Apical constriction requires patterned apical surface remodeling to synchronize cellular deformation

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Abstract   Apical constriction is a basic mechanism for epithelial morphogenesis, making columnar cells into wedge shape and bending a flat cell sheet. It has long been thought that an apically localized myosin generates a contractile force and drives the cell deformation. However, when we tested the heighten apical surface contractility in a cellular Potts model simulation, the constriction increased a pressure inside the cell and pushed its lateral surface outward, making the cell into not the wedge shape but instead a drop shape. To keep the lateral surface straight, we considered an alternative model in which the cell shape was determined by cell membrane elasticity and endocytosis, and the increased pressure is balanced among the cells. The cellular Potts model simulation succeeded in reproducing the apical constriction, and it also suggested that a too strong apical surface tension might prevent the tissue invagination.

INTRODUCTION

The apical constriction is a general and fundamental process to bend or invaginate a flat epithelial tissue (*Martin and Goldstein, 2014*). In the flat epithelium, cells are in a columnar shape with apical-basal polarity. During the apical constriction, the columnar cells decrease its apical surface and turn into wedge shape, inducing a small angle between its opposite lateral surfaces. This cellular deformation is synchronized in a cluster of cells, and the angle of the lateral surfaces is summed across the cells to a larger tissue scale curve. Studies in genetics and molecular biology had revealed an involvement of the motor protein non-muscle myosin II which crosslink and contract actin filament network. The actomyosin network is formed beneath the cell apical surface (cortical actomyosin) and lining apical-lateral cell-cell interface (circumferential actomyosin), linked with adherens junction and tight junction so that it makes a supracellular continuous structure (*Martin et al., 2010*; *Takeichi, 2014*).

In *Drosophila melanogaster* mesoderm invagination, ventral cells express transcription factors Twist and Snail which then induce expression of numerous genes including regulators of the actin and myosin, such as T48, Fog, Mist, RhoGEF2, Rho, and Rock (*Martin, 2020*). The cortical and the circumferential actomyosin network are formed in the cells, and the cells undergo the apical constriction, while down-regulation of those regulators impaired the localization and activation of the myosin, and delayed or prevented the apical constriction (*Dawes-Hoang et al., 2005*; *Xie et al., 2016*; *Kölsch et al., 2007*; *Mason et al., 2013*). Similar regulation of the myosin by Rho and Rock was found responsible for the apical constriction in other organs (*Brodu and Casanova, 2006*;

Chung et al., 2017; *Guru et al., 2022*; *Manning et al., 2013*; *Röper, 2012*) and other insects (*Münster et al., 2019*; *Benton et al., 2019*; *Urbansky et al., 2016*). In vertebrate, the cortical myosin was not observed but the circumferential myosin was enriched, recruited together with Rock by PDZ domain containing protein Shroom 3, which was also found responsible for the apical constriction (*Nishimura and Takeichi, 2008*; *Lang et al., 2014*; *Plageman et al., 2010*; *Lee et al., 2007*; *Chung et al., 2010*; *McGreevy et al., 2015*; *Haigo et al., 2003*; *Hildebrand, 2005*). Since the interruption to those pathways reduced the myosin accumulation and prevented the apical constriction, it had been modeled that the apical myosin generated a contractile force on the cell apical surface and drove the cell and tissue deformation.

The myosin-based model was validated by simulations (*Rauzi et al., 2013*) with finite element model (*Conte et al., 2012*; *Perez-Mockus et al., 2017*) or vertex model (*Sherrard et al., 2010*; *Polyakov et al., 2014*; *Inoue et al., 2016*; *Pérez-Vonzález et al., 2021*). In the finite element model, a cell was represented by few quadrilateral elements and its cytosol was assumed to be viscous. The contractile forces on the cell apical, lateral, and basal surfaces were inferred from deformation of cells of the same representation (*Brodland et al., 2010*). The forces were varied in the simulation to show that the surface tension of the invaginated cells were more influential than those of surrounding cells, and that the invagination was robust against the variance in the force. In the vertex model, a cell was represented by a quadrilateral, polygon, or polyhedron with a surface tension on each edge or face and constraint on an area or volume inside. The differential surface tension was assigned based on the observation of myosin distribution in a modeled tissue, and varied to investigate its contribution. These studies succeeded in reproducing the apical constriction by the heighten apical surface tension, and further investigated additional components, such as differential lateral and basal surface tension, boundary condition by an extra cellular matrix, and compression by the surrounding cells, to make the resultant simulated deformation more similar to the modeled actual tissue.

Along with the cell autonomous deformation with the apical constriction, many epithelial tissues bending and invagination were found mechanically promoted by surrounding cells. For example, actomyosin was localized to a boundary between amnioserosa and lateral epidermis (*Hutson et al., 2003*; *Hayes and Solon, 2017*) or neural plate and epidermal ectoderm (*Hashimoto et al., 2015*; *Galea et al., 2017*, *2018*). This supracellular cables draw arcs and merge at a midline like a canthus so that a contraction of the cable results in the decrease of the inner amnioserosa or neural plate apical surface. Similar supracellular myosin cables were found encircling salivary gland (*Röper, 2012*; *Chung et al., 2017*) and tracheal pit (*Nishimura et al., 2007*; *Ogura et al., 2018*) during their invagination. Also, a planar compression was proposed to be able to drive the epithelial folding (*Brezavšček et al., 2012*).

In this study, we revisited the driving mechanism of the apical constriction and investigated a cell physical property required for the cellular deformation. First we modified a cellular Potts model, which was another framework to simulate cellular tissue (*Graner and Glazier, 1992*), so that it could represent an epithelial tissue, and simulated an effect of the apically localized myosin. Contrary to expectations, the cellular Potts model did not reproduce the apical constriction but the cells were delaminated one-by-one from edges. By analysing a change in energy with respect to the cell shape, we found that a force to contract the cell apical surface was substantially affected by the entire cell shape. Next we considered an alternative driving mechanism for the apical constriction incorporating a cell apical surface elasticity, which might be regulated by endocytosis. It succeeded in reproducing the apical constriction and epithelial bending. In addition, the encircling supracellular myosin cable largely promoted the invagination by the apical constriction, suggesting that too high apical surface tension may keep the epithelium apical surface flat.

RESULTS

Extended cellular Potts model to simulate epithelial deformations

To model cellular and tissue deformations, we use the cellular Potts model (*Graner and Glazier, 1992*). The model was first proposed to simulate cell sorting by assigning differential cell adhesion, which was expressed by a contact energy. In the basic cellular Potts model, cells are represented by sets of sites in a 2D lattice (*Figure 1*a), and their deformation is simulated by repeatedly updating a label on a randomly chosen site (*Figure 1*b, b'). For a site i, the contact energy is calculated by $\sum_{j \in N(i)} J(\sigma(i), \sigma(j))$, where $N(i)$ is a neighborhood of i, $\sigma(i)$ represents a cell to which i belongs, and J defines an affinity between the cells (*Figure 1*c). The contact energy is summed among all sites to represent a surface contact energy. Later, in various studies of cell and tissue-scale phenomena, the cellular Potts model was employed with additional elements and was extended to 3D. For example, *Fortuna et al.* (*2020*) simulated a cell migration with material labels representing nucleus, cytoplasm, and lamellipodium within each cell. To simulate morphogenesis of an epithelial tissue with apico-basaly polarized cells, the intracellular material labels were utilized to represent cell apical surface, lateral cell-cell junction, and basal surface (*Belmonte et al., 2016*; *Adhyapok et al., 2021*).

(*Belmonte et al., 2016*; *Adhyapok et al., 2021*). Also, the cellular Potts model was combined with a finite element model so that it simulated how cells generated and were pulled by stress/strain on the extracellular matrix (ECM) (*van Oers et al., 2014*; *Rens and Merks, 2017*).

