Size photometry and fluorescence imaging of immobilized immersed extracellular vesicles

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ABSTRACT

Immunofluorescence analysis of individual extracellular vesicles (EVs) in common fluorescence microscopes is gaining popularity due to its accessibility and high fluorescence sensitivity, however, EV number and size are only measurable using fluorescent stains requiring extensive sample manipulations. Here we introduce highly sensitive label-free EV size photometry (SP) based on interferometric scattering (iSCAT) imaging of immersed EVs immobilized on a glass coverslip. We implement SP on a common inverted epifluorescence microscope with LED illumination and a simple 50:50 beamsplitter, permitting seamless integration of SP with fluorescence imaging (SPFI). We present a high-throughput SPFI workflow recording >10,000 EVs in 7 min over multiple fields of view, pre- and post-incubation imaging to suppress background, along with automated image alignment, aberration correction, spot detection, and

EV sizing. We identify the upper sizing limit of SP with 440 nm illumination and extend EVs sizing from ~35 nm in diameter to >200 nm with dual 440 nm and 740 nm illumination. We benchmark SP to flow cytometry using calibrated silica nanoparticles and demonstrate superior, label-free sensitivity. We showcase SPFI's potential for EV analysis by experimentally distinguishing surface and volumetric EV dyes, observing the deformation of EVs adsorbed to a surface, and by uncovering distinct subpopulations in <100 nm-in-diameter EVs with fluorescently tagged membrane proteins.

INTRODUCTION

Cells secrete membrane-bound structures called extracellular vesicles (EVs) that are fundamental to tissue homeostasis and cell signaling¹, and have garnered great interest as biomarkers and for therapeutics². Their sizes range from a few tens of nanometers to more than a micrometer, and they carry a large variety of bioactive molecules such as proteins, lipids, nucleic acids, and metabolites. Major efforts have been devoted to the development of single vesicle measurement techniques to characterize the heterogeneity of EVs and the features that govern cellular activation pathways, EV cargo sorting³ and how they change in health and disease⁴. Currently, there are no generally accepted classifiers to distinguish EV subpopulations such as exosomes or ectosomes due to various technical challenges. Single protein expression levels suggest EV subpopulations, but due to considerable overlap are insufficient for classification⁵. Multi-dimensional characterization will be required and could involve multiplex proteomic analysis or incorporation of orthogonal parameters such as EV size. It is furthermore understood that EV diameters in various samples follow inverse power-law distributions where small EVs with

diameter d < 100 nm are by far the most abundant^{6,7}. Consequently, the limit of detection (LOD) of any single EV technology profoundly influences the results. For instance, proteomic analysis significantly varies between measurements with LODs of 100 nm compared to 150 nm⁸, suggesting that techniques with LODs substantially below 100 nm are required for reproducible and unbiased sample characterization.

Several techniques show potential to meet these requirements and advance EV subclassification. For example, electron microscopy can be combined with immunogold labelling to both size and characterize EVs down to d < 30 nm but has exceedingly low throughput and is thus impractical. Optical methods that can combine EV immunofluorescence analysis with EV sizing based on light scattering and are preferred owing to their higher throughput. Indeed, flow cytometry (FCM) emerged as the current gold standard for single EV analysis owing to its comparative ease of use, low operating cost, and simultaneous measurement of both EV size (by side scattering) and protein expression (by immunofluorescence)⁸. Many common flow cytometers used to detect EVs have an LOD ranging between $d \approx 100 - 150$ nm, while specialized nanoflow cytometers can detect EVs as small as $d \approx 40$ nm, but at the expense of fluorescence multiplexing⁹. The LOD of FCM is commensurate to nanoparticle tracking analysis (NTA) with a label-free LOD of $d \approx 80$ nm and which is commonly used to assess EV size distributions. Lowering the detection threshold remains a challenge as nanoparticle scattering scales with d^6 , and lowering the LOD just threefold requires almost a thousand-fold gain in sensitivity. Recently, fluorescence imaging of EVs captured on a surface by widefield or TIRF microscopy^{10–12} has gained traction as alternative to FCM due to its accessibility, higher fluorescence sensitivity¹³, and the possibility to manipulate EVs following their immobilization¹⁴. The absence of a scatter signal for sizing is circumvented using fluorescent

labels (e.g. dyes), but at the cost of further sample manipulation and susceptibility to false positive detection events, e.g. due to label aggerates. Alternatively, super-resolution imaging can be used to size EVs down to $d \approx 30$ nm but suffers from low throughput and requires very high expression of target proteins (at least 20) in a single EV⁶. Label-free detection of surface bound particles by dark-field imaging is well established¹⁵, however, it requires specialized microscopes with bright light sources and strong stray light suppression, and the scattering signal also scales unfavorably $\propto d^6$, and has not been useful for EV analysis.

Interference imaging methods such as interference scattering microscopy (iSCAT)¹⁶, single particle Interference Reflection Image Sensing (SP-IRIS)¹⁷, and several others^{18–21} offer relief as they optically amplify signals with minimally added noise²², are robust against stray light, and benefit from a more favorable contrast scaling $\propto d^3$. iSCAT and SP-IRIS require the least complex optical setups as both generate the interferometric reference field by reflection of the illumination light from the imaging substrate in common-path geometry. SP-IRIS achieves this using custom-made silicon chips coated with a thin dielectric reflective layer. It employs imaging from the top, mostly though air, since immersing samples adds complications such as additional reflections from the top surface or minute changes in liquid depth that affect focusing. iSCAT on the other hand uses common glass coverslips without any modification. Imaging is conducted not through the medium, but through the coverslip, and immersed samples can readily be analyzed. iSCAT offers exceptional sensitivity including label-free single molecule detection and mass photometry²³, it can be combined with fluorescence²⁴ and has been used to detect $EVs^{25,26}$. However, reliance on specialized optical setups has so far limited both wider adoption for EV sizing and integration with EV fluorescence analysis.

