The role of active mRNA-ribosome dynamics and closing constriction in daughter chromosome separation in *Escherichia coli*

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Abstract

The mechanisms by which two sister chromosomes separate and partition into daughter cells in bacteria

remain poorly understood. A recent theoretical model has proposed that out-of-equilibrium central dogma

reactions involving mRNA and ribosomes play a significant role in this process. Here we test this idea in

the *Escherichia coli* model system using high-throughput fluorescence microscopy in microfluidic devices.

We compare our experimental observations with predictions from a reaction-diffusion model that includes

central dogma-related reactions and excluded volume interactions between ribosomal subunits, polysomes,

and chromosomal DNA. Our results show that the non-equilibrium reactions of ribosomes cause them to

aggregate at the midcell, and this process facilitates the separation of the two daughter chromosomes.

However, the observed effects are weaker in live cells than our one-dimensional reaction-diffusion model

predicts. Rather than relying solely on active mRNA–ribosome dynamics, our data suggest that the closing

division septum via steric interactions and potentially entropic forces between two DNA strands coupled to

cell elongation act as additional mechanisms to ensure faithful partitioning of the nucleoids to two daughter

cells.

Significance

The mitotic spindle separates chromosomes in eukaryotic cells, but bacteria lack this structure. It remains

unclear how bacterial chromosomes partition before cell division. It has been hypothesized that non-

equilibrium dynamics of polysomes, that is, mRNA-ribosome complexes, actively drive the separation of

bacterial chromosomes. Using quantitative microscopy combined with computational modeling, we show

that polysome dynamics facilitates the separation of daughter chromosomes in *Escherichia coli,* but this

process does not constitute the sole mechanism. Our findings suggest that the closing division septum via

steric interactions and potentially entropic forces between the two DNA strands act as additional

mechanisms.

Introduction

Chromosome replication and segregation are essential processes in the life cycle of a cell (1-3). Segregation can be divided into two stages: the individualization of daughter chromosomes and their further separation, which enables their partitioning into daughter cells during cytokinesis (4). Completion of both stages is required for the successful propagation of cells from one generation to the next.

In bacteria, the individualization stage, commonly referred to as DNA segregation, occurs for the most part concurrently with DNA replication (5,6). In *Escherichia coli*, the newly synthesized DNA loci near *oriC* begin to separate from each other with some delay after initiation of DNA replication (7). After this initial cohesion phase, the sister loci segregate from each other sequentially in the order in which they appear in the genome (8-10). Despite extensive previous research, the underlying molecular mechanism driving the segregation phase in bacteria remains poorly understood (2,3,11). Several active mechanisms for chromosome segregation have been investigated, including ParAB (ParA-ParB-*parS*) (2,11) and the homologous Spo0J-Soj system (12), but these are missing in many bacterial species, including *E. coli* (13,14). In addition to the Par-systems, structural maintenance of chromosomes (SMC) proteins, encoded by MukBEF in *E. coli*, have also been implicated in the active segregation (15,16). SMC involvement in segregation is postulated to occur via a successive condensation of DNA around *oriC* via loop extrusion and repulsion of condensed DNA masses away from each other due to excluded volume interactions and chain entropy (15,16). While this is a plausible scenario, MukBEF is not essential for DNA segregation in *E. coli*, except possibly at the fastest growth rates when multifork replication is present (17).

Rather than relying on dedicated protein machinery for segregation, which would expend energy through ATP hydrolysis, it is plausible that the free energy needed for the segregation of chromosomal DNA is derived mostly from the replication process itself. Replication creates two new proximal DNA strands, which, according to polymer physics, may experience a repulsive entropic force in confined cylindrical conditions (18,19). This scenario, which we refer to as the entropic segregation mechanism, has been supported by numerous computer simulations (20-27) but it has yet to be experimentally validated. However, the driving force from this so-called entropic mechanism vanishes once the spatial overlap between the two nucleoids becomes zero (27). Therefore, if only the entropic mechanism were present, the two individualized and segregated nucleoids would remain in contact rather than fully separating from each other. The entropic mechanism can, thus, explain the segregation but not the separation phase of the two nucleoids.

The question then arises: what drives chromosomes apart during the separation stage? In eukaryotes, the separation is carried out by the mitotic spindle, but no spindles are known to exist in bacteria. DNA translocase FtsK, which is a part of the bacterial cell division apparatus (divisome) (28-30), could be involved in this process (29). However, cells with inactivated FtsK (FtsK K997A) (31) appear to separate their nucleoids and partition them into two daughter cell compartments in the majority of cells (32).

Evidence for a different mechanism for separation was revealed by Wu *et al.* (33). These authors observed
that in filamentous cells (more than 10 µm) with two nonreplicating nucleoids, the nucleoids separated from
each other over large distances so that their centers stayed close to the cell's ¼ positions (33). Wu *et al.*
explained such positioning by osmotic pressure from the accumulation of macromolecular crowders, but
their explanation relied on the assumption that these crowders are produced in a specific way outside the
two nucleoids (33). An alternative explanation, which did not rely on this assumption, was proposed by
Miangolarra *et al.* (34). These authors developed a model that included reactions between chromosomal
DNA, mRNA, ribosomal subunits, and polyribosomes (polysomes), along with their diffusion within the cell.
This reaction-diffusion model, which we refer to as the active mRNA-ribosome dynamics model, included
reactions for mRNA synthesis, ribosome assembly and disassembly on mRNA, and mRNA degradation.
The diffusion of these species was affected by the steric (excluded volume) interactions between these
molecules. The model predicted that due to these interactions, the 30S and 50S ribosomal subunits can
diffuse relatively unhindered into the nucleoid region, consistent with single-molecule tracking experiments
(35,36), while 70S ribosomes, which are further assembled into polyribosomes (polysomes), cannot. The
key insight from the model was a finding that mRNA synthesis within the nucleoid and the assembly of 30S
and 50S ribosomal subunits into polysomes around these mRNAs create instability within the nucleoid. This
instability splits the nucleoid into two distinct nucleoid lobes at some critical cell length (34). The model also
predicted that the two daughter nucleoid centers would position at the quarter points of the cell once the
transition occurred and stay in these locations even when cells filamented, consistent with the experimental
observations by Wu *et al*. (33).

Here, we investigate the mechanisms that individualize chromosomes in the late stages of segregation and
drive their separation in normal growth conditions. In particular, we test the active mRNA-ribosome
dynamics model in non-filamenting cells growing in slow-growth conditions when chromosomal DNA has a
simple topology. To this end, we use high-throughput microscopy in microfluidic devices to determine the
distributions of ribosomes, nucleoids, and replication machinery during cell and replication cycles. We
quantitatively compare the acquired experimental data to our own reaction-diffusion model, which is
conceptually similar to the one proposed by Miangolarra *et al*. (34). We first find that the termination of DNA
replication does not have a significant effect on the nucleoid segregation dynamics. Instead, DNA expands
and daughter nucleoids separate in proportion to cell length increase during this stage of the cell cycle. Our
data, when compared to the model, supports the idea that active ribosome dynamics facilitates the
separation of two nucleoids in normal growth conditions, as recently reported (37). However, in cells where
active ribosome dynamics are halted, we observe that the separation of daughter nucleoids still increases
proportionally with cell length. In addition to this effect, we identify closing constriction, as a new mechanism
that has a significant effect in separating the two daughter nucleoids and partitioning them into two daughter
cells. Based on this evidence, we propose a multi-process model to explain the nucleoid segregation and
separation dynamics in *E. coli*.