To simulate the epithelial tissue morphogenesis, we developed a new cellular Potts model with the cell apical-basal polarity. The monolayer tissue was represented in a 2D lattice in order to save a computational cost, and while the general 2D cellular Potts model cast the epithelium plane to the lattice, our model cast the tissue section perpendicular to the plane. In this model, the cell was represented by a set of sites labeled with a cell identity σ (*Figure 1*a), a cell type, and a cytosol τ corresponding to the cell type (*Figure 1*d). The cell apical and basal surfaces, lateral cell-cell junction, and apico-lateral sites where the adherens junction might locate were marked (*Figure 1*e, f, f'). This cell representation also cast the vertical section of the cell. Outside the epithelial cells, sites were labeled as a medium (*Figure 1*a) and with materials representing an apical ECM and fluids in the apical and basal sides (*Figure 1*d). In this study, we assumed the cell surface tension consisting of contractility and elasticity, where the contractility is not dependent on a surface width nor the strain but constant and to decrease the surface, and the elasticity is proportional to the surface strain and to return the surface to its original width. In our cellular Potts model, the surface contractility was defined by the contact energy, and expressed by

$$\mathcal{H}_{sc} = \sum_i \sum_{j \in N(i)} J(\delta(\sigma(i), \sigma(j)), \tau(i), \tau(j)), \tag{1}$$

where δ is the Kronecker delta, σ includes the cells and medium, τ includes the cytosols, apical ECM, and fluids, and J defines the contact energy between the materials $\tau(i)$ and $\tau(j)$ when the sites i and j are in different cells or medium, i.e., $\delta(\sigma(i), \sigma(j)) = 0$. Then, the cell apical surface contractility J_a was defined by the contact energy between the cytosol and the apical side medium fluid, the lateral cell-cell junction J_l contractility was defined by the contact energy between the cytosols, and the basal surface contractility J_b was defined by the contact energy between the cytosol and the basal side medium fluid. For the surface elasticity, a perimeter length of the 2D cell was evaluated. In the same way with the surface contact energy, a perimeter length P of a cell and its apical, lateral, and basal parts P_a, P_l, and P_b were calculated by summing together the adjacency of the sites to different cells, and the surface elasticity was calculated by

$$\mathcal{H}_{set} = E_s (P - P_0)^2 \tag{2}$$

for the total perimeter, or by

$$\mathcal{H}_{sep} = E_s ((P_a - P_{a0})^2 + (P_l - P_{l0})^2 + (P_b - P_{b0})^2) \tag{3}$$

for the partitioned perimeters, where E_s denotes a surface elastic modulus and P_0, P_{a0}, P_{l0}, and P_{b0} represent reference values. To simulate the effect of the supracellular myosin cable encircling the invaginated cells, we added a potential energy to the model. The supracellular myosin cable is formed along the adherens junction and pulls it horizontally toward a center of the invagination in the tissue plane. The potential energy was defined by a scalar field and made a horizontal gradient, and the scalar values for the adherens junction sites were averaged in each cell:

$$\mathcal{H}_{\text{mc}} = \frac{1}{|X|} \sum_{i \in X} U(i), \tag{4}$$

where X denotes a set of the adherens junction sites in the cell, and U denotes the scalar field. The model also included an area constraint

$$\mathcal{H}_{\text{ac}} = \lambda(A - A_0)^2, \tag{5}$$

where λ denotes a bulk modulus, A represents an area of a cell, and A_0 denotes its reference value. By combining the terms (*Equation 1-Equation 5*), a tissue's total energy $\mathcal{H}$ was calculated. The cellular Potts model simulates the cell and tissue deformation by updating the label when it decreases $\mathcal{H}$, or when it increases $\mathcal{H}$, with a probability of $e^{-\Delta H/T}$, where T denotes a fluctuation allowance of the system (*Figure 1*g).

Before simulating the apical constriction, we first examined parameters for a stable epithelium. The energy $\mathcal{H}$ included only the terms of the contact energy (*Equation 1*) and the area constraint (*Equation 5*), and thus the surface tension was determined by the contact energy. When the cell basal and apical surface tensions were equal to the lateral cell-cell junction tension, cells were rounded and disconnected from each other, and the epithelium could not maintain the integrity (*Figure 1—figure Supplement 1*). Therefore, we assigned 2 or 4 times higher contact energy to the basal and apical surface than the lateral cell-cell junction in the following simulations.

Tissue deformation by heighten apical contractility simulated with cellular Potts model

To test whether the apically localized actomyosin could drive the apical constriction, we simulated the epithelial cell monolayer with heighten apical surface contractility. Some preceding studies assumed that the apical myosin generated the contractile force (*Sherrard et al., 2010*; *Conte et al., 2012*; *Perez-Mockus et al., 2017*; *Pérez-Vonzález et al., 2021*), while the others assumed the elastic force (*Polyakov et al., 2014*; *Inoue et al., 2016*; *Nematbakhsh et al., 2020*). It seems natural to consider that the myosin generates a force proportional to its density but not to the surface width nor the strain. Also, an experiment with a cell cultured on a micro pattern showed that the myosin activity was well corresponded by the contractility, and its effect on the elasticity, if any, was minor (*Labouesse et al., 2015*). Therefore, we assigned to cells an additional apical surface contractility by setting J_a larger than J_b for the center pale blue cells.

In contrast to the vertex model, our simulations could not reproduce the apical constriction, but cells with the heighten apical surface contractility were covered by surrounding cells and delaminated one by one from edges (*Figure 2*a, *Figure 2—video 1*). When the compression $A_0/\bar{A}$ and the entire surface contractility was small, the covering cells were horizontally spread and their surface was extended largely during the covering. To prevented the surface expansion, we included the surface elasticity term (*Equation 2*). It kept the cells surface perimeter in a narrower range as expected, but still the center cells were delaminated and resulted in similar shapes (*Figure 2*b). We also simulated deformation with the supracellular myosin cable but without the heighten apical surface contractility. Either with or without the surface elasticity, the center cells were covered by the pulled cells and delaminated (*Figure 2*c).

These results suggest that the heighten apical surface contractility and the supracellular myosin cable cannot drive the apical constriction. It was not limited to specific parameters but reproducible for various contractilities. It contradicts the current model that the apical constriction is driven by the heighten contractility generated by the apically localized myosin.

Figure 1. Cellular Potts model for an epithelial tissue. (**a**) Cells representation in the cellular Potts model. The cells were colored differently, and a white space represents the medium. (**b**, **b'**) An update of a label from **b** to **b'**. The panels show an area inside a black line in **a**. The label on a randomly selected site (yellow arrowhead) was replaced with one on a randomly selected neighboring site (magenta arrowhead). (**c**) Neighborhood of a site for the contact energy. The panel shows an area inside a blue line in **a**. A site marked with a circle is adjacent to three sites in a different cell (black star) and a site in the medium (white star). (**d**) Material labels in the model. Four types of cytosol are colored blue, green, red, and dark gray. Inner and outer body fluid are colored white and light gray, and the apical ECM is colored dark yellow. (**e**) Cell polarity markings. The apical, basal, lateral, and none-polar surfaces are colored blue, reg, green, and gray respectively. The adherens junction is colored yellow. (**f**, **f'**) Tissue and cell representations in our model. The panel **f'** shows an area inside a black line in **f**, and lightens an area around a cell. (**g**) Algorithm of the cellular Potts model simulation.

Figure 1—figure supplement 1. Epithelial cell surface tension and tissue integrity.

Figure 2. Simulations of epithelial tissue with the heighten contractility. (**a**) The ratio $A_0/\bar{A}$ indicates how much cells were apico-basally compressed by the lateral cell-cell junction contractility J_l. The horizontal axis J_b/J_l indicates a ratio between the basal surface and lateral cell-cell junction contractilities. The vertical axis J_a/J_b indicates a ratio between the apical and basal surface contractilities in center pale blue five cells. The other surrounding pale red and green cells were assigned the apical surface contractility equal to the basal surface contractility. (**b**) Simulations w/o the surface elasticity, where $A_0/\bar{A} \sim 1.01$, $J_b/J_l = 4$, and $J_a/J_b = 1.6$. The vertical axis represents time (1 time/1000 updates). (**c**) Simulations of the supracellular myosin cable. Cells marked with an asterisk were assigned a potential energy on their adherens junction so that they were pulled toward the midline. $A_0/\bar{A} \sim 1.01$ and $J_b/J_l = 4$.

Figure 2—video 1. Simulation of epithelial tissue with the heighten contractility. The center pale blue five cells were assigned a heighten apical surface contractility: $A_0/\bar{A} \sim 1.01$, $J_b/J_l = 4$, and $J_a/J_b = 1.6$, where $A_0/\bar{A}$ indicates how much cells were apico-basally compressed by the lateral surface contractility J_l, J_b/J_l represents a ratio between the basal and lateral surface contractilities, and J_a/J_b represents a ratio between the apical and basal surface contractilities in center pale blue five cells. There were 1000 updates between frames.

