corresponding to the PCA of a

reduced set of 3D coordinates, such as the Cα atoms of the helical domains, provide an

energetic criterium for the transitions among different conformations generated by

principal components. Figures 11 and 12 show the free energy projections of PC1-2, PC3-

4, and PC5-6 for FSHR-WT and I423T models, respectively. For comparisons, the free

energy projections for the LSRh were included in both figures. In Fig 8 we calculated

multipeak probability density distributions for the unrestrained runs mainly for PC1-3,

which produced several local free energy minima spread in the PC1-PC2 and PC3-PC4

projections (Figs 11 and 12). In the unimodal distributions of Fig 8, corresponding to

restrained runs, free energy minima in all projections were localized in a central region of

the plots (Figs 11 and 12). For the LSRh trajectory, projections of free energy showed

well defined minima in PC3-4 and PC5-6, whereas few local minima were detected in

PC1-2; transitions among configurations due to PC1-2 could be expected since no helical

restrains were imposed on the LSRh structure. Therefore, from projections of the free

energy on PC, we could determine the extent of the conformational dynamics of FSRH-

WT and I423T models in a fluid membrane environment, relative to the dynamics of the

LSRh. By using dihedral restraints in the FSRH-WT and I423T models, the modulation of

helical dynamics allowed to preserve interhelical contacts as well as the integrity of the

helicity of the TM domains, which are known to work as structural switches on the

activation process of GPCR. Our results suggest that including helical restraints could

improve the quality of a working model for membrane receptors, such as the FSHR-WT,

without any additional cost of computer time in comparison to the generation of

unrestrained trajectories.

Fig 12. Projection of the free energy landscape on principal components, vectors 1 to 6
(in nm), for LSRh at 300 K, and I423T at 310 K. Red dots cluster each other in the free
energy minimum. In the unrestrained run of I423T, the projection of PC1 and PC2 disclose
a broad distribution of local minima that suggests transitions among several
conformational states.
536
537 **Fig 11.** Projection of the free energy landscape on principal components, vectors 1 to 6
538 (in nm), for LSRh and FSHR-WT at 300 K. Red dots cluster each other in the free energy
539 minimum. In the unrestrained run of FSHR-WT, the projection of PC1 and PC2 disclose
540 a broad distribution of local minima that suggests transitions among several
541 conformational states.
542

Discussion

Understanding signal transduction mediated by the FSHR, a prototype member of the

glycoprotein hormones subfamily of GPCR, relies on the description of the conformational

free energy landscape in the membrane environment. In the process of signal

internalization, the intracellular receptor domains engage G proteins, or β-arrestins [6]

through conformational changes that expose particular intracellular sequences involved

in the activation of these pivotal, universal signal transducers [7, 14, 16]. In this study, we

used the covariance matrix to unveil the conformational dynamics of helical domains in

simulation trajectories at the $\sim 10^2$ ns time scale for both the FSHR-WT and the I423T

mutant at 300 and 310 K, respectively (Table 1). Further, to evaluate the significance of

conformational fluctuations in template-based models, the dynamics in a homologous

receptor (Figs 1 and 2) such as the LSRh were also analyzed. Relaxation of the initial

LSRh crystal structure in a fluid membrane environment, evolved to states compatible

with intermediates in early stages of activation. Therefore, we chose the LSRh structure

as a reference system because it is an early intermediate of the activation of visual

rhodopsin in invertebrates [39], with the retinal chromophore in all-*trans* configuration.

Hence, it was expected that fluctuations of the TM domains, helices 1 to 7, populate

configurational states in preparation for the conformational changes forming the binding

pocket for engaging intracellular transducers. Finding a dominant state in a relaxed crystal

structure or in a reoptimized model could suggest common features, assuming that TM

interactions emerge from: I) Sequence similarity, whereby side chain interactions could

be found in conserved motifs; II) Interhelical interactions that became permanent in the

simulation not necessarily present in the initial model; and III) Interhelical interactions in

the crystal structure that remain stable in the fluid membrane environment. By using

dihedral restraints in the FSRH-WT and I423T models, the modulation of helical dynamics

allowed to preserve interhelical contacts as well as the integrity of the helicity of the TM

domains, which are known to work as structural switches on the activation process of

GPCR. From total RMSF calculated for the Cα atoms, we found that 2^{16} time frames were

sufficient to include most of the variability in the restrained runs, while in the unrestrained

runs the total RMSF was not stable at the ~10^2 time scale. Our results suggest that

including helical restraints could improve the quality of a working model for membrane

receptors, such as the FSHR-WT, without any additional cost of computer time in

comparison to the generation of unrestrained trajectories.

Conclusions

In this study we explored the conformational dynamics of the LSRh relaxed in a membrane bilayer as well as a working model of the FSHR-WT and its I423T mutant. Since the homology modeling protocol employed was based on the GPCR-Tasser approach —which takes into account the TM-score as a parameter to evaluate the quality of a GPCR model— we tested the behavior of such parameter as a function of the simulation time for a set of MD trajectories at 300 K and 310 K, including helical restraints in the FSHR-WT and I423T models. We found that using helical restraints improves the TM-score with values close to one, without compromising the computer time required to generate simulation trajectories at $\sim 10^2$ ns time scale. The advantages of using helical restraints in working models included the predominant conformational state shown by the dynamics of the receptor, resembling the behavior also detected for the LSRh structure and the re-optimization of the interactions among the TM helices when the homology modeling misses relevant contacts in the helical bundle. Although the overall structure of the receptor is well preserved in the unrestrained runs and the first principal mode presented similar trends in the collective motion of the TM helices, we suggest that helical restraints may be part of the modeling refinement of a GPCR homology model.

599
600
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602 review & editing. **Alfredo Ulloa-Aguirr**e Conceptualization, investigation, writing -
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Supporting information

S1_info. Data describing the dynamics of the 3D structure of the membrane protein receptors: FSHR-WT, I423T and LSRh.

S1 Fig. Time series for the TM-score of FSHR-WT

S2 Fig. Root mean square displacement (RMSD) for FSHR-WT

S3 Fig. Root mean square displacement (RMSD) for I423T

S4 Fig. Root mean square displacement (RMSD) for LSRh

S5 Fig. Contact maps as function of simulation time for the TM-helices in the FSHR-WT

S6 Fig. Root mean square fluctuations of Ca atoms calculated as function of block size for trajectories

water

LSRh   FSHR

Phosphate
head groups

TM2

TM1

TM6

Lipid
bilayer

TM3

TM4   TM5

TM7

intracellular side

Figure2

Figure3

Figure4

Figure5

Figure6

Figure7

helical restraints   no helical restraints

Figure8

Figure10

Figure11

Figure12

Figure1

Figure9