Title page

Optimal cell length for exploration and exploitation in chemotactic planktonic bacteria

Authors: Òscar Guadayol[1*], Rudi Schuech[2], Stuart Humphries[3]

Affiliations:

[1] Mediterranean Institute for Advanced Studies, IMEDEA (UIB-CSIC), Miquel Marqués 21, 07190 Esporles, Balearic Islands, Spain

[2] Department of Mathematics and Center for Computational Science, Tulane University, New Orleans, LA 70118, USA

[3] School of Life Sciences, Joseph Banks Laboratories, University of Lincoln, Lincoln LN6 7DL, UK

* **Correspondence to**: Òscar Guadayol, Mediterranean Institute for Advanced Studies, IMEDEA (UIB-CSIC), Miquel Marqués 21, 07190 Esporles, Balearic Islands, Spain, oscar@guadayol.cat

Running head: Optimal bacterial length for chemotaxis

Author Contribution Statement

OG: Conceptualization, Methodology, Software, Formal analysis, Investigation, Data curation, Writing – Original Draft, Visualization. **RS**: Conceptualization, Methodology, Software, Writing-Review & Editing. **SH**: Conceptualization, Writing - Review & Editing, Supervision, Project Administration, Funding Acquisition.

Abstract

Elongated morphologies are prevalent among motile bacterioplankton in aquatic systems. This is often attributed to enhanced chemotactic ability, but how long is best? We hypothesized the existence of an optimal cell length for efficient chemotaxis resulting from shape-imposed physical constraints acting on the trade-off between rapid exploration versus efficient exploitation of nutrient sources. To test this hypothesis, we first evaluated the chemotactic performance of elongated cephalexin-treated *Escherichia coli* towards α-methyl-aspartate in an agarose-based microfluidic device that creates linear, stable and quiescent chemical gradients. Our experiments showed that cells of intermediate lengths aggregated most tightly to the chemoattractant source. We then replicated these experimental results with Individual Based Model (IBM) simulations. A sensitivity analysis of the IBM allowed us to gain mechanistic insights into which parameters drive this trend and showed that the poor chemotactic performance of very short cells is caused by loss of directionality, whereas long cells are penalized by brief, slow runs. Finally, we evaluated the chemotactic performance of cells of different length with IBM simulations of a phycosphere – a hotspot of microbial interactions in the ocean. Results indicated that long cells swimming in a run-and-reverse pattern with extended runs and moderate speeds are most efficient at harvesting nutrients in this microenvironment. The combination of microfluidic experiments and IBMs proves thus to be a powerful tool for untangling the physical constraints that motile bacteria are facing in aquatic systems.

Main text

Introduction

Shape is a functional trait that influences several important tasks in bacterial life, including motility, nutrient uptake, resistance to grazing, passive dispersion, and pathogenicity (Young 2006). Among the characters commonly used to parametrize microbial shape, cell length exhibits one of the highest degrees of variability. Not only is there considerable interspecific diversity in cell length but also intraspecific. Cell length may fluctuate throughout the cell cycle and depends on growth rate and nutritional environment (Vadia and Levin 2015; Willis and Huang 2017) and sometimes increases dramatically as stressed individuals stop dividing (Erill et al. 2007).

One of the most important behaviours that planktonic bacteria can undertake is chemotaxis, which allows a motile bacterium to detect chemical gradients (exploration) and to track these gradients and remain close to the nutrient source to maximize uptake (exploitation) (Stocker and Seymour 2012). Elongated cell morphologies are particularly prevalent among planktonic chemotactic bacteria (Dusenbery 1998) probably because they help microbial swimmers overcome the most insidious challenge they face: Brownian motion. Swimming bacteria live in a low Reynolds number regime and, due to their small size, cannot maintain a straight course as they are constantly deflected by colliding water molecules. To overcome this challenge, flagellated bacteria have evolved a characteristic swimming pattern, consisting of a sequence of runs and sudden reorientations. During a run, a cell swims in a relatively straight line, whereas in the reorientation phase the direction of motion changes in species-specific ways (Mitchell 2002; Mitchell and Kogure 2006; Grognot and Taute 2021). Elongation, by increasing rotational drag, reduces the effects of Brownian motion, thus allowing longer and straighter runs and increasing signal-to-noise ratios and chemotactic performance (Dusenbery 2009). Furthermore, the increase in rotational drag due to elongation also restricts reorientation angles (Guadayol et al. 2017), which helps cells to track gradients more rapidly (Locsei 2007; Nicolau et al.

2009).

Elongation also poses some challenges to efficient chemotaxis. First, the restriction of reorientation

angles leads to higher dispersal rates (Lovely and Dahlquist 1975) and so shortens the time spent close

to a nutrient source. Second, elongation increases translational friction, slowing the cell and decreasing

swimming efficiency (Dusenbery 1998; Schuech et al. 2019). And finally, it can slow down or weaken

the cell's response to environmental cues as, given the diffusive nature of intracellular transport of

signals, the time it takes for chemotactic signals to travel from the membrane receptors to the flagellar

motors increases quadratically with cell length (Segall et al. 1985). In bacteria with randomly

distributed multiple flagella, this signal delay can also induce desynchronization of flagella (Maki et al.

2000), ultimately leading to alterations in the motility pattern (Guadayol et al. 2017).

In summary, current theories predict both potential advantages and disadvantages of elongation for

chemotaxis (see *Supplementary material*, *Theoretical Background* and Table S1). We hypothesized

that optimal cell lengths for efficient chemotaxis exist, imposed by a trade-off between the need to run

straight, swim efficiently, and explore the environment quickly, and the need to reorient when desired,

respond promptly to environmental changes, and maximize the time spent near nutrient sources.

We sought to experimentally test our hypothesis by using *Escherichia coli* cultures treated with

cephalexin, an antibiotic that stops cells from dividing without otherwise physiologically impairing

them (Rolinson 1980; Ishihara et al. 1983; Maki et al. 2000). To assess the influence of cell length on

chemotactic performance, we subjected cultures at various stages of elongation to microfluidic

chemotactic assays with different gradients of concentration of the chemoattractant α-methyl-aspartate.

To gain a mechanistic understanding of the experimental results, we then built an individual based

model (IBM) of our experimental system in which the chemotactic behaviour was simulated with a

well-known model for the aspartate chemotactic pathway in *E. coli* (Tu et al. 2008). We used this IBM

to further explore the consequences of elongation for chemotaxis in the phycosphere, the chemically

distinct area surrounding phytoplanktonic cells where most interactions with bacteria, with global

significance, take place (Seymour et al. 2017; Raina et al. 2022).

Materials and methods

Cultures

Experiments were performed with the chemotactic *E. coli* strain AW405. Cultures were grown

overnight at 30ºC in 25 mL of minimum growth media (Maki et al. 2000) from frozen glycerol stock

and shaken at 200 rpm at an orbit of 19 mm in 100 mL Erlenmeyer flasks. The saturated culture was

diluted 100-fold in 25 mL of fresh medium. Growth was monitored hourly using a cell density meter

measuring at a wavelength of 600 nm. After ~2 h, once cultures had achieved exponential growth

(optical density ~0.2), a sample was taken to conduct a chemotactic assay on the non-treated

population, and then cephalexin was added to a final concentration of 60 μgL^{-1}. Cephalexin is a β-

lactam antibiotic that promotes the formation of increasingly elongated cells and filaments by

interfering with the formation of the septal ring and stopping cell division (Eberhardt et al. 2003).