Energy landscape with respect to cell shapes

To investigate what made the difference between the vertex model and the cellular Potts model results, we first looked into an early deformation of the cells in the cellular Potts model (*Figure 3*a, b). It was characterized by 1) that only the edge cell shrank its apical surface, 2) that inner cells remained in the columnar shape, and 3) that a lateral surface of the edge cell was curved outward. These could be explained with a balanced apical surface contractility. For the sake of simplicity, consider a system of vertices linked by edges in a line, and the edges bear differential contractilities. For a vertex, a force exerted on the vertex is a sum of the opposing contractile forces on the linking two edges, and it will be zero when the two edges bear an equal contractility. In the same way, a force exerted on an apical cell-cell junction by the apical surface contractility was a sum of them and it would be zero when two adjacent cells bore an equal apical surface contractility. Then, a practical pulling force on the cell-cell junction by the cell apical surfaces remained only on the edge of the cells with heighten apical surface contractility (*Figure 3*c, c'). The edge cell shrank its apical surface and consequently decreased its volume. The decrease in the volume increased the pressure, and a difference in the pressure between the cells resulted in the lateral surface curvature.

Since the cellular Potts model results seemed acceptable, we next examined why the inner cells could shrink their apical surface in the vertex model by analyzing the energy with respect to a cell shape. The cell shape was represented by vertices and edges linking the vertices, and its energy function was defined by a combination of the pressure and the surface contractility. While the conventional vertex model assumes only straight edges, we included also curved edges. *Figure 3*d shows cell shapes with which the energy of the cell was at a minimum for an apical surface width and a curvature of the lateral surface. When the cell lateral surface was restricted to be straight, a plot of energy with respect to the apical surface width was well fitted by a quadratic polynomial ($y = 0.0216x^2 - 4.8409x + 5062.8$, $R^2 = 0.9999$) (*Figure 3*e, f), indicating that the practical force to constrict the apical surface was attenuated by the deformation. Because of the attenuation, the constricting force was unbalanced between the shrunk edge cell and the inner columnar cell, and it allowed the inner cells to shrink the apical surface simultaneously. On the other hand, when the cell lateral surface was allowed to be curved and the cell changed its shape along a gradient of the energy landscape (*Figure 3*e), a plot of energy with respect to the apical surface was well fitted by a linear function ($y = -4.3135x + 5055.3$, $R^2 = 0.9992$) (*Figure 3*f), indicating that cell could shrink its apical surface without attenuating the constricting force. When the practical constricting force remained constant, the inner cell would remain in the columnar shape. These results indicate that the difference between the vertex model and the cellular Potts model results was due to the straight lateral surface, while it seems not natural for the epithelial cells. Also, we expect that the cellular Potts model simulation with the heighten apical surface contractility reproduced a phenotype of endocytosis down-regulation as discussed later.

Apical constriction by modified apical elasticity

Since the cellular Potts model simulation indicated that the heighten apical surface contractility could not drive the apical constriction, we revisited requirements for the apical constriction. It requires 1) a mechanism to simultaneously shrink the apical surface among cells in a cluster, and 2) a mechanism to keep the lateral surface straight. Actually, the condition 2) can be satisfied when the condition 1) is achieved, because the simultaneous apical shrinkage will increase the pressures in the cells, but the increase will also be synchronized and thus the pressure will be balanced. As suggested by the cell shape energy analysis (*Figure 3*f), the apical shrinkage can be synchronized when the energy function and its derivative are both increasing with respect to the apical width. The energy function of the partitioned perimeters elasticity (*Equation 3*) has a such feature, and it may correspond to the stretching elasticity of cell membrane (*Helfrich, 1973*), where the apical reference value P_{a0} can be modified by apical endocytosis/exocytosis. By decreasing P_{a0} with the endocytosis, the cell membrane would contribute to the apical surface tension proportionally to the change $P_a - P_{a0}$.

Figure 3. Cell shape and practical force to constrict the apical surface. (**a**) The edge cell at the early phase of the cellular Potts model simulation with the heighten apical surface contractility. (**b**) Illustration of the edge cell. (**c, c'**) Illustration of the apical surface contractile force acting on the cell-cell junction (**c**), and a summed result of the forces (**c'**). (**d**) Map of cell shapes. For the apical width and a curvature of the right side lateral surface, the energy of the cell is at minimum with the shape. The pressure and the surface contractility were set so that the cell took the columnar shape for the apical width w, apical-basal height h, and 0 curvature. (**e**) Energy landscape of the cell shapes for the apical width and the lateral curvature. Red line shows a path following a gradient of the energy. (**f**) Plots of energy with respect to the apical width. Blue plot shows the energy when the lateral surface was restricted to be straight. Red plot shows the energy along the path in **e**.

To validate our new model, we simulated the epithelial cell monolayer with the decreased elastic reference value (*Figure 4*a, *Figure 4—video 1*). The inner cells shrank the apical surface simultaneously and they were deformed into the wedge shape with fluctuating but almost straight lateral surfaces, succeeded in invaginating the tissue with various parameters. Note that the simulation results were not at a steady state but in the middle of the deformation. The simulations were paused at an arbitrary time point because it was hard to assess whether the tissue was in the steady state or not, and resultant configurations were checked whether the cells were delaminated or not, and evaluated the invagination by comparing the center pale blue cells and surrounding pale green cells in an average distance to the apical ECM. Interestingly, the tissue apical surface was kept relatively flat when the surface elastic modulus was high even though the cells became the wedge shape (*Figure 4*b).

Variation in configuration and resultant deformations

Next we simulated the apical constriction by the elasticity remodeling with various conditions. The apical constriction could bend the tissue with thinner and thicker cells (*Figure 5—figure Supplement 1*), demonstrating that the model was robust against cell shapes. When the surrounding cells were assigned an affinity to the apical ECM, the cells were spread on the ECM and pushed the center cells toward the midline, and the invagination was slightly promoted (*Figure 5—figure Supplement 2*).

We also simulated the apical constriction by the elasticity remodeling with additional supracellular myosin cable (*Figure 5*). The edge pale green cells adjacent to the center pale blue cells were assigned the potential energy so that their adherens junctions were pulled toward the midline (*Equation 4*). Surprisingly, the tissue invagination was largely promoted by the supracellular myosin cable with a wide range of the pulling force and also wide ranges of other parameters. Note that the supracellular myosin cable alone could not reproduce the apical constriction (*Figure 2*c), and that the cell surface elasticity, with which the tissue was kept almost flat unless the supracellular myosin cable was added, were sufficient to bend the tissue with enough high pulling force on the adherens junction.

Figure 4. Simulations of epithelial tissue with the modified surface elasticity. (**a**) Results of the simulations. The center pale blue five cells were assigned $P_{a0} = 0$, while the others were assigned P_{a0} equivalent to P_a initial value. The ratio $A_0/\bar{A}$ indicates how much cells were compressed, J_b/J_l indicates a ratio between the basal surface and lateral cell-cell junction contractility, and E_s denotes the surface elastic modulus for the inner pale blue and green 13 cells. The edge pale red cells were assigned 0.1 times smaller surface elastic modulus than the inner cells. (**b**) Plots showing difference in distance from the apical ECM between the constricting cells and the surrounding cells. Average distances were compared, and a larger difference indicates a deeper invagination. Magenta horizontal lines indicate 0.73, an average difference between the center cells and surrounding cells when all of the cells were assigned P_{a0} equivalent to the P_a initial value, as a control. Results of three simulations were averaged.

Figure 4—video 1. Simulation of epithelial tissue with the decreased elastic reference value. The center pale blue five cells were assigned a decreased apical surface elastic reference value $P_{a0} = 0$. The cells were assigned $A_0/\bar{A} \sim 1.01$ and $J_b/J_l = 2$, and the inner pale blue and green cells were assigned $E_s = 2.5$ while the edge pale red cells were assigned $E_s = 0.25$. There were 10000 updates between frames.

Figure 5. Simulation of apical constriction with the supracellular myosin cable. The center pale blue five cells were assigned $P_{a0} = 0$, and the siding pale green two cells adjacent to the center pale blue cells were assigned the potential energy on their adherens junction so that they were pulled toward the midline. The magnitude C_r indicates a gradient of the potential energy, $A_0/\bar{A}$ indicates how much cells were compressed, J_b/J_l indicates a ratio between the basal and lateral cell-cell junction contractility, and E_s denotes the surface elastic modulus for the inner pale blue and green 13 cells. The edge pale red cells were assigned 0.1 times smaller surface elastic modulus than the inner cells.

Figure 5—figure supplement 1. Simulation with various cell heights

Figure 5—figure supplement 2. Simulation with cell-ECM adhesion

Inferred pressure distribution around *Drosophila* embryo tracheal pit

In the new model, we proposed that the simultaneous increase in the pressure would keep the cell lateral cell-cell junction straight. However, it is also possible that other mechanisms may keep the lateral cell-cell junction straight, for example by decreasing cell volume so that the pressure is constant and balanced among the cells. To check if the cell pressure was increased in an actual tissue, we implemented a geometrical tension inference around an invaginated tracheal pit of *Drosophila* embryo. The geometrical tension inference is a method to measure relative junctional tensions and cell pressures from shapes of the cells (*Roffay et al., 2021*), and we extended 2D Bayesian method (*Ishihara and Sugimura, 2012*) to 3D for the measurement of invaginated tissue.