Here we introduce size photometry (SP) of immobilized EVs by differential iSCAT pre- and post-incubation imaging of a (capture) surface, together with concomitant SP and fluorescent imaging (SPFI) using a widefield fluorescence microscope. We describe iSCAT using LED light sources and cameras intended for fluorescence microscopy, and benchmark it against flow cytometry using calibrated silica particles. We introduce an automated image analysis pipeline that includes image pre-processing and background suppression steps and allows to analyze captured vesicles over multiple fields of view (FOVs) for high throughput SP analysis. Recognizing an upper sizing limit for SP, we introduce combined SP using two illumination wavelengths and extend the dynamic range of SP for polydisperse samples such as liposomes and EVs from $d \approx 35$ to > 200 nm. Next, we validate EV sizing in SPFI on EV samples stained with surface and intraluminal dyes as function of EV diameter and furthermore observe contrast behavior consistent with EV deformation on the surface. Finally, we analyze EVs expressing a fluorescent fusion protein and identify two distinct clusters of EVs with d < 100 nm, which are revealed by correlating EV sizes with the expression of the protein tag, and would be missed by techniques with LODs of d > 100 nm.

<u>RESULTS</u>

SPFI allows to combine sizing of immersed EVs inside a microfluidic flow cell with further characterization using fluorescent reporters or affinity binders. Particle sizing is enabled through label free iSCAT imaging on common glass coverslips which can be implemented in any widefield epifluorescence microscope by the addition of a 50:50 beamsplitter that is switched with the dichroic filter cube, shown in Fig. 1**a**. We use the terminology 'size photometry'

because nanoparticles are conventionally characterized by size, however, we note that iSCAT contrasts are principally dependent on both EV size and refractive index and only few implementations truly decouple them^{25–27}. A detailed description of the iSCAT imaging principle is given in Supplementary Section S1, describing the mathematical underpinning of its increased sensitivity for small particles compared to darkfield imaging. SP records pairs of pre- incubation (F_{pre}) and post-incubation (F_{post}) images over multiple FOVs to maximize imaging throughput while being selective to particles that are immobilized on the coverslip during incubation, Fig. 1b. Each F_{pre} and F_{post} contain only static features and can be averaged indefinitely, Fig. 1d, relaxing the need for high speed cameras which are otherwise common for dynamic iSCAT imaging²⁸.

We introduce an SPFI analysis pipeline (*https://github.com/junckerlab/SPFI*) to handle data preprocessing and image analysis starting with a step to correct chromatic aberrations between different imaging channels to maximize FOV sizes while ensuring correct spot matching between the channels (Fig. S2). Our calibration-free algorithm can map and correct distortions between any two SPFI channels directly on the respective post-incubation images. Next, the algorithm digitally aligns F_{pre} and F_{post} images (Supplementary Section S3) and detects large surface inhomogeneities on F_{pre} to be excluded from subsequent analysis. Our algorithm then localizes candidate spots in the image data as local contrast minima in $F_{diff} = F_{post} - F_{pre}$, thereby identifying captured EVs as well as sampling the image background (c.f. Fig. 1c, Fig. S3) to generate an experimental noise floor and inform thresholding (details in Supplementary Section S4). SP measurements are taken either in a single illumination wavelength or under dual illumination with a short and a long wavelength to maximise sizing range (described below) and candidate contrasts are measured as center pixel values. For SPFI, fluorescence images (Fig. 1e) are superposed with the iSCAT images and fluorescence intensities *I* for all candidate spots are measured (c.f. Fig. 1c, right panel). Taking differential fluorescence images is possible but often inconsequential due to overall low autofluorescence background in empty flow cells. The correlation between SP measurements and fluorescence signal intensities are reported as scatter plots of *I* and *C* (Fig. 1f) which are reminiscent of FCM data and led us to adopt data analysis workflows and thresholding protocols common to FCM. EVs are identified in them as candidate spots that exceed the SP background noise found in buffer control experiments, thus establishing a count of the total EV number and removing the ambiguity of thresholding on EV fluorescence alone¹².

Our analysis pipeline includes an iSCAT flatfield pre-processing step to correct setup-specific illumination inhomogeneities in the raw iSCAT images. The illumination background is identical for each FOV and can vary up to ~20 % within each image depending on the region of interest (ROI), see Fig. 2a, 2b. Flatfield processing is necessary because pre- and post-incubation images are often misaligned up to ~15 % of the FOV size (shifts are caused by manual pipetting and imperfect stage movements) and subtraction of the raw images can lead to residual background that overwhelms EV contrasts <1%. Flatfield images are collected as series of images from random positions over the sample which are averaged into a single flatfield image *FF*. The preand post-incubation images are then background corrected by division like $F_{pre} \rightarrow F_{pre}/FF$ and $F_{post} \rightarrow F_{post}/FF$.

iSCAT imaging parameters and image noise

The iSCAT images in this manuscript were taken by averaging 200 frames for both pre- and post-incubation images and normalizing them using flatfield images consisting of 960 individual frames (120 arbitrary positions on a slide to eliminate local interference by dust with 8× average each). We found those imaging parameters by considering the trade-off in imaging speed and detection sensitivity due to image noise as described below. The time it took to capture these images depended on the chosen FOV size since this limited camera acquisition speed and saving time. For example, in 1200×1200 pixel (88 × 88 μ m²) FOVs with 200 frames could be taken in ~12 s, split roughly equally between actual camera exposure and software overhead from computing the average and storing the data on the hard disk drive. Flatfield images needed to be taken only once per substrate regardless of the number of imaged FOVs and took about 180 s each. Imaging ten different FOVs in one iSCAT channel with 560 EVs each (Fig. S4, Supplementary Section S5) thus took about 5 min, or 20 FOVs with ~11,000 EVs total about 7 min.

The choice of imaging parameters presents a trade-off between imaging time and final image noise, and thus detection sensitivity. The image noise in iSCAT is commonly minimized using cameras with high dynamic range and well depth, exposing pixels close to saturation, as well as averaging of multiple individual frames. Fig. 2c shows an image noise characterization for a Prime 95B (Photometrics) with a well depth of 80,000 e^- and readout noise of 1.6 e^- (see methods). The solid black line indicates the theoretical noise floor consisting purely of photon shot noise (square root of the number of detected photons) whereas the dashed black line includes technical camera readout noise according to manufacturer specifications. This

prediction was validated by taking pairs of iSCAT images (each averaged *n* times), subtracting the two to remove fixed-pattern noise, and calculating the standard deviation σ_n (black diamonds). The noise in SP data includes a second contribution from flatfield processing which adds noise depending on the number of averages n_{FF} of the flatfield image (crosses, color coded). We found that the measured noise saturated close to the expected noise floor (dashed lines) for small n_{FF} but exceeded the prediction for larger n_{FF} slightly. For the imaging parameters used throughout this manuscript, we predicted an approximate image standard deviation $\sigma = 5 \times 10^{-4}$ where the flatfield contributed less than a third. The noise floor for SP contrasts C_{SP} reported in SPFI plots below exceeded this value (Fig. 2d, note third root scaling) as our candidate spot localization registers contrast minima and thus samples at the extreme of the noise distribution. Additional noise overhead stemmed from small focus offsets between the averaged flatfield images and the actual FOV which produced artifacts shown in the inset (see also Supplementary Section S6).

