Lipid gating of BK channels and mechanism of activation by negatively charged lipids

Andrei Mironenko[1], Bert L. de Groot[1], Wojciech Kopec[1,2]*

[1]Computational Biomolecular Dynamics Group, Max Planck Institute for Multidisciplinary Sciences, Am Fassberg 11, 37077 Göttingen, Germany

[2]Department of Chemistry, Queen Mary University of London, 327 Mile End Road, London E1 4NS, United Kingdom

*Corresponding author: Queen Mary University of London, Mile End Road, London E1 4NS, United Kingdom, +49 551 201-2306, w.kopec@qmul.ac.uk, wkopec@gwdg.de

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Abstract

BK channels are a class of K^+ channels that possess an unusually high conductance and are synergistically gated by intracellular Ca^{2+} and voltage. Despite the significant array of experimental and computational data, many aspects of their function and dynamics remain unclear - such as how ion permeation is halted in the closed state of the channel. Available CryoEM structures obtained in deactivating conditions capture the channel with a wide, unobstructed pore, in contrast to e.g. a helix bundle crossing observed in some K^+ channels. Several hypotheses of BK closure were proposed, including selectivity filter and hydrophobic gating. In this work, we expand on the model of hydrophobic gating and focus on the role of lipids. Based on our atomistic and coarse-grained molecular dynamics simulations, we propose lipid entrance into the pore - either with lipid tails or entire lipid molecules - through the membrane-facing fenestrations to be a critical determinant of BK conductivity. Furthermore, we elucidate the mechanism of BK activation by negatively charged lipids, and suggest that they act by a multi-modal mechanism, which encompasses lipid entry reduction, increase of the K^+ occupancy of the pore, and stabilization of the channel's open-state structure - in broad agreement with experimental data. This presents an example of the crucial role of lipids in regulating BK channel's activity, and paves the way to further understanding of BK function in such complex environments as cellular membranes.

Introduction

Potassium (K^+) channels enable permeation of K^+ ions across biological membranes with high conductance and high selectivity, which plays an important role in a variety of cellular processes [1–3]. In particular, BK channels ('big K^+', also known as maxiK or Slo1) stand out among other K^+ channels due to their particularly large single-channel conductance - 100-300 pS [4]. BK channels are activated synergistically by intracellular Ca^{2+} and membrane depolarization; this allows them to regulate Ca^{2+} homeostasis and membrane potential, and participate in the processes such as the termination of the action potential, neurotransmitter release and many others [5–7].

The first full-length CryoEM structures were obtained for the BK channel from *Aplysia californica* [8,9], followed by structures from other organisms (including human); some of the structures also featured additional regulatory subunits [10–12]. The structures revealed that BK channels, in addition to the pore domain, contain the regulatory transmembrane voltage and intracellular Ca^{2+}-sensing domains that enable the synergistic activation (*Fig. 1A*). In certain cases, CryoEM structures were obtained either in the presence, or absence of Ca^{2+}, allowing to isolate the putatively open and closed states of the channel, respectively [8–11]. However, whether these structures actually correspond to open (conductive) or closed (non-conductive) states has been unclear, especially considering the fact that neither structure had a physical constriction of the permeation pathway that would prevent ion permeation in the closed state [8,9,11] - contrary to e.g. Kv channels, where a helix-bundle crossing sterically blocks the pore [13,14]. In a recent CryoEM study, mutations of voltage-sensing arginines that promote deactivation made it possible to

obtain structures with the pore as narrow as ~8 Å in diameter [12], which nevertheless can still fit a hydrated K^+ ion [15].

While it is possible that BK conformations with a fully constricted pore might exist, albeit not captured by the currently available structures - mechanisms other than physical block by helix-bundle-like structures have also been proposed. One such possibility is gating at the selectivity filter (SF). The SF is the narrowest part of the ion permeation pathway in K^+ channels that largely determines the channel's conductance and selectivity [16]. In some K^+ channels, ion permeation can be stopped via conformational changes at the SF, as is the case in several K2P channels [17,18]. While in BK the SF is in the same (conductive) conformation in both Ca^{2+}-bound and Ca^{2+}-free CryoEM structures [8,9], it was shown that large organic molecules are able to block the pore not only in the open, but in the closed state as well, raising the possibility that the channel is gated at the SF [19–23]; to note, some data support such block only in the partially closed state, but not fully closed [19]. In another study, molecular dynamics (MD) simulations of the human BK channel showed destabilization of the SF in the Ca^{2+}-free-like state [24], although it should be noted that a homology model was used for the structure of the channel, which, together with an initial configuration containing water molecules in the SF, could promote destabilization of the SF structure [16].

An alternative hypothesis is hydrophobic gating. In this mechanism, water molecules leave the inner pore of the channel in the closed state, creating a hydrophobic barrier for ion permeation through the pore [25]. Hydrophobic gating was first proposed for model nanopores based on MD simulations [26,27]: even if the diameter of a narrow pore was bigger than that of a water molecule, such pores could still undergo spontaneous dehydration-hydration transitions. There's structural and computational evidence of hydrophobic gating in several biological ion channels as well (e.g. GLIC [28,29], MscS [30–32], or TWIK-1 [33]). For BK, several MD studies showed cavity dehydration in the Ca^{2+}-free state, and a hydrated pore in the Ca^{2+}-bound state [34–36]; additionally, the Ca^{2+}-free structure underwent further pore narrowing and rotation of pore-lining residues, proposed to increase the hydrophobicity of the pore [34]. The idea of pore dehydration in BK agrees well with the experimental mutagenesis data, where the voltage of half activation correlates with the hydrophobicity of pores of various mutants [35–39].

Despite the overall support of the hydrophobic gating model in BK channels by experimental and computational data, several questions remain. First, to our knowledge, the ion permeability of either a hydrated or dehydrated pore has been tested directly using MD simulations only for the isolated pore domain of BK, which was shown to be conformationally unstable [35] - but not in the full-length channel. Second, in several MD studies of BK channels pore dehydration was accompanied by entrance of lipid tails into the pore through the fenestrations between the pore-forming S6 helices [12,34–36] - including a recent study where free energy calculations suggested the necessity of lipid entry for pore dehydration [40] - calling for further clarification of the role lipids play in regulating the channel's function. This is especially interesting, as the mentioned fenestrations occur only in the CryoEM structures obtained in deactivating conditions, but not in activating [8,9,11]. Moreover, some CryoEM structures capture lipid tails in these fenestrations [11]. At the same time, the inaccessibility of the BK pore to blockers in the fully

closed state observed in an experimental study [19] could indicate the existence of a steric block of the pore, which could be posed either by the channel itself, or by lipids occupying the pore. Finally, the involvement of lipids in BK regulation is further suggested by the dependence of the channel's activity on the membrane lipid composition - it was shown to be affected by the lipid's tail length [41], as well as head group charge and size [42,43].

In this work, we study the role of lipids in BK gating using atomistic and coarse-grained MD simulations. First, we directly test ion permeabilities of the full-length BK CryoEM states using simulations with applied voltage. Then, we describe the various modes of lipid behavior we observed: in particular, partial and full entry of lipids into the pore, as well as the interplay between lipid entry and conformational changes in the protein and its effect on ion permeation. We then investigate the mechanism of action of negatively charged lipids that activate BK [42,43], and demonstrate the integral role of lipid head group and tail dynamics in the pore in the activation mechanism.

Results

All-atom simulations of Ca^{2+}-bound and Ca^{2+}-free states

Previously, it was shown that Ca^{2+}-free states of BK generally undergo cavity dehydration in MD simulations, and the pore in Ca^{2+}-bound channels remains hydrated [34–36]; additionally, a recent MD study used applied voltage to test the conductivity of the isolated pore domain of BK, although without special restraints to prevent lipids from entering the pore the channel's tetrameric structure collapsed [35]. Therefore, the ability of the CryoEM states of the full-length channel to conduct ions has not been confirmed directly. To bridge this gap, we first performed atomistic MD simulations of BK from *Aplysia californica* in equilibrium (i.e. at 0 mV), and then under applied voltage starting from various frames of the equilibrium runs, in both Ca^{2+}-bound and Ca^{2+}-free states (PDB ID 5tj6 [9] and 5tji [8], respectively; +300 mV, POPC membrane, 0.8 M KCl, CHARMM36m). The simulation durations for each condition tested are given in *SI Table 1*.

The pore of Ca^{2+}-bound stayed hydrated throughout most of the equilibrium and applied voltage simulations in agreement with the previous studies, with the exception of spontaneous dehydration events under applied voltage (2 out of 6 replicas), and a decrease in pore hydration in one equilibrium replica out of 5 (*Fig. 1B, SI Fig. 1B*). Conversely, the Ca^{2+}-free channel showed a dehydrated pore in all simulations except 1) one equilibrium replica out of five, and 2) all three applied voltage simulations started from this replica (*Fig. 1B, SI Fig. 1A, C*). Notably, pore dehydration was always accompanied by entrance of lipids into the pore. We observed two types of such events: in the first type, only lipid tails entered through the membrane-exposed fenestrations between S6 helices of neighboring subunits present in the Ca^{2+}-free structure ('partial' lipid block), which was the case for several replicas of the Ca^{2+}-free state (*Fig. 1C, SI Fig. 1D*). The second type was entrance of entire lipid molecules through the fenestrations into the pore ('full' lipid block), which occurred in one equilibrium replica of Ca^{2+}-free (full lipid block was also preserved in the simulations at +300 mV spawned from this replica), as well as in the

two events of spontaneous dehydration of the pore in Ca^{2+}-bound at +300 mV, described in more detail below (*Fig. 1C, SI Fig. 1A, Fig. 2*). In terms of the channel's conductance, in simulations with applied voltage Ca^{2+}-bound showed steady outward K^+ permeation as long as the pore remained hydrated and not blocked by lipids (24±11 pS), and Ca^{2+}-free didn't permeate as long as either partial or full block occurred. It should be noted that the observed conductance value is an order of magnitude lower than the experimental one (up to 300 pS) [4] - a known feature of force fields without explicit description of polarization, such as the commonly used in biomolecular simulations CHARMM36m, also employed in this study [16,44]. Notably, in the simulations of the Ca^{2+}-free state at +300 mV where the pore remained hydrated and not blocked by lipids, the channel showed steady outward K^+ permeation (8±3 pS, *Fig. 1D, SI Fig. 1C, D, SI Table 1*).