Cephalexin is not otherwise known to alter growth rate, physiology, or flagellar motility (Rolinson

1980; Ishihara et al. 1983; Maki et al. 2000). Once cephalexin was added, chemotactic performance

was evaluated every hour, as cells progressively elongated, up to five hours after the addition of

cephalexin.

Before performing any chemotactic assay, cultures were washed into random motility buffer (RMB,

Turner et al. 2000) by filtering and backwashing with RMB through a 0.2 µm sterile membrane three

times (Bell et al. 2005). This procedure resulted in highly motile suspensions of bacteria with limited

loss of cells (usually less than half) and was much quicker than the classical centrifugation technique.

Microfluidic device

To evaluate the chemotactic ability of *E. coli,* we used an agarose-based microfluidic device with three

channels (Fig. 1A, Diao et al. 2006; Cheng et al. 2007; Ahmed et al. 2010). This design allows the

generation of quiescent, steady linear chemical gradients. We used a design with three 600 µm wide

and 100 µm deep channels running parallel and separated by 200 µm wide walls and embedded in 3%

agarose (corresponding to 'design 2' in Ahmed et al. 2010). Details on the fabrication and calibration

of the devices are given in *Supplementary material, Microfluidic devices*.

Chemotactic assays

Chemical gradients were generated by pumping in the outer channels of the device RMB with different

concentrations of α-methyl-aspartate. In all experiments we kept the concentration of aspartate in the

source channel at 100 µM, while changing the concentration at the sink channel to achieve different

gradients. We performed experiments with three different sink concentrations (0, 25 and 50 µM)

which, after accounting for the differences from theory quantified in the device calibration, resulted in

cross-channel average relative concentration gradients of 1.20 mm^{-1}, 0.72 mm^{-1} and 0.40 mm^{-1}

respectively. Relative gradients are defined as $\frac{dC}{dx}/\overline{C}$, where C is the concentration of α-methyl-

aspartate, x is the distance along the gradient and $\overline{C}$ is the average concentration over the whole width

of the channel (see a list of all symbols used throughout the text in Table 1). These linear gradients

were similar to those considered in previous experimental and theoretical studies of bacterial

chemotaxis (e.g. Kalinin et al. 2009; Ahmed et al. 2010).

We imaged cells in a phase contrast inverted microscope (Zeiss AxioVert A1) equipped with a

temperature-controlled stage (PE100-ZAL System, Linkam Scientific Instruments Ltd., Tadworth, UK)

and using a 10X objective (0.55 NA), which gave a depth of field of 12.4 µm. Videos were recorded

using a sCMOS camera (Hamamatsu ORCAFlash2.8, Hamamatsu City, Japan). The resulting lateral

resolution was of 2.75 µm pixel^{-1}, and the field of view was 524 µm by 598 µm, allowing the

visualization of almost the entire width of the channel at once. Videos were taken at 30 fps which is

fast enough to efficiently track *E. coli* (Guadayol et al. 2017).

Figure 1*: Chemotactic responses in example populations of E. coli of different cell lengths in a microfluidic device creating a linear gradient of α-methyl-aspartate.* ***a)*** *Schematic of the microfluidic device used for the chemotactic assays.* ***b)*** *Epifluorescence microscope image of a gradient of fluorescein in the microfluidic device showing the entire width of the central channel where bacteria are located, and partially the lateral channels were buffers with different amount of chemoattractant are flowing. Vertical dashed red lines mark the limits of the channels. Superimposed in green is the profile of average fluorescein concentration (in arbitrary units of fluorescence intensity).* ***c)*** *Bacterial cell densities across the width of the central channel for populations of E. coli of the three example cell lengths exposed to a linear gradient of α-methyl-aspartate (* $\nabla C/\overline{C}=1.2\,mm^{-1}$ *). The shortest cells (1.7 ±0.3 µm, in purple) are untreated* E. coli *cells. The other two examples correspond to the average length of populations two hours (5.2±1µm, in red) and three hours (13.2±2.4 µm, in*

*orange) after treatment with cephalexin. Circles are empirical measurements; continuous lines are the nonlinear least squares fits of the exponential. **d**) Inverse cumulative frequency of run times for cells swimming upgradient (continuous lines) and downgradient (dashed lines) with a maximum deviation of 45 degrees from the direction of the gradient. **e**) Polar plot of the average run speeds per mean direction of the runs.*

In each experiment, several one-minute-long videos were taken at regular intervals starting 10 minutes

after bacteria were seeded in the centre channel, giving ample time for the chemical linear gradient to

re-establish (Ahmed et al. 2010), and ending 60 minutes after seeding. Motile bacteria were tracked

using in-house MATLAB code described elsewhere (Guadayol et al. 2017) and published under an

open-source license (Guadayol 2016). The code characterizes shapes and positions of individual cells

and registers frequency, duration, speed and direction of runs, as well as frequency, duration and angle

of reorientations.

A few minutes after injecting the bacterial culture into the central channel, cells visibly aggregated

towards the outer channel with the highest chemoattractant concentration, following an exponential

distribution. The inverse of the exponential decay constant, the "chemotactic precision length scale"

(L), was used to parameterize chemotactic performance (Kalinin et al. 2009; Son et al. 2016). Precision

length scales were estimated for each cell-length class and chemical gradient by fitting an exponential

function to the distribution of bacteria across the centre channel using the Least Absolute Residual

(LAR) robust curve fitting algorithm. Data points less than 50 µm from the walls were discarded

because light diffracted by the agarose walls interfered with the particle detection algorithm. Other

metrics of chemotactic response were considered but were ultimately disregarded because of large

errors associated to their estimation in the experimental setup, and because of their sensitivity to the

local gradient (*Supplementary material*, *Choice of response parameter*).

An occasional imperfect seal between the agarose and the PDMS resulted in shallower nutrient

gradients and lower bacterial densities than expected. We identified experiments with deficient seals as

those that met any of the following criteria: 1) residual flow in the centre channel was above 2 µm s^{-1}

(i.e., ~10% of the average speed for *E. coli*); 2) average number of cells detected was lower than 500;

3) distribution of cells was not monotonically decreasing away from the nutrient source, and 4)

coefficient of determination R^2 of the fitted exponential model was <0.1. Overall, we recorded 126

videos, of which 89 were successful and used in subsequent analyses. All experimental data, including

experiments not included in the final analyses, have been uploaded to a public data repository

(Guadayol et al. 2024).