Results

The midcell DNA density minimum evolves independently of replication termination

DNA segregation cannot occur without DNA being first replicated. We first hypothesized that the replication
solely drives the progression of segregation. In previous works on the segregation of nucleoids using
tracking of fluorescently labeled loci in the chromosome (5,7,10,38,39), following bulk segregation (26,37),
or labeling the replication terminus by MatP fusion (26,40), the segregation process had not been directly
tracked during the replication cycle at the single-cell level. To do so now, we labeled the beta-clamp protein,
DnaN, with a fluorescent fusion (Ypet-DnaN) (41). To follow the dynamics of the whole chromosome, we
used HupA-mCherry, which has been tested before and is known to bind uniformly throughout the
chromosome (42). We imaged cells in mother machine microfluidic devices (43,44), obtaining about 30-50
time points per cell cycle from several hundred cells (Movie S1-2, SI Fig. S1). We averaged these data over
the cell population and produced kymographs (Fig. 1A-B). We chose to carry out measurements in slow-
growth conditions where, at most, only a single replication cycle is ongoing (45) and where we could
accurately determine timings for the initiation and termination of the replication (SI Fig. S2).

To characterize the segregation process, we determined the total extent of all nucleoids in the cell, $L_{nuc,ext}$,
and the sum of nucleoid lengths, $L_{nuc,sum}$ (Fig. 1C; for details, see Materials and Methods). These
measurements showed that the outer edges of the nucleoids, as determined by $L_{nuc,ext}$, stayed at about a
constant distance (0.5/2 μm ≈ 0.25 μm) from the cell ends throughout the replication cycle (Fig. 1C).
Notably, the slope of $L_{nuc,ext}$ and $L_{nuc,sum}$ vs. time curves did not show any discernible change at the
termination of replication (Fig. 1C). We would have expected the slope of these curves to decrease at the
termination because the increase in the amount of chromosomal DNA stopped. To further verify this result,
we aligned the curves from individual cells relative to replication termination time, Trt , but still found only
a minor changes in the slope of the curves at the replication termination (Fig. S3). We also found that the
outer edges of the nucleoids remain at approximately constant distance from the cell poles ($i.e.\ L_{cell} -$
$L_{nuc,ext} \approx const$) not only throughout the replication cycle but also during the division cycle in both slow (SI
Fig. S4) and moderately fast growth conditions (SI Fig. S5). Note that the measured exclu

.sion distance of nucleoids from poles (0.25 μm) is slightly less than the length of the hemispherical cap
region (0.3-0.35 μm) (46). The data thus show that instead of nucleoid elongation and separation halting at
the replication termination, the length of nucleoids ($L_{nuc,ext}$ and $L_{nuc,ext}$) continues to increase and the
increase is in proportion to cell length. This dependence has also been noticed in some earlier studies
(26,33,47), although the replication period was not directly determined in these works.

In addition to $L_{nuc,sum}$ and $L_{nuc,ext}$, we also determined the DNA density at midcell during the replication
cycle as a different measure of DNA segregation and separation (Fig. 1D). In the early stages of replication,
there was a maximum DNA density at midcell (Fig. 1A). Upon completion of about 60% of the replication

cycle, the DNA amount at the midcell started to decrease (Fig.1D). This decrease corresponded to the

formation of a local DNA density minimum (Fig. 1E). As with our other nucleoid measurements, the midcell

minimum continued to decline further after replication termination without any distinct changes in the slope

of the curve. Altogether, these results, combined with the nucleoid length measurements, indicate that while

chromosome replication is a prerequisite for segregation, nucleoid segregation is not strongly driven by the

replication process itself at late stages of the replication cycle and after its completion. Instead, the

separation of two nucleoids follows overall cell growth.

Correlated and anti-correlated periods of DNA and ribosome density distributions during the cell

cycle

In balanced growth, cell size increases proportionally to ribosome numbers. We hypothesized that the DNA

segregation at late stages of the replication cycle could be driven by active ribosome dynamics that involve

mRNA synthesis and degradation and are associated with the assembly and disassembly of polyribosomes

(34). To test this hypothesis, we constructed a strain (for details, see SI Text, *Strain construction*) where the

nucleoids were labeled with HupA-mYpet fluorescent marker (26,42) and ribosomes with RpsB-mCherry;

the latter being part of the 30S ribosomal subunit (48,49) . Both labels had been tested in previous studies

(26,42,48,49). Since this strain lacked a marker for replication, we followed nucleoid segregation and

ribosome accumulations throughout the division cycle (Fig. 2A-C, SI Fig. 6, Movie S3-4), which we

customarily refer to as the cell cycle.

In the first half of the cell cycle, the population-averaged data show a unimodal distribution of DNA with a

maximum close to the cell center (Fig. 2D, Age 0.10, 0.30). The cell cycle age (age for short) here is

expressed as the fraction of a completed cell cycle. Notably, the ribosome density distribution is similar to

the DNA density distribution from the same cells, showing a broad density maximum at the midcell. Here,

the measured ribosome density is the sum of free 30S ribosome subunits and densities of 30S subunits

incorporated into free ribosome particles and polysomes. The co-localization of the two distribution maxima

appears inconsistent with equilibrium thermodynamic models (50,51), where ribosomes phase-separate

from DNA-rich regions. Note that there are no local maxima in the ribosome density at cell poles in

distributions in Fig. 2D. The excess ribosome density at cell poles is not expected to result in maxima in the

extreme ends of 1D density distributions, which are projections of the 3D density onto the long axis of the

cell. The maxima are missing because of the rounded shape of the poles and diffraction-limited imaging.

Also, there could be an excess ribosome density present in the vicinity of the inner membrane of the cell,

but this is not visible in microscopy images (SI Fig. S6) and in the density distributions along the short axes

of the cell (SI Fig. S7), likely due to projection effects.

As the cell cycle progresses, the nucleoid expands, and the population-average DNA density distribution

becomes bimodal (also referred to as bilobed nucleoid (7)) with a local minimum near its center (Fig.2D,

Age=0.50). This minimum is accompanied by a sharper ribosome density maximum at the same location.

In this case, the observed densities of ribosomes and chromosomal DNA at the midcell are anti-correlated

at age = 0.5, consistent with the phase separation of nucleoid and ribosomes, as predicted by some

equilibrium thermodynamic models (50,51). The excess ribosome density increases, and the DNA density

at midcell decreases further as the cell cycle progresses (Fig. 2D, Age=0.70). However, at the last stages

of the cell cycle, the ribosome density decreases instead of further increasing, while the DNA density

decreases too (Fig. 2D, Age=0.90). Again, this positively correlated behavior cannot be explained by

equilibrium considerations of volume exclusion between ribosomes and chromosomes.

To show these changes occurring in the cell cycle, we plot the Pearson correlation coefficient, R, between

the DNA and ribosome distributions along the long axis of the cell as a function of cell age (Fig. 2E). At the

beginning of the cell cycle, the two distributions show small but positive R values, as expected based on

the coincidence of their maxima. At about cell age 0.5, the correlation coefficient starts to decrease and

becomes negative. This decrease in R corresponds to a concurrent decrease in DNA and an increase in

ribosome amounts at midcell (Fig. 2F). The correlation coefficient reaches its minimum at cell age ~0.75,

then starts to increase and becomes positive by the end of the cell cycle. At cell cycle age ~0.75, the midcell

ribosome numbers reach their peak and start to decrease thereafter, while the amount of DNA at midcell

continues to decrease until the end of the cell cycle (Fig. 2F). To understand the cause of the decrease in

ribosome numbers at midcell at the end of the cell cycle, we compared these data with the midcell phase

contrast signal from the same cells (Fig. 2F, black line). There is an upturn in the phase contrast signal at

the time when the ribosome numbers start to decrease. The increase in the phase contrast signal signifies

the onset of constriction (51). We find that at single-cell level the midcell ribosome density starts to decrease

after the onset of constriction (SI Fig. S8A) and shortly (10 min) before cells divide (SI Fig. S8B). The

simplest explanation for the decrease in midcell ribosome numbers after the onset of constriction is the

decrease in cell volume in the constricted region.

Based on traces in Fig. 2D, the excess fraction of ribosomes in the midcell is small throughout the cell cycle.

