

## A direct fiber approach to model sclera collagen architecture and biomechanics

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**Short Title:** Directly modeling of sclera collagen fibers

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1 **Highlights**

2 Collagen fibers are the main load-bearing component of eye tissues.

3 Conventional sclera modeling ignores that fibers are long, interwoven and interact.

4 We demonstrate a direct fiber model with long, interwoven and interacting fibers.

5 Collagen fiber mechanical properties were estimated using inverse fitting.

6 The model captures simultaneously sclera fiber structure and macroscale mechanics.

## 7 **Abstract**

8 Sclera collagen fiber microstructure and mechanical behavior are central to eye physiology and  
9 pathology. They are also complex, and are therefore often studied using modeling. Most models  
10 of sclera, however, have been built within a conventional continuum framework. In this framework,  
11 collagen fibers are incorporated as statistical distributions of fiber characteristics such as the  
12 orientation of a family of fibers. The conventional continuum approach, while proven successful  
13 for describing the macroscale behavior of the sclera, does not account for the sclera fibers are  
14 long, interwoven and interact with one another. Hence, by not considering these potentially crucial  
15 characteristics, the conventional approach has only a limited ability to capture and describe sclera  
16 structure and mechanics at smaller, fiber-level, scales. Recent advances in the tools for  
17 characterizing sclera microarchitecture and mechanics bring to the forefront the need to develop  
18 more advanced modeling techniques that can incorporate and take advantage of the newly  
19 available highly detailed information. Our goal was to create a new modeling approach that can  
20 represent the sclera fibrous microstructure more accurately than with the conventional continuum  
21 approach, while still capturing its macroscale behavior. In this manuscript we introduce the new  
22 modeling approach, that we call direct fiber modeling, in which the collagen architecture is built  
23 explicitly by long, continuous, interwoven fibers. The fibers are embedded in a continuum matrix  
24 representing the non-fibrous tissue components. We demonstrate the methodology by modeling  
25 a rectangular patch of the posterior sclera. The direct fiber model presented incorporated  
26 specimen-specific fiber orientations derived from polarized light microscopy data of histological  
27 cryosections. The fibers were modeled using a Mooney-Rivlin model, and the matrix using a Neo-  
28 Hookean model. The fiber parameters were determined by inversely matching experimental equi-  
29 biaxial tensile data from the literature. After reconstruction, the direct fiber model orientations  
30 agreed well with the microscopy data both in the coronal plane (adjusted  $R^2=0.8234$ ) and in the  
31 sagittal plane (adjusted  $R^2=0.8495$ ) of the sclera. With the estimated fiber properties  
32 ( $C_{10}=5746.9$  MPa;  $C_{01}=-5002.6$ MPa, matrix shear modulus 200kPa), the model's stress-strain  
33 curves simultaneously fit the experimental data in radial and circumferential directions (adjusted  
34  $R^2$ 's 0.9971 and 0.9508, respectively). The estimated fiber elastic modulus at 2.16% strain was  
35 5.45GPa, in reasonable agreement with the literature. During stretch, the model exhibited  
36 stresses and strains at sub-fiber level, with interactions among individual fibers which are not  
37 accounted for by the conventional continuum methods. Our results demonstrate that direct fiber  
38 models can simultaneously describe the macroscale mechanics and microarchitecture of the  
39 sclera, and therefore that the approach can provide unique insight into tissue behavior questions  
40 inaccessible with continuum approaches.

## 41 **1. Introduction**

42 Collagen fibers are the principal load-bearing component of sclera (Boote et al., 2020; Girard  
43 et al., 2009b; Grytz et al., 2014a; Jan et al., 2017a; Pijanka et al., 2012), and thus play an  
44 important role on eye physiology and pathology (Coudrillier et al., 2012; Ethier et al., 2004; Pijanka  
45 et al., 2012; Summers Rada et al., 2006). This has motivated many studies aimed at  
46 understanding the role of sclera collagen microarchitecture on macroscale eye biomechanical  
47 behavior (Coudrillier et al., 2013; Girard et al., 2009a; Grytz et al., 2011).

48 Because of the complexity of sclera microstructure and the difficulty of accessing it directly for  
49 experimentation, numerical models have been widely developed and used for the studies  
50 (Coudrillier et al., 2013; Coudrillier et al., 2015b; Girard et al., 2009a; Girard et al., 2009b; Grytz  
51 et al., 2011; Hua et al., 2020; Sigal et al., 2004; Voorhees et al., 2018; Voorhees et al., 2017).  
52 The models have been formulated within a continuum mechanics framework in which collagen  
53 fiber architecture has been approximated using statistical distributions. For example, in a common  
54 approach, the scleral collagen microarchitecture is accounted for by collagen fiber “families”, each  
55 of which is described through the family preferred orientation and a “degree of alignment” around  
56 this preferred orientation (Coudrillier et al., 2013; Girard et al., 2009a; Grytz et al., 2011). The  
57 model parameters are often determined through an inverse fitting process by matching  
58 experimentally-determined sclera behavior under inflation tests. More advanced models have  
59 incorporated experimental data on collagen fiber orientations, obtained for example from small-  
60 angle light scattering or related scattering techniques (Coudrillier et al., 2013; Schwaner et al.,  
61 2020a; Zhang et al., 2015). Other recent models have focused on incorporating regional variations  
62 in fiber family characteristics (Kollech et al., 2019). The success of the conventional models at  
63 the macroscale encouraged their use to predict microstructural characteristics of sclera fibers,  
64 such as collagen fiber crimp (Girard et al., 2009a; Grytz et al., 2014a; Grytz et al., 2011). These  
65 models were helpful to understand the interplay between mechanics and structure, both micro  
66 and macro, and necessary because of the limitations of the experimental tools at the time.

67 Advances in the experimental tools have enabled a more comprehensive characterization of  
68 the sclera microstructure and mechanics (Behkam et al., 2019; Brown et al., 2007; Coudrillier et  
69 al., 2016; Hoerig et al., 2022; Jan et al., 2015; Lee et al., 2022a; Ling et al., 2019; Pijanka et al.,  
70 2019; Sigal et al., 2014; Winkler et al., 2010; Yang et al., 2018a; Yang et al., 2021). Studies using  
71 these tools have demonstrated structural and mechanical characteristics of the sclera that are  
72 potentially crucial yet are not accounted for by the conventional continuum approach for modeling,  
73 such as fiber interweaving, fiber-fiber interactions and the in-depth fiber orientation distributions

74 (Boote et al., 2020; Jan et al., 2017b; Lee et al., 2022a). Altogether this indicates that continuum  
75 models are limited in the ability to capture and describe sclera structure and mechanics at the  
76 smaller, fiber-level, scales. For instance, elsewhere we have shown that not accounting for  
77 interweaving and fiber-fiber interactions can introduce substantial errors when estimating sclera  
78 fiber mechanical properties using inverse modeling (Wang et al., 2020). This is also important  
79 because accurate predictions at the small scale are crucial if the intent is to use the models to  
80 understand effects at the scale of cells, axons and for studying mechanobiology.

81 To address the limitations of conventional models and better account for microstructure,  
82 models have been developed that explicitly incorporate collagen fiber networks (Hadi and  
83 Barocas, 2013; Islam and Picu, 2018; Licup et al., 2015; Picu et al., 2018; Zhang et al., 2013).  
84 The models, however, are limited in their ability to represent specimen-specific collagen  
85 architecture. In addition, the models were formed by short fibers, sometimes generated  
86 stochastically, and do not represent well the long fibers that form the sclera (Boote et al., 2020).  
87 We contend that long fibers can have fundamentally different mechanical behavior than short  
88 fibers (Voorhees et al., 2018), and thus that explicitly accounting for them is important in  
89 specimen-specific modeling of the eye. Altogether, the limitations of the current modeling tools  
90 highlight the need to develop more advanced modeling techniques that can incorporate detailed  
91 information on fibers.

92 Our goal was to demonstrate that it is possible to build a model of sclera that represents fiber  
93 microstructure better than the conventional modeling approaches and that also captures sclera  
94 macroscale behavior. In this manuscript we introduce a new modeling approach, that we call  
95 direct fiber modeling, in which the collagen architecture is accounted for by long, continuous,  
96 interwoven fibers. The fibers are embedded in a continuum matrix representing the non-fibrous  
97 tissue components. We demonstrate the methodology by modeling a rectangular patch of  
98 posterior pole sclera. First, we show that a specimen-specific direct fiber model can be built based  
99 on high resolution polarized light microscopy data of histological sections. We show that the fiber  
100 orientation distributions of the direct fiber model follow closely those of the histology  
101 simultaneously in both the coronal and sagittal planes. This is a more demanding requirement  
102 than in conventional models in which only the fiber orientations in coronal plane are accounted  
103 for (Girard et al., 2009b). Second, we used the model in an inverse modeling approach by fitting  
104 experimental biaxial stress-strain data from the literature. We show that the direct fiber model can  
105 match the experiment simultaneously in both radial and circumferential directions. Overall, direct  
106 fiber modeling can provide unique insight into the interplay between tissue architecture and

107 behavior and help answer questions that have been inaccessible with the conventional continuum  
108 approach.