The tracheal pit is formed in a tracheal placode by the apical constriction (*Nishimura et al., 2007*; *Kondo and Hayashi, 2013*). An initially flat apical surface of the tissue was bent to make the pit (*Figure 6* a). During the process, there were no obvious change in a distribution pattern of cellular junctional tension inside and around the tracheal pit (*Figure 6*b). On the other hand, the apical surface area decreased and the relative pressure increased among the invaginated cells (*Figure 6*c-e). Note that variances of the apical area and the relative pressure did not increase (*Figure 6*d, e), suggesting that an actual pressure value was also increased simultaneously by the cell deformation. These results were consistent with our model, indicating that the lateral cell-cell junction of the tracheal pit cell was kept straight by the increasing but balanced pressure.

DISCUSSION

In this study, we demonstrated that the heighten apical surface contractility could not drive the apical constriction, and proposed the alternative driving mechanism with the apical surface elasticity remodeling. When the cells were assigned the high apical surface contractility, the cells were rounded and delaminated one by one from edges in the cellular Potts model simulation. This simulation results disagree with the conventional model of the apical constriction (*Martin and Goldstein, 2014*; *Takeichi, 2014*; *Martin, 2020*) and its vertex model simulation results. However, the vertex model simulation was substantially affected by the straight lateral surface (*Figure 3*d-f), which was not natural for the epithelial cells. Therefore we concluded that the cellular Potts model better predicted a result of the heighten apical surface contractility than the vertex model. For the apical constriction, we considered the modified apical surface elasticity as a plausible underlying mechanism, where the synchronized decrease in cell volume and increase in pressure balances the pressure between the cells and keeps the lateral cell-cell interface straight. The requirement of simultaneous contraction is also consistent with the non-autonomous disruption of apical constriction (*Guglielmi et al., 2015*; *Galea et al., 2021*). With the elasticity-based apical constriction model, the cellular Potts model reproduced a tissue invagination. Also, we confirmed the increase in pressure among invaginated cells in *Drosophila* embryo tracheal pit. Interestingly, the tissue apical surface remained relatively flat when the cells were assigned higher surface elastic modulus. Though it might seem contradicting that the surface elasticity was the dominant factor for the apical constriction, it could be explained that a too high surface tension tried to keep the apical surface short and thus straight as a result.

If the apical surface elasticity was the determinant factor for the apical constriction, then how does the apical surface contractility generated by the myosin contribute to it, and how is the apical surface elasticity regulated by molecules? In many tissue, the apical myosin was found not permanently active but its flow, activation, and resultant apical surface contraction was pulsed, followed by a partial relaxation (*Martin et al., 2009*; *Solon et al., 2009*; *Martin et al., 2010*; *Mason et al., 2013*; *Booth et al., 2014*; *Kerridge et al., 2016*; *Mason et al., 2016*; *Krueger et al., 2020*; *Martin and Goldstein, 2014*; *Miao and Blankenship, 2020*). The contraction and partial relaxation is iterated to decrease the apical surface progressively, and it is called ratchet model or ratcheting mechanism. While it is yet to be clarified how the ratcheting prevents the apical surface from expanding back to its original width during the relaxation, some studies pointed out an involvement of the endocytosis

Figure 6. Change in junctional tension and cell pressure distribution during tracheal pit invagination. (**a**) Vertices and edges representation of adherens junction inside and around *Drosophila* embryo tracheal pit. The panels show two tracheal pits from the beginning of the invagination (0 min) and after 15 and 30 minutes. Black bars on top and at left side indicates positions of the y-z and x-z slices. (**b**) Heat maps showing cell junctional tensions at the three time points. The relative junctional tensions were averaged weighted with the edge lengths for each cell. (**c**) Heat maps showing relative cell pressure at the three time points. (**d**) Change in apical surface area among the invaginated constricting cells (lines) and surrounding cells (dotted lines). Colors indicate different embryos, and the values were averaged in each embryo. Error bars indicate standard deviations. (**e**) Change in relative cell pressure among the invaginated constricting cells (lines) and surrounding cells (dotted lines). Colors indicate different embryos, and the values were averaged in each embryo. Error bars indicate standard deviations. A scale bar in **a** represents 10 nm.

in the apical constriction (*Lee and Harland, 2010*; *Mateus et al., 2011*; *Ossipova et al., 2014*, *2015*; *Le and Chung, 2021*), and especially in its ratcheting (*Miao et al., 2019*; *Kowalczyk et al., 2021*; *Chen and He, 2022*). When the apical endocytosis was impaired but the myosin activity was retained, cells exhibited the pulsed apical contraction but fully turned back to its original shape in the following relaxation. Interestingly, *Miao et al.* (*2019*) downregulated endocytosis in a *Drosophila* gastrula and observed the apical contraction in the edge cells of a mesoderm which would be invaginated by the apical constriction in the wild type embryo. This edge specific contraction resembles the result of simulation with the heighten apical surface contractility (*Figure 2*), and the retained myosin activity might be consistent with the heighten contractility. Together, the myosin surely generates the pulsed contractility and also might be constantly generating contractility at some magnitude. Then, how the myosin and its pulsed and constant contractility contribute to the ratcheting and the apical constriction? Although most of the preceding studies supposed that the endocytosis regulated the ratcheting indirectly by facilitating apical actomyosin network remodeling, here we would propose that the cell membrane itself increases the apical surface tension by changing the reference value. Since the cell membrane is an elastic sheet, it bears the surface tension when extended, and its magnitude is proportional to the change from the reference value. The reference value is an area of the membrane when the membrane is not stretched nor compressed, and it can be decreased by the endocytosis. It might be also possible that the remodeling of the apical actomyosin ratchets the relaxation, as actin filament linker βH-spectrin was also reported regulating the ratcheting (*Krueger et al., 2020*). However, a resultant apical surface tension must be increasing with respect to the apical area, and an accountable mechanism is not known for the actomyosin network for such ratcheting and physical property. In summary, we propose that 1) the actomyosin generates pulsed contractility, and the cell apical surface is contracted temporally; 2) the contraction will be fully relaxed without the endocytosis; 3) the endocytosis decrease the reference value of the apical surface elasticity during the pulsed contraction; 4) after the pulsed contraction and endocytosis, the cell apical surface is partially relaxed (*Figure 7*a).

From the model of pulsed contractility and modified elasticity, we got expectations as described below. When the apical myosin is constantly active as in the conventional model, the heighten contractility is balanced among the cells and makes the edge cells delaminated (*Figure 7*b). Therefore the myosin activatoin must be pulsed and sporadic. When the myosin activity is pulsed and sporadic but not accompanied by the surface elasticity modification, the temporal contraction is fully relaxed and the tissue remains in the same shape (*Figure 7*c). When the sporadic pulse of myosin activation is followed by the endocytosis, the cells undergoes the simultaneous apical contraction, and then the pressure is balanced among the cells, keeping the cell lateral cell-cell junction straight and deforming the cells into the wedge shape (*Figure 7*d). In the actual tissue, the apical localization of the myosin and regulators of the endocytosis were both patterned. However, it may be possible in theory that the myosin activation is not patterned but is pulsed across the entier tissue, while the endocytosis is patterned, and the tissue will be bent at the site of the heighten endocytotic activity (*Figure 7*e). Another interesting notion was that too high myosin activation would prevent the apical constriction. Even if the apical contraction is synchronized and the lateral surface is kept straight so that the cells are deformed into the wedge shape, the too high apical surface tension will try to make the apical surface as short as possible, and the apical surface will be kept flat (*Figure 7*f). On the other hand, mechanisms which does not generate the apical surface tension, such like the zippering (*Bahri et al., 2010*; *Hashimoto et al., 2015*; *Galea et al., 2017*) or the supracellular myosin cable encircling the contracting cells (*Nishimura et al., 2007*; *Röper, 2012*; *Chung et al., 2017*), may promote the deep invagination. Note that we also demonstrated that the supracellular contractile ring alone could not drive the tissue invagination (*Figure 2*c). How actual individual tissue accomplish the synchronized apical contraction would be clarified by examining the physical property and dynamics of the cell membrane, the apical actomyosin network, the adherens junction, and the tight junction.