To quantify this fraction, we decompose the ribosomal density distribution into a midcell fraction and a

uniform background by fitting the data from individual cells (see SI Text, *Quantification of midcell amount*

and the excess midcell fraction). We find that the population-average excess fraction of ribosomes is only

4% of the cell total at its peak at cell age ~0.75 (Fig. 2G). At the same time, the excess fraction before the

DNA density minimum at midcell begins to form is 2.5%. The equilibrium thermodynamic models (50,51)

can qualitatively explain the increase in the excess fraction from 2.5% to 4%. However, they fail to explain

the 2.5% fraction present in the early cell cycle when there is no DNA density minimum at the center of the

nucleoid.

To verify the validity of the above findings, we constructed a strain in which the fluorescent mCherry fusion

was added to the C-terminus of the 50S ribosomal protein RplI (RplI-mCherry) instead of labeling the 30S

subunit (SI Fig. S9A-C). This label has also been previously characterized and tested (52). The cells with

RplI-mCherry label showed the same behavior as the one with RpsB-mCherry (SI Fig. S9D-G), including

co-localization of ribosome and DNA density distributions at the midcell at early cell cycle (SI Fig. S9D,

Age=0.15) and comparably small excess fraction of ribosomes at midcell at its peak value (3.5%) (SI Fig.
S9G).

We also investigated whether the same behaviors would occur in faster growth rates in M9 glucose + CAS
medium (SI Fig. S10). Under these conditions, the replication cycle initiates predominantly in the previous
cell cycle and completes at the early cell cycle (45). Accordingly, a minimum in the nucleoid density
distribution is already observed at the midcell at cell birth (SI Fig. S10C). Consequently, unlike in slow-
growth conditions, there is no positive correlation between DNA and ribosome density distribution during
the early stages of the cell cycle (SI Fig. S10D). Aside from this difference, the other behaviors (SI Fig.
S10E-F), including a small excess ribosome fraction at the cell center at its peak value (5.5%), are
consistent with those observed under slower growth conditions.

The DNA and ribosome correlation patterns in the early and midcell cycle are qualitatively
explained by the coupled DNA and ribosome dynamics model

To understand if the observed ribosome and nucleoid localization patterns can be explained by the active
dynamics of mRNA and ribosomes, we developed our own reaction-diffusion model (see SI Text, *The*
Coupled DNA and Ribosome Dynamics Model). We then compared the calculated cell cycle-dependent
DNA and ribosome densities from the model to the experimental data. Our 1D model is similar to the one
developed by Miangolarra *et al*. (34), considering mRNA synthesis and degradation, as well as the
assembly of polysomes (polyribosomes) from ribosomal subunits. Polysome and ribosomal subunits
interact with each other and with DNA via excluded volume interactions. However, unlike the previously
published model, ours uses a power-law term instead of a virial expansion to describe DNA (51,53) and
polysome (51) self-interactions, as the latter has been tested in experiments (53). As a simplification, our
model does not explicitly simulate bare mRNA and polysomes of varying sizes; instead, it considers only
an average-sized polysome. Unlike the previous model (34), ours includes cell growth, assuming DNA and
ribosome amounts increase with cell length (or volume).

To qualitatively compare the model to experimental data, we plotted the nucleoid and ribosome density
distribution along the cell length at different cell cycle times (Fig. 2H, SI Fig. S11A-B). The modeled
ribosome density represents the total density of ribosomal subunits (either 30S or 50S) and includes
contributions from both polysome-associated and free ribosomal subunits (SI Fig. S11C–D). As such, this
density is directly comparable to our experimental measurements. While the cell cycle times in Fig. 2H were
chosen to approximate those in Fig. 2D, they are not identical because the model does not exactly
reproduce the cell cycle timing when a single nucleoid starts to split into two lobes. Notably, however, the
ribosome distribution from the model shows a small local maximum at the cell middle, coinciding with the
maximum of the nucleoid distribution at the early stages of the cell cycle (Fig. 2H, Age 0.12, 0.22). This
maximum is qualitatively consistent with experimental data (Fig. 2D, Age 0.10, 0.30). The midcell maximum
in ribosome distribution arises in the model because 30S and 50S ribosomal subunits can diffuse into the

nucleoid region, where they assemble into polysomes. Since the rate of polysome formation is proportional
to DNA density in the model, the maxima of ribosome and DNA densities coincide.

The model predicts that excess ribosome density at the midcell cleaves the initial single nucleoid into two
nucleoid masses at some critical cell length due to excluded volume interactions (Fig. 2H, Age=0.32). After
this transition, ribosome and DNA density distributions become anti-correlated (Fig. 2H, Age=0.72; SI Fig.
S11E), which is consistent with experimental observations. Separation of the two DNA masses in the model
occurs rapidly, within 5% (0.05) of the cell cycle. During this transition, the excess fraction of ribosomes at
midcell reaches 48% and approximately maintains this value for the rest of the cycle (SI Fig. S11F). In
contrast, the experiment shows that the fraction reaches only 4.5% (Fig. 2G). The model thus correctly
predicts a local maximum in ribosome density at the cell middle in the early stages. However, it shows a
significantly higher accumulation of ribosomes in the cell middle in the later stages of the cell cycle
compared to the experimental data. We hypothesized that the substantially higher ribosome fractions
observed in the model are due to its 1D nature. In a 3D cell, ribosomes could form an outer shell around
the nucleoid, as has been proposed by earlier single-molecule tracking experiments of ribosomes (34).
Since this shell could not be derived from our 1D model, we added its effect *ad hoc* to our model (SI Fig.
S12A-C). Additionally, we assumed a uniform ribosome volume density in the shell that matches the density
predicted by the 1D model at the cell poles (for details, see SI Text, *The Concentric Shell Model*). With this
addition, the experimental (Fig. 2D) and modeling densities (Fig. 2I) showed similar values, with the model
yielding about 3% excess of ribosomes at midcell at its peak (SI Fig. S12E). However, to reach this
agreement, the inner radius of the added shell needed to be 50% of the cell radius, while it has been
estimated based on single-molecule tracking experiments that the inner radius is about 80% of the cell's
radius (35). These comparisons suggest that the coupled DNA and ribosome dynamics model could better
explain the data if it were implemented in 3D.

Constriction pushes aside daughter nucleoids

In the late stages of the cell cycle, the model and experimental data do not agree with each other. The data
reveal decreasing ribosome and nucleoid densities at the midcell (Fig. 2F), whereas the model maintains
almost a constant fraction (SI Fig. S11 A-B, F). The discrepancy appears to be due to the closing
constriction. To further understand the role of constriction closure in separating daughter nucleoids, we
carried out experiments in which the beta-lactam antibiotic cephalexin inhibited constriction formation. It is
well known that despite the inhibition of constriction, cephalexin-treated cells continue to grow in volume
and replicate DNA at the same rates as untreated cells (54). We then compared the changes in ribosome
and DNA densities between the cephalexin-treated and untreated cells. Since cephalexin prevents cell
division, cell age cannot be used to track the progression of DNA and ribosome densities, as done in the
previous measurements in Fig. 2. Instead, we determined the timing for the first appearance of the local
DNA density minimum at the cell center, T_{min}, following the procedure defined earlier (45) (for details see
SI Text, *Determining A_{min} and T_{min} from DNA and ribosomal distributions*). The appearance of this

minimum correlated with the replication termination timing, T_{rt}, and was detected on average shortly after

T_{rt} in the strain that carries DNA and replisome markers (SI Fig. S13A-B). We then aligned the time-

dependent data relative T_{min} (Fig. 3A, B) processing the data from untreated cells the same way (SI Fig.

S14A-B). At T_{min}, both the treated and untreated cells show an increase in midcell ribosome numbers (Fig.

3C) and a simultaneous decrease in the midcell DNA amount (Fig. 3D), consistent with the excluded volume

interactions between ribosomes and DNA. However, at later times ($time - T_{min} > 50$ min), the midcell

ribosome numbers stay constant in treated cells, while they decrease in untreated ones (Fig. 3C). This

comparison further confirms that the decrease in ribosome numbers in untreated cells is caused by the

closing constriction.