## 109 **2. Methods**

110 This section is organized in three parts following the same general order as the process for  
111 building and using a direct fiber model. First, experimental data on sclera fibers and orientations  
112 was obtained using established PLM imaging of histological cryosections (Jan et al., 2015).  
113 Second, a direct fiber finite element model of a patch of posterior sclera was built based on the  
114 PLM-derived orientation data. The fiber architecture of the model was first built from the collagen  
115 fiber orientation data in the coronal sections, and then iteratively optimized to also match the  
116 orientation data from sagittal sections. The direct fiber model was then embedded in a matrix  
117 representing non-collagenous components. Third, the combined fiber and matrix model was used  
118 in an inverse fitting optimization process to match the models' simulated stress-strain behaviors  
119 with equi-biaxial test data from the literature (Eilaghi et al., 2010). This process produced  
120 estimates of the fiber mechanical properties. Below we describe these parts in detail.

121 Modeling was done in Abaqus 2020X (Dassault Systemes Simulia Corp., Providence, RI, 171  
122 USA). Customized code and the GIBBON toolbox (Moerman, 2018) for MATLAB v2020  
123 (MathWorks, Natick, MA, USA) were used for model pre/post-processing and inverse modeling.

### 124 **2.1 Histology, polarized light microscopy and fiber orientation quantifications**

125 The study was conducted in accordance with the tenets of the Declaration of Helsinki and the  
126 Association of Research in Vision and Ophthalmology's statement for the use of animals in  
127 ophthalmic and vision research. For the fiber orientation distribution in the coronal plane, we used  
128 a porcine eye that was also part of a previous study on sclera architecture (Gogola et al., 2018b).  
129 Details of the eye preparation, histological processing, PLM imaging and post-processing  
130 methods are reported elsewhere (Jan et al., 2015; Jan et al., 2017b). Briefly, a healthy eye without  
131 known abnormalities was obtained from the local abattoir and processed within 24 hours of death.  
132 The episcleral tissues, fat and muscles were carefully removed, and the globe was perfusion and  
133 immersion fixed in 10% formalin for 24 hours at an IOP of 0 mmHg. After fixation, the optic nerve  
134 head region was isolated with an 11 mm circular trephine and serially cryo-sectioned coronally  
135 with a slice thickness of 30  $\mu\text{m}$  (Figure 1 A). Seventeen serial sections were imaged with PLM  
136 using the 0.8x objective (NA 0.12) of an Olympus SZX16 microscope, paired with a dual chip  
137 Olympus DP80 camera (4.25  $\mu\text{m}/\text{pixel}$ ).

138 For the fiber orientation distribution in the sagittal plane, we used a healthy sheep eye that  
139 was processed in the same way, except that the optic nerve head was sectioned sagittally. Three  
140 sections through the middle of the scleral canal were selected and imaged with PLM using an

141 Olympus IX83 microscope with 4x objective (1.49 $\mu$ m/pixel). The higher resolution in this  
142 orientation was selected to resolve better the in-depth fiber interweaving. Fiber orientation  
143 distributions were normalized for use, and therefore we do not expect this to affect the  
144 reconstructions. Please see the discussion for a discussion of the potential consequences of  
145 having used different species for the coronal and sagittal planes.

146 PLM images were processed to derive at each pixel the in-plane collagen orientation (in  
147 Cartesian coordinates) and a parameter which we have previously referred to as “energy” (Yang  
148 et al., 2018b). Energy helps identify regions without collagen, such as outside of the section, and  
149 regions where the collagen fibers are primarily aligned out of the section plane, so that they can  
150 be accounted for in the orientation distribution.

151 For the coronal fiber orientation, the coronal sections were stacked sequentially and  
152 registered (Gogola et al., 2018b). After registration, the PLM data was reprocessed to align  
153 orientation values (Gogola et al., 2018b). As target to build the direct fiber model we selected a  
154 rectangular block of sclera, 2.00 x 1.91 mm in size, in the temporal side of the optic nerve head  
155 (Figure 2). The location and shape were chosen to match the sample tested experimentally  
156 (Eilaghi et al., 2010). We then calculated the distribution of collagen fiber orientations from the  
157 PLM images. We used the pixel-level data from 17 sections, weighted by the local “energy”. This  
158 allowed us to build the in-plane distribution based on fibers in the same plane (Yang et al., 2018b).

159 For the sagittal fiber orientation, we analyzed the images of the 3 sagittal sections. Naturally,  
160 the sagittal sections were not from the same sample as the coronal sections, and therefore we  
161 needed a different method to assign fiber orientations. It seemed inadequate to arbitrarily pick the  
162 orientation distribution of a region. Thus, instead, we derived a representative sagittal orientation  
163 distribution. We did this by manually defining 236 rectangular areas, 715  $\mu$ m x 715  $\mu$ m, all in the  
164 sclera near the posterior pole in regions approximating the location selected in the coronal plane.  
165 The orientation distributions were calculated for each rectangular area, again limiting this to the  
166 in-plane collagen fibers by using the out-of-plane information derived from the energy parameter  
167 (Yang et al., 2018b). We then derived the representative sagittal orientation distribution by  
168 averaging the distributions across the 236 regions. In addition, we also calculated the sagittal fiber  
169 anisotropy from the median of the anisotropies among the 236 areas. The anisotropy indicates  
170 the degree of fiber alignment and was quantified as circular standard deviation (Gogola et al.,  
171 2018b). Perfectly aligned fibers have an anisotropy of 1 and evenly dispersed fibers have an  
172 anisotropy of 0. We reasoned that the average distribution and median anisotropy provide a better  
173 representation of the orientation distribution in this direction than any single distribution.



## 174 **2.2 Direct fiber model construction**

### 175 **2.2.1 Fibers**

176 Fibers were simulated using 3-dimensional linear truss elements (T3D2 in Abaqus). Fiber  
177 locations were defined by Cartesian coordinates (X, Y, Z) of element nodes. To define fibers, we  
178 sampled orientation values from PLM images at regularly spaced “seed” points with a spacing of  
179 272  $\mu\text{m}$ . At each seed point, a straight fiber 8.5  $\mu\text{m}$  in diameter was traced in the section plane at  
180 an angle equal to the orientation at the seed point, meanwhile, the fiber passed through the seed  
181 point and extended the full span of the region. The process was applied for each layer, and then  
182 the fibers were stacked, resulting in a stack of 2D layers of fibers, each with a large number of  
183 fibers crossing or “interpenetrating”. To resolve fiber interpenetrations, an algorithm was used to  
184 refine and displace fiber elements in the whole structure (Matuschke et al., 2021; Matuschke et  
185 al., 2019). Briefly, if the smallest distance between two elements was less than the fiber diameter,  
186 an interpenetration was detected. Interpenetrated elements were shifted apart iteratively until all  
187 interpenetrations were resolved. During the process, fiber elements were re-meshed so that the  
188 element length was kept between 17 $\mu\text{m}$  to 25.5 $\mu\text{m}$ . Fibers were smoothed by controlling the fiber  
189 minimum radius of curvature.

190 It is important to make a few important notes regarding our use of the term “fiber”. The collagen  
191 of the sclera has a complex hierarchical structure even more complex than that of the cornea and  
192 tendon (Boote et al., 2020; Jan et al., 2018). The models were reconstructed using a fiber  
193 diameter of 8.5  $\mu\text{m}$ , which means that what we refer to as a fiber most likely represents a group  
194 of fibers that elsewhere may instead be described as a fiber bundle. Nevertheless, because our  
195 intent with this work is to call attention to the power of incorporating detailed microstructural  
196 information on sclera, we decided to use the term fiber, while also being careful to acknowledge  
197 at several critical points that these may be understood as fiber bundles. This is further addressed  
198 in the discussion.

199 To evaluate the similarity of the model and the images in the coronal plane, we quantified the  
200 model’s coronal orientation distribution and compared it with the distribution of the PLM images.  
201 We counted the occurrences of element orientations, where the orientation of an element was the  
202 slope angle in the coronal section plane and the number of occurrences was estimated as the  
203 element volume. This allowed us to account for uneven element sizes and to properly compare  
204 with PLM pixel-based measurements.

205 To better account for the fiber distribution in 3D, including the through-depth fiber undulations,  
206 we did the following: we shifted the fiber elements locally in-depth according to the fiber relative  
207 positions observed in coronal section images. If a fiber element had an orientation that disagreed  
208 with the orientation in the image by more than 45 degrees, the element was “shifted” in the  
209 posterior direction. A fiber segment that was in agreement with the local orientation was kept at  
210 the original sclera depth. After applying this process throughout the volume, we did another  
211 iteration resolving fiber interpenetrations. We then compared the in-depth fiber anisotropy of the  
212 model with the representative anisotropy in the sagittal plane. If the model had a larger in-depth  
213 anisotropy than the images, we shifted the fibers, increasing the fiber undulation amplitudes and  
214 decreasing the degree of fiber alignment. Conversely, if the model had a lower anisotropy than  
215 the sagittal images, we decreased the amplitude of fiber undulations. Each of these steps was  
216 followed by an iteration resolving fiber interpenetrations and a check in the agreement between  
217 the model and sample anisotropies in the coronal plane. Adjusted R-square (adjusted  $R^2$ ) values  
218 were used to evaluate the fitness of orientation distributions. This process converged to a set of  
219 continuous fibers in excellent agreement with the PLM-derived orientations simultaneously in both  
220 coronal and sagittal planes. Adjusted  $R^2$  is similar to the conventional  $R^2$  in that it indicates fit to  
221 a curve, but it adjusts for the number of points considered, avoiding potentially misleading  
222 excellent fits that are the result of many points of comparison (Miles, 2005).