In conclusion, we elucidated how distinctively the surface contractility and the surface elasticity

Figure 7. Hypothetical model of endocytosis-based apical constriction. (**a**) Flow diagram of the ratcheting by endocytosis. The cell apical surface was contracted by the pulsed myosin activation (1). Without the endocytosis, the apical surface would be fully relaxed (2). By the endocytosis, the apical surface reference value was decreased (3). Because of the modified reference value, the cell apical surface was partially relaxed (4). (**b**) Expected deformation by the heighten apical surface contractility. (**c**) Expected deformation by the sporadic apical surface contractility. (**d**) Expected deformation by the patterned sporadic contractility and endocytosis. (**e**) Expected deformation by the general sporadic contractility and the patterned endocytosis. (**f**) Expected deformation with the heightened apical surface tension.

355 would regulate tissue morphogenesis. Our simulation results were consistent with the studies
356 of the endocytosis in apical constriction, and suggested why and how the endocytosis and the
357 ratcheting was required for the apical constriction.

METHODS
Cellular Potts model with subcellular components
To simulate epithelial cell and tissue deformation, we developed a cellular Potts model with sub-
cellular components, subcellular locations, and a potential energy. The cellular Potts model is a
stochastic simulation framework representing a cell with a set of sites in a lattice, and the defor-
mation is simulated by updating a label on the site in Metropolis-style. In the update, a change
in energy $\Delta\mathcal{H}$ is calculated between the original and a proposed new configuration, and accepted
if $\Delta\mathcal{H} \leq 0$ or with probability $e^{-\Delta\mathcal{H}/T}$ if $\Delta\mathcal{H} > 0$, where T represents a system temperature. With
the subcellular components, the energy $\mathcal{H}$ is defined by a contact energy between the subcellular
components inside and across cells, and an area constraint of the cell. The contact energy is defined
by

$$\mathcal{H}_{\text{sc}} = \sum_{i} \sum_{j \in N(i)} J(\delta(\sigma(i), \sigma(j)), \tau(i), \tau(j)), \tag{6}$$

where each of i and j denotes a site, $N(i)$ is a neighborhood of i, $\sigma(i)$ represents a cell to which i belongs, $\tau(i)$ represents a subcellular component at i, δ is the Kronecker delta, and J define the affinity between the subcellular components either inside or across cells. The area constraint is defined by

$$\mathcal{H}_{\text{ac}} = \sum_{\sigma} \lambda (A(\sigma) - A_0(\sigma))^2, \tag{7}$$

where σ denotes a cell, λ denotes a bulk modulus, $A(\sigma)$ represents an area of σ, and $A_0(\sigma)$ denotes a reference value for σ.

A ratio $A_0/\bar{A}$ is a parameter outside the simulation, and $A_0(\sigma)$ for a cell σ with a width w and a height h at an initial configuration is set $whA_0/\bar{A}$. The contact energy between cytosols across cells, which defines the cell lateral surface contractility J_l, is also dependent on $A_0/\bar{A}$ and set $\lambda w^2 h(A_0/\bar{A} - 1)$. The apical surface contractility J_a and the basal surface contractility J_b are defined by the contact energy for the cytosol to '*outer body fluid*' and '*inner body fluid*' across cells respectively, and set proportional to J_l.

In the update, first a site i is randomly selected, and then another site j is randomly selected from the neighborhood $N(i)$. New labels representing cell and subcellular component for the site i is determined based on the labels on i and j. For example, if the site i was labeled with a cell σ_1 and j was labeled as '*medium*' for cell identity and with '*apical ECM*' for the subcellular component, the new label for i was '*medium*' and '*outer body fluid*'. If the site i was labeled with '*medium*' and '*apical ECM*', then the update was rejected so that a cell would not intrude into the apical ECM. The update is iterated for an arbitrary duration, and results were visually checked if the tissue integrity was retained, the cells were delaminated, or the tissue was bent or invaginated.

Cellular Potts model with partial surface elasticity

To simulate the differential phyisical properties of the apical, lateral, and basal surfaces, the subcellular locations are marked automatically. Sites adjacent to different cells but not to the medium are marked as lateral. Sites which are adjacent to medium and were adjacent to the apical ECM at the initial configuration are marked as apical. Sites which are adjacent to medium and not apical are marked as basal. In addition, sites are marked as adherens junction if they are not apical, adjacent to an apical site of any cell, and adjacent to a lateral site of different cell.

For the surface elasticity, the total, apical, lateral, and basal perimeter lengths of a cell σ $P(\sigma)$, $P_a(\sigma)$, $P_l(\sigma)$, and $P_b(\sigma)$ are calculated in the same way with the contact energy, i.e., by summing numbers of adjacent sites in different cells for all sites, all sites marked apical, lateral, and basal in σ respectively. The surface elasticity is calculated by

$$\mathcal{H}_{\text{set}} = \sum_{\sigma} E_{s\sigma}(P(\sigma) - P_0(\sigma))^2 \tag{8}$$

for the total perimeter, or by

$$\mathcal{H}_{\text{sep}} = \sum_{\sigma} E_{s\sigma}((P_a(\sigma) - P_{a0}(\sigma))^2 + (P_l(\sigma) - P_{l0}(\sigma))^2 + (P_b(\sigma) - P_{b0}(\sigma))^2) \tag{9}$$

for the partitioned perimeters, where $E_{s\sigma}$ denotes a surface elastic modulus of the cell σ and $P_0(\sigma)$, $P_{a0}(\sigma)$, $P_{l0}(\sigma)$, and $P_{b0}(\sigma)$ represents reference values.

Cellular Potts model with potential energy

To simulate an effect of supracellular myosin cable encircling the invaginated cells, we assigned a potential energy to the adherens junction. Tension on a circular actomyosin cable would shrink the circle, and the cable would be pulled toward the center. Quantitatively, the cable is pulled by γ/R, where γ is the tension and R is the radius of the cable curvature as determined from the Young-Laplace equation. The pulling force was simulated by a gradient of the potential energy decreasing toward the middle line, constant in each half side and taking V shape. The potential

energy was defined by a scalar field, and the scalar values corresponding to the adherens junction sites were averaged in each cell by

$$\mathcal{H}_{\mathrm{mc}} = \sum_{\sigma} \frac{1}{|X_{\sigma}|} \sum_{i \in X_{\sigma}} U(i), \tag{10}$$

where X_{σ} denotes a set of the adherens junction sites in the cell σ, and U denotes the scalar field. When the adherens junctions sites moved toward the middle line, $\mathcal{H}_{\mathrm{mc}}$ would decrease, and the movement would be accepted if the increase of other terms in $\mathcal{H}$ was smaller than the decrease. The pulling force γ/R is equivalent to a rate of the gradient change.

Implementation of the simulations

The simulations were implemented by Matlab custom scripts which are available at GitHub (https://doi.org/10.5281/zenodo.8260142).

Energy landscape with respect to cell shape

To analyze the energy with respect to cell shape, we prepared a simple 2D model of single cell represented by four vertices and four edges. Apical and basal edges were straight, while a lateral edge was straight but another lateral edge was assigned a constant curvature. Let v_1, v_2, v_3, and v_4 denote the vertices, where an edge between v_1 and v_4 is the apical, an edge between v_2 and v_3 is the basal, an edge between v_1 and v_2 is the curved lateral, and an edge between v_3 and v_4 is the straight lateral edge. For a radius r of the curvature, an angle θ can be determined as $2r\sin\theta = ||v_2 - v_1||$ and so a length of the curved edge is $2r\theta$. Also, an area between a chord and an arc for v_1 and v_2 is $r^2\theta - 1/2||v_2 - v_1||r\cos\theta$, and an area A in the cell is

$$A = r^2\theta - \frac{1}{2}||v_2 - v_1||r\cos\theta + \frac{1}{2}(v_1 \times v_2 + v_2 \times v_3 + v_3 \times v_4 + v_4 \times v_1). \tag{11}$$

In the same way with the cellular Potts model and the vertex model, the energy $\mathcal{H}$ of the cell was defined with the surface contractility and the area constraint,

$$\mathcal{H} = J_a||v_1 - v_4|| + J_l(2r\theta + ||v_4 - v_3||) + J_b||v_3 - v_2|| + \lambda(A - A_0)^2, \tag{12}$$

where J_a, J_l, and J_b denote the apical, lateral, and basal surface contractilities, λ denotes the bulk modulus, and A_0 is the reference value. For an apical width, v_1 and v_4 positions were set, and for the lateral curvature, v_2 and v_3 positions were searched for a minimum energy. For the cell to be stable in a rectangular shape with the apical width $w = 13$, lateral height $h = 18$, 0 curvature, $\lambda = 1$, and the area reference value $A_0 = 1.02wh$, the surface contractilities were set as $J_b = \lambda(A_0 - wh)h$ and $J_l = \lambda(A_0 - wh)w$. The apical surface contractility J_a was set $J_a = 1.6J_b$. The apical width and the curvature were varied from w to $w/4$ and 0 to $0.5/h$ in 100 bins respectively, and the energy landscape was plotted.