The data also indicates that, in addition to reducing the number of ribosomes at midcell, the constriction

affects the separation of the two daughter nucleoids. Comparison of DNA density at the midcell reveals a

clear decrease in DNA in untreated cells relative to treated ones for $time - T_{min} > 50$ min (Fig. 3D). The

effect of constriction is also noticeable in the measured distance between two nucleoid lobes, increasing

more significantly in untreated cells than in treated ones (Fig. 3E). Altogether, these data show that the

closing constriction separates two chromosomes in the late stages of the cell cycle. The effect is not caused

by FtsK, since FtsK is recruited to the divisome in both cephalexin-treated and untreated cells (55). The

constriction-driven separation appears, thus, to be a steric effect: the closing constriction physically pushes

aside the daughter nucleoids.

In untreated cells in slow-growth conditions, the onset of constriction occurs effectively at the time of the

DNA density minimum, T_{min} (45). Therefore, the effects arising from the active ribosome dynamics and

constriction on the separation of two daughter nucleoids are superimposed for $time > T_{min}$. However, the

data from cephalexin-treated cells can be directly compared with our model (for $time > T_{min}$) because it

excludes effects arising from the constriction process. The comparison shows that the DNA and ribosome

density distributions predicted by the model (Fig. 3 F-G) are at a qualitative level similar to the experimental

data from cephalexin-treated cells (Fig. 3A-B). In quantitative comparisons, the model reproduced some

trends in the experimental data but missed others. The model quantitatively correctly predicts the distance

between the peaks of two nucleoid lobes (Fig. 3H) but fails to predict the time-dependent decrease in the

amount of DNA at midcell (Fig. 3I). In the experiments, the DNA amount changes only 10% over 60 min,

while in the model it changes 50% within about 5 min. However, after this initial 5 min period, the model

and experiment show comparable rates of DNA density decrease at the midcell (Fig. 3I). The reason for

the ten-fold slower segregation dynamics observed in the experiments could be related to the choice of

model parameters. Reducing the diffusion coefficient of DNA 10-fold considerably slowed down the rate of

the change in DNA amount at the midcell, although the cell cycle time for the transition for DNA density

from high to low values shifted to later times (SI Fig. S15). Further adjustments of model parameters may

overcome this shift. Altogether, the model shows qualitative agreement with the experimental data. Its

quantitative accuracy could be potentially improved by extending the model to three dimensions and refining

its unknown parameters through additional experiments and optimization.

The effect of disruption of polysome formation on nucleoid separation

To perturbatively test the effects of active ribosome dynamics on the separation of two nucleoids, we treated

E. coli cells with a lethal concentration of rifampicin (300 μg/ml), which inhibits the initiation of RNA synthesis

(56) but does not stop ongoing DNA replication (6). As in the previous measurements, we also added

cephalexin to the culture medium at the same time as rifampicin to prevent cells from constricting. Since

rifampicin inhibits RNA synthesis, no new polysomes can form after the drug reaches the cells. The existing

transcripts can still be completed, but this occurs on the timescale of tens of seconds for a single gene and

about a minute for longer operons. These timescales are much shorter than our measurement frame rate

(4 min). On a slightly longer timescale (3-5 min (57)), the existing polysomes disassemble as their

associated mRNAs degrade, releasing ribosomal subunits.

To understand the effects arising from these processes, we compared data from cells treated

simultaneously with rifampicin and cephalexin (SI Fig. S16A–C) to cells treated only with cephalexin (SI

Fig. S16D–F). In this comparison, we align the time-dependent data to the antibiotic treatment start time,

T_{ab}, instead of T_{min} because most Rif-treated cells lack a detectable midcell DNA density minimum.

Additionally, we excluded from the analysis cells with visible constriction at the start of the treatment. As

expected, the localized accumulation of ribosomes at the midcell is not present in cells following rifampicin

treatment (SI Fig. S16B, C) due to polysome dissociation and subsequent diffusion of free ribosomal

subunits onto the nucleoid region.

Following rifampicin addition, the midcell DNA density minimum almost immediately ceases to deepen,

although it persists rather than disappearing (Fig. 4A). If steric interactions between DNA and polysomes

alone created this minimum, we would expect its complete disappearance on a timescale comparable to

polysome dissociation. In contrast, cephalexin-only treated cells continue to show a gradual decrease in

midcell DNA density (Fig. 4A). This difference between treatment conditions indicates that active mRNA-

ribosome dynamics contribute significantly to establishing the midcell DNA density minimum but are not

solely responsible.

We next examined nucleoid segregation by measuring the distance between nucleoid centers (D_{peak}) in

rifampicin-treated cells (Fig. 4B, SI Fig. S17A). This distance continues to increase for approximately 50

minutes, long after active ribosome dynamics ceases (within minutes of rifampicin addition) and polysomes

dissociate (3–5 min). Under normal growth conditions, nucleoids separate proportionally to cell elongation

(Fig. 1C). To determine if a similar relationship exists following rifampicin treatment, we compared nucleoid

center distances to cell length increases (SI Fig. S17), normalizing both metrics at the time of antibiotic

addition (Fig. 4B). Indeed, even without active ribosome dynamics, nucleoid centers continue separating

proportionally to cell elongation, which continues for about 50 min after treatment starts. This finding
indicates that cell growth alone is sufficient to drive nucleoid segregation (Fig. 4C). However, rifampicin
treatment also results in nucleoid expansion, causing increased midcell DNA density when two nucleoids
are present (Fig. 4A). These findings show that chromosome segregation progresses despite active
ribosome dynamics and polysome dissociation, while the separation of daughter chromosomes at the
midcell is affected.

To compare these data to our model, we divide the rifampicin-treated cells into three age groups based on
their nucleoid segregation stages. This grouping was necessary because rifampicin treatment starts at
different cell cycle ages for the asynchronous cell population in Fig. 4. The first group, representing 12% of
cells, includes those cells where T_{min} was detected before treatment began ($T_{min} < T_{ab}$, $N = 92$) (SI Fig.
S18). The second group, comprising 17% of cells, includes cells in which nucleoid separation was detected
after treatment started ($T_{min} > T_{ab}$, $N = 138$). The largest group, accounting for 71% of cells ($N = 556$),
shows no detectable nucleoid separation either before or after treatment. Our following analysis focuses
on the first two groups, as the third group shows minimal changes in ribosome and nucleoid distributions
after the treatment. To model these two groups of cells, we set the polysome formation rate to zero at the
start of the treatment. We then evolve the model so that cell length changes the same way as in the
experiments (see SI Text, *Rifampicin Treatment in the Model*, for details).

The group of cells that were late in their replication cycle at the onset of antibiotic treatment ($T_{min} < T_{ab}$),
show a local DNA density minimum (Fig. 5A) and a midcell ribosome accumulation at the start of rifampicin
treatment (Fig. 5B). The midcell ribosome accumulation vanishes within about 16 min from the start of the
treatment (Fig. 5B-C). The same occurs in the model, although at a faster rate (Fig. 5E-F). At the same
time, the model and the data qualitatively differ in how DNA distributions evolved. The experimental data
show that the midcell DNA density minimum persists, albeit becoming shallower over a 50-minute period,
before plateauing (Fig. 5A, C, SI Fig. S19A). In contrast, in the model, the two nucleoid masses merge so
that the local DNA density minimum at the midcell completely vanishes (Fig. 5D-F, SI Fig. S19B-C). The
coupled DNA ribosome dynamics model is thus unable to explain the persistence of the experimentally
observed nucleoid minimum after treatment.