## 223 **2.2.2 Matrix**

224 The model used for inverse fitting consisted of the direct fiber model and a coincident matrix  
225 with an overall dimension of 2.00 x 1.91 x 0.35 mm. The model of fibers and matrix had the same  
226 dimensions, and fiber end-nodes resided on the matrix surfaces.

## 227 **2.3 Model inverse fitting**

### 228 **2.3.1 Meshing and material properties**

229 Fibers were modeled as a hyperelastic Mooney-Rivlin material (Holzapfel, 2001):

$$230 \quad W = C_{10}(I_1 - 3) + C_{01}(I_2 - 3) + \frac{1}{2}K(J - 1)^2 \quad (1)$$

231 where  $W$  was the strain energy density,  $C_{10}$  and  $C_{01}$  were the material constants, restricted by  $C_{10}$   
232 +  $C_{01} > 0$  and would be determined by inverse modeling,  $I_1$  and  $I_2$  were the first and second  
233 invariants of the right Cauchy-Green deformation tensor,  $K$  was the bulk modulus and  $J$  was the  
234 determinant of the deformation gradient. As noted above, the fiber element diameter was  
235 assumed to be 8.5 $\mu$ m, and the element size was between 17 $\mu$ m to 25.5 $\mu$ m. The matrix was

236 meshed with linear eight-noded, hybrid hexahedral elements (C3D8H in Abaqus) and modeled  
237 as a neo-Hookean material with a shear modulus of 200 kPa (Coudrillier et al., 2015a; Girard et  
238 al., 2009b).

### 239 **2.3.2 Interactions**

240 Fiber-fiber interactions were simulated by preventing fiber interpenetrations using Abaqus'  
241 general contact with no friction. Preventing fiber interpenetration is computationally expensive  
242 and therefore commercial implementations are often directed at maintaining small  
243 interpenetrations rather than avoiding them altogether. We reasoned that a small interpenetration,  
244 in low single digits percent of fiber diameter, will likely not represent a major deviation from fiber  
245 physical behavior. We calculated the extent of the interpenetration in the following way: we  
246 defined a deformed model based on the original model and the final nodal displacements. A  
247 customized MATLAB program was then used to find all existing interpenetrated element pairs.  
248 For each of these pairs, we quantified the percentage of interpenetration as:

$$249 \text{percentage of interpenetration} = \frac{d - x_{\min}}{d} \quad (2)$$

250 where  $d$  was fiber diameter, and  $x_{\min}$  was the minimum distance between two elements.  
251 Interpenetration thus varied between 0 when there is no interpenetration and 1 when two  
252 elements are perfectly overlapping. We found that 99.4% of the interpenetrations were smaller  
253 than 5%. This is further addressed in the discussion.

254 Fiber-matrix interactions were ignored, as is usual in biomechanical models of the eyes (Grytz  
255 et al., 2014a; Grytz and Meschke, 2009; Petsche and Pinsky, 2013).

### 256 **2.3.3 Finite element analysis procedure**

257 The fiber-matrix assembly was subjected to a quasi-static process of equi-biaxial stretch to  
258 match the experimental results and boundary conditions reported elsewhere (Eilaghi et al., 2010).  
259 The matrix was simulated using Abaqus standard implicit procedure. Due to the presence of  
260 complicated fiber contacts, the direct fiber model was simulated using Abaqus dynamic explicit  
261 procedure, meant to improve the convergence and computational efficiency. The model fiber  
262 volume fraction (VF) is 7%. The resulting stresses  $\sigma$  along radial and circumferential directions  
263 were contributed by both matrix and fibers, where the matrix stress was weighted by fiber VF:

$$264 \sigma = (1 - VF)\sigma_{\text{matrix}} + \sigma_{\text{fibers}} \quad (3)$$

265 Because the process was quasi-static, the dynamic explicit procedure would require an  
266 excessive number of small-time increments, which is computationally impractical. Therefore, to

267 run the dynamic analysis efficiently, mass scaling was implemented. We achieved a  $1e-5$  stable  
268 time increment, where we modeled the process in the shortest time period in which inertial forces  
269 remain insignificant. During the simulation, we assured that the inertial effects were negligible by  
270 keeping the kinetic energy less than 5% of the internal energy.

#### 271 **2.3.4 Boundary conditions and inverse modeling procedure**

272 The model was simulated iteratively to derive fiber material properties ( $C_{10}$  and  $C_{01}$ ) by  
273 inversely matching published experimental stress-strain data (Eilaghi et al., 2010).

274 **Selecting the target experimental set:** The fibers in the direct fiber model were not equally  
275 distributed in all directions. It seems reasonable to expect that this structural anisotropy will result  
276 in mechanical anisotropy (Coudrillier et al., 2015b). It was therefore important to select for inverse  
277 fitting, among the multiple experimental results reported in the literature (Eilaghi et al., 2010), the  
278 one that exhibited a mechanical anisotropy matching the structural anisotropy of our model. To  
279 do this, we first quantified the mechanical anisotropy of our model (Figure 4). We did this by  
280 subjecting the model to displacement-controlled equi-biaxial stretch tests. At each strain level we  
281 quantified the model's stress anisotropy as the ratio of stresses in the orthogonal directions  
282 ( $S_{11}/S_{22}$ ). The test was repeated several times with different fiber material property values. The  
283 results showed that the ratio of stresses had a fairly constant distribution and was essentially  
284 independent to the material properties changes within the ranges tested. With this, we were able  
285 to select 76940-TS as the case with similar anisotropy as our model from the several experimental  
286 results reported by Eilaghi et al..

287 To be consistent with the selected experimental data, we assigned an equi-biaxial stretch of  
288 2.16% as the displacement boundary condition to our fiber-matrix assembly. To optimize the fiber  
289 material properties ( $C_{10}$  and  $C_{01}$ ), we used the simplex search method of Lagarias et al. (Lagarias  
290 et al., 1998). The algorithm sought to identify the two parameters that yielded the closest match  
291 between the simulated and experimental stress-strain curves, simultaneously in both directions.  
292 We optimized the fitness by minimizing the cost function, defined as the residual sum of squared  
293 (RSS) between the model and experimental curves at 21 strain states. The optimization was  
294 terminated when the cost function value was smaller than 0.01. After optimization we also  
295 computed the adjusted  $R^2$  between the curves to assess curve similarity. The complete inverse  
296 fitting process was repeated 11 times with various starting parameters to test the consistency and  
297 "uniqueness" of the results.

298 For interpretation of the results and to compare with other studies, it is useful to “convert” the  
299 optimal  $C_{10}$  and  $C_{01}$  parameters of the hyperelastic Mooney-Rivlin model into more intuitive  
300 representations of fiber mechanical properties. Specifically, we estimated the fiber shear modulus  
301 as  $2(C_{10} + C_{01})$ . We also derived a fiber elastic modulus by simulating uniaxial stretch of a single  
302 straight fiber with the optimal  $C_{10}$  and  $C_{01}$  values and the hyperelastic Mooney-Rivlin material. The  
303 fiber elastic modulus was then obtained as the slope of the stress-strain curve at 2.16% strain.

### 304 3. Results

305 After construction, the model consisted of 1016 fibers (or fiber bundles). Fiber orientation  
306 distributions and overall anisotropies were similar between the model and the tissue, as measured  
307 using PLM (Figure 5). In the sagittal plane, the anisotropies were 0.6031 and 0.5976 for the model  
308 and tissue, respectively (a difference smaller than 1%); In the coronal plane, the anisotropies  
309 were 0.1777 and 0.1626 for the model and tissue, respectively (a difference smaller than 10%).

310 Figure 5 shows the fiber orientation distributions of the direct fiber model and PLM images.  
311 The match of distributions was achieved simultaneously in both coronal and sagittal planes. By  
312 Wilcoxon rank sum tests, there were no significant differences in fiber orientation distributions  
313 between the model and the PLM images, in either coronal ( $p > 0.7$ ) or sagittal ( $p > 0.6$ ) planes.