IBM model

We developed an individual-based model (IBM) to gain a mechanistic insight into how chemotactic

responses depend on cell length. The IBM was written in MATLAB code and published under GNU

license (Guadayol 2023). The IBM simulates cells behaving independently (i.e. without interacting or

communicating with other cells) in a 2D domain bound by parallel reflecting walls 600 µm apart,

which is the width of the channel in the experimental microfluidic device. The dimensionality of the

model not only replicates that of the microfluidic device, in which nutrient concentration is invariant in

the vertical, but also of the motility parameters, which were estimated from 2D videos. The

chemotactic behaviour of the cells in the IBM was simulated with a coarse-grained chemotactic

pathway model described extensively elsewhere (Tu et al. 2008; Kalinin et al. 2009; Shimizu et al.

2010). See *Supplementary material, Chemotaxis signalling pathway model* for details on our

implementation. Both the number of cells (10^5) and the time step (0.1 s) were selected to ensure errors

below 10% (*Supplementary material, Convergence tests*). Each simulation was randomly populated

with cells of a given cell length. The number of cells remained constant throughout the simulation.

Their "swimming speed" (*v*), "tumble bias" (*TB,* the proportion of time spent tumbling) and

"directional persistence" (α_p, defined here as the average of the cosine of the reorientation angles) were

drawn from published empirical functional relationships of these parameters versus cell length for *E.*

coli swimming in homogeneous conditions (Guadayol et al. 2017, *Supplementary material* Fig. S1). All

parameters used in the IBM are listed in Table 1.

	Symbol	**Parameter**	**Definition**	**Units**
Population-level parameters	C	Chemoattractant concentration	Environmental concentration of aspartate or α-methyl-aspartate	µM
	L	Chemotactic precision length-scale	Inverse of the exponential decay constant of bacteria distribution vs distance to nutrient source	µm
	t_{ss}	Time to steady state	Time it takes for IBM simulation to reach steady state, which is defined as the moment when the relative rate of change in L is lower than 10% per hour	min
	B	Normalized cell density	Density of bacteria per distance to source boundary, normalized to number of bacteria in the simulation	-
	N	Normalized nutrient exposure	Population-averaged, time-integrated exposure to chemoattractant for the entire population in IBM simulations, normalized by integrated exposures for non-chemotactic populations distributed uniformly along the gradient	-
	μ	Bacterial random motility	Coefficient of diffusion of a bacterial population in the absence of chemical gradients, calculated fitting Taylor's equation (Taylor 1922) to tracks from IBM simulations of bacteria swimming in homogeneous conditions (Schuech and Menden-Deuer 2014)	µm^2 s^{-1}
	v_c	Chemotactic speed	Mean speed at which bacteria move up a chemoattractant gradient	µm s^{-1}
Individual motility parameters	f_t	Translational friction coefficient	Translational friction coefficient along the long semi-axis of a prolate ellipsoid of revolution	kg s^{-1}
	f_r	Rotational friction coefficient	Rotational friction coefficient about the short semiaxes of a prolate ellipsoid of revolution	kg µm^2 s^{-1}
	D_r	Rotational diffusivity	Rate of rotational diffusion about the short semiaxes of a prolate ellipsoid of revolution: $D_r = \frac{kT}{f_r}$	rad^2 s^{-1}
	τ_r	Run time	Average duration of runs	s
	τ_t	Tumble time	Average duration of tumbles	s
	TB	Tumble bias	Proportion of time spent tumbling: $TB = \frac{\tau_t}{\tau_t + \tau_r}$	-
	α_p	Directional persistence parameter	Average of the cosine of the reorientation angles	-
	v	Swimming speed	Average swimming speed during runs	µm s^{-1}

Table 1: Summary of the symbols used for each parameter and their definition.

During a run, each cell was assumed to move at a constant run speed, whereas its swimming direction

was affected by rotational Brownian motion. To account for this effect, after each time step, the run

direction was changed by adding a normally distributed stochastic component with standard deviation

$\sigma = \sqrt{2 D_r \Delta t}$, where D_r is the size- and shape-dependent rotational diffusivity perpendicular to the long

axis and Δt is the time step of the simulation. Rotational diffusivity can be calculated as $D_r = \dfrac{kT}{f_r}$, where

k is the Boltzmann constant, T is the absolute temperature and f_r is the rotational friction coefficient

about the short axes. We modelled bacteria as rigid prolate ellipsoids of revolution (Dusenbery 1998),

and used well-known theoretical solutions for the friction coefficients (Perrin 1934).

The duration of both runs and tumbles was modulated stochastically from Poisson processes whose

frequencies are the inverse of the expected durations and which generate the exponential distributions

observed experimentally (Berg and Brown 1972). The expected duration of a tumble (τ_t) changed only

with cell length, whereas the expected duration of a run (τ_r) depended also on environmental

chemoattractant concentrations (see following subsections).

At the start of a simulation, cells were distributed uniformly and oriented randomly within a 600 by 600

µm square domain. A reflecting boundary condition was imposed at the two edges of the domain

perpendicular to the nutrient gradient, analogous to the walls of the centre channel in the chemotactic

device; our results were insensitive to this choice (*Supplementary material, Boundary behaviour, Fig.*

S2). Cells bumping into one of the two edges parallel to the gradient were translated to the other

boundary without a change in course or speed (i.e. a periodic boundary condition).

Phycosphere model

The phycosphere was simulated with the same IBM model on an annular 2D domain, in which the

inner circle approximated the wall of a spherical phytoplankton cell of radius 25 µm. The outer

boundary was set to a radius of 500 µm, corresponding to half the average distance between

neighbouring microalgae at a concentration of 10^3 cells·mL^{-1}. Bacteria encountering a boundary

bounced away conserving the same angle of incidence and speed. A 2D domain was necessary to

match the dimensionality of the measurements of size-dependent motility parameters (*Supplementary*

material Fig. S1), which were obtained from 2D videos.

The chemoattractant concentration gradient around a spherical phytoplankton cell at steady state,

assuming background concentrations to be 0, follows $C(r)=C_0 r_0 r^{-1}$ (e.g. Jackson 1987), where r is the

radial distance from the centre of phytoplankton cell, r_0 is the radius of the cell, and C_0 is the

concentration of the chemoattractant at the cell wall. C_0 was set to be constant (corresponding to a

situation where the cell is releasing the chemoattractant at a constant rate) with an arbitrary value of 10

µM that ensured concentrations of aspartate within the sensitivity range of *E. coli*, which is a poor

chemotacter in comparison with free-living planktonic bacteria (Stocker 2011). In this scenario, at

steady state the shape of the chemical gradient is independent of molecular diffusivity.

Phycosphere simulations were run for 10 minutes to match reported temporal spans of formation and

dissipation of bacteria aggregates around nutrient hotspots (Blackburn et al. 1998; Smriga et al. 2016).