Size photometry of silica nanoparticles

We characterized the sensitivity of SP using silica nanoparticles of known sizes and compared it to nanoparticle FCM using a CytoFLEX-S instrument (Beckman Coulter Life Sciences). We measured violet side scattering (VSSC using a 405 nm laser) and side scattering (SSC using a 488 nm laser) of aminated silica particles with 50 nm, 100 nm and 200 nm diameters in Fig. 3a. We found that the instrument was able to resolve 100 nm and 200 nm particles which is in line with an ~84 nm detection limit obtained by the manufacturer on an optimized instrument²⁹.

Notably, even for the 200 nm particles the signal was within the instrument noise level and could only be distinguished at higher sample concentrations and high detection rates.

We next took SP measurements under 440 nm and 740 nm iSCAT illumination wavelengths after adsorbing the particles onto plasma-activated glass coverslips. We show SP scatterplots for 50 nm, 100 nm, and 200 nm particles in Fig. 3b and representative post-incubation images for these particles as well as for 20 nm and 80 nm ones in Fig. 3c. We found that particles of all sizes could be detected (Fig. 3c) and that even 50 nm particles cluster separately from the background noise floor (Fig. 3b). The scatterplot for the 50 nm particles revealed large, noncircular iSCAT spots that are consistent with aggregates of two or more particles which thus likely dominated the corresponding FCM plot. Similar aggregates were also visible for 100 nm and 200 nm particles but produced excessively bright iSCAT contrasts and were discarded by the localization algorithm. For the 200 nm particles (rightmost panel in Fig. 3c), we observed a contrast inversion under 440 nm illumination which requires special attention for SP sizing. The basic model of interferometric imaging predict that contrasts C scale with particle diameters d to the third power²⁵ which we refer to as d^3 model. This behavior is expected to break for particles >> 100 nm when Mie scattering causes imbalance of forward and back scattering and increasing particle sizes requires incorporation of additional interferometric phase factors. In this manuscript, we limit ourselves to particle sizing using the d^3 model and hence report SP contrasts in the scatterplots (e.g. Fig. 3b) scaled to their third root $\sqrt[3]{C_{\lambda}}$ which are then proportional to the particle sizes. We combined measurements C_{λ} of 20 nm, 50 nm, 80 nm, 100 nm, and 200 nm particles for both wavelengths λ in Fig. 3d and fitted the first three samples to d^3 curves (solid lines, color-coded for wavelength). Using the fit, we determined an SP calibration for sizing silica particles using $d = \kappa_{\text{Silica}} \sqrt[3]{C_{440 \text{ nm}}}$ with proportionality factor

 $\kappa_{\text{Silica}} = 256 \text{ nm}$. This calibration can be used to estimate sizes of other particles with known refractive indices in the d^3 regime (see below) and, by assuming minimally resolvable contrasts of $\sqrt[3]{C_{440 \text{ nm}}} \ge 0.11$, determines an SP detection limit for silica spheres of $d_{\min} \approx 28 \text{ nm}$.

We continued to investigate the breakdown of the d^3 model and contrast inversion of $C_{440 \text{ nm}}$ using finite element (FEM) simulations in the COMSOL optics module (Fig. S5, Supplementary Section S7). The simulation data, plotted as dashed lines in Fig. 3d, reproduced the contrast inversion observed for 200 nm particles with 440 nm illumination and the absence of contrast inversion for 740 nm. While these simulations improve on the simple d^3 model, they still miss important effects such as the Gouy phase shift³⁰ (see also Fig. S7). Modeling of the full interferometric point spread function³¹ in future work will be instrumental to extend the SP sizing range by using information in the rings around large particles to alleviate ambiguity of contrasts around the inversion point.

SPFI of liposomes with dual wavelength iSCAT imaging

Fig. 4 shows SPFI data on fluorescently stained liposomes which were prepared from DOPC lipid incorporating 1% DiI dye by extrusion through track-etched membrane filters with a 400 nm pore size (representative images in Fig. 4a). Extrusion is known to produce liposomes with sizes both smaller and larger than the filter pore, however, the addition of the lipid dye allowed for an independent measure of total lipid content within each liposome. Fig. 4c shows SPFI scatterplots that were generated from the image data under 440 nm iSCAT illumination, and third root SP contrasts and DiI intensities were found to correlate well as shown by a linear fit to the data above $\sqrt[3]{C_{440 \text{ nm}}} > 0.1$ (dashed line). For many bright liposomes ($\sqrt[3]{I_{\text{DII}}} > 12$), SP

measurements underestimated liposome sizes due to the contrast inversion and many spots landed well above the fit line. Intriguingly, we found a maximum $\sqrt[3]{C_{440 \text{ nm}}} \approx 0.4$ which is even larger than the $\sqrt[3]{C_{440 \text{ nm}}} \approx 0.35$ that we measured for 100 nm silica beads. This indicated that liposomes flattened upon adsorption on the glass surface which would result in a decrease of the interferometric phase factors and thus a delay of the contrast inversion. Flattening has been documented, and moreover liposomes are often used to form supported lipid bilayers³². We performed COMSOL simulations in Fig. 4b showing that flattening of solid particles while maintaining the total mass indeed results in a delay of the iSCAT contrast inversion without affecting the initial d^3 scaling.