314 Figure 6 illustrates fiber displacements and stresses at half and full stretch. The large  
315 differences in stress between fibers were consistent with the anticipated process of stretch-  
316 induced recruitment (Grytz and Meschke, 2010; Jan and Sigal, 2018). All 11 optimization runs led  
317 to curves that were in good agreement with the experiment. They all led to fairly consistent  
318 estimates of fiber shear modulus, with an average of 1522.8 MPa and a standard deviation  
319 (STDEV) of 38.9 MPa (Supplementary Table 1). Interestingly, these were obtained with  $C_{10}$  and  
320  $C_{01}$  coefficients that varied substantially. For example,  $C_{10}$  ranged from 118.9 to 7881.8 MPa.

321 Figure 7 shows that the stress-strain curves of the optimal model fit very well the experimental  
322 data in both radial (adjusted  $R^2 = 0.9971$ ; RSS = 0.00095) and circumferential directions (adjusted  
323  $R^2 = 0.9508$ ; RSS = 0.0098) simultaneously. The material parameters for the inverse fitting run  
324 that led to the closest model and experiment fit were:

$$325 \quad C_{10} = 5746.9 \text{ MPa}$$

$$326 \quad C_{01} = -5002.6 \text{ MPa}$$

327 The estimated fiber elastic modulus at 2.16% strain was 5.45GPa. The estimated fiber shear  
328 modulus is 1488.6 MPa.

## 329 4. Discussion

330 Our goal was to develop a direct fiber modeling approach that captures both the microscale  
331 collagen fiber architecture and the macroscale mechanical behavior of the sclera. We  
332 demonstrated our approach by modeling a rectangular patch of the posterior sclera. The model  
333 incorporated several fiber characteristics ignored by most previous models of posterior sclera,  
334 such as fiber interweaving, fiber-fiber interactions, long fibers, and a physiologic experimentally-  
335 derived in-depth fiber orientation distribution. Below we discuss in detail the importance of  
336 considering these characteristics for the study of sclera biomechanics.

337 A common preconception about the effects of fiber interweaving, and the resulting fiber-fiber  
338 interactions, seems to be that interweaving increases the stiffness of a material, or tissue.  
339 However, by comparing models with interweaving and non-interweaving fibers, we have shown  
340 that an interwoven architecture is, as a structure, more compliant than a non-interwoven  
341 architecture (Wang et al., 2020). This is consistent with the literature on textile mechanics (Saiman  
342 et al., 2014; Stig and Hallström, 2019). This effect can be understood by noting that the fibers of  
343 the interweaving architecture are undulated, whereas those from the non-interweaving models  
344 are straight. Straight fibers aligned with the load carry the forces more efficiently, and are shorter,  
345 and thus, the overall model is stiffer. Since fiber interweaving and the resulting fiber-fiber  
346 interactions play an important role in determining the structural stiffness of the sclera, it should be  
347 considered when modeling sclera biomechanics, as did our direct fiber model. Note that we are  
348 not the first to recognize the importance of fiber interweaving and fiber-fiber interactions in soft  
349 tissue biomechanics (Elliott and Setton, 2001; Guerin and Elliott, 2007; Nerurkar et al., 2011;  
350 Wagner and Lotz, 2004; Zhang et al., 2013). For example, in annulus fibrosus of intervertebral  
351 disc, interlamellar shearing can account for nearly 50% of the total stress associated with uniaxial  
352 extension. Therefore, interweaving collagen fiber layers may play an important role in annulus  
353 fibrosus tissue function.

354 It is important to note that the scale of the undulations of interweaving is different from that of  
355 the collagen fiber waviness referred to as crimp by us, (Gogola et al., 2018a; Jan et al., 2018; Jan  
356 et al., 2017a; Jan and Sigal, 2018) and others (Grytz et al., 2014a; Grytz et al., 2020a; Grytz and  
357 Meschke, 2009). For example, ocular collagen fibers crimp has a period typically under 20  $\mu\text{m}$ ,  
358 (Gogola et al., 2018a; Jan et al., 2018; Jan et al., 2017a; Jan and Sigal, 2018) much smaller than  
359 the estimates of interweaving undulations on the order of 100 to 300  $\mu\text{m}$  (Wang et al., 2020).  
360 From a mechanical perspective, the undulations of crimp affect directly the biomechanical  
361 behavior of a fiber under load, with the crimp “disappearing” as the fiber is loaded or stretched,

362 and eventually recruited (Jan et al., 2022; Jan and Sigal, 2018). The interactions between  
363 interwoven fibers mean that these undulations do not “disappear” under load, and are limited by  
364 the interlocking (Lee et al., 2022b).

365 The collagen fibers of the sclera are long and continuous; thus, they can transfer forces over  
366 a long distance (Boote et al., 2020; Voorhees et al., 2018). However, the conventional continuum  
367 approach to modeling sclera only considers the local or regional orientation distribution of these  
368 long fibers and maps such information into fairly small “finite elements” (Coudrillier et al., 2013;  
369 Grytz et al., 2011). As a result, the force transmission is continuous across elements rather than  
370 along fibers, which alters the predicted mechanical behavior locally and potentially at the  
371 macroscale (Campbell et al., 2015; Coudrillier et al., 2013; Kollech et al., 2019; Roberts et al.,  
372 2010; Voorhees et al., 2018; Zhang et al., 2015; Zhou et al., 2019). We are not the first to  
373 recognize the importance of accounting for long fiber mechanics (Grytz et al., 2020b; Huang et  
374 al., 2017; Lanir, 2017). One approach proposed to address the problem has been through the use  
375 of symmetry boundaries across elements. In this manner it is possible to simulate fibers that are  
376 continuous across elements. The approach, however, substantially limits the type of models that  
377 can be considered compared with the direct fiber model we present.

378 The collagen architecture of the sclera varies in depth (i.e., the direction perpendicular to the  
379 scleral surface) (Jan et al., 2017b; Pijanka et al., 2015). Such variations are crucial in the load-  
380 bearing capacity of the sclera (i.e., bearing shear stresses), and may even have clinical  
381 implications (Danford et al., 2013). Unfortunately, most fiber-aware models of the sclera only  
382 account carefully for the fiber orientations within the scleral plane, while the in-depth orientations  
383 are ignored or are modeled in a much more simplified manner than the out of plane orientations  
384 (Coudrillier et al., 2013; Voorhees et al., 2017; Zhang et al., 2015). Our direct fiber model  
385 incorporates both in-plane and in-depth specimen-specific fiber orientation distributions that  
386 match those measured using PLM, and thus, has a higher fidelity for representing the sclera and  
387 its mechanics.

388 Ultimately, a model is only as good as the predictions that can be made using it. On this, we  
389 point out that the fiber elastic modulus estimated with our direct fiber model is well within the range  
390 of values reported in the literature. For example, we estimated a fiber elastic modulus of 5.45 GPa,  
391 compared with those of 2-7 GPa of the bovine Achilles tendon (Van Der Rijt et al., 2006; Yang et  
392 al., 2007) and 5-11.5 GPa of the rat tail (Wenger et al., 2007). In contrast, the estimates of fiber  
393 elastic modulus derived using continuum models, which used constitutive model formulations  
394 based on highly simplified assumptions of fiber architecture and behavior at the microscale level,



395 are several orders of magnitude smaller, between 1 and 200 MPa (Coudrillier et al., 2015b; Grytz  
396 et al., 2014b; Schwaner et al., 2020a; Schwaner et al., 2020b). Hence, whilst both continuum and  
397 direct fiber models can closely approximate the macroscale sclera behavior, the estimates of fiber  
398 mechanical properties derived from both types of models can be substantially different. This may  
399 not be a problem when the intent is limited to describing tissue mechanical response. However,  
400 an important application of continuum models is to “infer” characteristics of the underlying fibers.  
401 The large differences in fiber properties estimated by the continuum models and the experimental  
402 measurements already suggest that these should be interpreted very carefully as the values  
403 inferred may be inaccurate. A common application of this approach is to compare the fiber  
404 properties inferred from healthy and unhealthy tissues, or from young and old donors. The  
405 argument in this case is based on the idea that the methods may not produce accurate fiber  
406 estimates, but that the comparison remains valid. This may indeed be the case. But it is also  
407 possible that if the changes with pathology or age involve aspects of the fibers that are not  
408 accounted for in the continuum model formulation, then the origin of the changes will end up  
409 artefactually attributed to another tissue characteristic. Thus, while all models involve important  
410 approximations, if the goal is to derive estimates at the fiber scale, we posit that direct fiber models  
411 are preferable over continuum models.

412 The methodology to build the direct fiber model structure can be applied to other collagenous  
413 tissues. The reconstruction method is not tied to PLM and can instead be done using second  
414 harmonic imaging or confocal microscopy, as long as they provide information on the density and  
415 orientation of collagen fibers. The level of detail necessary from the images can depend on the  
416 complexity of the tissue in question and the accuracy needed from the reconstruction and  
417 simulations. Because the sclera has a complex 3D architecture, we combined histological  
418 information from two sectioning directions. This may not be necessary for other tissues, or  
419 perhaps it may be possible to obtain depth information by taking advantage of the confocal nature  
420 of the imaging or 3D PLM (Yang et al., 2018b).