Such diverse behavior of the Ca^{2+}-free state - a hypothetical closed state of the channel - not only reflects the earlier free energy calculations which showed that Ca^{2+}-free can possess a hydrated pore [40], but also raises a question of the relationship between the pore conformation, lipid block and the conductive state of the channel. Generally, we observed a variety of pore domain conformations in the simulations of the Ca^{2+}-free state. These conformations often differed from the initial CryoEM conformation, which specifically manifested in F304 side chains rotating into the pore, as well as overall pore narrowing (*Fig. 1E, F, SI Fig. 2*) - similar to observations in previous work [34]. To note, in one simulation replica without applied voltage, the pore became effectively sealed at the residue I308, indicating that full pore closure can in principle be achieved - however, it was concurrent with lipid entry, thus it is unclear whether full pore closure (and thus a possible steric block of ion permeation posed by the channel itself) is possible without lipid block (*SI Fig. 1A, SI Fig. 2A*). Overall, conformational changes in the Ca^{2+}-free state showed a certain degree of correlation with the conductive state of the channel; specifically, the aforementioned simulation replicas of Ca^{2+}-free under voltage which did not undergo pore dehydration and lipid block and were conductive (*SI Fig. 1C*), showed a smaller degree of pore narrowing and F304 rotation relative to the CryoEM structure, compared to the replicas which did undergo pore dehydration/lipid block and were not conductive (*SI Fig. 2B*). To note however, the equilibrium replica which underwent full lipid block, as well as the consecutive simulations under voltage where the channel was not conductive - and lipids continued to reside in the pore - did not show deviation in the pore width compared to the Ca^{2+}-free CryoEM value (and showed a smaller degree of F304 rotation) (*SI Fig. 1A, SI Fig. 2A - 'rep4', SI Fig. 2C*), suggesting dependence of conformational behavior on the mode of lipid block (partial or full). Thus, to summarize, based on our equilibrium and applied voltage simulations of the full-length BK, we confirm that the Ca^{2+}-bound structure corresponds to an open, conductive state of BK. At the same time, the interplay between further conformational changes in the channel's pore relative to the Ca^{2+}-free CryoEM structure and the obligatory pore dehydration/lipid block was needed to stop ion permeation in the Ca^{2+}-free state.

Spontaneous closure of the open state

As we mentioned earlier, the Ca^{2+}-bound state remained hydrated and conductive throughout most of our simulations; however, in two out of six replicas at +300 mV the pore of the open channel underwent spontaneous dehydration, accompanied by a full lipid block of the pore in both cases (*Fig. 2A, SI Fig. 3*) - in addition to one equilibrium replica with decreased pore hydration and increased lipid presence in the pore (*SI Fig. 1B*). Generally, the entrance of an entire lipid molecule was made possible by widening of one of the fenestrations between S6 helices of neighboring subunits (*Fig. 2B, SI Fig. 4B*), accompanied by motions of bulky hydrophobic residues obscuring the fenestrations in the Ca^{2+}-bound structure, such as F307 and I308 (*Fig. 2C*). Interestingly, opening of the fenestrations and motions of these residues resemble the respective differences between Ca^{2+}-bound and Ca^{2+}-free CryoEM structures, supporting the notion that we observed transitions from the open to the closed state of the pore (*Fig. 2D*). Additionally, in a previous study that combined MD simulations and mutagenesis experiments, the positively charged residue clusters at the end of S6 helices (corresponding to ^{317}KRQK^{320} in the BK from *Aplysia californica* studied here) and S6 glutamates of the neighboring subunits ('S6Glu') were shown to play a role in the stabilization of the open state of the channel, and proposed to form salt bridges in the closed, but not the open state of the channel [24]. In our simulations, we observed clear formation of such salt bridges before fenestration widening in one of the two replicas with a spontaneous lipid block (*Fig. 2E, SI Fig. 4A*). In the other replica with a full lipid block, however, the salt bridges were forming only transiently - similar to several other replicas with no lipid block. Additionally, one of the replicas with no lipid block showed formation of a relatively long-lived salt bridge (*SI Fig. 4A*). Taken together, this could indicate the existence of mechanisms that induce fenestration widening (and thus increased propensity for a lipid block) with some degree of independence of salt bridge formation - or, alternatively, the possibility of pore closure and lipid block happening at longer simulation times in these replicas as well.

Effect of membrane composition on lipid block

Membrane lipid composition plays a significant role in regulating the behavior of BK channels [45,46]. In light of the model of lipid block as the closure mechanism in BK, it is thus interesting to not only test the physical basis of action of different lipids broadly, but also elucidate whether it might be related to different propensities of lipids to block the channel's pore. Here, we focus on the negatively charged POPS lipid, the presence of which - and negatively charged lipids in general - was shown to increase both the open probability (Po) of the channel and its single-channel current [43]. To tackle this issue computationally, we first performed coarse-grained simulations of Ca^{2+}-bound and Ca^{2+}-free states of BK using the Martini3 force field [47] in a POPE:POPS (3:1 mol.) membrane - the ratio used in the experimental study [43] - with control simulations in a POPC membrane. The switch from all-atom to coarse-grained simulations was done to increase the accessible simulation times and extensively sample the relevant configurations of lipids around the protein. After 40 μs-long simulations for each combination of the channel state and membrane composition, we observed a significant decrease in the pore occupancy by lipids in POPE:POPS compared to POPC (*Fig. 3A*). Interestingly, the pore stayed hydrated in the Ca^{2+}-free state, (*SI Fig. 5A*), opposite to what we observed in most of the

atomistic simulations of Ca^{2+}-free. To note, in the Martini3 variant with elastic networks that we used [47,48], the conformational changes of the Ca^{2+}-free we observed in the atomistic simulations were restricted, which might have increased the barrier for pore dehydration and/or lipid entry (reflecting the previous observation of a certain degree of correlation between the propensity for lipid block and further conformational changes relative to the Ca^{2+}-free CryoEM structure, *SI Fig. 2B*). In line with the hydrophilic environment of the pore, the number of the Martini lipid beads in the pore in Ca^{2+}-free was rather low for both POPE:POPS and POPC membranes - on average less than ~3 out of 12 beads in a single residue for every lipid type, which would correspond to ~8-12 out of ~50 non-hydrogen atoms in a single lipid in atomistic representation. Reflective of the hydrophilic environment, it was lipid headgroups which occupied the pore most frequently as opposed to tails (*SI Fig. 6*). When water was artificially excluded from the pore using restraints in the Ca^{2+}-free state, POPE:POPS and POPC membranes showed practically no difference in lipid entry, with a drastically increased average number of lipid beads in the pore for both membrane compositions (*Fig. 3B, SI Fig. 5B*). Given the absence of the effect of POPS on the fully closed state characterized by a hydrophobic pore and substantial lipid entry, this suggests that the presence of POPS in the membrane increases the barrier for lipid entry only in the open and intermediate closed states, as long as the pore is still hydrated.

Interestingly, not only did the average number of lipid beads in the pore decrease in presence of POPS, but so did the average number of individual lipid residues simultaneously occupying the pore (*SI Fig. 6C*). Notably, the S6 helices contain two negatively charged glutamates per chain - E310 and E313, the same glutamates that form salt bridges in the closed state with the KRQK cluster at the lower ends of S6 helices mentioned above (*Fig. 3C*). Here however, the presence of eight negatively charged residues close to the fenestrations between S6 helices could in principle provide a barrier for entry of similarly charged POPS molecules. This could be especially relevant for lipid entry into the pore, as even though POPS comprises only a quarter of the membrane, the area around the protein is actually enriched with POPS, including the region near the pore (*Fig. 3D, SI Fig. 7A*). Indeed, the average electrostatic potential calculated at the level of E313, which coincides to a large extent with the position of lipid headgroups near fenestrations (*SI Fig. 7B*), is more negative both in the fenestrations and the pore relative to the regions of the membrane close to the fenestrations in Martini simulations (*Fig. 3E*). Thus, we suggest the electrostatic repulsion between E310/E313 and the negatively charged POPS headgroups as the driving force behind reduced lipid entry, explaining the increased open probability of the channel in presence of POPS observed in experiments [42,43]; mutual repulsion between the headgroups of different POPS residues that occupy the fenestrations could also in principle contribute to this effect.

Previously, we mentioned that in simulations of Ca^{2+}-bound and Ca^{2+}-free (hydrated pore) states primarily the headgroups contributed to the transient occupancy of the pore by lipids. Aligning with the larger POPS occupancies in the membrane regions surrounding the protein, the pore in POPE:POPS was enriched by negatively charged beads (measured as the occupancy difference between the negatively and positively charged beads), compared to POPC in our Martini3 simulations, with the biggest increase in the simulations of the Ca^{2+}-free state with a hydrated pore (*Fig. 4A, middle*). In the Ca^{2+}-bound state this effect was less pronounced, likely

due to the tightly closed fenestrations (*Fig. 4A, middle*). At the same time, in the Ca^{2+}-free state with a hydrated pore the average number of K^+ residing in the pore increased, whereas it remained practically unchanged in the Ca^{2+}-bound state (*Fig. 4A, left*). It is possible that a similar increase in local K^+ concentration in the pore could generally happen in open, Ca^{2+}-bound-like conformations due to the channel's intrinsic flexibility (and e.g. some degree of fenestration opening that would allow the negatively charged headgroups to be closer to the pore), which was however restricted by the elastic networks in the Martini force field. In our control atomistic simulations of Ca^{2+}-bound, we did observe an increase in the average number of K^+ ions in the pore in one of the two POPE:POPS simulation sets, compared to the other with a different initial membrane configuration (as well as relative to POPC simulations) - while POPS headgroups were present in the pore in this simulation set and absent in the other. It should be noted that the number of K^+ ions in the pore correlated with the pore occupancy by lipid tails as well, which was different in both POPE:POPS simulation sets, as well as in POPC, possibly reflecting the limited sampling in the atomistic simulations (*SI Fig. 8*). Taking these force field- and sampling-related limitations into account, and drawing from the mechanism of activation proposed for small-molecule negatively charged activators of K^+ channels (NCAs) [49], we cautiously propose the explanation for the increased single channel current by negatively charged lipids: the headgroups of negatively charged lipids occupy the fenestrations (or transiently the pore), thereby electrostatically attracting and increasing the local concentration of K^+ in the pore, leading to an increased probability of a permeation event, via multi-ion knock-on mechanism (*Fig. 4A, right*).