We confirmed that many larger EVs that underwent the contrast switch at 440 nm maintained dark iSCAT contrast when imaged under a longer 740 nm wavelength, as shown in Fig. 4d. Expectedly, this came at the cost of reduced sensitivity overall which we visualized by a SP scatterplot of both contrast measurements in Fig 4e. We found that in the region above the noise floor of the 740 nm channel and below the contrast inversion point of the 440 nm channel (purple dots in Fig. 4e) both measurements are roughly proportional to each other. Using a linear fit to the data (forcing a zero *y*-intersect), we determined a contrast proportionality factor $\eta = 2.6$ between the two wavelengths which was reasonably close to the theoretically predicted value of 2.8 based on wavelength scaling. Both measurements were combined to $\sqrt[3]{C_{SP}}$ in Fig. 4f, where $\sqrt[3]{C_{440 \text{ nm}}}$ was used for small liposomes (blue dots), $\eta \sqrt[3]{C_{740 \text{ nm}}}$ for large ones (red dots), and intermediately sized ones were averaged between the two measurements (purple spots). Combining measurements in two wavelengths using a SP correlation plot (Fig. 4e) can thus maximize the dynamic range for sizing polydisperse samples at the expense of dual measurements.

SPFI of extracellular vesicles

Small extracellular vesicles (sEV) from the supernatant of several cell lines were purified by size exclusion chromatography (SEC) and imaged using combined 440 nm and 740 nm illumination wavelengths. EV sizes were estimated by $d = \kappa_{EV} \sqrt[3]{C_{SP}}$ where the calibration factor $\kappa_{EV} =$ 357 nm was calculated from κ_{Silica} (c.f. Fig. 3) by assuming a uniform refractive index for EVs of $n_{EV} = 1.40$ (Supplementary Section S1). With $\sqrt[3]{C_{440 nm}} = 0.11$ as estimate for the noise floor (Fig. 2d), the calibration gives a lower detection threshold of SP for EVs with $0.11\kappa_{EV} \approx$ 35 nm diameters using our setup. We performed numerical calculations (Fig. S1) and FEM simulations (Fig. S6) comparing this solid sphere equivalent refractive index calibration to coreshell models^{25,33} which showed that core-shell geometries led to only minor corrections of EV sizes and led us to adopt a single valued $n_{EV} = 1.40$ for simplicity. Additionally, we found that experimentally measured SP contrasts at 440 nm illumination before contrast inversion can exceed the prediction of either solid or core-shell EV models, suggesting that EVs, just as liposomes, can flatten upon surface adsorption³⁴.

Fig. 5a shows an SPFI scatterplot of EVs from HT29 cells which were biotinylated via their membrane proteins using the membrane impermeable Sulfo-NHS-Biotin and then incubated with Alexafluor 674 (AF647)-labelled streptavidin. EVs were immobilized on glass coverslips precoated with poly-L-lysine (PLL) which is positively charged at neutral pH and captures EVs by electrostatic interaction. The scatterplot was divided into four quadrants (Q1 to Q4, labelled in Fig. 5a) by applying SP contrast and fluorescence intensity thresholds found in control experiments (Fig. S8). We note that biotinylation is not specific to EVs and spots in Q1 on the top/left include many biotinylated soluble proteins as well as streptavidin molecules that were not washed off. However, by correlating the fluorescence data to label-free iSCAT, these spots 13 can be efficiently removed prior to further analysis. Assuming a constant membrane protein density for EVs of different sizes, the fluorescence signal intensity is expected to correlate with EV surface area. We tested this prediction by a fit to $\sqrt[3]{I_{AF647}} \propto d^{2/3}$ for $\sqrt[3]{C_{SP}} > 0.1$, shown as solid line, which resulted in a correlation coefficient of $R^2 = 0.77$. Any linear fit to the data would result in non-negligible *y*-axis intercepts which cannot be explained purely by detection sensitivities.

To further corroborate these finding, a second sample of HT29 derived EVs was stained with Cell Tracker Deep Red (CTR) and captured on a plasma cleaned glass surface, Fig. 5b. Since CTR accumulates in the EV lumen, the fluorescence intensity is expected to scale with the EV diameter and hence $\sqrt[3]{I_{\text{CTR}}} \propto \sqrt[3]{C_{\text{SP}}}$. This scaling behavior was confirmed using a linear fit (solid line in Fig. 5b) to the datapoints above $\sqrt[3]{C_{\text{A}}} > 0.1$ which resulted in a correlation coefficient $R^2 = 0.81$. Similar to Fig. 5a, spots above $\sqrt[3]{C_{\text{SP}}} > 0.4$ showed systematic deviation from the fit, indicating potential underestimation of EV sizes by SP due to limited sizing range at 740 nm (see also Fig. S9 for SPFI plots of the individual wavelengths). Fig. S10 shows additional control experiments where EVs were co-stained with CTR and Carboxyfluorescein succinimidyl ester (CFSE) which both showed similar correlation with third root SP contrasts and similar detection sensitivities.

Next, we show a comparison of SPFI versus other single EV methods in Figs. 5c-5g on EVs from A431 cells expressing CD63 fused to green fluorescent protein (CD63-GFP) in addition to their wildtype CD63. CD63 is a transmembrane protein from the tetraspanin family and a common marker of small EVs due to its high abundance and role in exosome biogenesis. Fig. 5c shows FCM data of this sample which had a LOD of $d \approx 100$ nm and includes considerable

technical noise as seen from the buffer control (Fig. S11, size calibration in Fig. S13). Fig. 5d shows a negative stain TEM image which confirmed the presence of many EVs of 50 nm and below. Fig. 5e shows an SPFI scatterplot of the data which detected EVs down to \sim 35 nm in our calibration and shows non-uniform GFP fluorescence with two clear clusters. The first cluster comprises bright GFP⁺ spots ($\sqrt[3]{I_{GFP}} > 5$) on EVs with diameters up to ~100 nm with linear correlation between contrast and fluorescence. The second cluster consists of a much larger number of dim GFP⁺ spots ($\sqrt[3]{I_{GFP}} < 4$) on EVs up to only ~80 nm (see Fig. 5f for images of the EVs highlighted with arrowheads in Fig. 5e). The SPFI plot visually resolves the two clusters by combining EV size and CD63 expression data and the mismatch in size range between the clusters suggests that they represent different subpopulations. The same measurement in FCM, due to its higher detection threshold, could not resolve the two clusters. The iSCAT data furthermore allows to count the total EV number and we found that 5291 out of 10801 spots above the gating thresholds were positive for GFP, or in other words the overall EV labeling efficiency of the CD63-GFP transfected cells within our iSCAT detection range was 49 %. Fig. S11 shows a comparison of the same sample captured on a plasma activated glass and on PLL coated glass, and we found that the overall shape of the SPFI plot was largely comparable and both GFP+ clusters are readily identifiable on either surface. SP contrast measurements differed minorly, which will be further investigated in future work to improve sizing and include capture bias of different surfaces. Fig. S12 furthermore shows the same EVs purified via a 10,000 g ultracentrifugation protocol which led to the inclusion of much larger EVs in the sample. We found that our current single pixel contrast measurements indeed limit SP sizing at 740 nm to EVs of ca. 200 nm. However, the data is suggestive that improved sizing including a full interferometric PSF model will increase the range in the future (see Fig. S12).