421 A common concern with inverse fitting is the uniqueness of the so-called optimal parameters  
422 (Girard et al., 2009a; Girard et al., 2009c; Zhou et al., 2019). When we repeated the optimization  
423 with various starting parameters, we obtained fairly consistent fiber shear modulus predictions.  
424 However, the material model parameters  $C_{10}$  and  $C_{01}$  varied substantially. This was not a surprise  
425 given the material formulation we used and the clear potential for interactions between the  
426 parameters. Altogether, this reinforces the importance of focusing on parameters with a clear

427 physical interpretation, in this case fiber shear modulus. Such parameters are more likely to have  
428 better stability.

429 From a numerical perspective, prevention of fiber penetrations is computationally intensive.  
430 We took advantage of highly mature general contact tools implemented in commercial code to  
431 keep solution time reasonable. Ensuring absolutely no interpenetration would have required an  
432 impractical number of small elements to represent the undulating fibers. To avoid this problem,  
433 we decided to quantify and track the interpenetrations and consider a solution valid if these  
434 remained below 5% of the fiber diameter. In a worst-case scenario this would mean that two  
435 interacting fibers have “pushed” into each other such that the distance between their centers is  
436 only 95% of twice their radii. We assumed that the fibers are all circular and remain circular,  
437 despite pushing into each other.

### 438 **Limitations**

439 When interpreting the results from this work, it is important to consider also the following  
440 limitations. First, we considered highly simplified fiber-fiber interactions and ignored fiber-matrix  
441 interactions. We are aware that the forces between fibers may be much more complex, including  
442 friction, crosslinks and several other physical processes that are challenging to simulate. Some  
443 of these may have major impact on the results, and others may not. This likely depends on the  
444 specific structure and loading conditions of the tissue. We carried out a friction sensitivity test in  
445 which we repeated the inverse fitting while using several friction coefficients. The results were  
446 very similar in stress-strain (less than 5% difference). We thus conclude that in the particular test  
447 reported herein ignoring friction did not adversely affect the accuracy of the model estimates  
448 substantially. Fiber-matrix interactions are potentially even more complex than fiber-fiber  
449 interactions given the wide diversity of components that form what we describe simply as “matrix”.  
450 These include non-fibrous components, such as GAG chains, that may act as lubricants and affect  
451 sliding, in which case their presence could make the tissue more compliant (Hatami-Marbini and  
452 Pachenari, 2020; Murienne et al., 2016). Other components include elastin fibers and cells. While  
453 we acknowledge the highly simplified fiber-fiber and fiber-matrix interactions we used in this study,  
454 we argue that the methodology we have described can be extended to incorporate much more  
455 complex interactions. Conventional continuum modeling techniques of the eye and other tissues  
456 simply ignore fiber-fiber and fiber-matrix interactions, yet it is not obvious to readers that these  
457 assumptions have been made, or how to relax them.

458       Second, we did not account for sub-fiber level features, such as collagen fiber crimp and  
459 various fiber or bundle diameters. Future work would benefit from introducing collagen fiber crimp  
460 (Grytz et al., 2014a) and various fiber diameters (Komai and Ushiki, 1991).

461       Third, our direct fiber model does not account for the very large number and complexity of  
462 fibers of the sclera. We posit that the fiber-level mechanics presented using the direct fiber model  
463 with reduced fiber density will likely be similar to the actual tissue. Therefore, our objective of  
464 introducing direct fiber model for simulating sclera is fulfilled. Meanwhile, it's beneficial to use  
465 fewer fibers to reduce computational cost. We agree that in future work, it will be worthwhile to  
466 work on improving computational efficiency and building the model with more fibers.

467       Fourth, we did not account for regional variations. In our model structure we did not observe  
468 the strong well-aligned region of fibers in the radial direction that have been reported to take about  
469 10% of the sclera nearest to the choroid. This could be due to variations between eyes, regional  
470 inhomogeneities (the radial fibers appear to be more readily distinguished closer to the optic nerve  
471 head) or due to the specific coronal sections used for the reconstruction. It is likely that other  
472 regions and other eyes will produce different models that respond differently to mechanical  
473 loading. We posit that our goal in this manuscript was modest, aiming to demonstrate it is possible  
474 to build a fiber-based model that produces a macroscopic mechanical behavior matching  
475 experimental tests. It will be beneficial to consider more features and variations in future work.

476       Fifth, the model used to demonstrate the direct fiber method was reconstructed from PLM and  
477 experimental data from different species. We posit that the impact was minor because our goal  
478 was to demonstrate the reconstruction and modeling methodology. It should be clear to readers  
479 that the process could be redone in other species. In this sense the key result from this work, that  
480 the direct fiber modeling technique can work still stands. In future work, it would be beneficial to  
481 have the image and experimental data consistent, i.e., from the same species, to better estimate  
482 the fiber mechanical properties of a certain species.

483       Sixth, the matrix mechanical properties were kept constant at literature values and not  
484 optimized iteratively like the fiber properties. This was for the sake of simplicity. Introducing matrix  
485 properties changes to the model may be necessary to properly account for the effects of age and  
486 interweaving, given the evidence that changes in matrix properties could result in age-related  
487 changes in sclera properties (Grytz et al., 2014b).

488       Seventh, although our model has similar boundary conditions to the experiment it is  
489 impossible to match the experiment precisely. Holding the tissues with clamps, rakes or hooks

490 alters the boundary conditions, in ways that are extremely complicated to replicate  
491 computationally. Moreover, the experiments to which we fit the model were biaxial stretch tests.  
492 This is not the physiologic mechanical condition of sclera. The mechanical behavior of sclera  
493 under inflation or more complex modes could be different.

## 494 **5. Conclusion**

495 We have shown the possibility of developing specimen-specific direct fiber model of sclera  
496 that can represent the sclera fibrous microstructure better than the previous continuum modeling  
497 approaches and allow accurate capture of sclera mechanics. We successfully built the model with  
498 long, continuous, interwoven fibers that takes into account the effects of fiber interweaving and  
499 fiber-fiber interactions. Our results have demonstrated that the direct fiber model can match the  
500 fiber orientations measured in high-resolution PLM images simultaneously in coronal and sagittal  
501 planes. The model properties can be optimized through inverse fitting to match experimental  
502 stress-strain responses. The estimated fiber elastic modulus is in good agreement with the  
503 literature. The direct fiber modeling methodology potentially has broad application to simulate  
504 other fiber-based tissues. Overall, the direct fiber modeling technique in this study is important for  
505 characterizing sclera collagen architecture at the fiber level, analyzing microstructural responses  
506 to macroscale mechanical loadings, and for understanding the scleral biomechanical environment.