The effects we observed so far relate primarily to the propensity of different lipids to enter the pore, which can be both inhibitory, as is the case for lipid block, and activatory, in case of the transient pore occupancy by POPS headgroups. However, another possible dimension of the activatory effect of negatively charged lipids is stabilization of the open state conformation of the channel. In Martini simulations we observed disproportionately frequent interaction of POPS (i.e. with a frequency bigger than the fraction of POPS in the membrane, or 1/4) with the KRQK cluster, and, conversely, rare interactions with the S6 glutamates, reflective of the electrostatic attraction between POPS molecules and KRQK, and repulsion from E310/E313 (*Fig. 4B*). This, in principle, could interfere with the formation of the S6 glutamates-KRQK salt bridge characteristic of the closed state, an effect proposed recently for the by phosphatidylinositol 4,5-bisphosphate (PIP_2, charge -3) activation of the SK2 channel, in which similar salt bridges can form [50]. Indeed, in our control atomistic simulations, we observed a higher frequency of states with larger KRQK-glutamate distances in POPE:POPS compared to POPC, indicating stabilization of the pore domain's structure (*SI Fig. 9A*). Similar to the Martini simulations, S6Glu was depleted of POPS relative to POPE in both simulation sets, and KRQK enriched with POPS in one simulation set; in the other set, KRQK did not show disproportionately frequent interactions with POPS relative to its molar fraction in the membrane, nevertheless the existing POPS-KRQK contacts and the repulsion of POPS from S6Glu appeared to be sufficient to stabilize the channel's conformation in this case (*SI Fig. 9B*). Additionally, we performed atomistic simulations of an isolated pore domain in POPE:POPS and POPC as an approach to increase sampling due to a largely reduced size of the system (*SI Table 1*). However, despite longer simulations and a larger number of initial membrane configurations tested, we observed

very low conformational stability of the pore in both membrane compositions (*SI Fig. 10*) - indicating caution that needs to be exercised when using the isolated pore domain as a model of BK in MD simulations. To summarize, based on our simulations we propose a multi-modal mechanism of activation by negatively charged lipids, involving reduced lipid entry, increased K^+ occupancy of the pore and stabilization of the open state conformation.

Discussion

BK channels present an interesting case where CryoEM structures, despite being obtained in activating and deactivating conditions, do not provide a clear definition of open (conductive) and closed (non-conductive) states, particularly due to the absence of a physical constriction in the permeation pathway in deactivating conditions, such as a helix-bundle crossing [51,52], or a collapsed SF [53,54] occurring in some other K^+ channels. Later, computational and experimental studies suggested a model where the channel's pore in the closed state undergoes dehydration, thus establishing a hydrophobic barrier to permeation of K^+ [34–36,39]. However, to our knowledge, the ion permeability of the BK CryoEM states has not been tested directly in the full length BK, such as with MD simulations with applied voltage. At the same time, the pore-lining S6 helices of neighboring subunits form membrane-exposed fenestrations in Ca^{2+}-free structures (deactivating conditions), which in some CryoEM structures contain lipid tails [11], in addition to the occupancy of the pore by lipid tails in several MD studies [34,35,40], proposed to increase the hydrophobicity of the pore in the closed state. In this work, we focused on the role of lipids in the BK channel's gating. Our MD simulations of the full-length channel in equilibrium and under applied voltage allowed us to test the conductivity of the full length channel in Ca^{2+}-bound and Ca^{2+}-free states directly, and show that the pore block by lipids is a defining characteristic of the closed, non-conductive state of the channel. Pore block occurred in two modes: in one mode, only tails of lipids can enter the pore through the S6 fenestrations, likely not only increasing the hydrophobicity of a dewetted closed-state pore, but also posing a steric barrier to ion permeation ('partial block', *Fig. 1C, left*). Second, in certain cases (including the events of spontaneous pore closure in the Ca^{2+}-bound state) entire lipid molecules were able to enter into the pore through the fenestrations ('full block', *Fig. 1C, right*). Both modes of lipid block were able to stop ion permeation on MD timescales (microseconds). Generally, both mechanisms align with experimental data, including some of the CryoEM structures of BK containing lipid tails in fenestrations mentioned earlier (supporting partial block) [11], inaccessibility of the pore in the closed state for large pore blockers (supporting both full and partial block, with lipid tails providing a steric barrier in the latter case) [19], and dependence of Po and single-channel currents on the membrane composition (such as activation of the channel by lipids with bigger headgroups, which can be explained by a higher energetic barrier for full lipid entry through the fenestrations between the S6 helices) [45,46,55]. Additionally, in a CryoEM structure of a K^+ channel from the K2P family (TREK-1), which has similar membrane-exposed fenestrations in the closed state, it was possible to resolve a lipid molecule occupying the pore [56], not only confirming the results of an earlier MD study showing lipid entry into the pore in another K2P family member TWIK-1 [57], but also indicating that a full lipid block can generally serve as a pore closure mechanism. Interestingly, in previous MD studies of BK it was possible to achieve dehydration

of the pore without lipid entry using restraints that prevented lipids from entering the pore, which could, in principle, be sufficient to prevent ion permeation [35,36]. However, in our simulations where the motions of lipids were unrestrained, pore dehydration was generally accompanied by lipid entry, and, given the mentioned experimental data that broadly agrees with the model of lipid block, it is unclear how frequent in reality would be the events of dehydration without the presence of lipids in the pore.

The observation of two modes of lipid block - partial and full - also raises a question, whether they could correspond to different functional states of BK. In electrophysiological recordings, BK can occupy several open (conductive) and closed (non-conductive) states that have different life times [58]. Naively it can be assumed that the barrier to relieve the partial block (move only the lipid tails out of the pore) is lower than that for the full block, thus in the framework of lipid block as the closing mechanism, the partial block would correspond to the short-lived closed states, and full block to the long periods of closure.

One hypothesis regarding the Ca^{2+}-free state lacking a physical constriction of the pore is that it does not represent the final closed state, but rather an intermediate one, and further closure of the pore is possible [8], as is the case for a prokaryotic homolog of the BK channel, MthK, which shows a helix bundle-like structure in its closed state [59]. Indeed, later cryoEM structures and MD simulations of BK with mutations in the voltage sensor allowed to capture a state with a narrower pore (albeit wider than a hydrated K^+ ion) [12,34]. Further pore narrowing also occurred in practically all of our simulations of the Ca^{2+}-free state where the partial lipid block happened. Additionally, in Martini3 simulations where such changes were limited by elastic networks, the pore remained hydrated and with a relatively low degree of lipid entry in the Ca^{2+}-free state, unless restraints preventing pore hydration were used, presenting further evidence of the importance of additional conformational changes in achieving the fully closed state. However, it should be noted that the mode of pore closure here differed from the helix bundle mechanism - the pore closed only partially in all replicas except one with full pore closure accompanied by lipid block - and at least in all cases we observed it appeared to be only accessory to lipid block.

Finally, we investigated the mechanism of BK activation by negatively charged lipids. It has been shown that such lipids can increase both the Po of the channel, as well as single-channel current [42,43]. In light of our observation of lipid block characterizing the closed state of the channel, it was interesting to test whether activation by negatively charged lipids is related to different lipids having different propensities for lipid block. Based on our atomistic and coarse-grained simulations of a membrane containing the negatively charged POPS lipid, we suggest a multi-modal mechanism of its action, with 1) reduced lipid entry in open and intermediate closed states when the pore is still hydrated, and thus increased Po, 2) increased local K^+ concentration in the pore, which explains the increase in the single-channel current, and 3) stabilization of the open-state conformation of the channel via disruption of the ^{317}KRQK^{320}-E310/E313 salt bridge formation, also explaining the increased Po. Considering the latter mechanism, a mutagenesis study featuring another negatively charged lipid - PIP_2 - showed that in the human BK channel, it interacts specifically with the RKK motif, equivalent to KRQK in the *Aplysia* BK [42]. Additionally, for another K^+ channel with similar positively/negatively

charged residues in the lower S6 helices - SK2 - MD simulations and mutagenesis also suggested disruption of the salt bridge and activation by PIP_2 [50]. Even though experimentally this effect was studied specifically for PIP_2, we hypothesize that the channel's open-state structure might be stabilized via this mechanism by negatively charged lipids in general.

In terms of the physical basis of activation by negatively charged lipids, electrostatic interactions played a key role: eight E310/E313 residues on S6 helices contribute to regions of negative electrostatic potential near and in the pore that increases the barrier for POPS molecules to enter it (*Fig. 3C, E*). At the same time, the general enrichment of the areas near the channel by POPS enhances this effect on lipid entry, despite POPS constituting only 25% of the membrane in our case (*Fig. 3D, SI Fig. 7A*). In experiments, the open probability of the channel was shown to be higher with a bigger negative charge on the lipid's headgroup [42], which is in agreement with the idea of electrostatic repulsion of the negative charges on the lipids from the pore by S6 glutamates. Conversely, positively charged lipids were shown to decrease the open probability [60] - in the framework of the model we propose, their effect could be explained by negatively charged E310/E313 attracting the positively charged lipids into the pore, although this would require further validation (such as by studying the effect of E310/E313 mutants on the channel's activity in positively and negatively charged lipid-containing membranes; it should be noted that such mutants might affect not only lipid entry, but also the conformational stability of open and closed states, given the role that E310/E313 also play in the latter [24]). Similarly, POPS headgroups residing close to the pore attracted K^+ into it (*Fig. 4A*) - we propose that the electrostatic attraction between the negatively charged headgroups and K^+ lies behind the increase in the single channel current, akin to the mechanism proposed for small-molecule negatively charged activators of K^+ channels [49].

Interestingly, some previous studies suggested a dual effect of the negatively charged PIP_2 - not only activatory, but also inhibitory at high enough voltages [61]. Such inhibition was hypothesized to originate from PIP_2 binding to an intersubunit cleft between two voltage sensors, thereby augmenting the negative charge on the cytosolic side of the membrane and effectively pulling the voltage sensors into the 'down' (inactivated) configuration, as evidenced by suppression of inhibition in presence of voltage-sensor-binding β subunits, or via shielding of charges by intracellular divalent cations [61]. In our Martini3 simulations, POPS forms a distinct binding site in the outer voltage sensor region (*Fig. 3D, SI Fig. 7A, SI Fig. 11*). This generally supports the notion of involvement of negatively charged lipids in voltage activation, however the associated conformational changes in the voltage sensor and their coupling to the pore domain are beyond the scope of this work.