DISCUSSION

We described how to upgrade standard epifluorescence microscopes for SP using iSCAT imaging with the addition of a broadband 50:50 beamsplitter for ~ US\$ 150. We were able to detect EVs down to $d \approx 35$ nm using high throughput pre- and post-incubation imaging over multiple FOVs, exceeding the detection sensitivities of commercial tools as shown in Fig. 6. We furthermore introduced SPFI which leverages FCM workflows using SP as orthogonal detection to size EVs and remove false fluorescence detection events from unbound labels or substrate autofluorescence. Our SPFI pipeline includes fully automated image pre-processing including iSCAT flatfield corrections and aberration corrections, automated FOV alignment, candidate spot detection, and contrast measurements using two SP illumination wavelengths. SPFI imaging consumes only 10 µl of sample per flow cell and more than 10,000 EVs and nanoparticles could be imaged within ~7 min for pre- and post-incubation images per iSCAT wavelength over 20 FOVs; more FOVs can be acquired if desired.

We validated SP on silica particles of various sizes and confirmed that SP contrasts scale approximately with d^3 for silica nanospheres up to about ~100 nm diameter using an imaging wavelength of 440 nm and about 150 nm under longer 740 nm wavelength illumination. Accurate sizing of EVs on the other hand is expected to be more challenging since iSCAT contrast values are dependent on both particle size and refractive index and EVs can vary in density due to different amounts of cargo. However, using fluorescent dyes as proxies for EV size, SPFI was able to distinguish surface and volumetric stains via the relationship between size and fluorescence intensity *I* that was proportional to d^2 and d^3 , respectively, which we believe to be the first direct experimental demonstration. EV sizing was furthermore established using a simple $C \propto d^3$ model based on scaling the silica reference data for an EV solid sphere equivalent refractive index of $n_{\rm EV} = 1.40$. We found that our SP sizing range was between around 35 nm to 200 nm which is a competitive dynamic range compared to NTA and FCM (see also Fig. S12). The measured sizes of SEC purified EVs up to 200 nm diameters were consistent between SP, NTA and FCM and we propose that the upper limit for SP can likely be increased through image analysis in which the microscope point spread function is explicitly modeled and ambiguity in EV size due to double-valued contrasts around the inversion point can be resolved. We note that reliable EV detection is possible for any EV larger than 35 nm with dual wavelength imaging since the contrast switches occur at different EV sizes and near invisible spots at either wavelength are readily registered by the other one (c.f. Fig. S12). Finally, we demonstrated the utility of SPFI for EV subpopulation discovery on CD63-GFP tagged EVs by comparing to other state of the art techniques where SPFI revealed that GFP was present in two clusters that differed in both GFP brightness and EV sizes. Notably, clear separation between the clusters was only seen for EVs below 100 nm in size and FCM thus only picked up diffuse expression on larger EVs.

A corollary of the above discussion is that the capture conditions including buffer, pH, salt concentration, and surface and EV charge could influence EV physisorption and introduce capture bias. We mainly captured EVs by electrostatic interaction using positively charged PLL coated coverslips. We found, however, that EVs could also be immobilized on negatively charged surfaces in PBS buffer, suggesting screening of the EV charge and major contribution by Van der Waals forces to EV adsorption. We found that electrostatic attraction on PLL increased capture rates but could also increase EV flattening on the surface. Flattening without material loss results in a reduction of the interferometric phase resulting from the out-of-plane

material. This extends the range in which $C \propto d^3$ as it delays the contrast inversion, which could increase the range of SP. Strong deformations, however, might lead to fluid loss from the EV lumen and thus additional complications for EV sizing. Future work will investigate whether EV deformability in different capture environments could open the door to all optical single vesicle stiffness or density measurements using SP.

Conclusions

SPFI of EVs is readily accessible by the addition of a 50: 50 beam splitter to any modern epifluorescence microscope and runs on an openly available image processing pipeline (*https://github.com/junckerlab/SPFI*). SP on pre- and post-incubation images can detect and size EVs from $d \approx 35$ nm to > 200 nm, and its combination with fluorescence imaging generates rich multimodal single EV datasets that could help identify EV characteristics and clusters in > 500 EVs per FOV, and > 10,000 in each microfluidic flow cell in our experiments. SPFI can distinguish membrane from intraluminal dyes in EVs, captures deformation of EVs, and can reveal subpopulations that are only distinguishable in EVs with d < 100 nm. Finally, SPFI should be broadly applicable to EVs and other rigid or soft, synthetic or biological nanoparticles.



Figure 1. Overview of combined EV size photometry (SP) and fluorescence imaging (SPFI). a Schematic of the microscopy setup including a microfluidic flow cell with EVs immobilized on a coverslip in an epifluorescence microscope with LED illumination, a high sensitivity sCMOS camera, and an off-the-shelf 50:50 beamsplitter for

iSCAT imaging. A motorized filter turret automatically switches between the dichroic mirror for fluorescence imaging, and the 50:50 beam splitter for iSCAT. **b** Workflow to suppress background inhomogeneities and surface bound crud. The background is registered in a pre-incubation iSCAT image, then EVs are immobilized, unbound EVs are washed, and a post incubation iSCAT image is taken. The difference of the two images efficiently removes small surface inhomogeneities and allows to measure accurate contrast values of small particles within an undulating background. **c** Workflow to generate SPFI scatterplots. Typically, the iSCAT image is used to register candidate spots and contrasts or fluorescence intensities are subsequently measured on all channels. **d** differential iSCAT of pre- and post-incubation image which is used for EV localization. pre: image before particle incubation; post: image

after particle incubation; SP: pre-image subtracted from the post-image. Three particles are highlighted: i is a particle that is already present before incubation and is excluded from the rest of the analysis; ii is a particle that is detected in fluorescence, but not in iSCAT, and iii is a particle that is detected both in iSCAT and in fluorescence. **e**

Fluorescence pre- and post-incubation images that were taken concomitantly to the iSCAT in **d**. The generation of a difference image in fluorescence is optional, and mostly inconsequential for substrates with low autofluorescence. **f** SPFI scatterplot of the processed data shown in **d** and **e**, taken on EVs from A431 cells that were stably transfected to express CD63 fused with green fluorescence protein.