507

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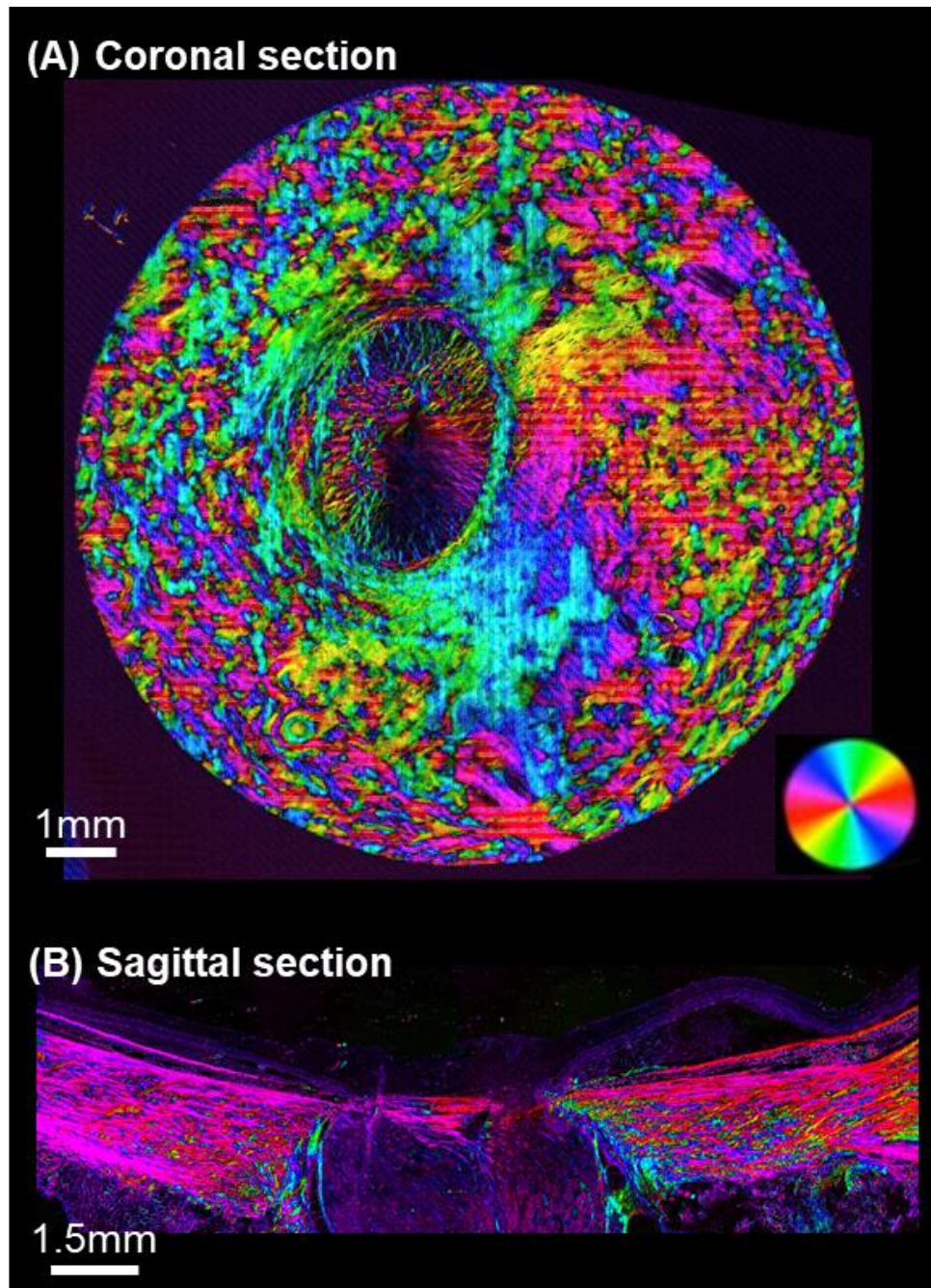
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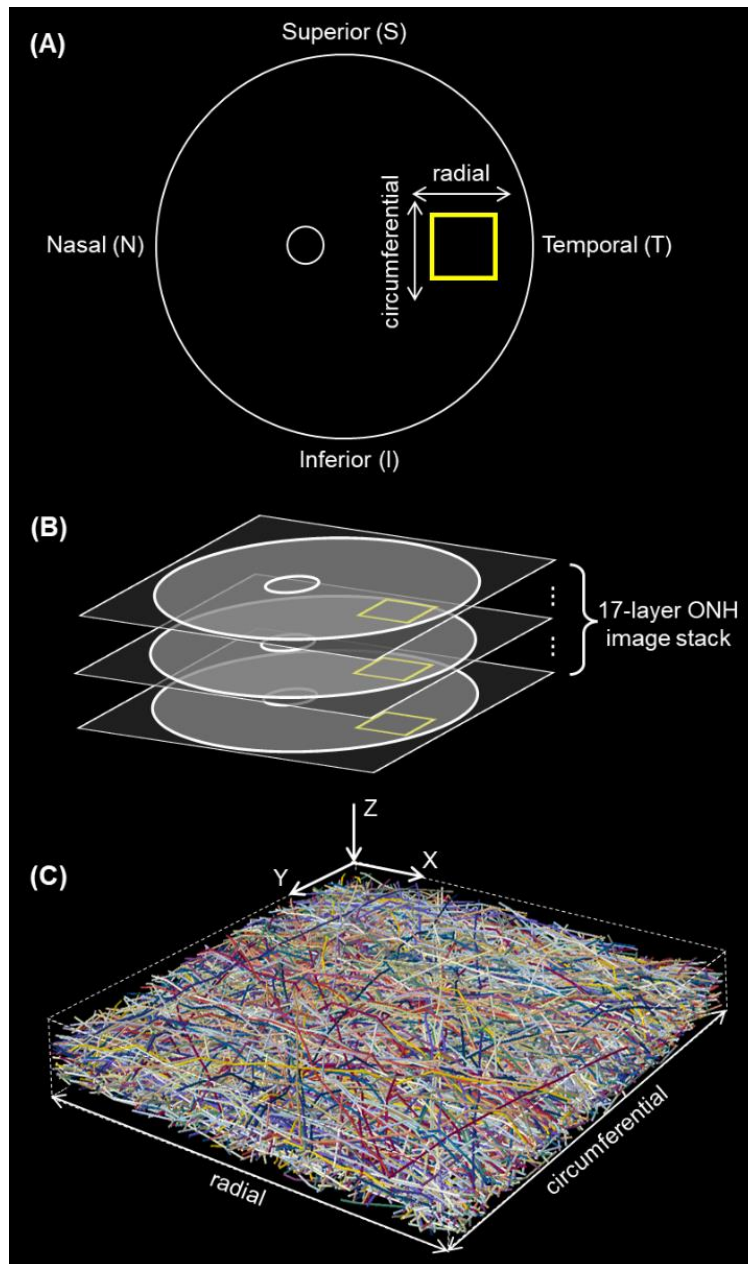
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740 **Figure 1. (A)** Example PLM image of a coronal section of a pig eye through the lamina cribrosa.

741 **(B)** Example PLM image of a sagittal section of a sheep eye through the optic nerve head (ONH).

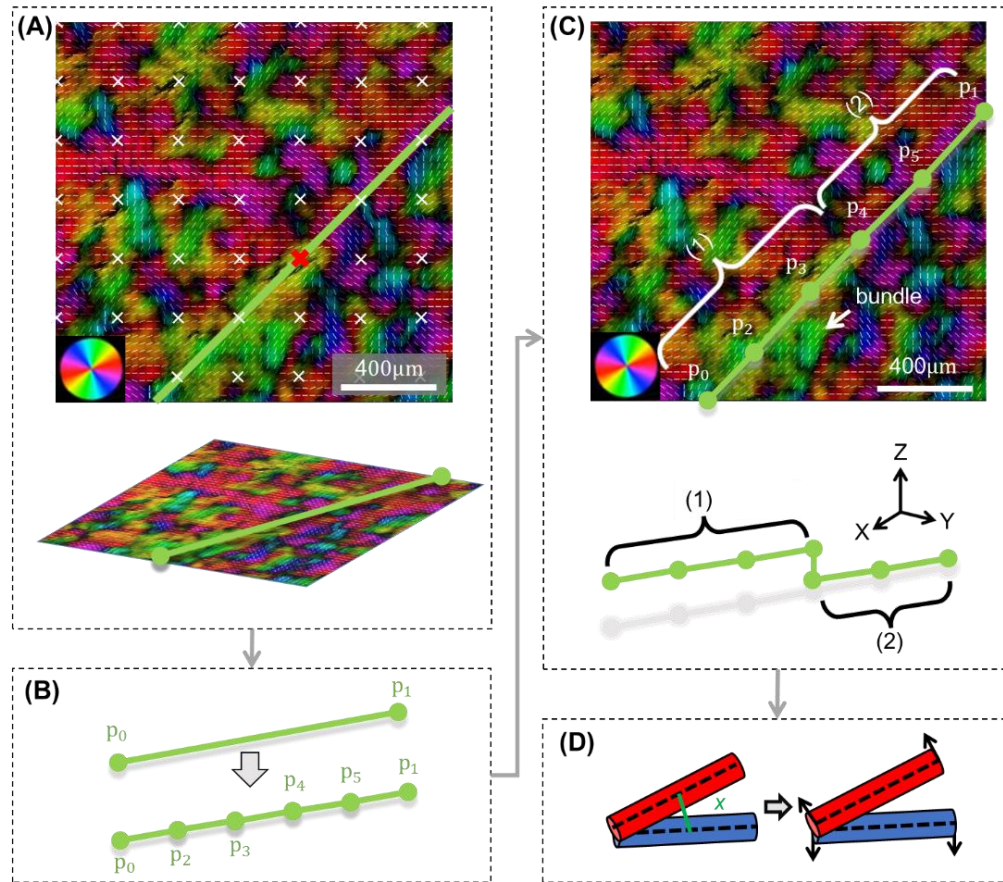
742 The colors indicate the local fiber orientation in the section plane, and the brightness the “energy”

743 parameter (see main text).