A gradient map was obtained from the energy landscape by Prewitt filter, and a path from the apical width w and 0 curvature following the gradient was calculated numerically.

The cell shape searching and the path tracking were implemented by Matlab custom scripts which are available at GitHub (https://doi.org/10.5281/zenodo.8354144).

Image acquisition

Drosophila eggs were collected at 25 °C, and dechorionated embryos were mounted on a glass-bottom dish with glue and submerged under water. The lateral epidermis were imaged with a confocal microscope (Olympus FV1000) equipped with a ×60 oil immersion objective lense (PLAPON 60X0) at 25 °. Z stack of 0.7 μm interval with slices of 0.125 μm x-y resolution was taken every 3 minutes.

Cell tracking

The cells were manually tracked with Fiji plugin TrackMate manual tracking.

Conversion of adherens junction into graph

For the geometrical tension inference, the cells were segmented manually and converted to the vertices and edges representation. The 3D image of cadherin was projected onto the x-y plane, and the adherens junction was traced manually. The trace was narrowed to 1 pixel boundary by Fiji plugin MorphoLibJ Segmentation Interactive Marker-controlled Watershed. In the segmented image, the boundary pixels were scanned and the vertices were allocated for pixels incident to more than two cells, i.e., the boundary pixel was next to pixels in the cells. Next, all pairs of the cells were enumerated and checked if there were two vertices incident to both of the cells, and edges were allocated to link those vertices. Pixels incident to both of the cells were assigned to the edge, and a depth of the brightest pixel in the stack for x and y of each pixel was read. When the depth differed by more than 3 from an average depth of adjacent edge pixels, the depth was masked as a noise. For each vertex, a plane was fitted by the least square method to x, y, and depths of the incident edges pixels, and depth of the plane at the vertex x and y position was assigned to the vertex. Vertices closer than 2 to each other were merged into a vertex.

3D Bayesian tension inference

The balance equation was obtained in a tangent plane for each vertex. For a vertex v_0, its adjacent vertices v_1, v_2, ... were collected and a plane was fitted to them by the least square method. The vertices v_0, v_1, ... and edges between them e_1, e_2, ... were projected on the plane, and the projected edge e'_1, e'_2, ... were measured their lengths and angles. An axis was set parallel with e'_1, and another axis was also set orthogonal to the first axis. The balance equation was obtained with respect to the axes. A total force f_1 to move v_0 along the first axis was represented by an equation

$$f_1 = \sum_{i=1}^{k} -\cos(\theta_i)T_i + \frac{1}{2}\sum_{j=1}^{k}(|e'_{j+1}|\sin\theta_{j+1} - |e'_j|\sin\theta_j)P_j, \tag{13}$$

where k is a number of the adjacent vertices, θ_i denotes an angle between the edge e'_1 (and thus the first axis) and the edge e'_i, $|e'_i|$ represents the length of e'_i, $e'_{k+1} = e'_1$, T_i denotes a tension exerted on e'_i, and P_j denotes a pressure in a cell containing edges e'_j and e'_{j+1}. Note that the symbols T_i and P_j follows the *Ishihara and Sugimura* (*2012*) original work, and do not confused with the symbols which denoted the perimeter length and temperature. In the same way, a total force f_2 to move v_0 along the second axis was represented by an equation

$$f_2 = \sum_{i=1}^{k} -\sin(\theta_i)T_i - \frac{1}{2}\sum_{j=1}^{k}(|e'_{j+1}|\cos\theta_{j+1} - |e'_j|\cos\theta_j)P_j. \tag{14}$$

The coefficients for T_is and P_js were combined in an $n \times m$ matrix **A** so that **Ap** represents the forces, where n is a number of the equations, m is a number of the tensions and pressures, and **p** is an m-dimensional vector composed of T_i and P_j.

The equations above are based on an assumption that an adherens junction is responsible for the edge tension and it makes an arc so that its angle of incidence to the vertex is parallel with the plane tangent at the vertex. We also approximated the tangent plane with the plane fitted to the adjacent vertices.

The Bayesian tension inference is to find the most likely distribution of tensions and pressures which makes **Ap** reasonably small. It was numerically estimated as below. Let a parameter $\mu = \sigma^2/\omega^2$, where σ represents errors of the equations and ω is a variance of the tensions and pressures distribution. The parameter μ was searched for a minimization of Akaike's Bayesian information criterion (ABIC). To calculate ABIC, we prepared $\tau = \mu^{1/2}$ and

$$\mathbf{S} = \begin{pmatrix} \mathbf{A} & \mathbf{0} \\ \tau\mathbf{B} & \tau\mathbf{g} \end{pmatrix}, \tag{15}$$

where

$$\mathbf{g} = \begin{pmatrix} 1 \\ \vdots \\ 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} \begin{array}{l} \left.\vphantom{\begin{matrix}1\\\vdots\\1\end{matrix}}\right\}\text{number of edges with tension, } n_e \\[2em] \left.\vphantom{\begin{matrix}0\\\vdots\\0\end{matrix}}\right\}\text{number of cells with pressure, } n_c \end{array} \quad , \tag{16}$$

and $\mathbf{B}$ was a diagonal matrix of $\mathbf{g}$. Then $\mathbf{S}$ was decomposed by a QR factorization, $\mathbf{S} = \mathbf{QR}$. Let $\mathbf{R}$ be expressed as

$$\mathbf{R} = \begin{pmatrix} \mathbf{H} & \mathbf{h} \\ \mathbf{0} & h \end{pmatrix}, \tag{17}$$

and $\mathbf{d}_{H1}$, $\mathbf{d}_{H2}$, ... be an array of absolute value of non-zero diagonal elements of $\mathbf{H}$. The ABIC was calculated by

$$\text{ABIC} = n_{km} + n_{km}\log\frac{2\pi h^2}{n_{km}} + 2\sum_i \log\mathbf{d}_{Hi} - n_e\log\mu + 2n_p, \tag{18}$$

where $n_{km} = n - n_c + 1$ and n_p is a number of parameters, here $n_p = 2$. After the search of μ which minimizes ABIC, $\mathbf{p}$ was evaluated by calculating

$$\mathbf{p} = \mathbf{H}^{-1}\mathbf{h}, \tag{19}$$

where $\mathbf{H}^{-1}$ is the Moore-Penrose pseudo inverse matrix of $\mathbf{H}$.

For the basis of the calculation, see *Ishihara and Sugimura* (*2012*).

Cell apical area calculation

For a cell, a plane was fitted to incident vertices of the cell. The vertices and edges between them were projected on the plane, and an area of the polygon was calculated as the cell apical surface area.

Implementation of the tissue apical surface analysis

The conversion from the segmented image to the vertices and edges representation, the geometrical tension inference, and the apical area calculation was implemented by Matlab custom scripts which are available at GitHub (https://doi.org/10.5281/zenodo.8320862).

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AUTHOR CONTRIBUTIONS

Cell and tissue modeling, S.Y. and F.G.; simulation design, S.Y.; geometrical tension inference, S.Y., S.I., and F.G.; programming, S.Y.; writing - original draft, S.Y.; writing - review & editing, S.Y. and F.G.; funding acquisition, S.Y.; supervision, F.G.

Competing interests

The authors declare no competing financial interests.

REFERENCES

References

Adhyapok P, Piatkowska AM, Norman MJ, Clendenon SG, Stern CD, Glazier JA, Belmonte JM. A mechanical model of early somite segmentation. iScience. 2021 Apr; 24(4):102317.

Bahri S, Wang S, Conder R, Choy J, Vlachos S, Dong K, Merino C, Sigrist S, Molnar C, Yang X, Manser E, Harden N. The leading edge during dorsal closure as a model for epithelial plasticity: Pak is required for recruitment of the Scribble complex and septate junction formation. Development. 2010 Jun; 137(12):2023–2032.

Belmonte JM, Clendenon SG, Oliveira GM, Swat MH, Greene EV, Jeyaraman S, Glazier JA, Bacallao RL. Virtual-tissue computer simulations define the roles of cell adhesion and proliferation in the onset of kidney cystic disease. Mol Biol Cell. 2016 11; 27(22):3673–3685.