To conclude, we present a model of BK gating where lipids play an integral role in determining the conductive state of the channel, and suggest a multi-modal mechanism of how the negative charge of lipid headgroups can activate the channel (*Fig. 5*). BK channels exist in complex multicomponent lipid environments, such as plasma, mitochondrial and nuclear membranes, also in various cell types [7]. At the same time, BK channel activity was shown to be regulated, in addition to the lipid headgroup charge, by headgroup size, tail length, as well as specific protein-lipid interactions [41–43,45,46,55,60]. This likely leads to intricate patterns of regulation in

multicomponent membranes, not only presenting further opportunities to study and validate the role of lipids in the function of BK channels, but also help inspire the design of similarly acting new small molecule modulators of BK channels.

Methods

System preparation

CryoEM structures of BK from *Aplysia californica* in Ca^{2+}-bound (PDB ID 5TJ6) and Ca^{2+}-free (PDB ID 5TJI) states [8,9] were used throughout this study. To build the all-atom systems, the CryoEM structures of the full length channel were embedded in a 1-palmitoyl-2-oleoyl-glycero-3-phosphocholine (POPC) lipid bilayer using CHARMM-GUI [62,63]. The non-terminal loops missing from the structures were added, and N- and C-termini were acetylated and methylamidated, respectively. The system with the Ca^{2+}-bound structure contained 1523 POPC molecules, 3641 K^+, 3649 Cl^-, 247063 water molecules, as well as 4 Mg^{2+} and 8 Ca^{2+} ions in their CryoEM binding sites. The system with the Ca^{2+}-free contained 1503 POPC, 3629 K^+, 3609 Cl^- and 245986 water molecules. This includes 4 K^+ added to the SF (binding sites S1-S4), as the CryoEM structure of the Ca^{2+}-free state, opposite to Ca^{2+}-bound, doesn't have any K^+ resolved inside the SF. Additionally, 2 water molecules per subunit were added to the cavities behind the SF for each CryoEM structure, as water molecules in these positions are usually present in high-resolution structures of other channels, and were shown to have a stabilizing effect on the SF structure [64–67]. The number of K^+ ions was chosen to yield a concentration of 0.8 M, and the number of Cl^- to yield systems with a total charge of zero.

The coarse-grained structures of the protein for Martini3 [47] simulations were obtained using martinize2 [68], and the protein-membrane systems were built using insane [69]. For the simulations with POPC membranes, the system with the channel in the Ca^{2+}-bound state contained 1602 POPC molecules, 3400 K^+ beads, 3384 Cl^- beads and 53638 water beads (regular); the system with the channel in the Ca^{2+}-free state contained 1546 POPC molecules, 3407 K^+, 3387 Cl^- and 53722 water beads. The systems with 1-palmitoyl-2-oleoyl-sn-glycero-3-phosphoethanolamine (POPE) and 1-palmitoyl-2-oleoyl-sn-glycero-3-phospho-L-serine (POPS) and a 3:1 ratio contained, in case of the Ca^{2+}-bound state, 400 POPS, 1201 POPE molecules, and 3600 K^+, 3184 Cl^- and 53642 water beads; in case of the Ca^{2+}-free state, 385 POPS, 1159 POPE molecules, 3599 K^+, 3194 Cl^- and 53716 water beads.

Systems containing the full-length channel for the control atomistic simulations in the POPE:POPS membrane were built similarly to the simulations of the full-length channel in POPC, as was the case for the isolated pore domain in POPE:POPS and POPC. To mitigate the limited membrane configuration sampling due to slower lipid dynamics compared to the Martini3 simulations [47], two systems with different initial configurations of lipids in the membrane were built for each state of the full-length channel in POPE:POPS. For the Ca^{2+}-bound state of the protein, one system contained 1200 POPE, 400 POPS, 3367 K^+, 2975 Cl^-, 226858 water molecules, as well as 4 Mg^{2+} and 8 Ca^{2+} in their binding sites; the other contained 1200 POPE,

400 POPS, 3368 K^+, 2976 Cl^-, 226877 water molecules, 4 Mg^{2+} and 8 Ca^{2+}. For the Ca^{2+}-free state, one system contained 1125 POPE, 375 POPS, 3191 K^+, 2796 Cl^- and 215037 water molecules; the other contained 1125 POPE, 375 POPS, 3191 K^+, 2796 Cl^- and 215033 water molecules.

For the isolated pore domain, one system in the POPC membrane, and four systems in the POPE:POPS membrane with different initial configurations of lipids in the membrane were built for each state of the channel. The POPC system with the Ca^{2+}-bound state contained 127 POPC, 106 K^+, 94 Cl^- and 7167 water molecules; with Ca^{2+}-free 126 POPC, 114 K^+, 102 Cl^- and 7676 water molecules. In POPE:POPS, for the Ca^{2+}-bound state each system contained 99 POPE, 33 POPS and, correspondingly, 142 K^+, 97 Cl^-, 9644 water molecules for systems 1 and 2; 142 K^+, 97 Cl^-, 9619 water molecules for system 3; and 141 K^+, 96 Cl^-, 9613 water molecules for system 4. For the Ca^{2+}-free state, they contained 99 POPE, 33 POPS, 147 K^+, 102 Cl^- and 1) 9969, 2) 9973, 3) 9979, 4) 9959 water molecules.

MD simulations

The systems containing the full-length channel were first equilibrated using the default 6-step protocol with the gradual release of restraints on the protein and lipids provided by CHARMM-GUI [62,63] to reach the temperature of 323 K and the pressure of 1 bar. CHARMM36m/TIP3P parameters were used for all atomistic simulations in this work. All simulations throughout this work were carried out using GROMACS, versions 2020 to 2024 [70]. Simulations without applied voltage were then run in the NPT ensemble using the semi-isotropic Parrinello-Rahman barostat [71] and velocity rescale (v-rescale) thermostat [72]. Periodic boundary conditions (xyz) were used. Hydrogen-containing bonds were constrained using the LINCS algorithm [73]; simulation timestep was 2 fs. Lennard-Jones interactions were treated using a cutoff scheme, with forces switched to 0 between 1.0 nm and 1.2 nm. Coulomb interactions were treated using PME [74] with a 1.2 nm cutoff. Five simulation replicas were run for each system without voltage. Simulations with applied voltage of the full-length channel in POPC were started from frames of equilibrium simulation replicas at 0.2-0.5 μs, isolating the functional state of interest (full, partial or no lipid block), and run for up to 1 μs/replica, 3-6 replicas per system. To reach the desired transmembrane voltage ΔV of +300 mV, a constant electric field was applied, calculated as $E = \Delta V / d_z$, where d_z is the size of the box along the *z*-axis [75]. All simulations are listed in *SI Table 1*.

The coarse-grained systems were energy minimized and equilibrated for 500 ns to reach the temperature of 323 K and pressure of 1 bar; for equilibration, v-rescale thermostat [72] and Berendsen barostat [76]. Martini3 parameters with elastic networks between backbone beads were used (500 $kJ*mol^{-1}*nm^{-2}$, with the lower and upper elastic bond cut-off lengths 0.5 and 0.9 nm, respectively, including intersubunit elastic bonds) [47,48]. During the equilibration, restraints with a force constant of 1000 $kJ*mol^{-1}*nm^{-2}$ on the backbone beads of the protein were used; for production runs it was decreased to 1 $kJ*mol^{-1}*nm^{-2}$. Lennard-Jones interactions were treated with a cutoff such that the potential was 0 at 1.1 nm. Electrostatic interactions were treated

using reaction field with a 1.1 nm cut-off, with a dielectric constant of 15 before the cut-off, and infinity beyond it. V-rescale thermostat and Parrinello-Rahman barostat were used for production runs. The simulation timestep in Martini3 simulations was 20 fs. For each combination of bilayer composition and channel state, a 40 μs-long simulation was run (*SI Table 1*).

Data analysis

Pore hydration, lipid entry into the pore, lipid interaction with the protein residues and conformational changes were analyzed using Python3 and mdanalysis [77,78]. For the pore occupancy analysis, the pore was defined as a cylinder with a 10 Å radius in the membrane plane, and 20 Å depth relative to the center of mass of T276 which form the lower S4 K^+ binding site in the SF. To count permeation events and calculate conductances, a custom Fortran script was used [79]. An outward permeation event was defined as a transition of a K^+ from the S2 to the S3 binding site in the SF (S3 to S2 for inward permeations). To calculate conductances, the difference between the number of outward and inward permeation events was used. Errors were calculated as 95% confidence intervals using Student's *t*-distribution. Electrostatic potential maps were obtained using the g_elpot tool [80]. The data were plotted using matplotlib and seaborn [81,82]. Molecular visualizations were rendered using VMD [83].

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Competing interests

The authors declare no competing interests.

Data availability

Molecular dynamics simulations parameter files, initial structures and topologies used in this work are available at https://doi.org/10.5281/zenodo.15519972.

Author contributions

B.L.d.G. and W.K. designed and supervised the project. A.M. performed molecular dynamics simulations and analyzed the data. A.M, B.L.d.G. and W.K. wrote the manuscript.