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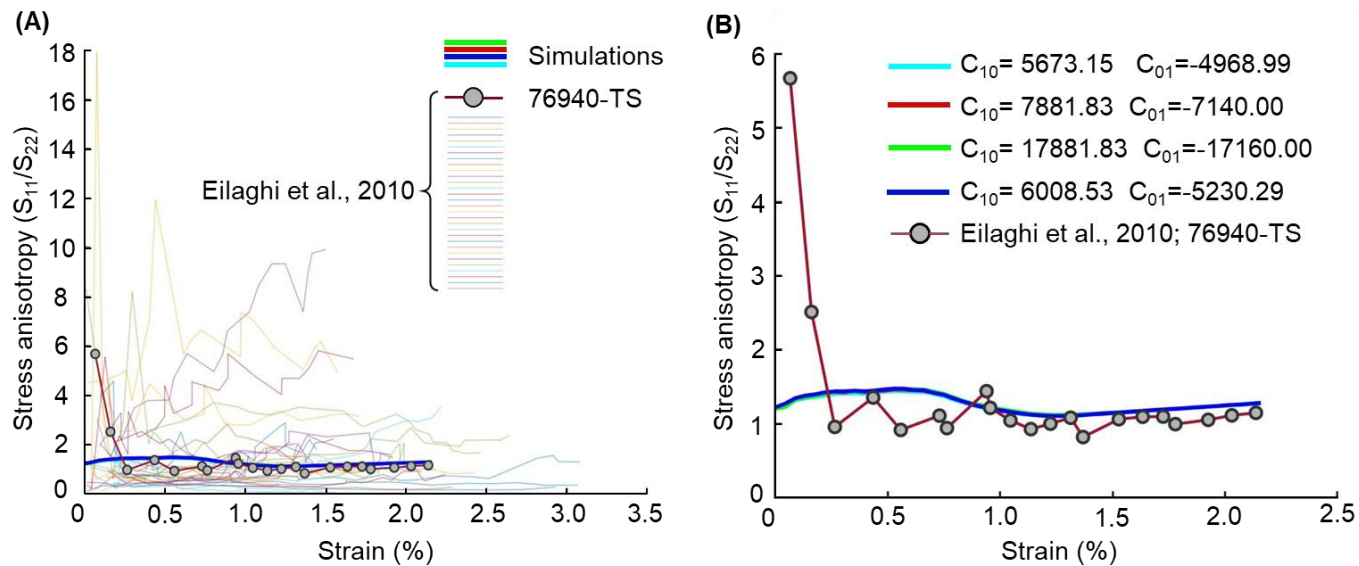
745 **Figure 2. (A)** Schematic coronal view illustrating the location of the scleral region modeled (yellow  
746 box, 2.00 x 1.91 mm). The region was located at the temporal sector. Also shown are the radial  
747 and circumferential directions, which correspond with the directions used for the equi-biaxial  
748 testing in the experiment and simulation. **(B)** A set of 17 coronal sections were stacked and  
749 registered. Fiber orientation data was extracted from the selected rectangular patch of sclera  
750 region (yellow box in panels A and B) in the image stack and used to build the direct fiber model  
751 **(C)**. Fibers, or fiber bundles, are shown in random colors to simplify discerning the complex  
752 interwoven architecture.



753

754 **Figure 3.** Workflow for creating and processing fibers. **(A)** The process begins with the pixel-level  
755 data from the coronal section PLM process (both orientation and energy). For clarity here we  
756 illustrate the process with a small square region. The colors of the PLM image represent the local  
757 collagen orientation. To help discerning orientations the images are shown with overlaid short  
758 white line segments representing the mean orientation over a small square region. (top panel) A  
759 regular grid of "seed" points was defined (white x marks). At these seed points the local fiber  
760 orientation was sampled, averaged as per the white lines to avoid noise. The energy information  
761 was used to skip defining fibers in regions without a reliable well-defined orientation, and to give  
762 preference to the in-plane fiber orientations over the out-of-plane ones. The in-plane orientation  
763 was then used to define a straight fiber in the image plane. In the example case, a fiber was traced  
764 at the grid point (red x mark) at an angle around 45 degrees. The bottom panel shows an isometric  
765 view of the image with the fiber overlaid. **(B)** (top) The fiber was initially defined by one element  
766 and two nodes,  $p_0$  and  $p_1$ . (bottom) The fiber was then refined into five elements and six nodes.  
767 **(C)** The average orientations over the fiber elements were then computed and compared with  
768 local image orientations to determine their agreement element by element. In the example shown,  
769 the element group (1),  $p_0$ - $p_2$ ,  $p_2$ - $p_3$ ,  $p_3$ - $p_4$ , agreed well with the direction of a yellow-green bundle.

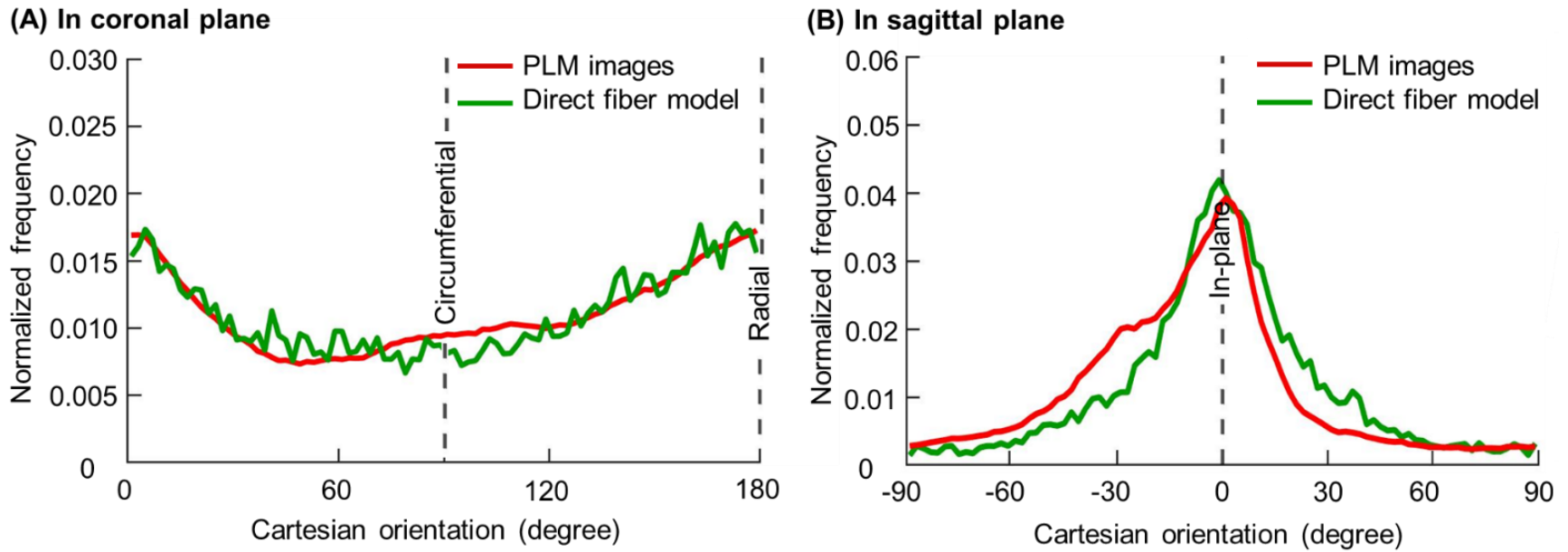
770 Accordingly, the group was accepted at this section. In contrast, the element group (2)  $p_4$ - $p_5$ ,  $p_5$ -  
771  $p_1$ , had poor agreement between element and image orientations. This group was then assigned  
772 a lower depth, effectively “pushing” the group or fiber segment to the depth of another section.  
773 Fiber connectivity was ensured by adding an element to connect the two element groups ((1) and  
774 (2)) at different depths. The fiber smoothness was restored when resolving fiber collisions.  
775 Meanwhile, elements were combined or split to maintain all element lengths within a pre-defined  
776 range. **(D)** Elements were moved apart if the shortest distance  $x$  between them was smaller than  
777 fiber diameter, indicating a collision.