Figure 2. SP noise characterization. a Image flatfield processing demonstrated on a raw post-incubation image and b flatfield corrected image, which is necessary to remove illumination inhomogeneities but adds image noise.
Shown is a 1200×1200 pixel ROI of a raw 16-bit iSCAT image which has close to 20% illumination inhomogeneity.

The color bar for the raw image is scaled to the bit depth of the image and dashed lines indicate limits used for plotting. The processed image is shown on a contrast greyscale of 1.5% around the image median. Dark spots are EVs from HT29 cells and circles indicate matching spots as a guide to the eye. **c** iSCAT image noise for a series of images collected at 40 FPS as function of number of images averaged. Black solid line represents an estimate of the photon shot noise and constitutes the lower theoretical noise limit whereas the black dashed line also includes the estimated camera readout noise. The black diamonds are measured SD of the image noise in the averaged timeseries, and crosses are for measured image noise after normalizing the timeseries with flatfield (FF) images consisting of $n_{\rm FF}$ frames. The latter is compared to the respective estimated noise floor (horizontal dashed lines) and used to set imaging parameters for the iSCAT endpoint measurements. **d** SPFI plot generated in a buffer control experiment over 20 FOVs using iSCAT at 440 nm wavelength and a red fluorescence channel (used in Fig. 5a) showing zero

events in the upper right quadrant. The inset shows a differential iSCAT image which is largely background-free

except several image artifacts indicated by white arrowheads.



Figure 3. Calibration, validation, and benchmarking of SP. a Nanoparticle flow cytometry (FCM) and **b** Size photometry (SP) scatterplots of a buffer control and silica nanoparticles with 50 nm, 100 nm, and 200 nm diameters. For FCM, the samples were diluted to matching concentrations based on the manufacturer specification, measured for 60 s each, and 5 % of total events are shown by their violet side scattering signal (VSSC-H at 405 nm) and side scattering signal (SSC-H at 488 nm). SP plots show iSCAT contrasts to the third root for 440 nm and 740 nm

illumination wavelength, respectively. The 50 nm particle signal in SP is clearly separated from the noise blob whereas for FCM no cluster is visible. **c** Representative SP images of particles of different sizes. All color scales are set from -1.0 % to +0.5 %, scalebar is 500 nm. Interference contrast inversion for 200 nm nanoparticles is observed for 440 nm illumination. **d** Theoretical (solid lines: fit to analytical model; dashed lines: calculated by finite element modeling) and experimental (dots) SP contrasts as function of particle size. Horizontal and vertical error bars are SD of the particle sizes specified by the vendor and SD of particle contrast, respectively.

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Figure 4. iSCAT at multiple wavelengths and fluorescence imaging of fluorescently stained liposomes. a Representative SP images of liposomes imaged at 440 nm (1, note white center in large dark spot indicated by a white arrowhead) and 740 nm (2) with aligned color scales (-1% to +0.5 %) and in fluorescence (3) to visualize their DiI stain. **b** FEM simulation of iSCAT contrasts at 440 nm of solid particles that remain spherical or flatten to 2/3 or 1/3 of the original height while maintaining a constant volume. Flattening allows larger particles to be measured by suppressing contrast inversion while increasing the maximal contrast. **c** SPFI scatterplot of DiI stained liposomes imaged with a 440 nm and **d** 740 nm iSCAT illumination, both showing a linear relationship albeit with different sensitivity (slope). **e** Scatter plot correlating the SP contrast at the two wavelengths. The purple dots indicate the size range at which the two contrast values are proportional to each other and is used to determine the proportionality factor η by a linear fit (solid black line). **f** Sizing of EVs using both wavelengths, with blue dots measured at 440 nm, red dots measured at 740 nm and purple dots averaged between the two measurements using their proportionality factor η. Gray dots fall below the iSCAT detection threshold using either of the wavelengths.

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Figure 5. SPFI scatterplots correlating EV size with fluorescence intensity. a SPFI scatterplot of biotinylated EVs that were stained with Streptavidin-AF647. Dashed lines are contrast and intensity thresholds separating real detection events from image noise. Quadrants Q1 to Q4 apply to all scatterplots. The solid line is a model assuming homogenous labelling of the EV membrane proteins. b SPFI scatterplots of EVs stained using cell tracker deep red (CTR) where the solid line is a model assuming dye uptake in the lumen of the EVs. c FCM measurements of EVs from A431 cells expressing CD63-GFP measured in the fluorescein (FITC) channel. d Negative stain TEM image of the CD63-GFP EVs. e SPFI scatterplot of the CD63-GFP EVs showing two distinct clusters of GFP+ spots, one group with $\sqrt[3]{I_{GFP}} > 4$ and EV sizes up to ~100 nm and a smaller group with $\sqrt[3]{I_{GFP}} < 4$ on EVs up to only ~80 nm, either of which could not be resolved in FCM (see c). f Post-incubation images of selected EVs identified by star symbols in f. Top row: GFP fluorescence; middle row: 740 nm iSCAT; bottom row: 440 nm iSCAT. Panels can be matched to the spots in f by the colors of their borders. Intensity bars show raw iSCAT and fluorescence contrasts

and scale bars are 500 nm.



Figure 6. Comparison of SPFI against established label-free detection techniques. SPFI sizing range and throughput are plotted for imaging with a 100× as well as a 60× objective (see Fig. S14). The SPFI throughput is calculated by assuming a capture density of 400 EVs / (100 µm)² together with the imaging parameters used in this manuscript (see methods, Fig. S4). FCM and NTA lines represent the commercial instruments used in this manuscript, whereas the dashed FCM line indicates the highest reported sensitivity on any FCM to our knowledge⁹. Shaded areas indicate commonly assumed size ranges for large EVs (IEVs), small EVs (sEVs) and smaller, non-

membranous nanoparticles (NPs).