Results

Experimental results

As cells elongated, their motility pattern changed: swimming speeds decreased, tumble angles were

narrower, runs shorter and tumbles longer (Fig. 2, *Supplementary material* Fig. S1). We detected

chemotactic responses in populations of *E. coli* of all cell lengths. As is commonly observed in *E. coli*,

these responses were caused by a directional bias in run times rather than in swimming speed (Fig. 1c,

d). After a few minutes, cells distributed along the chemical gradient following an exponential

relationship in all experiments (Fig. 1b) of the form $B(x)=B(0)e^{-x/L}$, where $B(x)$ is the cell density at

distance x from the wall closest to the most concentrated outer channel, and L is the chemotactic

precision length-scale that parametrizes how tightly cells aggregate near the source of chemoattractants

Figure 2: Example trajectories of cells of different length. **a**, **c** *and* **e** *are randomly selected experimental trajectories of cells exposed to a linear gradient of α-methyl-aspartate.* **b**, **d** *and* **f** *are randomly selected trajectories from the corresponding IBM simulations.* **a** *and* **b** *show trajectories of cells 1,7 ±0.3 µm long (i.e. untreated* E. coli *cells);* **c** *and* **d** *show trajectories of 5.2±1µm long cells (i.e. ~two hours after cephalexin treatment);* **e** *and* **f** *show trajectories of 13.2±2.4 µm long cells (i.e. ~three hours after cephalexin treatment).*

(Son et al. 2016). The observed exponential distribution is the steady-state solution of the Keller-

Segel's bacterial transport model for a linear chemical concentration profile (Kalinin et al. 2009; Son et

al. 2016). In this solution, at steady state $L = \mu / v_c$, where μ is the random motility coefficient that

parametrizes the diffusivity of bacteria in the absence of chemical gradients, and v_c is the chemotactic

drift velocity (i.e. the average speed at which bacteria swim towards the source). Thus, L can be

interpreted as the ratio of bacterial diffusivity to chemotactic ability. The shorter L is, the better bacteria

are at exploiting a nutrient source.

As previously observed (Kalinin et al. 2009), L decreased with increasing gradient steepness (Fig. 3).

Regardless of the steepness, the functional response of L to cell length was U-shaped, with minimum L

occurring for cells 2.5 to 5 µm long. This supports our hypothesis that some trade-off for optimal

Figure 3: Experimental and simulated precision length scales vs. cell length of E. coli *populations exposed to linear gradients of α-methyl-aspartate of different steepness for a period between 10 and 60 minutes. Colours represent the three gradients used in the experiments, blue being the shallowest gradient , red the intermediate and purple the steepest. Data points are experimental results, with error bars showing the standard errors of the exponential fits. Coloured areas show the IBM model outputs between 10 (upper limits) and 60 (lower limits) minutes of simulation time.*

chemotaxis is operating in *E. coli*. However, our experimental system alone does not allow identifying

which mechanisms are driving the observed pattern for two reasons. First, L is an emerging population

parameter that ultimately depends on at least four shape-dependent individual-based motility

parameters. These parameters are rotational friction coefficient f_r, swimming speed v, tumble bias TB

(that is the proportion of time spent tumbling), and directional persistence α_p (defined here as the

average of the cosine of the reorientation angles). All these parameters have non-linear functional

relationships with cell length (*Supplementary material* Fig. S1, Guadayol et al. 2017) and therefore

their individual effect cannot be isolated. Thus, the observed pattern has multiple possible non-

mutually exclusive explanations. For example, high L could be equally explained by conditions that

increase μ, such as high v or α_p, or by conditions that decrease v_c, such as low f_r. A summary of these

possible explanations based on current theories is given in the *Supplementary material*, *Theoretical*

background and Table S1.

The second reason that prevents identifying underlying mechanisms is that we manipulated cells into

adopting potentially unnatural shapes, which could induce experimental artefacts. For example, flagella

could become increasingly desynchronized as cells elongate, because of the longer distance that

chemical signals need to diffuse and because of the difficulty of forming a single bundle when flagella

are far apart (Lee et al. 2021). This might lead to changes in TB that may partly explain the increase in

L for the longest cells. However, such experimental artefacts do not necessarily reflect the biophysical

constraints on chemotaxis over evolutionary time and make it difficult to generalize results.

In order to tease apart which parameters are driving the observed pattern, we developed an individual-

based model (IBM) that simulates a population of chemotactic bacteria of a given cell length

swimming along a linear gradient of aspartate. Model simulations between 10 and 60 minutes

(corresponding to the experimental time span) show good fits with the experimental data for the three

chemoattractant gradients explored (Fig. 3).

IBM dynamics

In all IBM simulations, L monotonically decreased over time, approaching a horizontal asymptote (*Supplementary material* Fig. S3). We defined an operational steady state as the point at which the relative rate of change in L was lower than 10% per hour. The time t_{ss} it took for simulations to reach such steady state (grey circles in Fig. 4b) replicated the U-shaped relation with cell length and ranged between 31 and 107 minutes. This indicates that experiments were not long enough for populations to achieve steady state distributions, particularly for cells longer than 10 µm. Allowing simulations to proceed until steady state revealed a trend different to what was observed experimentally: L decreased almost monotonically with cell length (Fig. 4a) instead of showing a local minimum at intermediate cell lengths (Fig. 3). This divergence between long term simulations and experiments and simulations shorter than 60 minutes highlights the transient nature of the measurements. A more ecologically relevant parameter, the population-averaged time-integrated exposure to nutrients N, shows the relative advantage of cells of intermediate length over the longest cells to extend for simulations as long as 2 hours (Figs. 4c, 5a). This advantage, however, diminishes as simulation time increases (Fig. 5a).

The random motility coefficient (μ), which parametrizes the diffusivity of bacteria in the absence of gradients (and is therefore a proxy for exploration), shows also a maximum at intermediate lengths (Fig. 4d), although this occurs at shorter cells and is narrower than for nutrient exposure.

Sensitivity analysis

To understand the individual role of each motility parameter in the emerging pattern, we performed a sensitivity analysis of the model (Fig. 4) using a one-at-a-time (OAT) approach: we ran a series of simulations to steady state in which we changed each one of the relevant shape-dependent parameters (f_r, v, TB, α_p) whilst keeping the other parameters constant at the empirical average values for a wild-type cell 1.7 µm long and 0.75 µm wide at a temperature of 30°C swimming in a homogeneous environment ($f_r=0.4 \cdot 10^{-20}$ kg m^2 s^{-2}, $v=18$ µm s^{-1}, $TB=0$, $\alpha_p=0$, *Supplementary material* Fig. S1,

Figure 4: Sensitivity analyses at steady state of the IBM. a) Precision length scales L at steady state; b) time t_{ss} it takes for simulations to reach steady state. c) Normalized exposures to α-methyl-aspartate N after 30 minutes of simulation time. d) Bacterial diffusivities in homogeneous conditions. Grey circles represent simulations where all parameters changed with cell length according to the empirical relationships for cephalexin-treated E. coli *shown in Supplementary material Fig. S1. Coloured lines show the results of the one-at-a-time sensitivity analysis, where one parameter varies with cell length while the rest are fixed to the values corresponding to normal-size wild-type* E. coli.

Figure 5: Time-integrated nutrient exposure per cell vs cell length in simulations of A) a linear gradient of aspartate , and B) a phycosphere of 25 µm radius and a concentration of Asp at the surface of the microalga of 10 µM. Values are normalized by the average nutrient exposure of a non-chemotactic population uniformly distributed across simulation space.