778

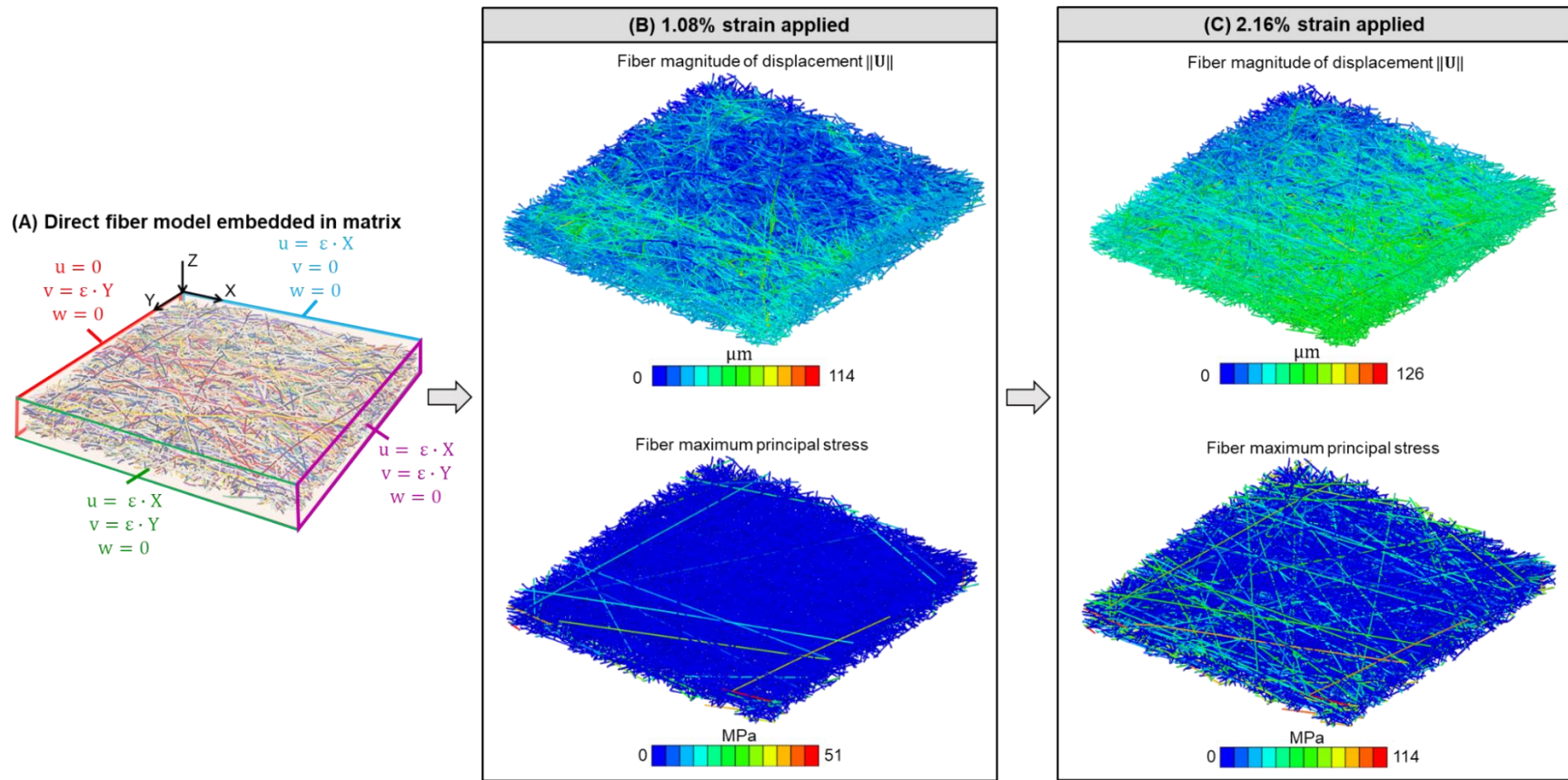
779 **Figure 4. (A)** To properly fit the model to an experimental curve it was necessary to select out of all the experimental curves the set  
 780 with a similar anisotropy. See text for more details. To do this we first assigned four sets of random material parameters,  $C_{10}$  and  $C_{01}$ ,  
 781 to the fiber model. The intent was to generate a set of model behaviors spanning the range of anisotropies to be expected over a wide  
 782 range of material properties. Equi-biaxial stretch was applied to the fiber-matrix assembly and the stress anisotropy ( $S_{11}/S_{22}$ ) of the  
 783 model computed at many steps through the stretching process. We then also computed the stress anisotropy of the experimental  
 784 curves reported by Eilaghi et al. There were substantial variations in their anisotropies. Out of all of the curves we selected 76940-TS  
 785 because it had the closest agreement with the model anisotropy. The plot shows all experimental curves with 76940-TS highlighted.  
 786 There was good agreement in anisotropy after about 0.25% strain. At smaller strains the experiment observed higher anisotropy,  
 787 potentially due to the challenge of balancing the initial loads in the clamping when using rakes. **(B)** In this plot only the four models and  
 788 the selected experimental curve. All  $C_{10}$  and  $C_{01}$  values are in MPa. The results revealed that the model stress anisotropy was  
 789 essentially independent of the fiber material properties. This can be discerned from the observation that the lines representing the four  
 790 simulations are almost indistinguishable. This means that the process of fitting fiber material properties will preserve the stress  
 791 anisotropy, allowing a close match of the stress-strain data simultaneously in both radial and circumferential directions.

## Fiber orientation distributions



792

793 **Figure 5.** Fiber orientation distributions of the direct fiber model (green lines) and the PLM images (red lines), in the **(A)** coronal and  
794 **(B)** sagittal planes. For the coronal plane the PLM orientation was derived from the stack of 17 images. In the coronal plane the radial  
795 direction corresponds to 0 and 180 degrees and the circumferential direction with 90 degrees. For the sagittal plane, the PLM orientation  
796 shown is the average orientation distribution over the 236 regions analyzed. See main text for details. Frequencies were normalized  
797 by the total sum of frequencies. Overall these results show that fiber orientation distributions of the direct fiber model agreed well with  
798 those from the PLM images in both coronal (adjusted  $R^2 = 0.8234$ ) and sagittal (adjusted  $R^2 = 0.8495$ ) planes.



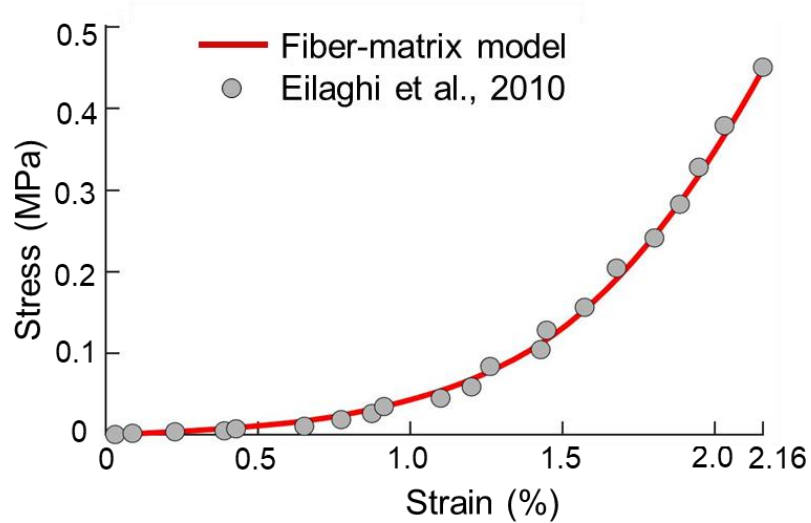
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800 **Figure 6. (A)** Schematic of the displacement boundary conditions for the fibers and matrix to represent scleral tissue subjected to 2.16%  
 801 equi-biaxial strain ( $\epsilon$ ).  $\mathbf{U}$  is the nodal displacement vector of fiber with components  $u$ ,  $v$  &  $w$  representing displacement in X, Y and Z  
 802 direction, respectively.  $\|U\|$  is the fiber magnitude of displacement. In panel A the fibers are randomly colored to facilitate discerning  
 803 their architecture. Panels **B** and **C** show isometric views of the direct fiber model with the fibers colored according to the magnitudes  
 804 of displacement (top row) or maximum principal stress (bottom row). The model is shown when subjected to **(B)** 1.08% strain or **(C)**  
 805 2.16% strain. From the images it is clear that the model fibers exhibit complex loading patterns.

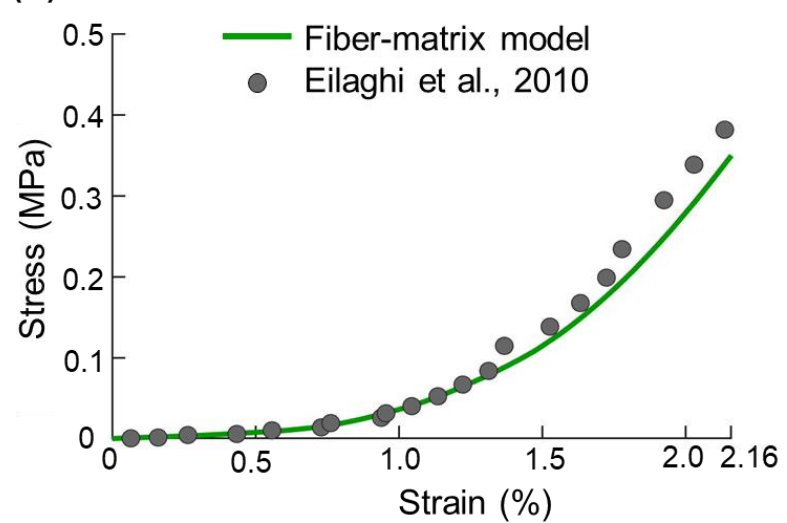


## Fit of stress-strain curves

(A) Radial direction



(B) Circumferential direction



806

807 **Figure 7.** Fit of stress-strain response was achieved between the model and the experiment simultaneously in both the (A) radial  
808 direction (adjusted  $R^2 = 0.9971$ ;  $RSS = 0.00095$ ) and (B) circumferential directions (adjusted  $R^2 = 0.9508$ ;  $RSS = 0.0098$ ).

809 **Supplementary Table 1**

Set #	$C_{10}$ (MPa)	$C_{01}$ (MPa)	Shear modulus (MPa)
1	5732.2650	-4989.8000	1484.93
2	5929.2900	-5169.5800	1519.42
3	5971.4700	-5230.2900	1482.36
4	5722.4110	-4981.2290	1482.36
5	881.8312	-100.0000	1563.66
6	119.4267	662.3430	1563.54
7	119.3580	662.0210	1562.76
8	5746.9000	-5002.6000	1488.60
9	119.4156	662.3429	1563.52
10	7881.8000	-7140.0000	1483.60
11	118.9000	659.5000	1556.80
Max	7881.8000	662.3430	1563.66
Min	118.9000	-7140.0000	1482.36
Average	3485.7334	-2724.2993	1522.87
STDEV	3142.4550	3160.9742	38.94

810 **Table 1.** 11 sets of  $C_{10}$  and  $C_{01}$  hyperelastic Mooney Rivlin material parameters that led to stress-  
811 strain curves in good agreement with the experimental data from the literature (see main text).  
812 The resultant fiber shear modulus is with an average of 1522.87 MPa and a standard deviation of  
813 38.94 MPa.