Figures

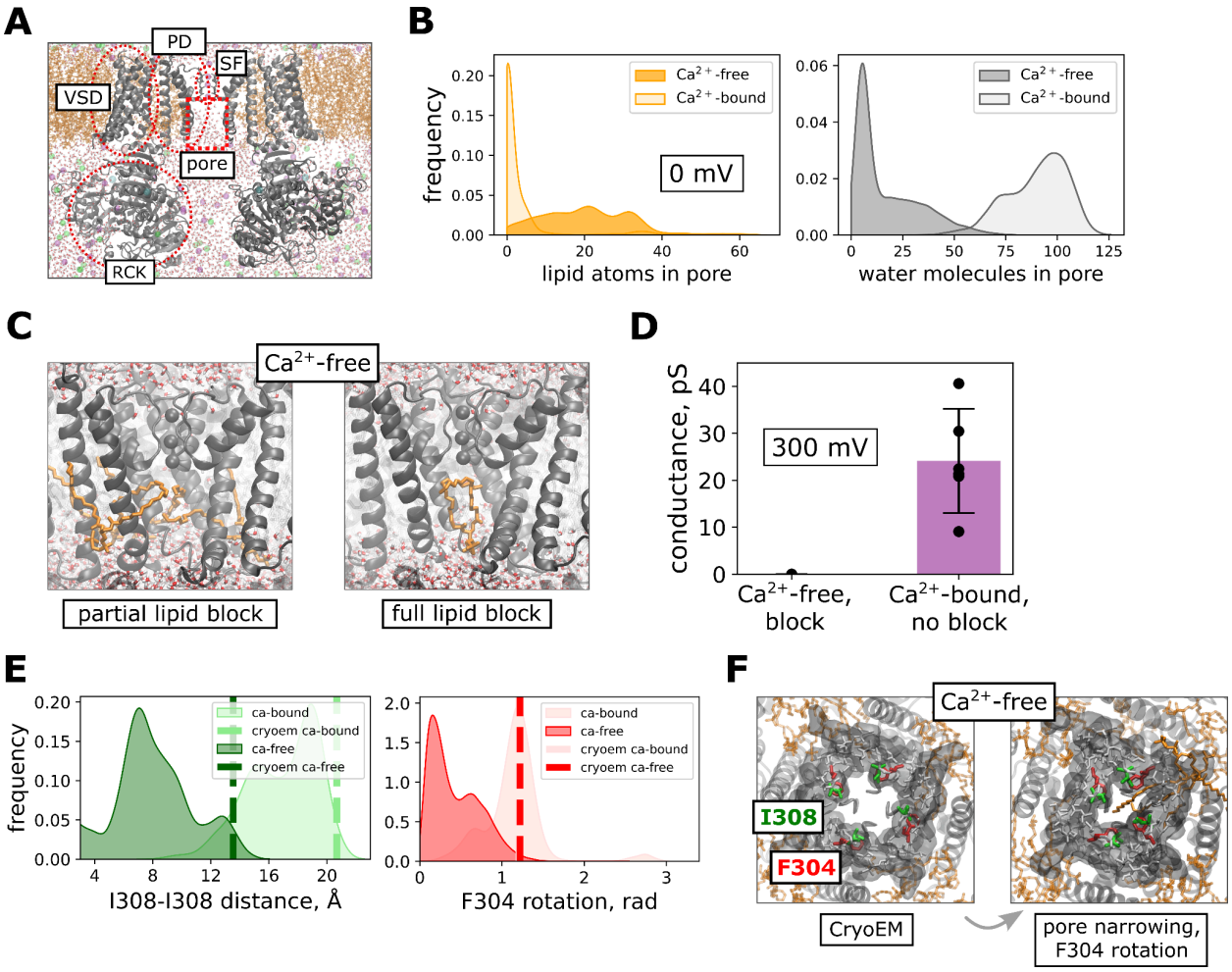

Fig. 1. (*A*) BK channels feature the pore domain (PD), the voltage-sensing domain (VSD), and the cytosolic Ca^{2+}-sensing domain (RCK), with the pore cavity below the selectivity filter highlighted. A snapshot from one of our atomistic simulations of BK, with the channel in the lipid bilayer, water and ions included in simulations is shown. (*B*) Distribution of the number of lipid atoms (non-hydrogen, left), and water molecules in the pore for all-atom simulations of Ca^{2+}-free and Ca^{2+}-bound without applied voltage. (*C*) Snapshots showing the two types of lipid block we observed in our simulations of the Ca^{2+}-free state: partial (only lipid tails are in the pore) and full lipid block (entire lipid molecules occupy the pore). (*D*) Conductances of Ca^{2+}-free and Ca^{2+}-bound under applied voltage (+300 mV, data for Ca^{2+}-free includes simulations started with either partial or full lipid block). (*E*) Distributions of pore widths, defined as distance between I308 residues of opposite subunits, and rotation of F304 into the pore as the angle between the center of mass of F304 from all subunits, and CA and CZ of a F304 from each subunit. Dashed lines represent values in respective CryoEM structures; F304 angle values are similar in both CryoEM structures. (*F*) Snapshots from simulations of Ca^{2+}-free showing the conformational changes relative to the CryoEM structures, including narrowing of the pore at I308 and inward rotation of F304 into the pore, necessary to achieve the non-conductive state of the channel.

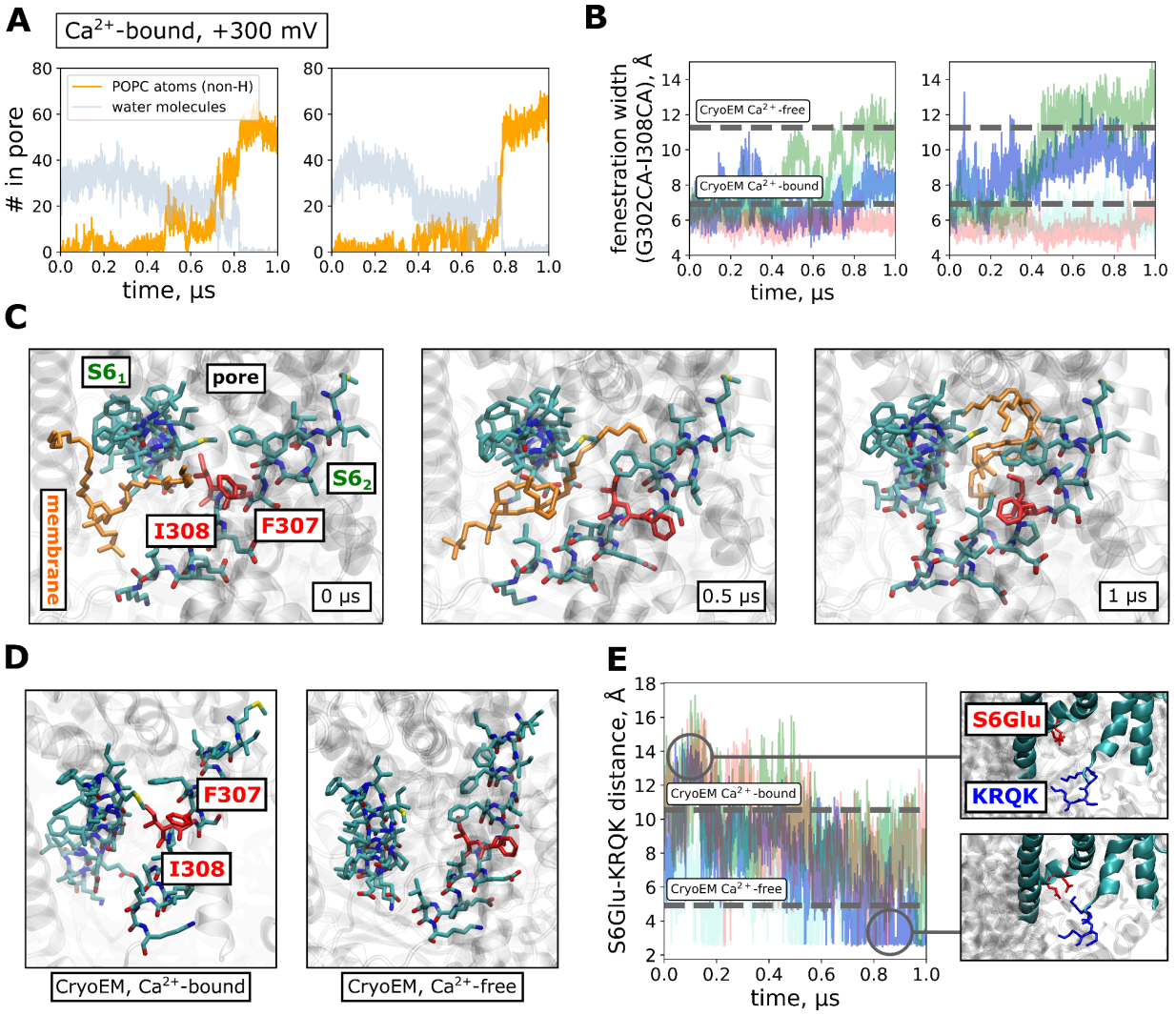

Fig. 2. Simulation replicas of Ca^{2+}-bound at +300 mV where spontaneous lipid block occurred. (*A*) Time traces of lipid entry and pore dehydration. (*B*) Time traces of fenestration width, defined as the distance between CA atoms of G302 residues of opposing subunits. (*C*) Simulation snapshots showing lipid entry into the pore and the accompanying opening of a fenestration between the pore-lining S6 helices, with motions of residues I308 and F307. (*D*) CryoEM structures of Ca^{2+}-free and Ca^{2+}-bound showing conformational differences similar to those observed during the spontaneous lipid block of Ca^{2+}-bound. (*E*) Formation of salt bridges between E310/E310 and ^{317}KRQK^{320} during the spontaneous lipid block, shown for one of the simulation replicas of Ca^{2+}-bound where the spontaneous lipid block occurred, indicative of an open to closed state transition; different colors represent different E310/E313-KRQK pairs (four in total).