Methods:

Microscopy setup

The iSCAT system was built on an inverted fluorescence microscope (Nikon TI2) using a 100× plan apochromat (λ series, Nikon) of NA 1.45 with a 1.5× intermediate magnification lens (inbuilt feature of the Ti2). We performed additional characterization of SP detection sensitivity using a 60× objective shown in Fig. S14, however, due to reduced sensitivity all experimental data shown was taken with the 100× objective. Imaging was done one a Prime95B (Photometrics) with 118×118 µm² FOV size on the full 1608×1608 pixel sensor. Illumination was provided via a standard fixed large FOV epifluorescence illuminator (Ti2-LAPP *fixed main branch*, Nikon) with built in fly-eye lens. The illuminator has an adjustable aperture which affected the achievable iSCAT contrasts (data not shown). All data for this manuscript were taken at the same setting which was optimized for highest contrast of 50 nm silica particles. iSCAT using multiple FOVs furthermore relied on an LED based reflectance autofocus (*Perfect Focusing System*, Nikon).

The fluorescence measurements were taken using a pentapass dichroic mirror (FF409/493/573/652/759-SP01, Semrock) and pentapass emission filter (FF01-432/515/595/681/809, Semrock). The emission was additionally filtered using 515/30 nm (GFP), 595/40 nm (SPI) or 680/42 nm (AF647) filter sets (Semrock) in an automated filter turret in front of the camera. Maximum illumination power of the Spectrapad was used in all channels and images were averaged 8× with 100 ms exposure time each.

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iSCAT imaging used the same light source (Spectra X Light Engine, Lumencor) which provides 370 mW output power in the 440/20 nm channel and 80 mW in the 740/20 nm channel at the maximum power setting according to manufacturer specifications. A filter cube with a widefield 50:50 beam splitter (BSW10R, Thorlabs) was added in the filter turret to switch between fluorescence and iSCAT modes. The frame exposure was maximised within the dynamic range of the camera, which resulted in exposure times around 4 ms at 440 nm and 25 ms at 740 nm. This allowed to image ROI of 512×512 pixels at 95 FPS, or 1200×1200 pixels at 40 FPS.

iSCAT noise characterization

The noise characterization in Fig. 2 was performed using the same settings as for regular imaging in FOVs of 512×512 pixels with camera exposure of 4.8 ms. The expected noise was calculated using the camera bit depth of $\delta = 2^{16}$, well depth of $\gamma = 80,000$ e-, quantum efficiency of $\eta =$ 95 % and readout noise $v = 1.6 e^{-}$. Due to illumination inhomogeneities, the maximal exposure is only achieved in a small region in the center of each frame while it drops off substantially towards the edges. We adjusted camera exposure time to achieve highest single pixel exposure of 98.5 % saturation which resulted in an average pixel exposure of $\rho = 75.4$ %. This corresponds to $n = \frac{\gamma}{\eta} \times \frac{\rho}{\delta} = 57.600$ captured photons per pixel for single frame averaging and allows to calculate the photon shot noise as $\sigma = \sqrt{n}$, shown as solid line in Fig. 2c. Since the camera readout noise was limiting the dynamic range to 50,000, we also calculated an adjusted noise floor by accounting for the readout noise as $\sigma' = \sigma_{sn}/\sqrt{v}$.

All points in Fig. 2c were generated using a dataset containing N = 5000 images (taken at a frame rate of ~500 images/s) of an empty coverslip with an average pixel saturation of ~90 % (similar to Fig. 2a). For each amount of averaging, we computed two images, I_U and I_E , where I_U

is the average of the first *n* uneven frames, and I_E is the average of the first *n* even frames of the original time series. This interleaved averaging of the frame made the analysis robust against slow illumination power fluctuations. We then calculated the difference $I_{diff} = I_E - I_U$ to eliminate all spatial information of the individual frames such as illumination inhomogeneities. We calculated the noise in I_{diff} using the standard deviation σ_{diff} . Since both individual images I_E and I_U contribute to σ_{diff} in equal parts, we can finally estimate the image noise for an iSCAT image averaging 2n frames as $\sigma_n = 1/\sqrt{2} \sigma_{diff}$.

To estimate the image noise after flatfield correcting the raw iSCAT frames, we generated two extra images, FF_U and FF_E , to mimic the real flatfield images. The stand-in flatfield images were summed from $n_{FF} = \{1, 10, 100, 1000\}$ additional frames of the timeseries which were not already used to generate F_E and F_U . F_U and F_E were then divided and subtracted as $F_{diff}^* = F_E/FF_E - F_U/FF_U$ such that the resulting noise can be calculated via the standard deviation of F_{diff}^* .

We note that additional background was present in real experiments which is described in Supplementary Section S6.

Flow cell preparation and sample incubation

All flowcells used in this manuscript used #1.5 glass coverslips (Fisher Scientific) as substrate. Coverslips were sonicated in acetone and isopropyl alcohol for 5 min each and rinsed in distilled water before blow-drying with nitrogen. The coverslips were plasma treated for 1 min and sealed into a custom hybriwell chamber (Grace bio-labs). The flow cells have a volume of 10 μ l with a rectangular center area of 6 mm × 8 mm with conical taper regions at the short sides (15 deg half angle) which leads to 1.5 mm diameter in- and outlets. The height of the flow cells is 150 μ m and a total of five such flow cells are contained on each hybriwell chamber. Initial filling of the flow cells is aided by capillary pressure after manual pipetting into the flow cell inlet. Subsequent sample filling or washing of the flow cell was done by pipetting onto the inlet while gently pressing a folded lint-free tissue (Kimwipe, Kimberly-Clark Professional) to the outlet. Flow cells were filled with the respective buffer (see sample descriptions) before taking the pre-incubation images. 10 μ l of sample was flushed into the flow cell and EV capture was monitored using iSCAT imaging (without flatfield processing). After the appropriate density was achieved (typically 1 – 3 mins), the flow cell was flushed with 100 μ l of buffer and the post-incubation images were taken.