Benton MA, Frey N, Nunes da Fonseca R, von Levetzow C, Stappert D, Hakeemi MS, Conrads KH, Pechmann M, Panfilio KA, Lynch JA, Roth S. Fog signaling has diverse roles in epithelial morphogenesis in insects. Elife. 2019 Oct; 8:e47346.

Booth AJR, Blanchard GB, Adams RJ, Röper K. A dynamic microtubule cytoskeleton directs medial actomyosin function during tube formation. Dev Cell. 2014 Jun; 29(5):562–576.

Brezavšček AH, Rauzi M, Leptin M, Ziherl P. A model of epithelial invagination driven by collective mechanics of identical cells. Biophys J. 2012 Sep; 103(5):1069–1077.

Brodland GW, Conte V, Cranston PG, Veldhuis J, Narasimhan S, Hutson MS, Jacinto A, Ulrich F, Baum B, Miodownik M. Video force microscopy reveals the mechanics of ventral furrow invagination in Drosophila. Proc Natl Acad Sci U S A. 2010 Dec; 107(51):22111–22116.

Brodu V, Casanova J. The RhoGAP crossveinless-c links tracheanless and EGFR signaling to cell shape remodeling in Drosophila tracheal invagination. Genes Dev. 2006 Jul; 20(13):1817–1828.

Chen W, He B. Actomyosin activity-dependent apical targeting of Rab11 vesicles reinforces apical constriction. J Cell Biol. 2022 Jun; 221(6).

Chung MI, Nascone-Yoder NM, Grover SA, Drysdale TA, Wallingford JB. Direct activation of Shroom3 transcription by Pitx proteins drives epithelial morphogenesis in the developing gut. Development. 2010 Apr; 137(8):1339–1349.

Chung S, Kim S, Andrew DJ. Uncoupling apical constriction from tissue invagination. Elife. 2017 Mar; 6.

Conte V, Ulrich F, Baum B, Muñoz J, Veldhuis J, Brodland W, Miodownik M. A biomechanical analysis of ventral furrow formation in the Drosophila melanogaster embryo. PLoS One. 2012; 7(4):e34473.

Dawes-Hoang RE, Parmar KM, Christiansen AE, Phelps CB, Brand AH, Wieschaus EF. folded gastrulation, cell shape change and the control of myosin localization. Development. 2005 Sep; 132(18):4165–4178.

Fortuna I, Perrone GC, Krug MS, Susin E, Belmonte JM, Thomas GL, Glazier JA, de Almeida RMC. CompuCell3D Simulations Reproduce Mesenchymal Cell Migration on Flat Substrates. Biophys J. 2020 06; 118(11):2801–2815.

Galea GL, Cho YJ, Galea G, Molè MA, Rolo A, Savery D, Moulding D, Culshaw LH, Nikolopoulou E, Greene NDE, Copp AJ. Biomechanical coupling facilitates spinal neural tube closure in mouse embryos. Proc Natl Acad Sci U S A. 2017 Jun; 114(26):E5177–E5186.

Galea GL, Maniou E, Edwards TJ, Marshall AR, Ampartzidis I, Greene NDE, Copp AJ. Cell non-autonomy amplifies disruption of neurulation by mosaic Vangl2 deletion in mice. Nat Commun. 2021 Feb; 12(1):1159.

Galea GL, Nychyk O, Mole MA, Moulding D, Savery D, Nikolopoulou E, Henderson DJ, Greene NDE, Copp AJ. Vangl2 disruption alters the biomechanics of late spinal neurulation leading to spina bifida in mouse embryos. Dis Model Mech. 2018 Mar; 11(3):dmm032219.

Graner F, Glazier JA. Simulation of biological cell sorting using a two-dimensional extended Potts model. Phys Rev Lett. 1992 Sep; 69(13):2013–2016.

Guglielmi G, Barry JD, Huber W, De Renzis S. An Optogenetic Method to Modulate Cell Contractility during Tissue Morphogenesis. Dev Cell. 2015 Dec; 35(5):646–660.

Guru A, Saravanan S, Sharma D, Narasimha M. The microtubule end-binding proteins EB1 and Patronin
modulate the spatiotemporal dynamics of myosin and pattern pulsed apical constriction. Development. 2022
Nov; 149(22).

Haigo SL, Hildebrand JD, Harland RM, Wallingford JB. Shroom induces apical constriction and is required for
hingepoint formation during neural tube closure. Curr Biol. 2003 Dec; 13(24):2125–2137.

Hashimoto H, Robin FB, Sherrard KM, Munro EM. Sequential contraction and exchange of apical junctions
drives zippering and neural tube closure in a simple chordate. Dev Cell. 2015 Jan; 32(2):241–255.

Hayes P, Solon J. Drosophila dorsal closure: An orchestra of forces to zip shut the embryo. Mech Dev. 2017 Apr;
144(Pt A):2–10.

Helfrich W. Elastic properties of lipid bilayers: theory and possible experiments. Z Naturforsch C. 1973;
28(11):693–703.

Hildebrand JD. Shroom regulates epithelial cell shape via the apical positioning of an actomyosin network. J
Cell Sci. 2005 Nov; 118(Pt 22):5191–5203.

Hutson MS, Tokutake Y, Chang MS, Bloor JW, Venakides S, Kiehart DP, Edwards GS. Forces for morphogenesis
investigated with laser microsurgery and quantitative modeling. Science. 2003 Apr; 300(5616):145–149.

Inoue Y, Suzuki M, Watanabe T, Yasue N, Tateo I, Adachi T, Ueno N. Mechanical roles of apical constriction, cell
elongation, and cell migration during neural tube formation in Xenopus. Biomech Model Mechanobiol. 2016
Dec; 15(6):1733–1746.

Ishihara S, Sugimura K. Bayesian inference of force dynamics during morphogenesis. J Theor Biol. 2012;
313:201–211.

Kerridge S, Munjal A, Philippe JM, Jha A, de las Bayonas AG, Saurin AJ, Lecuit T. Modular activation of Rho1
by GPCR signalling imparts polarized myosin II activation during morphogenesis. Nat Cell Biol. 2016 Mar;
18(3):261–270.

Kölsch V, Seher T, Fernandez-Ballester GJ, Serrano L, Leptin M. Control of Drosophila gastrulation by apical
localization of adherens junctions and RhoGEF2. Science. 2007 Jan; 315(5810):384–386.

Kondo T, Hayashi S. Mitotic cell rounding accelerates epithelial invagination. Nature. 2013 Feb; 494(7435):125–
129.

Kowalczyk I, Lee C, Schuster E, Hoeren J, Trivigno V, Riedel L, Görne J, Wallingford JB, Hammes A, Feistel K.
Neural tube closure requires the endocytic receptor Lrp2 and its functional interaction with intracellular
scaffolds. Development. 2021 Jan; 148(2):dev195008.

Krueger D, Pallares Cartes C, Makaske T, De Renzis S. H-spectrin is required for ratcheting apical pulsatile
constrictions during tissue invagination. EMBO Rep. 2020 Aug; 21(8):e49858.

Labouesse C, Verkhovsky AB, Meister JJ, Gabella C, Vianay B. Cell shape dynamics reveal balance of elasticity
and contractility in peripheral arcs. Biophys J. 2015 May; 108(10):2437–2447.

Lang RA, Herman K, Reynolds AB, Hildebrand JD, Plageman TF. p120-catenin-dependent junctional recruitment
of Shroom3 is required for apical constriction during lens pit morphogenesis. Development. 2014 Aug;
141(16):3177–3187.

Le TP, Chung S. Regulation of apical constriction via microtubule- and Rab11-dependent apical transport during
tissue invagination. Mol Biol Cell. 2021 May; 32(10):1033–1047.

Lee C, Scherr HM, Wallingford JB. Shroom family proteins regulate gamma-tubulin distribution and microtubule
architecture during epithelial cell shape change. Development. 2007 Apr; 134(7):1431–1441.

Lee JY, Harland RM. Endocytosis is required for efficient apical constriction during Xenopus gastrulation. Curr
Biol. 2010 Feb; 20(3):253–258.

Manning AJ, Peters KA, Peifer M, Rogers SL. Regulation of epithelial morphogenesis by the G protein-coupled
receptor mist and its ligand fog. Sci Signal. 2013 Nov; 6(301):ra98.

Martin AC. The Physical Mechanisms of Drosophila Gastrulation: Mesoderm and Endoderm Invagination.
Genetics. 2020 Mar; 214(3):543–560.

Martin AC, Gelbart M, Fernandez-Gonzalez R, Kaschube M, Wieschaus EF. Integration of contractile forces during tissue invagination. J Cell Biol. 2010 Mar; 188(5):735–749.