Guadayol et al. 2017). Because tumble time τ_t is insensitive to changes in chemoattractant in the

chemotactic pathway model, responses to *TB* mostly reflect responses to run time τ_r.

This analysis showed that, at steady state, changes in *v*, *TB* and, most noticeably, f_r drove decreases in *L*

as cells elongated (Fig. 4a). On the other hand, α_p drove increases in *L* that may moderate the decay in

L as cells elongate. The operating trade-offs are more clearly revealed in the steady state time (t_{ss}, Fig.

4b) and, ultimately, in the integrated exposure to nutrients (*N*, Fig. 4c) and random motility (μ, Fig.

4d): slow performances of short cells are explained by f_r, whereas those of long cells are driven by *v*

and *TB*. Steady state time remained insensitive to changes in α_p.

Chemotaxis in the phycosphere

One of the most important hotspots of microbial activity in aquatic systems is the phycosphere, the

volume influenced by the metabolic activity of a phytoplankton cell that shows a chemical composition

distinct of background seawater and where many of the interactions with other microorganisms occur

(Seymour et al. 2017; Raina et al. 2022). To challenge our results for cell length dependent

performance against an ecologically relevant scenario, we further used our model to explore which

shape may be optimal for chemotactic bacteria seeking phycospheres.

*Figure 6: Modelled chemotactic performances in a phycosphere. **a**) Normalized integrated nutrient exposure (N) after 10 minutes of simulation time, against cell length for simulated populations of E. coli around a spherical phytoplankton cell of radius 25 µm releasing Asp at a constant rate. **b**) bacterial densities normalized by background densities vs distance to phytoplankton cell wall. Colour-coded dashed lines represent populations of increasing cell length. The black line represents simulation for a population of wild-type E. coli cells displaying a run-and-reverse motility pattern, common in marine planktonic motile bacteria. Run-and-reverse pattern was modelled enforcing a constant tumble angle of 180 degrees for all reorientations. Inset shows the distribution of the chemoattractant Asp used in the simulations. **c-e**) integrated nutrient exposure N as a function of rotational friction coefficient f_r and swimming speed v (**c**), tumble bias TB (**d**), and directional persistence α_p (**e**). Although v, TB, and α_p can all be treated as independent of f_r in our simulations, the red lines show the observed dependences derived from experimental work with* E. coli *(Guadayol et al. 2017). Grey line shows the maximum nutrient exposures per rotational diffusion within each two-parameter space.*

We considered the diffusion of a chemoattractant from a 25 µm radius phytoplankton cell with a

constant rate of release and assumed steady state following previous models (Stocker 2011). In this

scenario, concentration does not decrease linearly as in our chemotactic device but as a function of $1/r$,

where r is the radial distance to the centre of the phytoplankton cell. We then ran IBM simulations

within a 2D plane to model the clustering of bacteria around phytoplankton. The results show bacterial

distributions consistently peaking some micrometres away from the phytoplankton cell surface (Fig.

6b). This is a behaviour commonly observed in IBMs of chemotaxis and is called the "volcano

effect"(e.g. Bray et al. 2007; Simons and Milewski 2011). It results from bacteria swimming past the

point of maximum concentration of attractant because it takes some time for the cell to detect a

decrease in chemoattractant concentration and initiate a tumble. This reaction time is often

parameterized by the latency time, which is a measure of the time needed for signals generated by

external stimuli to be processed through the complete signal transduction pathway. A sensitivity

analysis indicated that although the magnitude of this local peak was sensitive to changes in all motility

parameters, its location was most sensitive to swimming speed (*Supplementary material* Fig. S4)

because, given chemotactic response time, faster swimming translates into longer reaction distances.

Thus, for example, in our modelling framework the distribution of a population of bacteria swimming

at 20 µm s^{-1} peaked 40 µm away from the phytoplankton cell wall, whereas that of a population of

bacteria swimming at 100 µm s^{-1} peaked 200 µm away.

The time-integrated exposure to nutrients N is again a better measure of the ecological relevance of

these results. As in the linear gradient case, N increases rapidly with cell elongation up to a length of

around 3.3 µm and then decreases more gradually (Fig. 5a). To understand this trend, we explored the

2-D spaces defined by f_r on one axis and v, TB, and α_p on the other (Fig. 6c-e). This exercise shows that

as cells elongate, N is maximized by low TB (i.e. long runs), an α_p close to -1 (characteristic of a run-

and-reverse pattern), and a speed v of around 20 µm s^{-1}. Slow swimming leads to slow aggregation

around the phytoplankter, but swimming too fast amplifies the volcano effect (*Supplementary material*

Fig. S4).

Discussion

Physical constraints and biological trade-offs

The time-integrated exposure to chemoeffectors N, which we consider a measure of fitness of a

chemotactic strategy, depends on two factors: the precision and the speed of chemotaxis. All shape-

dependent motility parameters affect these two factors in different ways (Fig. 4). Chemotactic

precision, which we parameterize with the chemotactic precision length-scale L, improves sharply with

increasing f_r, and to a lesser extent, with decreasing v and increasing TB. The result is that, at steady

state, precision improves as cells elongate. On the other hand, the speed of the chemotactic process,

which is inversely related to the time t_{ss} to reach steady state, is positively related to v and f_r and

negatively related to TB. The result is that chemotaxis is fastest for *E. coli* cells of 3 to 7 µm in length.

Thus, the U-shaped response of L to elongation that we observed empirically likely resulted from

shape-dependent motility parameters acting oppositely on the precision versus speed of the chemotactic

process. Thus, alongside the response of random motility (μ) in the absence of gradients (Fig. 4d), our

results reflect the well-known trade-off between exploitation and exploration operating in chemotactic

bacteria in aquatic systems (Clark and Grant 2005; Altindal et al. 2011).

The general trends observed in L at steady state in response to changes in f_r, v, and α_p conform to

predictions from theoretical models. For example, the response of L to elongation-driven changes in f_r

(Fig. 4a) is consistent with the monotonic improvement in gradient detection predicted by a previous

study based purely on length-dependent changes in translational and rotational drag in ellipsoids of

revolution (Dusenbery 1998). The decrease in L in response to decreases in v as cells elongate is

consistent with the Keller-Segel model, which predicts a positive, linear relationship between these

parameters (Rivero et al. 1989). Finally, the increase in L with increasing α_p as cells elongate is

consistent with predictions from a previous study (Lovely and Dahlquist 1975) that established a

positive relation between α_p and the random motility parameter μ (and hence, also L).

The apparently small effect that α_p has on L and t_{ss} (Fig. 4a,b) is inconsistent with a strong positive

relationship between α_p and chemotactic velocity v_c predicted by a previous analytical model (Locsei

2007). This strong relationship should induce decreases in precision length (since $L=\mu/v_c$) and in t_{ss} as

cells elongate. Instead, L increases and t_{ss} remains invariant, suggesting that the effect of α_p on μ is

overpowering its expected effect on v_c. Alternatively, the marginal significance of α_p in our results may

also be reflecting its poor performance as a descriptor of reorientation behaviour, at least in peritrichous

bacteria such as *E. coli*. Motility patterns are defined by the mode of reorientation between runs. The

changes in motility pattern imposed by elongation in *E. coli* are poorly characterized by α_p because it

does not discriminate between 180º reorientations and true reversals (where a leading pole in the

previous run becomes the trailing pole, as the bundle is reformed in the opposite pole). Thus, for

example, α_p=0 could equally result from a truly uniform distribution of tumble angles and from a

pattern of angles in which half are 0º and half are 180º. In any case, the relationship between α_p and L

needs to be further explored.