We next examine the second group of cells ($T_{min} > T_{ab}$) where the DNA density minimum develops after
the start of rifampicin treatment (Fig. 5G-I, SI Fig. S20A). Although the observed DNA density minimum is
small in this group, its presence also contradicts the active mRNA-ribosome dynamics model. In the model,
the midcell DNA density minimum does not form without active ribosome dynamics (Fig. 5J-L, SI Fig. S20B-
C). The lack of a minimum is a qualitative feature of the model, independent of parameter selection. This
results from the overdamped nature of diffusive processes. Once the driving force is removed, that is, the
formation of polysomes is halted, then the separation of the two nucleoid lobes stops immediately (SI Fig.
S20B-C). The qualitative discrepancy between the model and the data for cells with $T_{min} > 0$ shows that
the active dynamics of mRNA and ribosomes is not responsible for DNA separation in this group of cells.

Discussion

Active mRNA-ribosome dynamics

This study used experiments and modeling to examine the role of mRNA-ribosome dynamics in the
segregation and separation of chromosomal DNA in *E. coli* cells. A characteristic feature of this mechanism,
as also observed in our measurements (Fig. 2), is the accumulation of polyribosomes (polysomes) at the
nucleoid (and cell) center before it starts to split into two lobes. According to our model, the accumulation
of polyribosomes results from strong excluded volume interactions between polysomes and chromosomal
DNA (51), leading to the exclusion of DNA from the midcell region. These qualitative features of the model
are consistent with our data in steady-state growth conditions. Perturbative measurements with rifampicin
treatment furthermore show that the nucleoid split does not develop in the majority of cells (77%) when
polysome formation is inhibited due to the lack of mRNA substrates. However, these measurements also
show that in cells where nucleoid lobes have already formed at the start of the treatment, the nucleoid
centers continue to move apart as the cells continue to elongate (Fig. 4H). The latter finding indicates that
the active mRNA-ribosome dynamics is not the sole process driving the late segregation of nucleoids.
Similar conclusions were recently reported by Papagiannakis *et al.,* although based on different evidence
(37).

Constriction effect

Our measurements also indicate that, in addition to active mRNA-ribosome dynamics, a closing constriction
plays an important role in separating the two daughter nucleoids and partitioning them into daughter cell
compartments (Fig. 3). This likely occurs via steric interactions between the constriction and DNA. These
steric interactions could potentially be mediated by crowders, ensuring that the inner membrane and
chromosomal DNA do not come into direct contact. Interestingly, earlier research has shown that the onset
of constriction is blocked until DNA segregation has progressed to the point where there is a local DNA
density minimum at midcell/mid-nucleoid (45). If the constriction were to start closing earlier, then it could
likely result in the trapping of incompletely replicated DNA under the constriction. Altogether, DNA
segregation to the point where there is a local DNA density minimum at the midcell appears to trigger the
onset of constriction, which completely separates one daughter chromosome from another.   The
segregation process can thus be self-amplifying through the formation of the constriction.

Evidence for additional mechanisms

In addition to the active mRNA-ribosome dynamics and closing constriction it is known from previous
research that DNA translocase FtsK (58,59) provides a DNA partitioning mechanism (29,32). FtsK usually
acts on about a 400 kb region near the replication termination region around the *dif*-site (29), but it can
directionally pump more than 3 Mb of chromosomal DNA across a closing constriction (32). This activity
mostly occurs at the end of the constriction phase when the pore separating the two daughters has a
diameter of less than 250 nm (32). The pore does not close fully until all the DNA is pumped across it.

Altogether, FtsK effectively guarantees a fail-safe completion of DNA partitioning. However, in most cell
divisions, FtsK is dispensable, as mutant cells are viable even when FtsK translocase activity has been
inhibited (32,60).

Our data suggest that, beyond active mRNA-ribosome dynamics, closing constriction, and FtsK-driven
translocation, additional mechanism(s) drive nucleoid segregation and separation. The main evidence for
this hypothetical mechanism comes from our finding that the two nucleoid centers move apart in proportion
to cell elongation. This movement also occurs during cephalexin and combined rifampicin and cephalexin
treatment. What could be this mechanism(s)? First, it can be ruled out that this process is due to the
equilibrium macromolecular crowding via excluded volume effects, which remains present after rifampicin
treatment. Indeed, passive crowding without reactions could lead to aggregation of DNA at high crowding
levels or simply hinder DNA segregation and separation due to increased viscosity. Consistent with this
view, a fusion of nucleoids can be observed in *E. coli* cells treated with chloramphenicol (61) – an antibiotic
which stops protein synthesis but not RNA synthesis. In this case, the high level of macromolecular
crowding appears to result from an increased level of RNA species (rRNA, tRNA, mRNA) in the cytosol
(62). Alternatively, it has been proposed that cell elongation could drive nucleoid segregation via transertion
linkages (63). However, this hypothesis can also be ruled out by our rifampicin measurements because the
linkages would disappear within a few minutes from the start of the treatment, whereas the daughter
nucleoids continue to move apart for a much longer period (50 min).

At the same time, the cell elongation-dependent movement of nucleoid lobes apart is potentially consistent
with the entropic segregation mechanism (17,18,25) under certain conditions. In particular, the entropic
segregation mechanism could explain that the distance between nucleoid lobes (D_{peak}) continues to
increase in proportion to cell length during rifampicin treatment (Fig. 4B, SI Fig. S17). During rifampicin
treatment, macromolecular crowding effects weaken because of the dissociation of polysomes to ribosomal
subunits (64). As a result, the chromosomal DNA expands and fills almost the whole cell volume (50,51).
According to the entropic segregation mechanism, one chromosome will occupy one half of the cell, and
the other chromosome the other half. The centers of each nucleoid are then at the ¼ position of the cell,
irrespective of cell length (33), and the distance between the lobes is proportional to cell length as observed
in our experiments. In our experiments, the nucleoids do not strictly fill the whole cell volume but remain
excluded from the poles (46). However, this exclusion does not significantly alter the above arguments
because the exclusion from poles is small, and the daughter DNA strands overlap at cell center (Fig. 4A).
As emphasized before, the overlap condition is strictly needed for the entropic mechanism to play any role
in driving the daughter nucleoids apart, and it cannot explain nucleoid separation when there is no overlap.

In summary, our measurements and modeling, combined with findings from previous research, bring out
the involvement of several partially overlapping processes in the segregation and separation of daughter
chromosomes in *E. coli*. Based on these findings, we propose the following sequence of events unfolding
during the cell cycle (SI Fig. S21): The entropic segregation mechanism may be involved during the

segregation stage when two chromosomes overlap (18,19). When about half of the DNA is replicated in

slow growth conditions, active mRNA-ribosome dynamics via excluded volume interactions become a

sufficiently strong driving force to lead to a midcell DNA density minimum. At this point, the entropic

mechanism of segregation still remains important, and, in fact, the associated force increases as more DNA

is replicated (26,27). However, the entropic force vanishes once two DNA molecules no longer overlap. In

parallel, the formation of the midcell DNA density minimum via the active mRNA-ribosome dynamics

triggers the onset of constriction in slow-growth conditions (45). The constriction further drives the two

daughter chromosomes apart and contributes to their separation. The final stages of DNA partitioning are

further aided by FtsK if DNA becomes trapped by the closing constriction.

Further work is needed to validate all the proposed processes mentioned above. In particular, the existing

evidence for the entropic mechanism remains indirect. Significant effort is also required to model more

faithfully the active mRNA-ribosome dynamics implementing the model in a 3D setting and including the

chain connectivity for DNA.

Materials and Methods

Details of bacterial growth conditions, microchip design, fluorescence microscopy setup, image analysis,

and modeling are given in the SI Text. The list of all strains and plasmids used in this study, including their

sources, is presented in the SI, Table S1. All experimental and modeling data are used in the main text, and

the SI figures are available from Zenodo (65). The Jupyter Notebooks used in our modeling are available

in reference (66).

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References

1. Bloom K, Joglekar A (2010) Towards building a chromosome segregation machine. *Nature* **463**:446-456. doi:10.1038/nature08912.

2. Wang X, Llopis PM, Rudner DZ (2013) Organization and segregation of bacterial chromosomes. *Nat. Rev. Genet.* **14**:191-203. doi:10.1038/nrg3375.