Martin AC, Goldstein B. Apical constriction: themes and variations on a cellular mechanism driving morphogenesis. Development. 2014 May; 141(10):1987–1998.

Martin AC, Kaschube M, Wieschaus EF. Pulsed contractions of an actin-myosin network drive apical constriction. Nature. 2009 Jan; 457(7228):495–499.

Mason FM, Tworoger M, Martin AC. Apical domain polarization localizes actin-myosin activity to drive ratchet-like apical constriction. Nat Cell Biol. 2013 Aug; 15(8):926–936.

Mason FM, Xie S, Vasquez CG, Tworoger M, Martin AC. RhoA GTPase inhibition organizes contraction during epithelial morphogenesis. J Cell Biol. 2016 Aug; 214(5):603–617.

Mateus AM, Gorfinkiel N, Schamberg S, Martinez Arias A. Endocytic and recycling endosomes modulate cell shape changes and tissue behaviour during morphogenesis in Drosophila. PLoS One. 2011 Apr; 6(4):e18729.

McGreevy EM, Vijayraghavan D, Davidson LA, Hildebrand JD. Shroom3 functions downstream of planar cell polarity to regulate myosin II distribution and cellular organization during neural tube closure. Biol Open. 2015 Jan; 4(2):186–196.

Miao H, Blankenship JT. The pulse of morphogenesis: actomyosin dynamics and regulation in epithelia. Development. 2020 Sep; 147(17):dev186502.

Miao H, Vanderleest TE, Jewett CE, Loerke D, Blankenship JT. Cell ratcheting through the Sbf RabGEF directs force balancing and stepped apical constriction. J Cell Biol. 2019 Nov; 218(11):3845–3860.

Münster S, Jain A, Mietke A, Pavlopoulos A, Grill SW, Tomancak P. Attachment of the blastoderm to the vitelline envelope affects gastrulation of insects. Nature. 2019 Apr; 568(7752):395–399.

Nematbakhsh A, Levis M, Kumar N, Chen W, Zartman JJ, Alber M. Epithelial organ shape is generated by patterned actomyosin contractility and maintained by the extracellular matrix. PLoS Comput Biol. 2020; 16(8):e1008105.

Nishimura M, Inoue Y, Hayashi S. A wave of EGFR signaling determines cell alignment and intercalation in the Drosophila tracheal placode. Development. 2007 Dec; 134(23):4273–4282.

Nishimura T, Takeichi M. Shroom3-mediated recruitment of Rho kinases to the apical cell junctions regulates epithelial and neuroepithelial planar remodeling. Development. 2008 Apr; 135(8):1493–1502.

van Oers RF, Rens EG, LaValley DJ, Reinhart-King CA, Merks RM. Mechanical cell-matrix feedback explains pairwise and collective endothelial cell behavior in vitro. PLoS Comput Biol. 2014 Aug; 10(8):e1003774.

Ogura Y, Wen FL, Sami MM, Shibata T, Hayashi S. A Switch-like Activation Relay of EGFR-ERK Signaling Regulates a Wave of Cellular Contractility for Epithelial Invagination. Dev Cell. 2018 Jul; 46(2):162–172.

Ossipova O, Chuykin I, Chu CW, Sokol SY. Vangl2 cooperates with Rab11 and Myosin V to regulate apical constriction during vertebrate gastrulation. Development. 2015 Jan; 142(1):99–107.

Ossipova O, Kim K, Lake BB, Itoh K, Ioannou A, Sokol SY. Role of Rab11 in planar cell polarity and apical constriction during vertebrate neural tube closure. Nat Commun. 2014 May; 5:3734.

Perez-Mockus G, Mazouni K, Roca V, Corradi G, Conte V, Schweisguth F. Spatial regulation of contractility by Neuralized and Bearded during furrow invagination in Drosophila. Nat Commun. 2017 Nov; 8(1):1594.

Pérez-Vonzález C, Ceada G, Greco F, Matejčić M, Gómez-González M, Castro N, Menendez A, Kale S, Krndija D, Clark AG, Gannavarapu VR, Álvarez Varela A, Roca-Cusachs P, Batlle E, Vignjevic DM, Arroyo M, Trepat X. Mechanical compartmentalization of the intestinal organoid enables crypt folding and collective cell migration. Nat Cell Biol. 2021 Jul; 23(7):745–757.

Plageman TF, Chung MI, Lou M, Smith AN, Hildebrand JD, Wallingford JB, Lang RA. Pax6-dependent Shroom3 expression regulates apical constriction during lens placode invagination. Development. 2010 Feb; 137(3):405–415.

Polyakov O, He B, Swan M, Shaevitz JW, Kaschube M, Wieschaus E. Passive mechanical forces control cell-shape change during Drosophila ventral furrow formation. Biophys J. 2014 Aug; 107(4):998–1010.

Rauzi M, Brezavšček AH, Ziherl P, Leptin M. Physical models of mesoderm invagination in Drosophila embryo. Biophys J. 2013 Jul; 105(1):3–10.

Rens EG, Merks RMH. Cell Contractility Facilitates Alignment of Cells and Tissues to Static Uniaxial Stretch. Biophys J. 2017 Feb; 112(4):755–766.

Roffay C, Chan CJ, Guirao B, Hiiragi T, Graner F. Inferring cell junction tension and pressure from cell geometry. Development. 2021 Mar; 148(18).

Röper K. Anisotropy of Crumbs and aPKC drives myosin cable assembly during tube formation. Dev Cell. 2012 Nov; 23(5):939–953.

Sherrard K, Robin F, Lemaire P, Munro E. Sequential activation of apical and basolateral contractility drives ascidian endoderm invagination. Curr Biol. 2010 Sep; 20(17):1499–1510.

Solon J, Kaya-Çopur A, Colombelli J, Brunner D. Pulsed forces timed by a ratchet-like mechanism drive directed tissue movement during dorsal closure. Cell. 2009 Jun; 137(7):1331–1342.

Takeichi M. Dynamic contacts: rearranging adherens junctions to drive epithelial remodelling. Nat Rev Mol Cell Biol. 2014 Jun; 15(6):397–410.

Urbansky S, González Avalos P, Wosch M, Lemke S. Folded gastrulation and T48 drive the evolution of coordinated mesoderm internalization in flies. Elife. 2016 Sep; 5:e18318.

Xie S, Mason FM, Martin AC. Loss of $G_{\alpha 12/13}$ exacerbates apical area dependence of actomyosin contractility. Mol Biol Cell. 2016 Nov; 27(22):3526–3536.

Figure 1—figure supplement 1. Epithelial cell surface tension and tissue integrity. The energy $\mathcal{H}$ included the terms of surface contact energy and area constraint. The constant $\bar{A}$ represents a target average area of cells in the simulation. The ratio $A_0/\bar{A}$ indicates how much the cells are compressed, and thus how strong the surface tension is. Here the lateral surface tension was defined by the contact energy between the cytosols in different cells, the basal surface tension was defined by the contact energy between the cytosol and the inner body fluid, and the apical surface tension was defined by the contact energy between the cytosol and the apical ECM or the outer body fluid. The lateral surface tension was determined based on the compression $A_0/\bar{A}$. The vertical axis J_b/J_l indicates a ratio between the basal and apical surface contact energy J_b and the lateral surface contact energy J_l.

Figure 5—figure supplement 1. Simulation of apical constriction with various cell heights. The cells were 2/3 lower (first and second columns) or 4/3 higher (third and fourth columns) in the apical-basal axis. The ratio $A_0/\bar{A}$ indicates how much cells were compressed, J_b/J_l indicates a ratio between the basal and lateral surface contractility, and E_s denotes the surface elastic modulus for the inner pale blue and green 13 cells. The edge pale red cells were assigned 0.1 times smaller surface elastic modulus than the inner cells. The center pale blue five cells were assigned $P_{a0} = 0$, while the others were assigned P_{a0} equivalent to P_a initial value.

Figure 5—figure supplement 2. Simulation of apical constriction by the elasticity remodeling with cell-ECM adhesion. The five center cells were assigned $P_{a0} = 0$, and the other surrounding cells were assigned an affinity to the apical ECM indicated by a ratio between cytosol-ECM and cytosol-outer body fluid contact energies J_s/J_a. The ratio $A_0/\bar{A}$ indicates how much cells were compressed, J_b/J_l indicates a ratio between the basal and lateral surface contractility, and E_s denotes the surface elastic modulus for the inner pale blue and green 13 cells. The edge pale red cells were assigned 0.1 times smaller surface elastic modulus than the inner cells.