Ecological and evolutionary implications

The optimal values for chemotaxis in our simulations may not be directly ecologically representative

because the IBM is based on *E. coli*'s chemosensory circuit for aspartate and uses species-specific

motility patterns and physiological responses to elongation. We therefore expect our results to be

somewhat dependent on species-specific sensitivities to different chemoattractants and to the

magnitude and shape of the chemical gradient, as well as the more generic aspects of the geometry and

boundary conditions of the model domain, and the simulation time. However, we suggest that our

results still reflect underlying trade-offs and evolutionary pressures acting on bacterial morphology.

Our results indicate that an optimal value for cell length arises from the well-known trade-off between

the need for rapidly tracking nutrient patches (exploration) and the need for remaining close to them

(exploitation) (Clark and Grant 2005; Altindal et al. 2011). Long cells will eventually maintain a tighter

distribution near a nutrient source whereas short cells will track gradients faster. Thus, the optimal cell

length of bacteria realized in nature will likely depend on the particulars of their environment (Stocker

and Seymour 2012). In environments where gradients are long and stable over time (relative to

microbial scales), such as at the sediment-water interface, we may expect selective pressure for

remaining close to the nutrient source and therefore for increasing directional control by increasing f_r

and α_p. This is consistent with observations of extremely long and large bacteria near and within

sediments (Fenchel 1994; Guerrero et al. 1999; Jørgensen 2006). In the planktonic environment, with

its patchy, dynamic and ephemeral nutrient landscape, speed is likely more important (Xie and Wu

2014). This may partly explain why the median aspect ratio of planktonic bacteria is relatively small

(around three:Dusenbery 1998), and why flagellated bacteria in oligotrophic environments tend to

swim faster (Mitchell et al. 1995). This argument, of course, does not reflect the fact that tasks and

costs other than chemotaxis may be important for bacterial fitness (Young 2006; Schuech et al. 2019),

nor their interaction with flow (Luchsinger et al. 1999; Taylor and Stocker 2012).

This need for speed in the planktonic environment may be moderated, as our simulations show, by the

occurrence of a volcano effect that penalizes very fast swimmers (Fig 6, *Supplementary material* Fig.

S4). The volcano effect results from cells overshooting the phycosphere as it takes some time (the

"latency time") for the bacterial chemotactic pathway to perceive a decrease in chemoattractant

concentration and induce reorientation (Bray et al. 2007; Simons and Milewski 2011). Thus, the

importance of the volcano effect depends on three factors: the shape and length of the chemoattractant

gradient (which will themselves depend on the size of the phytoplankton cell), the latency time, and the

swimming speed. The slower the cells are at reacting to a change in concentration and the faster they

swim, the further from the maximum concentration will their population distribution peak. In our

simulations this translates into optimal nutrient exposures at swimming speeds remarkably close to the

average value for wild-type *E. coli* (Fig. 6c). In general, the latency time will add a constraint on the

maximum swimming speed that is ecologically viable in a given planktonic system.

If the volcano effect proves to be an ecologically relevant mechanism, there are several factors that

could be reducing its impact in planktonic bacteria in aquatic systems. First, their latency times are

thought to be much faster (Stocker 2011) than those of *E. coli* (~0.2 s Block et al. 1982). Second,

speed-dependent responses such as chemokinesis or changes in motility pattern (Son et al. 2016),

which we have not considered here, may help bacteria remain longer around nutrient patches even

when their swimming speeds are high. Third, the adoption of swimming patterns with a reversal phase

and long runs, which are prevalent among marine bacteria (Johansen et al. 2002) can also maximize

nutrient exposure around a phycosphere (Fig. 6). And finally, the existence of a viscous layer of

exopolymers around a phytoplankton cell may minimize the volcano effect (Guadayol et al. 2021).

Otherwise, our simulations of the phycosphere depict an optimal strategy that is largely consistent with

observations of marine bacteria. The simulations predict that elongated cells with a run-and-reverse

pattern (i.e. $\alpha_p \to$ -1) and long runs (i.e. $TB \to 0$) will be most efficient in harvesting nutrients at

timescales of tens of minutes (Fig. 6). Among marine bacterial populations, the median cell aspect ratio

is close to 3 (Dusenbery 1998), the prevalent motility pattern is run-and-reverse (Johansen et al. 2002)

or has some sort of reversal phase (Xie et al. 2011), and the most common type of flagellation is a

single polar flagellum (Grognot and Taute 2021), which favours long runs at the cost of increasing

chemotactic response times (Sneddon et al. 2012).

We have focused on individual motility parameters (rotational friction, run speed, tumble and run

times, directional persistence) that are dependent on the shapes of the organisms in our experimental

system, and that therefore may be expected to change in natural populations displaying morphological

variability. However, cell shape is not the only factor sensitive to selective pressures that can influence

these motility parameters. For example, high f_r can also be achieved by means of long flagella, which

can act as stabilizers of the running trajectory (Mitchell 2002; Schuech et al. 2019). Such stabilization

would allow cells to be smaller while remaining chemotactically efficient. Similarly, changes in v or α_p

can be achieved by means other than elongation, for example by alterations in the number or positions

of flagella (Grognot and Taute 2021). Nevertheless, given the morphological diversity often seen in

natural populations, and the ease with which some cell populations appear to change shape, and in

particular aspect ratio, varying length may well be a response to an ever-changing nutrient landscape.

Chemotaxis is but one of several activities that may exert selective pressures on bacterial shape in

general and cell length in particular (Young 2006). Some of these pressures may in fact act in

opposition to the chemotactic mechanisms we reveal here. For example, specific nutrient uptake is

highest in very small and filamentous cells. Protozoan grazers predate most efficiently upon bacteria of

intermediate sizes and thus may select for small cells and filaments (Pernthaler 2005). And finally,

passive dispersal by Brownian motion is most efficient for small spherical cells (Dusenbery 1998).

Regardless, the fact that slightly elongated rods are the prevalent morphology among marine motile

bacterioplankton hints at the important role that motility and chemotaxis play in shaping natural

populations of aquatic bacteria.

In summary, our work dissects the main physical mechanisms that explain the prevalence of slightly

elongated forms among chemotactic bacterioplankton and showcases the interplay between them. Both

in a linear gradient replicating the long, stable gradients near static boundaries, and in a non-linear

gradient replicating a phycosphere in a much more dynamic planktonic setting, we reveal the existence

of an optimal cell length for bacterial chemotaxis and show how it arises from a tradeoff between the

needs for exploration and exploitation.