3. Gogou C, Japaridze A, Dekker C (2021) Mechanisms for chromosome segregation in bacteria. *Frontiers in Microbiology* **12**:685687. doi:10.3389/fmicb.2021.685687.

4. Batty P, Gerlich DW (2019) Mitotic chromosome mechanics: how cells segregate their genome. *Trends in Cell Biology* **29**:717-726. doi:10.1016/j.tcb.2019.05.007.

5. Woldringh CL, Hansen FG, Vischer NOE, Atlung T (2015) Segregation of chromosome arms in growing and non-growing *Escherichia coli* cells. *Frontiers in Microbiology.* **6**. doi:10.3389/fmicb.2015.00448.

6. Nielsen HJ, Li Y, Youngren B, Hansen FG, Austin S (2006) Progressive segregation of the Escherichia coli chromosome. *Mol Microbiol* **61**:383-393. doi:10.1111/j.1365-2958.2006.05245.x.

7. Bates D, Kleckner N (2005) Chromosome and replisome dynamics in *E. coli*: Loss of sister cohesion triggers global chromosome movement and mediates chromosome segregation. *Cell* **121**:899-911. doi:10.1016/j.cell.2005.04.013.

8. Viollier PH, Thanbichler M, McGrath PT, West L, Meewan M, McAdams HH, Shapiro L (2004) Rapid and sequential movement of individual chromosomal loci to specific subcellular locations during bacterial DNA replication. *Proc. Natl. Acad. Sci. U. S. A.* **101**:9257-9262. doi:10.1073/pnas.0402606101.

9. Nielsen HJ, Ottesen JR, Youngren B, Austin SJ, Hansen FG (2006) The *Escherichia coli* chromosome is organized with the left and right chromosome arms in separate cell halves. *Mol. Microbiol.* **62**:331-338. doi:10.1111/j.1365-2958.2006.05346.x.

10. Cass JA, Kuwada NJ, Traxler B, Wiggins PA (2016) *Escherichia coli* chromosomal loci segregate from midcell with universal dynamics. *Biophys J* **110**:2597-2609. doi:10.1016/j.bpj.2016.04.046.

11. Badrinarayanan A, Le TBK, Laub MT (2015) Bacterial chromosome organization and segregation. *Annu Rev Cell Dev Biol* **31**:171-199. doi:10.1146/annurev-cellbio-100814-125211.

12. Lin DC, Grossman AD (1998) Identification and characterization of a bacterial chromosome partitioning site. *Cell* **92**:675-685. doi:10.1016/s0092-8674(00)81135-6.

13. Livny J, Yamaichi Y, Waldor MK (2007) Distribution of centromere-like parS sites in bacteria: insights from comparative genomics. *J Bacteriol* **189**:8693-8703. doi:10.1128/JB.01239-07.

14. Cornet F, Blanchais C, Dusfour-Castan R, Meunier A, Quebre V, Sekkouri Alaoui H, . . . Bouet JY (2023) DNA segregation in enterobacteria. *EcoSal Plus* **11**:eesp00382020. doi:10.1128/ecosalplus.esp-0038-2020.

15. Brahmachari S, Marko JF (2019) Chromosome disentanglement driven via optimal compaction of loop-extruded brush structures. *Proc. Natl. Acad. Sci. U. S. A.* **116**:24956-24965. doi:10.1073/pnas.1906355116.

16. Harju J, van Teeseling MCF, Broedersz CP (2024) Loop-extruders alter bacterial chromosome topology to direct entropic forces for segregation. *Nat Commun* **15**:4618. doi:10.1038/s41467-024-49039-w.

17. Makela J, Sherratt D (2020) SMC complexes organize the bacterial chromosome by lengthwise compaction. *Curr Genet* **66**:895-899. doi:10.1007/s00294-020-01076-w.

18. Jun S, Mulder B (2006) Entropy-driven spatial organization of highly confined polymers: lessons for the bacterial chromosome. *Proc. Natl. Acad. Sci. U. S. A.* **103**:12388-12393. doi:10.1073/pnas.0605305103

19. Jun S, Wright A (2010) Entropy as the driver of chromosome segregation. *Nat. Rev. Microbiol.* **8**:600-607. doi:10.1038/nrmicro2391

20. Jun S, Arnold A, Ha BY (2007) Confined space and effective interactions of multiple self-avoiding chains. *Phys. Rev. Lett.* **98**:128303. doi:10.1103/PhysRevLett.98.128303.

21. Jung Y, Ha BY (2010) Overlapping two self-avoiding polymers in a closed cylindrical pore: implications for chromosome segregation in a bacterial cell. *Phys. Rev. E* **82**:051926. doi:10.1103/PhysRevE.82.051926.

22. Liu Y, Chakraborty B (2012) Segregation of polymers in confined spaces. *Phys. Biol.* **9**:066005. doi:10.1088/1478-3975/9/6/066005.

23. Dorier J, Stasiak A (2013) Modelling of crowded polymers elucidate effects of double-strand breaks in topological domains of bacterial chromosomes. *Nucleic Acids Res.* **41**:6808-6815. doi:10.1093/nar/gkt480.

24. Shin J, Cherstvy AG, Metzler R (2014) Mixing and segregation of ring polymers: spatial confinement and molecular crowding effects. *New Journal of Physics* **16**:053047. doi:10.1088/1367-2630/16/5/053047.

25. Chen YH, Yu WC, Wang JJ, Luo KF (2015) Polymer segregation under confinement: Influences of macromolecular crowding and the interaction between the polymer and crowders. *J. Chem. Phys.* **143**:134904. doi:10.1063/1.4932370.

26. Männik J, Castillo D, Yang D, Siopsis G, Männik J (2016) The role of MatP, ZapA, and ZapB in chromosomal organization and dynamics in *Escherichia coli*. *Nucleic Acids Res.* **44**:1216-1226. doi:10.1093/nar/gkv1484.

27. Polson JM, Kerry DRM (2018) Segregation of polymers under cylindrical confinement: effects of polymer topology and crowding. *Soft Matter* **14**:6360-6373. doi:10.1039/c8sm01062e.

28. Männik J, Bailey MW (2015) Spatial coordination between chromosomes and cell division proteins in *Escherichia coli*. *Frontiers in Microbiology* **6**:306. doi:10.3389/fmicb.2015.00306.

29. Stouf M, Meile J-C, Cornet F (2013) FtsK actively segregates sister chromosomes in *Escherichia coli*. *Proc. Natl. Acad. Sci. U. S. A.* **110**:11157-11162. doi:10.1073/pnas.1304080110.

30. Galli E, Midonet C, Paly E, Barre FX (2017) Fast growth conditions uncouple the final stages of chromosome segregation and cell division in *Escherichia coli*. *PLoS Genet.* **13**:1006702. doi:10.1371/journal.pgen.1006702.

31. Barre FX, Aroyo M, Colloms SD, Helfrich A, Cornet F, Sherratt DJ (2000) FtsK functions in the processing of a Holliday junction intermediate during bacterial chromosome segregation. *Genes Dev.* **14**:2976-2988. doi:10.1101/gad.188700.

32. Männik J, Bailey MW, O'Neill JC, Männik J (2017) Kinetics of large-scale chromosomal movement during asymmetric cell division in *Escherichia coli*. *PLoS Genet.* **13**:e1006638. doi:10.1371/journal.pgen.1006638.

33. Wu FB, Swain P, Kuijpers L, Zheng X, Felter K, Guurink M, . . . Dekker C (2019) Cell boundary confinement sets the size and position of the *E. coli* chromosome. *Curr. Biol.* **29**:2131-2144. doi:10.1016/j.cub.2019.05.015.

34. Miangolarra AM, Li SHJ, Joanny JF, Wingreen NS, Castellana M (2021) Steric interactions and out-of-equilibrium processes control the internal organization of bacteria. *Proc. Natl. Acad. Sci. U. S. A.* **118**: e2106014118. doi:10.1073/pnas.2106014118.