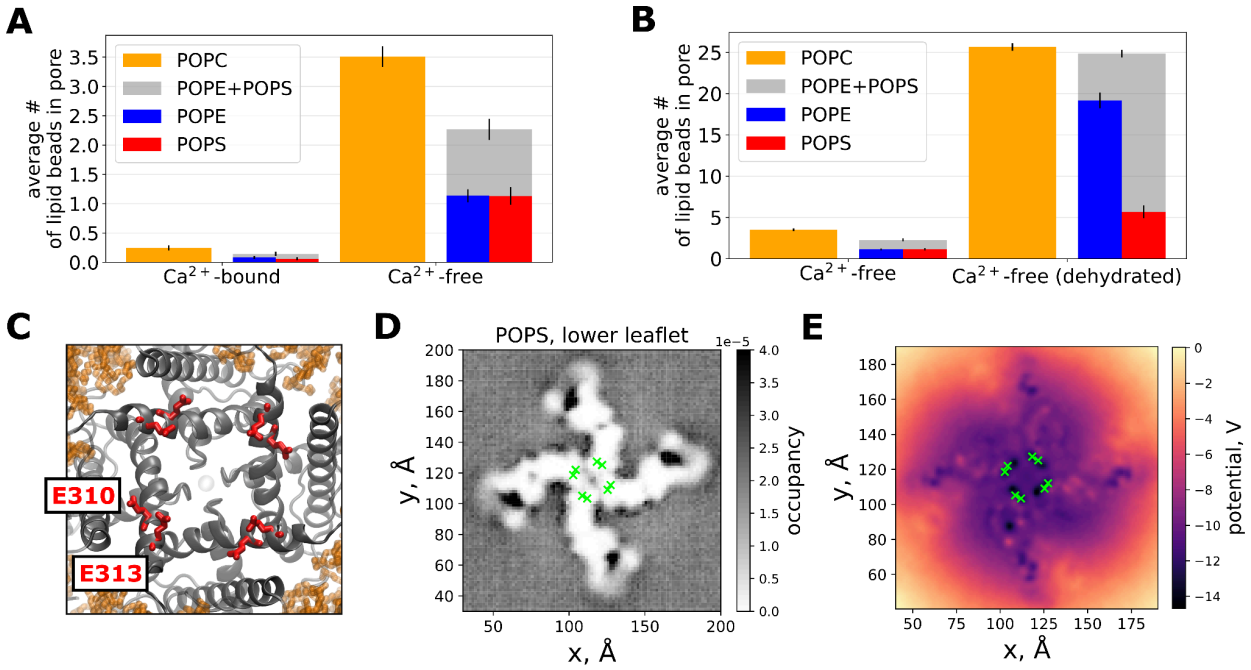

Fig. 3. (*A*) Lipid entry in Ca^{2+}-bound and Ca^{2+}-free states, in POPC and POPE:POPS membranes in Martini3 simulations with a hydrated pore in both systems, and (*B*) with restraints preventing water beads from entering the pore in the Ca^{2+}-free state. (*C*) E310/E313 residues on the pore-lining S6 helices, (*D*) example occupancies of POPS headgroups around the channel (Ca^{2+}-free state is shown) with E310/E313 depicted in green, and (*E*) average electrostatic potential in Martini3 simulations at the level of E310/E313.

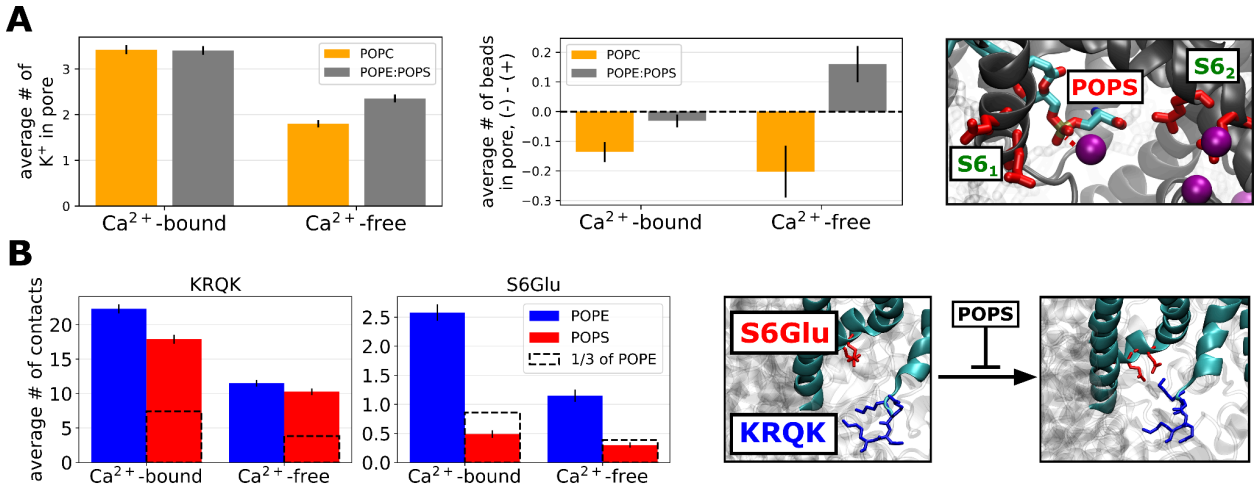

Fig. 4. (*A*) Difference in pore occupancy by negatively charged beads and positively charged beads in Martini3, showing increase in the average number of K^+ beads in pore in the Ca^{2+}-free state with a hydrated pore (*left*), accompanied by enrichment of the pore by negatively charged beads for that system in a POPE:POPS membrane (*middle*), explained by electrostatic attraction of K^+ into the pore by negatively charged headgroups of POPS transiently occupying the fenestrations between S6 helices and the pore; E310/E313 residues are shown for reference (*right*). (*B*) Average number of contacts of each lipid type with either the positively charged ^{317}KRQK^{320} cluster, or the negatively charged E310/E313 ('S6Glu', *left*), indicating that POPS binds to KRQK and is repelled from the E310/E313, thus possibly perturbing the formation of a salt bridge between ^{317}KRQK^{320} and E310/E313 characteristic of the closed state.

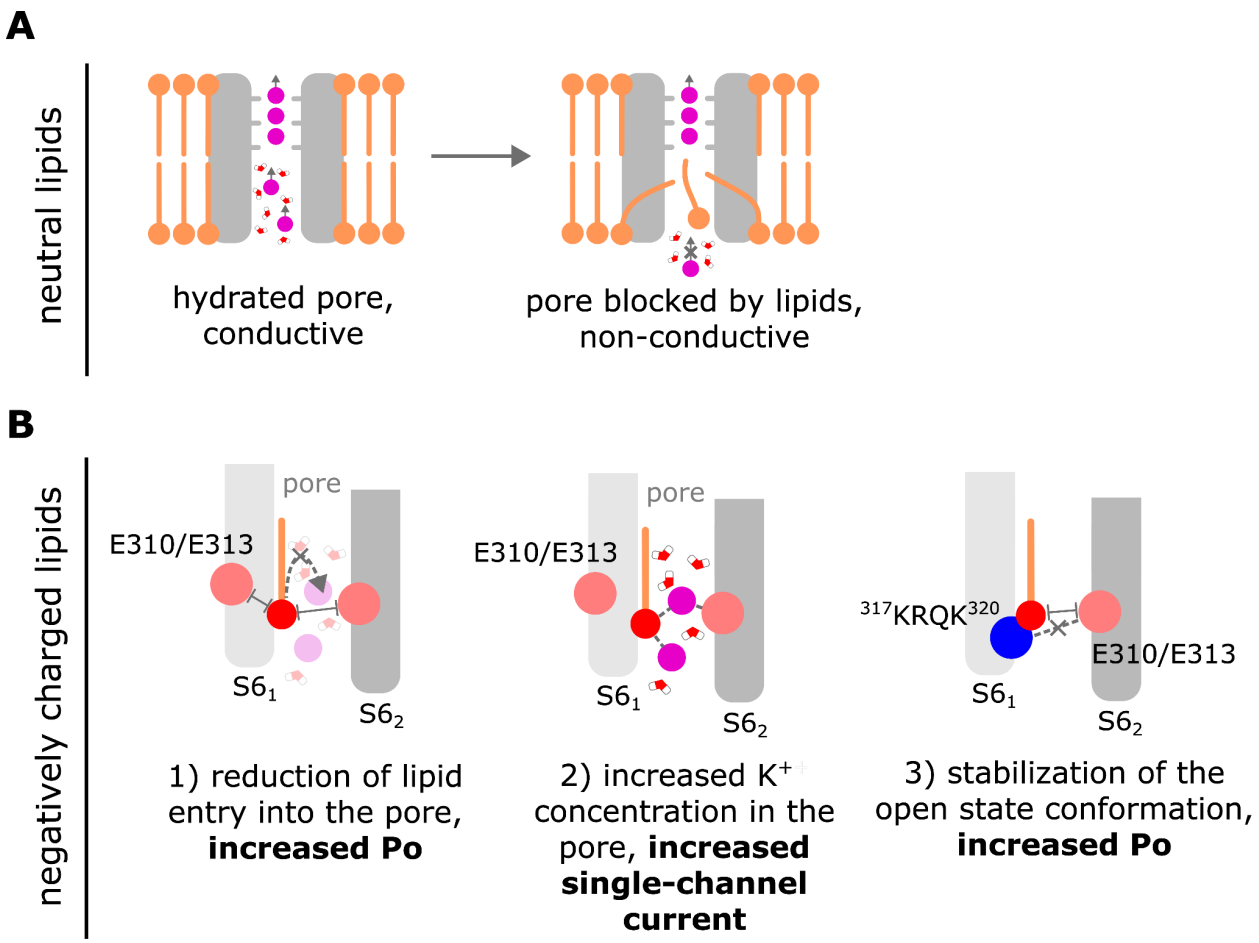

Fig. 5. A scheme illustrating the main hypotheses of this work. (*A*) While the pore of the channel in the open, conductive state of the channel remains hydrated, the closed state is characterized by either lipid tails, or entire lipid molecules entering the pore through membrane-facing fenestrations, and blocking ion permeation. (*B*) Presence of negatively charged lipids in the membrane leads to 1) reduced lipid entry into the pore due to the electrostatic repulsion from the fenestration-lining glutamate residues in S6 helices, 2) increased K^+ concentration in the pore and thus increased probability of a permeation event due to the electrostatic attraction between K^+ and negatively charged headgroups spontaneously occupying the fenestrations and the pore, and 3) disruption of the ^{317}KRQK^{320}-E310/E313 salt bridge formation, proposed to be a crucial step in channel closure [24], via interaction of the negatively charged lipids with ^{317}KRQK^{320} and repulsion from E310/E313. These effects explain the experimentally observed increase in the open probability and the single channel-current in negatively charged lipid-containing membranes [42,43].