Acknowledgements

We thank H. Fu for comments on previous versions of this manuscript. This study was funded by

Leverhulme Trust project RLA RL-2012-022, "Form and function in a microbial world", granted to

S.H., and by a Gordon and Betty Moore Foundation Marine Microbiology Initiative award, grant 6852

to S.H. and O.G. O.G. was supported by a María Zambrano grant from the University of the Balearic

Islands. No conflicts of interest are declared.

447   **References**

Ahmed, T., T. S. Shimizu, and R. Stocker. 2010. Bacterial chemotaxis in linear and nonlinear steady microfluidic gradients. Nano Lett. **10**: 3379–3385. doi:10.1021/nl101204e

Altindal, T., L. Xie, and X.-L. Wu. 2011. Implications of three-step swimming patterns in bacterial chemotaxis. Biophys. J. **100**: 32–41. doi:10.1016/j.bpj.2010.11.029

Bell, C. H., B. S. Arora, and T. A. Camesano. 2005. Adhesion of *Pseudomonas putida* KT2442 is mediated by surface polymers at the nano- and microscale. Environ. Eng. Sci. **22**: 629–641. doi:10.1089/ees.2005.22.629

Berg, H. C., and D. A. Brown. 1972. Chemotaxis in *Escherichia coli* analysed by three-dimensional tracking. Nature **239**: 500–504. doi:10.1038/239500a0

Blackburn, N., T. Fenchel, and J. G. Mitchell. 1998. Microscale nutrient patches in planktonic habitats shown by chemotactic bacteria. Science **282**: 2254–2256. doi:10.1126/science.282.5397.2254

Block, S. M., J. E. Segall, and H. C. Berg. 1982. Impulse responses in bacterial chemotaxis. Cell **31**: 215–226.

Bray, D., M. D. Levin, and K. Lipkow. 2007. The chemotactic behavior of computer-based surrogate bacteria. Curr. Biol. **17**: 12–19. doi:10.1016/j.cub.2006.11.027

Cheng, S.-Y., S. Heilman, M. Wasserman, S. Archer, M. L. Shuler, and M. Wu. 2007. A hydrogel-based microfluidic device for the studies of directed cell migration. Lab. Chip **7**: 763. doi:10.1039/b618463d

Clark, D. A., and L. C. Grant. 2005. The bacterial chemotactic response reflects a compromise between transient and steady-state behavior. Proc. Natl. Acad. Sci. U. S. A. **102**: 9150–9155. doi:10.1073/pnas.0407659102

Diao, J., L. Young, S. Kim, and others. 2006. A three-channel microfluidic device for generating static linear gradients and its application to the quantitative analysis of bacterial chemotaxis. Lab. Chip **6**: 381–388. doi:10.1039/B511958H

Dusenbery, D. B. 1998. Fitness landscapes for effects of shape on chemotaxis and other behaviors of bacteria. J. Bacteriol. **180**: 5978–5983.

Dusenbery, D. B. 2009. Living at micro scale: the unexpected physics of being small, Harvard University Press.

Eberhardt, C., L. Kuerschner, and D. S. Weiss. 2003. Probing the catalytic activity of a cell division-specific transpeptidase in vivo with β-lactams. J. Bacteriol. **185**: 3726–3734.

Erill, I., S. Campoy, and J. Barbé. 2007. Aeons of distress: an evolutionary perspective on the bacterial SOS response. FEMS Microbiol. Rev. **31**: 637–656. doi:10.1111/j.1574-6976.2007.00082.x

Fenchel, T. 1994. Motility and chemosensory behaviour of the sulphur bacterium *Thiovulum majus*. Microbiology **140**: 3109–3116. doi:10.1099/13500872-140-11-3109

Grognot, M., and K. M. Taute. 2021. More than propellers: how flagella shape bacterial motility behaviors. Curr. Opin. Microbiol. **61**: 73–81. doi:10/gjwcbr

Guadayol, Ò. 2016. trackbac.doi:10.5281/zenodo.3452818

Guadayol, Ò. 2023. chemotaxisIBM.doi:10.5281/ZENODO.8360667

Guadayol, Ò., T. Mendonca, M. Segura-Noguera, A. J. Wright, M. Tassieri, and S. Humphries. 2021. Microrheology reveals microscale viscosity gradients in planktonic systems. Proc. Natl. Acad. Sci. **118**. doi:10/ghsc3s

Guadayol, Ò., R. Schuech, and S. Humphries. 2024. Experimental data from "Optimal cell length for exploration and exploitation in chemotactic planktonic bacteria." doi:10.6084/m9.figshare.25452355

Guadayol, Ò., K. L. Thornton, and S. Humphries. 2017. Cell morphology governs directional control in swimming bacteria. Sci. Rep. **7**: s41598-017-01565-y-017. doi:10.1038/s41598-017-01565-y

Guerrero, R., A. Haselton, M. Solé, A. Wier, and L. Margulis. 1999. *Titanospirillum velox*: A huge, speedy, sulfur-storing spirillum from Ebro Delta microbial mats. Proc. Natl. Acad. Sci. **96**: 11584–11588. doi:10.1073/pnas.96.20.11584

Ishihara, A., J. E. Segall, S. M. Block, and H. C. Berg. 1983. Coordination of flagella on filamentous cells of *Escherichia coli*. J. Bacteriol. **155**: 228–237.

Jackson, G. A. 1987. Simulating chemosensory responses of marine microorganisms. Limnol. Oceanogr. **32**: 1253–1266. doi:10.4319/lo.1987.32.6.1253

Johansen, J. E., J. Pinhassi, N. Blackburn, U. L. Zweifel, and A. Hagström. 2002. Variability in

motility characteristics among marine bacteria. Aquat. Microb. Ecol. **28**: 229–237. doi:10.3354/ame028229

Jørgensen, B. B. 2006. Bacteria and marine biogeochemistry, p. 169–206. *In* H.D. Schulz and M. Zabel [eds.], Marine Geochemistry. Springer.

Kalinin, Y. V., L. Jiang, Y. Tu, and M. Wu. 2009. Logarithmic sensing in *Escherichia coli* bacterial chemotaxis. Biophys. J. **96**: 2439–2448.

Lee, W., Y. Kim, C. S. Peskin, and S. Lim. 2021. A novel computational approach to simulate microswimmers propelled by bacterial flagella. Phys. Fluids **33**: 111903. doi:10.1063/5.0069343

Locsei, J. T. 2007. Persistence of direction increases the drift velocity of run and tumble chemotaxis. J. Math. Biol. **55**: 41–60. doi:10.1007/s00285-007-0080-z

Lovely, P. S., and F. W. Dahlquist. 1975. Statistical measures of bacterial motility and chemotaxis. J. Theor. Biol. **50**: 477–496. doi:10.1016/0022-5193(75)90094-6

Luchsinger, R., B. Bergersen, and J. Mitchell. 1999. Bacterial swimming strategies and turbulence. Biophys. J. **77**: 2386. doi:10.1016/S0006-3495(99)77075-X

Maki, N., J. E. Gestwicki, E. M. Lake, L. L. Kiessling, and J. Adler. 2000. Motility and chemotaxis of filamentous cells of *Escherichia coli*. J. Bacteriol. **182**: 4337–4342. doi:10.1128/JB.182.15.4337-4342.2000

Mitchell, J. G. 2002. The energetics and scaling of search strategies in bacteria. Am. Nat. **160**: 727–740. doi:10.1007/BF02540225

Mitchell, J. G., and K. Kogure. 2006. Bacterial motility: links to the environment and a driving force for microbial physics. FEMS Microbiol. Ecol. **55**: 16.