35. Sanamrad A, Persson F, Lundius EG, Fange D, Gynna AH, Elf J (2014) Single-particle tracking reveals that free ribosomal subunits are not excluded from the *Escherichia coli* nucleoid. *Proc. Natl. Acad. Sci. U. S. A.* **111**:11413-11418. doi:10.1073/pnas.1411558111.

36. Mohapatra S, Weisshaar JC (2018) Functional mapping of the *E. coli* translational machinery using single-molecule tracking. *Mol. Microbiol.* **110**:262-282. doi:10.1111/mmi.14103.

37. Papagiannakis A, Yu Q, Govers SK, Lin W-H, Wingreen NS, Jacobs-Wagner C (2025) DNA/polysome phase separation and cell width confinement couple nucleoid segregation to cell growth in *Escherichia coli*. *eLife*. doi:10.7554/elife.104276.1.

38. Youngren B, Nielsen HJ, Jun S, Austin S (2014) The multifork *Escherichia coli* chromosome is a self-duplicating and self-segregating thermodynamic ring polymer. *Genes Dev.* **28**:71-84. doi:10.1101/gad.231050.113.

39. Makela J, Uphoff S, Sherratt DJ (2021) Nonrandom segregation of sister chromosomes by *Escherichia coli* MukBEF. *Proc. Natl. Acad. Sci. U. S. A.* **118**:e2022078118. doi:10.1073/pnas.2022078118.

40. Sadhir I, Murray SM (2023) Mid-cell migration of the chromosomal terminus is coupled to origin segregation in Escherichia coli. *Nat Commun* **14**:7489. doi:10.1038/s41467-023-43351-7.

41. Reyes-Lamothe R, Sherratt DJ, Leake MC (2010) Stoichiometry and architecture of active DNA replication machinery in *Escherichia coli. Science* **328**:498-501.

42. Wery M, Woldringh CL, Rouviere-Yaniv J (2001) HU-GFP and DAPI co-localize on the *Escherichia coli* nucleoid. *Biochimie* **83**:193-200.

43. Wang P, Robert L, Pelletier J, Dang WL, Taddei F, Wright A, Jun S (2010) Robust growth of *Escherichia coli*. *Curr. Biol.* **20**:1099-1103.

44. Yang D, Jennings AD, Borrego E, Retterer ST, Männik J (2018) Analysis of factors limiting bacterial growth in PDMS mother machine devices. *Frontiers in Microbiology* **9**:871. doi:10.3389/fmicb.2018.00871.

45. Tiruvadi-Krishnan S, Männik J, Kar P, Lin J, Amir A, Männik J (2022) Coupling between DNA replication, segregation, and the onset of constriction in *Escherichia coli*. *Cell Reports* **38**:110539. doi:10.1016/j.celrep.2022.110539.

46. Yang D, Männik J, Retterer ST, Männik J (2020) The effects of polydisperse crowders on the compaction of the *Escherichia coli* nucleoid. *Mol Microbiol* **113**:1022-1037. doi:10.1111/mmi.14467.

47. Gray WT, Govers SK, Xiang YJ, Parry BR, Campos M, Kim S, Jacobs-Wagner C (2019) Nucleoid size scaling and intracellular organization of translation across bacteria. *Cell* **177**:1632-1648. doi:10.1016/j.cell.2019.05.017.

48. Bakshi S, Choi H, Mondal J, Weisshaar JC (2014) Time-dependent effects of transcription- and translation-halting drugs on the spatial distributions of the *Escherichia coli* chromosome and ribosomes. *Mol. Microbiol.* **94**:871-887. doi:10.1111/mmi.12805.

49. Bakshi S, Siryaporn A, Goulian M, Weisshaar JC (2012) Superresolution imaging of ribosomes and RNA polymerase in live *Escherichia coli* cells. *Mol. Microbiol.* **85**:21-38. doi:10.1111/j.1365-2958.2012.08081.x.

50. Mondal J, Bratton BP, Li YJ, Yethiraj A, Weisshaar JC (2011) Entropy-based mechanism of ribosome-nucleoid segregation in *E. coli* cells. *Biophys. J.* **100**:2605-2613. doi:10.1016/j.bpj.2011.04.030.

51. Chang MH, Lavrentovich MO, Mannik J (2024) Differentiating the roles of proteins and polysomes in nucleoid size homeostasis in *Escherichia coli*. *Biophys J* **123**:1. doi:10.1016/j.bpj.2023.11.010.

52. Chai Q, Singh B, Peisker K, Metzendorf N, Ge XL, Dasgupta S, Sanyal S (2014) Organization of ribosomes and nucleoids in *Escherichia coli* cells during growth and in quiescence. *Journal of Biological Chemistry* **289**:11342-11352. doi:10.1074/jbc.M114.557348.

53. Cunha S, Woldringh CL, Odijk T (2001) Polymer-mediated compaction and internal dynamics of isolated *Escherichia coli* nucleoids. *J. Struct. Biol.* **136**:53-66. doi:10.1006/jsbi.2001.4420.

54. Gelber I, Aranovich A, Feingold M, Fishov I (2021) Stochastic nucleoid segregation dynamics as a source of the phenotypic variability in *E. coli*. *Biophys. J.* **120**:5107-5123. doi:10.1016/j.bpj.2021.10.004.

55. Yu XC, Tran AH, Sun Q, Margolin W (1998) Localization of cell division protein FtsK to the Escherichia coli septum and identification of a potential N-terminal targeting domain. *J Bacteriol* **180**:1296-1304. doi:10.1128/JB.180.5.1296-1304.1998.

56. Wlodarski M, Raciti B, Kotar J, Cosentino Lagomarsino M, Fraser GM, Cicuta P (2017) Both genome and cytosol dynamics change in *E. coli* challenged with sublethal rifampicin. *Phys Biol* **14**:015005. doi:10.1088/1478-3975/aa5b71.

57. Bernstein JA, Khodursky AB, Lin PH, Lin-Chao S, Cohen SN (2002) Global analysis of mRNA decay and abundance in Escherichia coli at single-gene resolution using two-color fluorescent DNA microarrays. *Proc. Natl. Acad. Sci. U. S. A.* **99**:9697-9702. doi:10.1073/pnas.112318199.

58. Demarre G, Galli E, Barre FX (2013) The FtsK family of DNA pumps. *Adv Exp Med Biol* **767**:245-262. doi:10.1007/978-1-4614-5037-5_12.

59. Besprozvannaya M, Burton BM (2014) Do the same traffic rules apply? Directional chromosome segregation by SpoIIIE and FtsK. *Mol. Microbiol.* **93**:599-608. doi:10.1111/mmi.12708.

60. Kennedy SP, Chevalier F, Barre F-X (2008) Delayed activation of Xer recombination at *dif* by FtsK during septum assembly in *Escherichia coli*. *Mol. Microbiol.* **68**:1018-1028. doi:10.1111/j.1365-2958.2008.06212.x.

61. Van Helvoort J, Huls PG, Vischer NOE, Woldringh CL (1998) Fused nucleoids resegregate faster than cell elongation in *Escherichia coli* pbpB(Ts) filaments after release from chloramphenicol inhibition. *Microbiology* **144**:1309-1317. doi:10.1099/00221287-144-5-1309.

62. Xiang YJ, Surovtsev IV, Chang YJ, Govers SK, Parry BR, Liu J, Jacobs-Wagner C (2021) Interconnecting solvent quality, transcription, and chromosome folding in *Escherichia coli*. *Cell* **184**:3626-3642. doi:10.1016/j.cell.2021.05.037.

63. Woldringh CL (2002) The role of co-transcriptional translation and protein translocation (transertion) in bacterial chromosome segregation. *Mol. Microbiol.* **45**:17-29. doi:10.1046/j.1365-2958.2002.02993.x.