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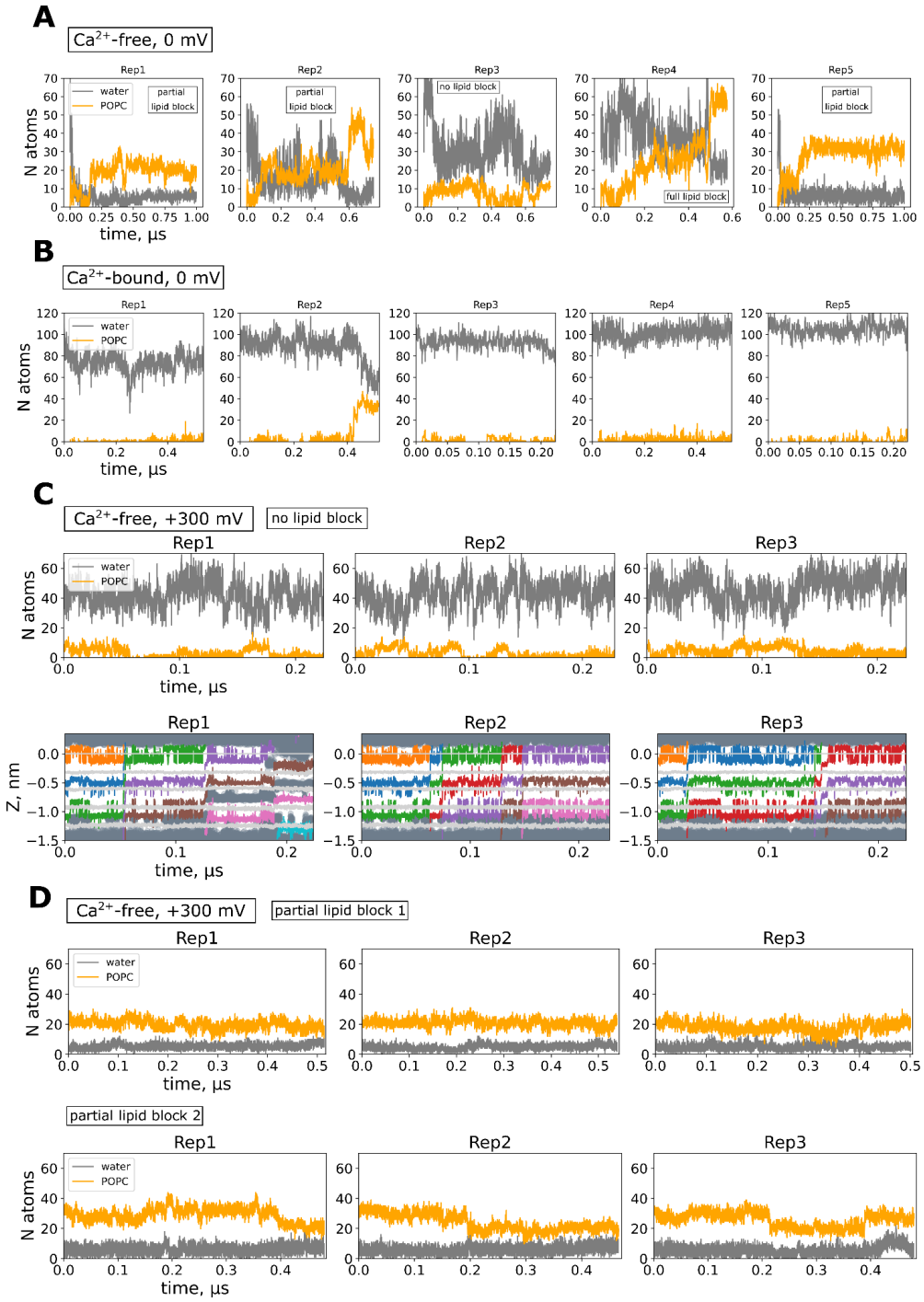
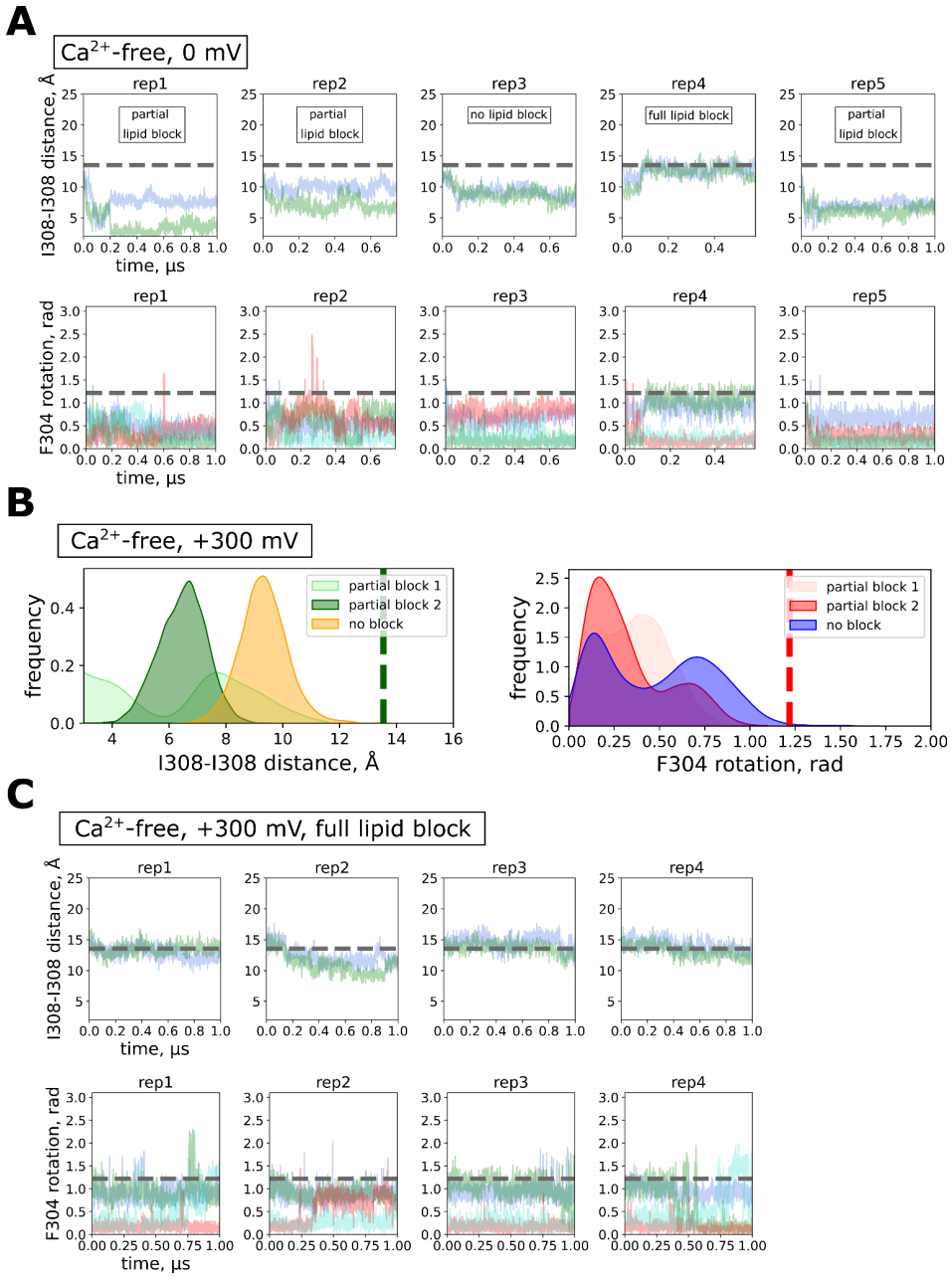
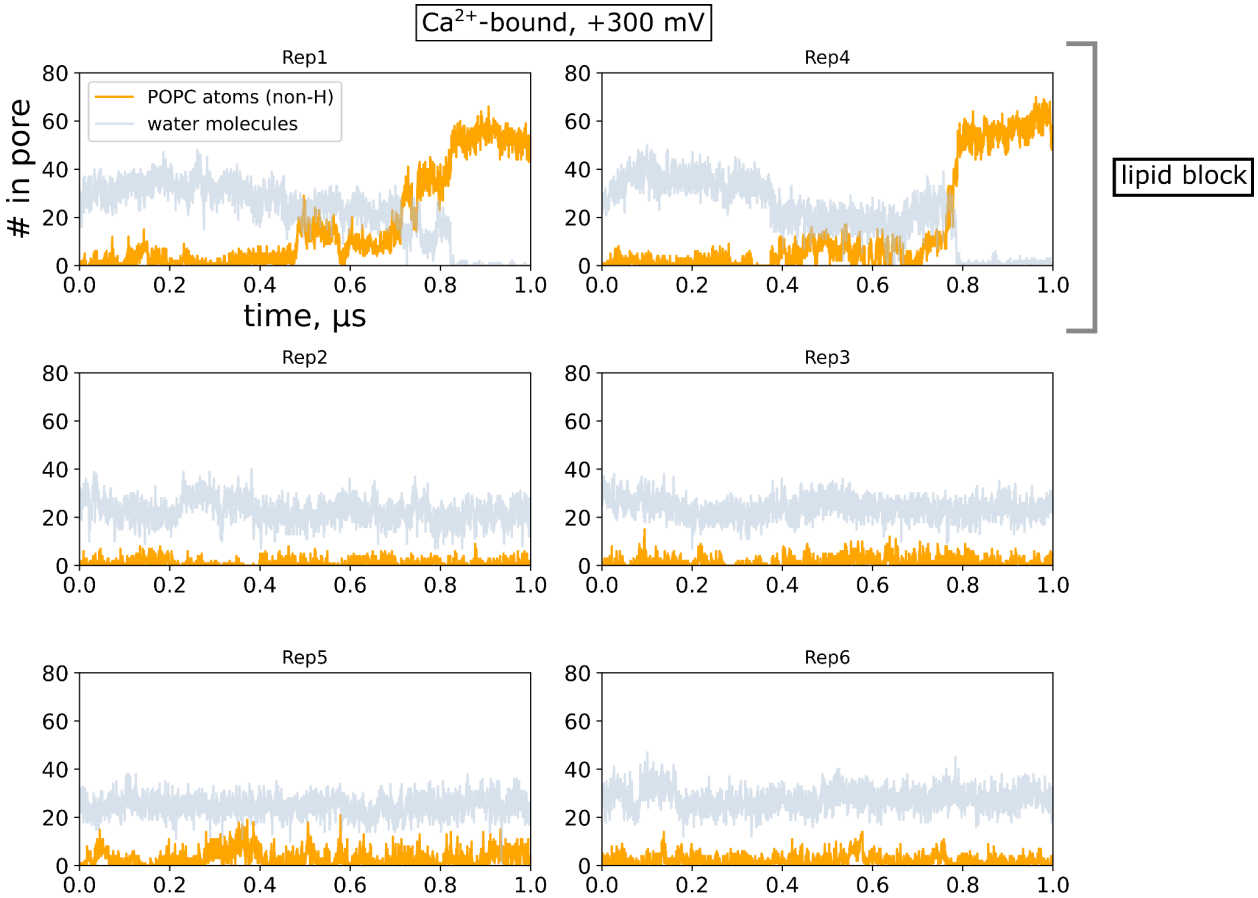