Calibration samples

Aminated silica nanoparticles 20, 50, 80, 100, and 200 nm in diameter were purchased from nanocomposix (SIAN20-25M, SIAN50-25M, SIAN80-25M, SIAN100-25M, SIAN200-25M). Prior to use, the particles were sonicated and then diluted to approximate concentrations of 1×10^9 particles per ml in 2 mM acetate buffer at a pH of 5 where they attained positive surface charge.

Liposomes were prepared from DOPC lipid (Avanti lipids, cat. no. 850375) in chloroform solution which was spiked at a 1 % molar ratio with DiI dye in ethanol. The solvents were evaporated in a test tube under vacuum overnight. The lipid film was resuspended in PBS to a molar concentration of 10 mM and allowed to swell and form large multilamellar bubbles for 30 mins. The solution was then extruded by passing 21 times through a filter with pore size of 400 nm (Avanti lipids) and diluted further for the incubation.

Cell culture and EV isolation

Human colorectal cancer cells HT-29 (ATCC) and CD63-GFP modified A431 cells (CD63 with C-terminal fusion of GFP, a gift from the lab of Janusz Rak³⁵) were cultured in Dulbecco's modified Eagle Medium (DMEM) culture medium with 10 % FBS and 1 % (vol/vol) penicillinstreptomycin. The cells were incubated in a humidified incubator at 37 °C and 5 % CO₂ and grown to about 60 % confluency. The medium was then replaced by DMEM with 5 % EVdepleted FBS and the cell supernatant containing the EVs was harvested after ~48 hrs. The supernatant was spun at 400 g for 15 min to pellet cell debris and filtered by 450 nm syringe filters to remove larger EVs. For each EV sample, 5 ml supernatant was concentrated to 500 µl using 100 kDa Amicon Ultra-0.5 centrifugal filters (Millipore) for 5 mins at 5,000 g. The concentrated supernatant was purified using single use qEV.70nm size exclusion columns (IZON) with 0.8 ml void volume after sample loading and 0.6 ml sample collection volume.

CTR staining of EVs

EV samples were stained twice with CTR after it was found that the correlation of size to fluorescence intensity in the SPFI plots could be considerably increased compared to single rounds of staining (data not shown). Single staining resulted similar high staining efficiencies as reported for double staining (i.e. the number CTR+ EVs) but many large EVs were only relatively dim and variations between EVs of the same SP contrast were large. For the first staining round, 2 µl of Cell Tracker Deep Red dye (Invitrogen) at 1 mM concentration was added to 500 µl concentrated cell supernatant and incubated for 2 hrs at 37 °C. EVs were then purified using IZON qEV.70nm size exclusion columns as described above. Afterwards, a second round of staining was done using 2 µl of CTR in 500 µl purified EVs for 2 hrs at 37 °C. Excess dye was 31

removed in six washes with PBS using 100 kDa Amicon Ultra-0.5 centrifugal filters (Millipore) for 5 mins at 5,000 g (each spin concentrates 500 μ l sample to ~15 μ l).

Biotinylation and AF647-Streptavidin staining of EVs

Biotinylation was done using the EZ-Link Sulfo-NHS-LC-LC-Biotinylation Kit (Thermo Scientific). 0.5 mg of Sulfo-NHS-LC-LC-Biotin was added into 500 μ l of purified EVs (typically ~10⁹ EVs per ml measured by NTA) and incubated for 30 min at room temperature. Unreacted biotin was removed in six washes with PBS using 100 kDa Amicon Ultra-0.5 centrifugal filters (Millipore) for 5 min at 5,000 g. Subsequently, 1 μ l of Alexa Fluor 647 labelled streptavidin at 1 μ g/ml was added and incubated for 20 min. Excess streptavidin was removed in six washes with PBS using 100 kDa Amicon Ultra-0.5 centrifugal filters (Millipore) for 5 min at 5,000 g.

Flow cytometry measurements

All FCM data presented here was taken using a CytoFLEX-S Flow Cytometer (Beckman Coulter) equipped with 405 nm, 488 nm and 638 nm lasers, and operated using CytExpert Software. The 405 nm violet laser was used for side scattering (VSSC) and the 488 nm laser for forwards scattering (FSC). VSSC gain was set to 200, SSC gain to 500 and FITC gain to 1000, the VSSC threshold was set to 1400. Before each sample measurement, the system was washed using Azide buffer at pH 5 for silica nanoparticles, ultrapure water for PS calibration beads and PBS for EV samples. An event rate below 1000 per second was ensured for these buffer controls. Samples were loaded and after the event rate became stable and measurements were taken for 60 s each.

The size calibration (Fig. S13) was done in FCMPass³³ v. 4.2.14 using 81 nm, 100 nm, 152 nm, 203 nm, 240 nm, 303 nm, 345 nm, 401 nm and 453 nm polystyrene beads (ThermoFisher Scientific Cat no. 3080A, 3100A, 3150A, 3200A, 3240A, 330A, 3350A, 3400A, 3450A). The calibration was done using the default *Average RI* core-shell model with a 5 nm thick shell of refractive index 1.4863 and core refractive index of 1.3859.

NTA measurements

NTA measurements were performed on a Nanosight NTA 3.4 using a 488 nm laser. Five measurements of 60 s each were taken per sample at a syringe pump speed of 40 at 25 frames per second and detector threshold of 5. The camera level was set to 14 for the SEC sample and 10 for the 10,000 g UC sample to keep overexposure below 10 % of traces.

Negative stain TEM imaging

Copper grids were negatively charged for 20 sec at 20 mA and 5 µL of EV sample was incubated without dilution after SEC purification for 10 min. The grids were washed three times in ultrapure water for 30 s each and fixed in 2 % glutaraldehyde solution for 2 min followed by additional three washes in ultrapure water for 30 s each. Staining was performed with a 2% uranyl acetate solution for 45 sec, the grids were blotted dry using filter paper and residual moisture was left to evaporate for 1 hr. TEM images were taken on a FEI Tecnai 12 at a working voltage of 120 kV.

Data availability

Matlab scripts for evaluating SPFI data, including all image pre-processing, alignment and spot localization algorithms detailed in Supplementary Sections 2, 3 and 4 are available via https://github.com/junckerlab/SPFI.

All raw data supporting the main findings of this manuscript can be found at https://doi.org/10.5281/zenodo.10208005.

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Declaration of Interest Statement

The authors declare that the research was conducted in the absence of any commercial or

financial relationships that could be construed as a potential conflict of interest.

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