64. Männik J, Männik J, Amarasinghe C, Chang M-H, Lavrentovich MO (2025) The role of macromolecular crowders in the formation and compaction of the *Escherichia coli* nucleoid. *EcoSal Plus* **13**:eesp-0002-2024. doi:10.1128/ecosalplus.esp-0002-2024.

65. Amarasinghe C, Chang M-H, Lavrentovich MO, Männik J (2025) Source data for "The role of active mRNA-ribosome dynamics and closing constriction in daughter chromosome separation in *Escherichia coli*" (Zenodo). doi:10.5281/zenodo.15866580.

66. Chang M-H, Amarasinghe C, Lavrentovich MO, Männik J (2025) Modeling codes for "The role of active mRNA-ribosome dynamics and closing constriction in daughter chromosome separation in *Escherichia coli*" (Zenodo). doi:10.5281/zenodo.15866629.

Figures

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Fig. 1. Segregation of nucleoids as a function of DNA replication time (age). (A) The population average
DNA density distribution along the long axis of the cell as a function of replication cycle age. Time zero
corresponds to replication initiation and one to termination (also indicated by a vertical dashed line). HupA-
mCherry labeling is used to infer DNA density distribution. Red corresponds to high- and blue to low-
intensity values. The dashed horizontal lines indicate a midcell region. (B) The same for the average
replisome density distribution. Ypet-DnaN labeling is used to infer this distribution. (C) Left: definition of
parameters characterizing the dimensions of nucleoids: $L_{nuc,sum}$ and $L_{nuc,ext}$. Right: change of cell
length L_{cell}, $L_{nuc,sum}$ and $L_{nuc,ext}$ as a function of replication cycle time. The dashed vertical line marks
replication termination. (D) Left: schematics showing the region where HupA-mCherry and Ypet-DnaN
midcell amount is measured (indicated also by the horizontal dashed lines in panels A and B). The midcell
amount is based on the integrated intensity of fluorescent labels in this band. Right: the normalized midcell
amount for nucleoid (red) and replisome (green) as a function of replication age. The curves for individual
cells are normalized by the maximum value during the whole measurement period and then averaged over
the cell population. (E) Left: schematics showing the DNA density distribution for a single cell at a single
time point and how the depth of the local minimum, A_{min}, is defined. Right: A_{min} as a replication cycle time.
Error bars in panels C-E are s.e.m. N=438

Figure 2. Cell-cycle dependent changes in DNA and ribosome density distributions and their comparison to model predictions. (A) Population averaged kymograph of DNA density distribution (HupA-mYpet) along the long axis of the cell as a function of cell age. Zero corresponds to cell birth and one to division. Red corresponds to high- and blue to low-intensity values. The dashed horizontal lines indicate a midcell region. (B) The same for ribosome density (RpsB-mCherry) distribution. (C) Phase contrast kymograph for the same cell population. (D) The normalized population-averaged density of nucleoid (green) and ribosomes

(red) along the long axis of the cell at different cell ages. The curves are normalized by their peak values.

(E) Pearson correlation coefficient, R, between nucleoid and ribosome densities as a function of cell age.

The coefficient is calculated for each cell at different time points and then averaged over the cell population.

(F) The integrated intensity within the 0.75 μm wide band around the cell middle is shown for the nucleoid

(green), ribosome (red) and phase contrast signals (black). The curves from individual cells are normalized

by their peak values and then averaged over the cell population. The left vertical dashed line corresponds

to cell age of 0.40 and the right one to 0.68. (G) The relative excess fraction of ribosomes in the midcell

band as a function of cell cycle time. The curve is derived from a fitting of the ribosome signal (see Methods).

N=486 for panels A-G. The error bars are s.e.m. for panels E-G. (H) The normalized DNA and ribosome

densities along the long axis of the cell at different cell ages, as predicted by the coupled DNA and ribosome

dynamics model. The curves are normalized by their peak values. Note that the DNA (left) and ribosome

(right) signals have different scales, and the cell ages are comparable but not identical as those in panel D.

(I) The same curves that are recalculated, assuming a concentric shell of ribosomes around the nucleoid

(for details, see the Concentric Shell Model in Methods). The shell's inner radius is $0.5 \times R_{cell}$.

Figure 3. Segregation of nucleoids with and without closing constriction. (A) A population-averaged kymograph of DNA density along the long axis of the cell as a function of absolute (unscaled) time from the formation of local DNA density minimum, T_{min}, for cephalexin-treated non-constricting cells. The vertical axis shows normalized cell length, which remains unchanged throughout the experiment. Note that cells filament in this measurement. (B) The same for ribosome density distribution in this cell population. (C) The average number of ribosomes at midcell as a function of time from T_{min} for cephalexin-treated (red) and untreated cells (violet). The curves are normalized so that the amounts are one at time T_{min}. (D) The normalized DNA amount at midcell is presented the same way as in the previous panel. (E) The average distance between nucleoid centers (peak distance) as a function of time from T_{min} for cephalexin-treated (red) and untreated cells (violet). The distances are normalized at T_{min}. The schematic on the top shows how the peak distance is defined. (F) Kymograph of nucleoid density from the coupled DNA and ribosome dynamics model presented the same way as the experimental data in panel A. (G) Kymograph of ribosome density from the model. (H) The peak distance as a function of time relative to T_{min} from cephalexin-treated cells (red) and from the model (black). (I) Normalized midcell DNA amount for cephalexin-treated cells (red) and model (black). The linear fits to both curves are shown by dashed lines. The model is fitted for $\text{time} - T_{min} > 5 \text{ min}$. (J) The fit coefficients for panel I are listed in Table S3. All error bars are s.e.m. N=235.

Figure 4. The effects from the disruption of active mRNA-polysome dynamics with rifampicin (Rif) on nucleoid separation. (A) The depth of the local DNA density minimum at midcell as a function of time relative to the start of antibiotic treatment with Rif + Cef (blue) and Cef only (pink). The times are relative to the start of antibiotic treatment, T_{ab}, with Rif and Cef (cephalexin), and with only Cef, respectively. All analyzed cells in this figure are non-constricting. (B) The average cell length (green) and the nucleoid peak distance (black) as a function of time from the start of the treatment for Rif + Cef treated cells. The curves from individual cells are normalized to one at T_{ab}. Error bars in panels G and H are s.e.m. (C) Schematic summarizing the expansion and movement of nucleoids and ribosomes during Rif treatment. The top corresponds to a cell just before and bottom after Rif (+Cef) treatment. $N = 786$ for Rif+Cef and $N = 310$ for only Cef treated cells.

Figure 5. Comparing DNA and ribosome dynamics in Rif+Cef treated cells to the coupled DNA and ribosome dynamics model. Cells from two age groups are compared. (A-F) Data and model for cells that have formed a local DNA density minimum before the antibiotic treatment, $T_{min} < T_{ab}$ ($N = 92$). (A, B) The kymographs of DNA and ribosome density, respectively. The times are relative to the start of antibiotic treatment. (C) The average DNA (green) and ribosome (red) distributions along the long axis of the cell at different time points of experiment. (D, E) The DNA and ribosome density kymographs, respectively, from the model corresponding to group of cells with $T_{min} < T_{ab}$. The cell length in the beginning of the treatment is 2.87 μm. Note shorter times in the model compared to experimental data. (F) The average DNA (green) and ribosome (red) distributions along the long axis of the cell at different time points from the model. (G-

L) Data and model for cells that formed a local DNA density minimum after the start of the antibiotic

treatment, $T_{min} > T_{ab}$ ($N = 138$). (G,H) The kymographs of DNA and ribosome density, respectively. (I) The

average DNA (green) and ribosome (red) distributions along the long axis of the cell at different time points

of the experiment. (J, K) The DNA and ribosome density kymographs, respectively, from the model

corresponding to group of cells with $T_{min} > T_{ab}$. The cell length in the beginning of the treatment is 2.24 μm.

Note shorter times in the model compared to experimental data. (F) The average DNA (green) and

ribosome (red) distributions along the long axis of the cell at different time points from the model.