Mitchell, J. G., L. Pearson, A. Bonazinga, S. Dillon, H. Khouri, and R. Paxinos. 1995. Long lag times and high velocities in the motility of natural assemblages of marine bacteria. Appl. Environ. Microbiol. **61**: 877–882.

Nicolau, D. V., J. P. Armitage, and P. K. Maini. 2009. Directional persistence and the optimality of run-and-tumble chemotaxis. Comput. Biol. Chem. **33**: 269–274. doi:10/chfkjq

Pernthaler, J. 2005. Predation on prokaryotes in the water column and its ecological implications. Nat.

Rev. Microbiol. **3**: 537–546.

Perrin, F. 1934. Mouvement brownien d'un ellipsoide - I. Dispersion diélectrique pour des molécules ellipsoidales. J Phys Radium **5**: 497–511. doi:10.1051/jphysrad:01934005010049700

Raina, J.-B., B. S. Lambert, D. H. Parks, and others. 2022. Chemotaxis shapes the microscale organization of the ocean's microbiome. Nature **605**: 132–138. doi:10.1038/s41586-022-04614-3

Rivero, M. A., R. T. Tranquillo, H. M. Buettner, and D. A. Lauffenburger. 1989. Transport models for chemotactic cell populations based on individual cell behavior. Chem. Eng. Sci. **44**: 2881–2897. doi:10.1016/0009-2509(89)85098-5

Rolinson, G. N. 1980. Effect of β-lactam antibiotics on bacterial cell growth rate. Microbiology **120**: 317–323. doi:10.1099/00221287-120-2-317

Schuech, R., T. Hoehfurtner, D. J. Smith, and S. Humphries. 2019. Motile curved bacteria are Pareto-optimal. Proc. Natl. Acad. Sci. 201818997. doi:10.1073/pnas.1818997116

Schuech, R., and S. Menden-Deuer. 2014. Going ballistic in the plankton: anisotropic swimming behavior of marine protists. Limnol. Oceanogr. Fluids Environ. **4**: 1–16. doi:10.1215/21573689-2647998

Segall, J. E., A. Ishihara, and H. C. Berg. 1985. Chemotactic signaling in filamentous cells of *Escherichia coli*. J. Bacteriol. **161**: 51–59.

Seymour, J. R., S. A. Amin, J.-B. Raina, and R. Stocker. 2017. Zooming in on the phycosphere: the ecological interface for phytoplankton–bacteria relationships. Nat. Microbiol. **2**: 17065. doi:10.1038/nmicrobiol.2017.65

Shimizu, T. S., Y. Tu, and H. C. Berg. 2010. A modular gradient-sensing network for chemotaxis in *Escherichia coli* revealed by responses to time-varying stimuli. Mol. Syst. Biol. **6**: 382. doi:10.1038/msb.2010.37

Simons, J. E., and P. A. Milewski. 2011. The volcano effect in bacterial chemotaxis. Math. Comput. Model. **53**: 1374–1388. doi:10.1016/j.mcm.2010.01.019

Smriga, S., V. I. Fernandez, J. G. Mitchell, and R. Stocker. 2016. Chemotaxis toward phytoplankton drives organic matter partitioning among marine bacteria. Proc. Natl. Acad. Sci. 201512307.

doi:10.1073/pnas.1512307113

Sneddon, M. W., W. Pontius, and T. Emonet. 2012. Stochastic coordination of multiple actuators reduces latency and improves chemotactic response in bacteria. Proc. Natl. Acad. Sci. **109**: 805–810. doi:10.1073/pnas.1113706109

Son, K., F. Menolascina, and R. Stocker. 2016. Speed-dependent chemotactic precision in marine bacteria. Proc. Natl. Acad. Sci. **113**: 8624–8629. doi:10.1073/pnas.1602307113

Stocker, R. 2011. Reverse and flick: Hybrid locomotion in bacteria. Proc. Natl. Acad. Sci. **108**: 2635–2636. doi:10.1073/pnas.1019199108

Stocker, R., and J. R. Seymour. 2012. Ecology and physics of bacterial chemotaxis in the ocean. Microbiol. Mol. Biol. Rev. **76**: 792–812. doi:10.1128/MMBR.00029-12

Taylor, G. I. 1922. Diffusion by continuous movements. Proc. Lond. Math. Soc. **s2-20**: 196–212. doi:10.1112/plms/s2-20.1.196

Taylor, J. R., and R. Stocker. 2012. Trade-offs of chemotactic foraging in turbulent water. Science **338**: 675–679.

Tu, Y., T. S. Shimizu, and H. C. Berg. 2008. Modeling the chemotactic response of *Escherichia coli* to time-varying stimuli. Proc. Natl. Acad. Sci. **105**: 14855–14860. doi:10.1073/pnas.0807569105

Turner, L., W. S. Ryu, and H. C. Berg. 2000. Real-time imaging of fluorescent flagellar filaments. J. Bacteriol. **182**: 2793–2801. doi:10.1128/JB.182.10.2793-2801.2000

Vadia, S., and P. A. Levin. 2015. Growth rate and cell size: a re-examination of the growth law. Curr. Opin. Microbiol. **24**: 96–103. doi:10.1016/j.mib.2015.01.011

Willis, L., and K. C. Huang. 2017. Cell size: fat makes cells fat. Curr. Biol. **27**: R592–R594. doi:10.1016/j.cub.2017.05.017

Xie, L., T. Altindal, S. Chattopadhyay, and X.-L. Wu. 2011. Bacterial flagellum as a propeller and as a rudder for efficient chemotaxis. Proc. Natl. Acad. Sci. **108**: 2246–2251. doi:10.1073/pnas.1011953108

Xie, L., and X.-L. Wu. 2014. Bacterial motility patterns reveal importance of exploitation over exploration in marine microhabitats. Part I: Theory. Biophys. J. **107**: 1712–1720. doi:10.1016/j.bpj.2014.07.058

Young, K. D. 2006. The selective value of bacterial shape. Microbiol. Mol. Biol. Rev. **70**: 660–703.

doi:10.1128/MMBR.00001-06

Data availability statement

The dataset of videos generated and analysed during the current study is not publicly available due to

storage restrictions but is available from the corresponding author on reasonable request. Bacterial

trajectories and cell lengths extracted from the videos are published in a public repository (DOI:

10.6084/m9.figshare.25452355), alongside the details of the experimental conditions needed to

reproduce all analyses and figures. The code developed and used in this study is available under an

open-source license at Zenodo.org (DOI: 10.5281/zenodo.3452818; DOI: 10.5281/zenodo.8360667).