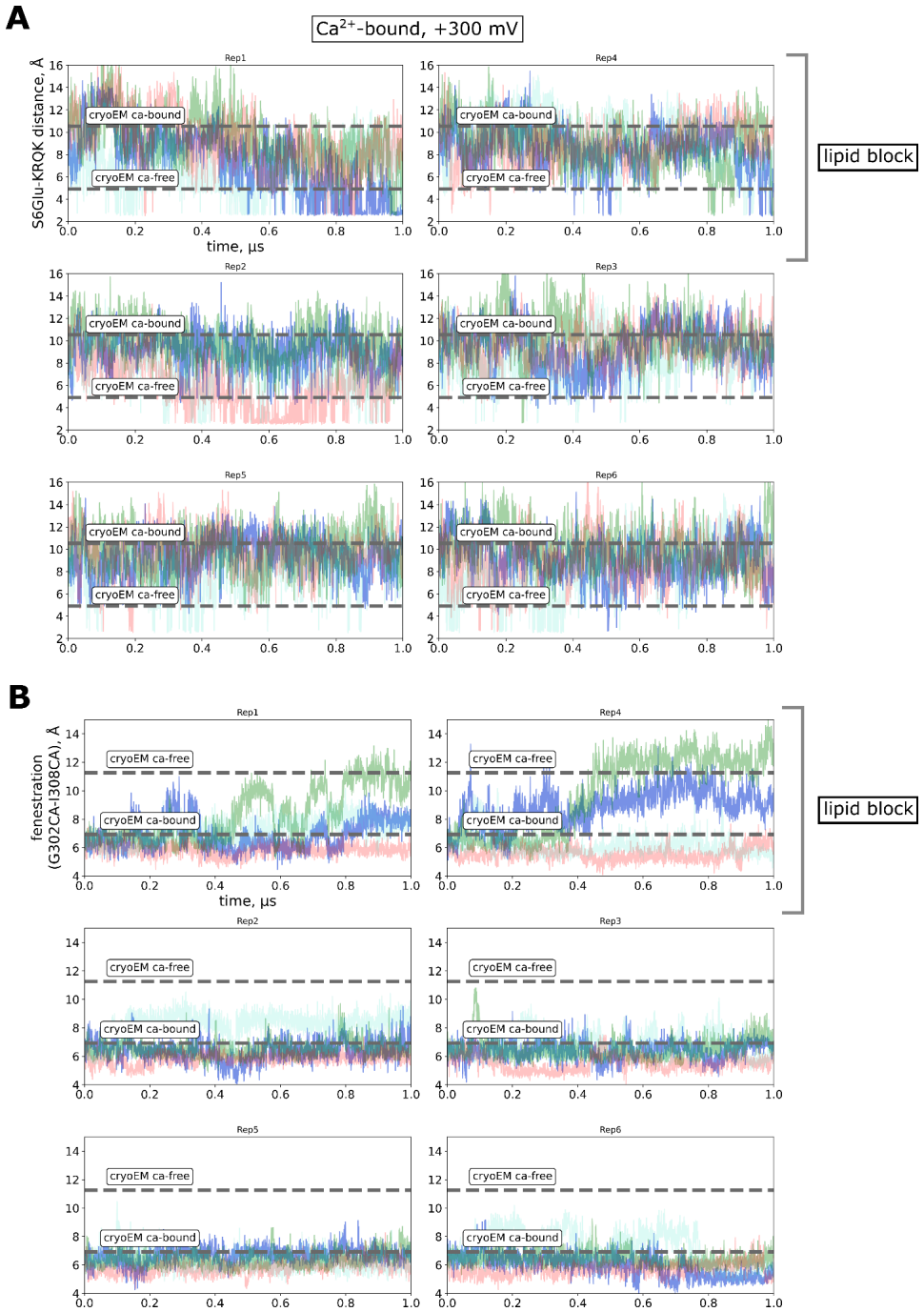
SI Fig. 1. Time traces of the number of lipid (non-hydrogen) and water oxygen atoms in the pore in atomistic simulations without applied voltage of (*A*) Ca^{2+}-free, (*B*) Ca^{2+}-bound, and (*C, upper*) replicas of Ca^{2+}-free under applied voltage, where the pore remained hydrated, not blocked by lipids, and conductive, as indicated by time traces of K^+ (colored lines), water (dark gray scatter) in the SF (carbonyls and hydroxyls of SF-forming K^+ binding sites are shown as light gray lines) (*C, lower*). (*D*) Time traces of lipid and water oxygen atoms in the pore, Ca^{2+}-free +300 mV simulation sets with partial lipid block.



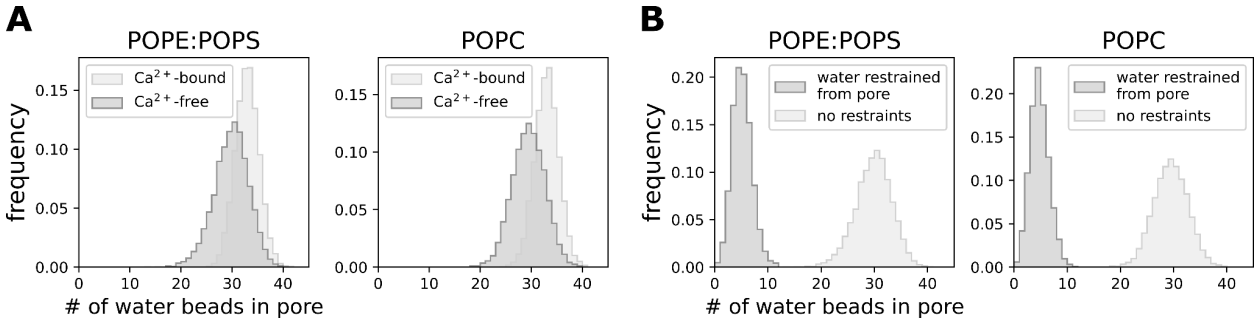
SI Fig. 2. (*A*) Time traces of pore opening (distance between CA atoms of I308 residues of opposing subunits) and inward rotation of F304 (defined as the angle between center of mass of all F304 and CA, CZ atoms of F304 for each subunit) in the Ca^{2+}-free state, in atomistic simulations without applied voltage. (*B*) Pore closure and F304 rotation in +300 mV simulations of Ca^{2+}-free in the two simulation sets with a partial lipid block, and the simulation set where the lipid block did not occur. (*C*) Time traces of pore opening and inward rotation of F304 in the Ca^{2+}-free state, in atomistic simulations at +300 mV where full lipid block occurred. Dashed lines represent the CryoEM structure values for the Ca^{2+}-free state.



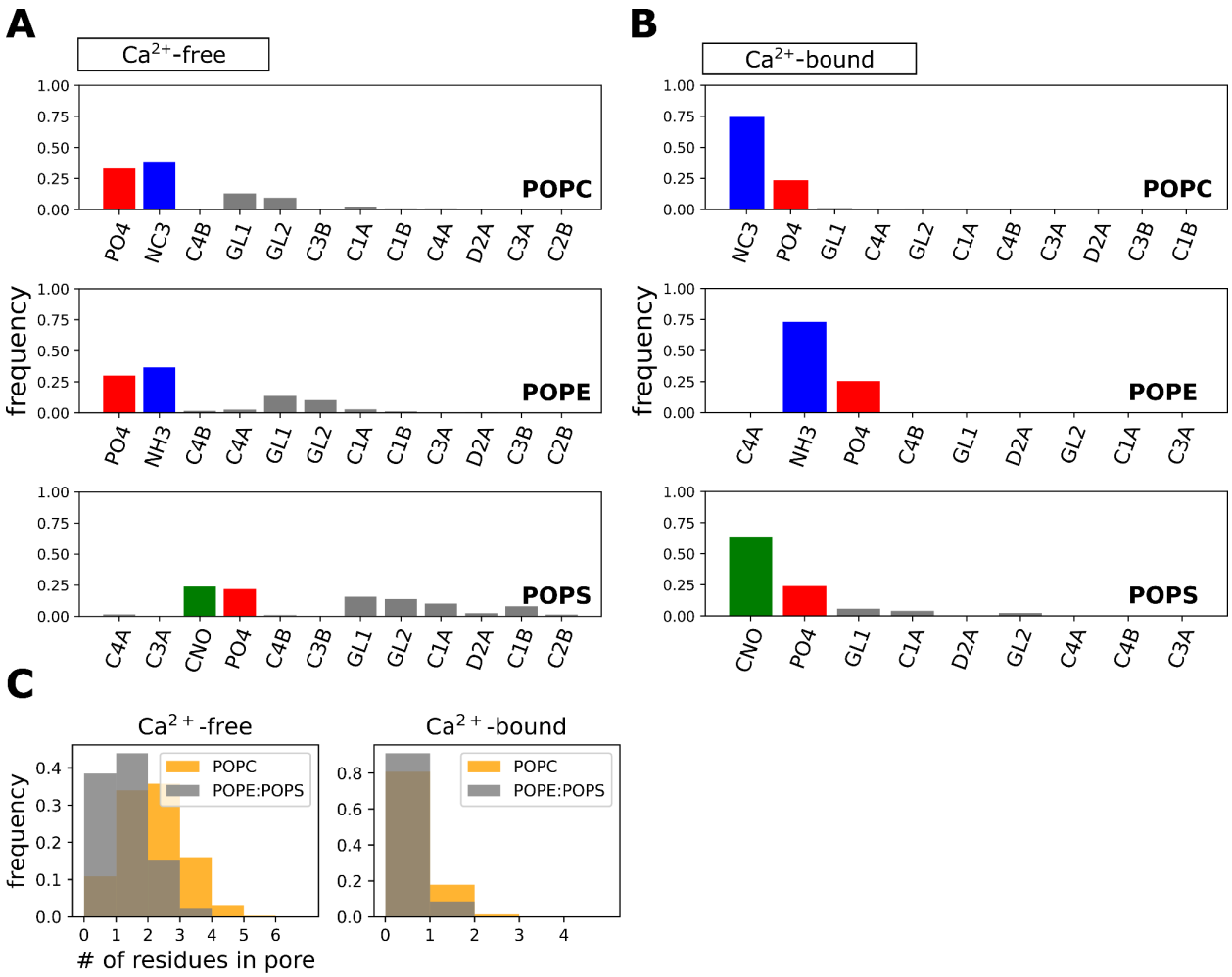
SI Fig. 3. Traces of the number of water molecules and POPC atoms (non-hydrogen) in the pore, atomistic simulations of the Ca^{2+}-bound state at +300 mV. The replicas where the lipid block of the pore occurred are highlighted.



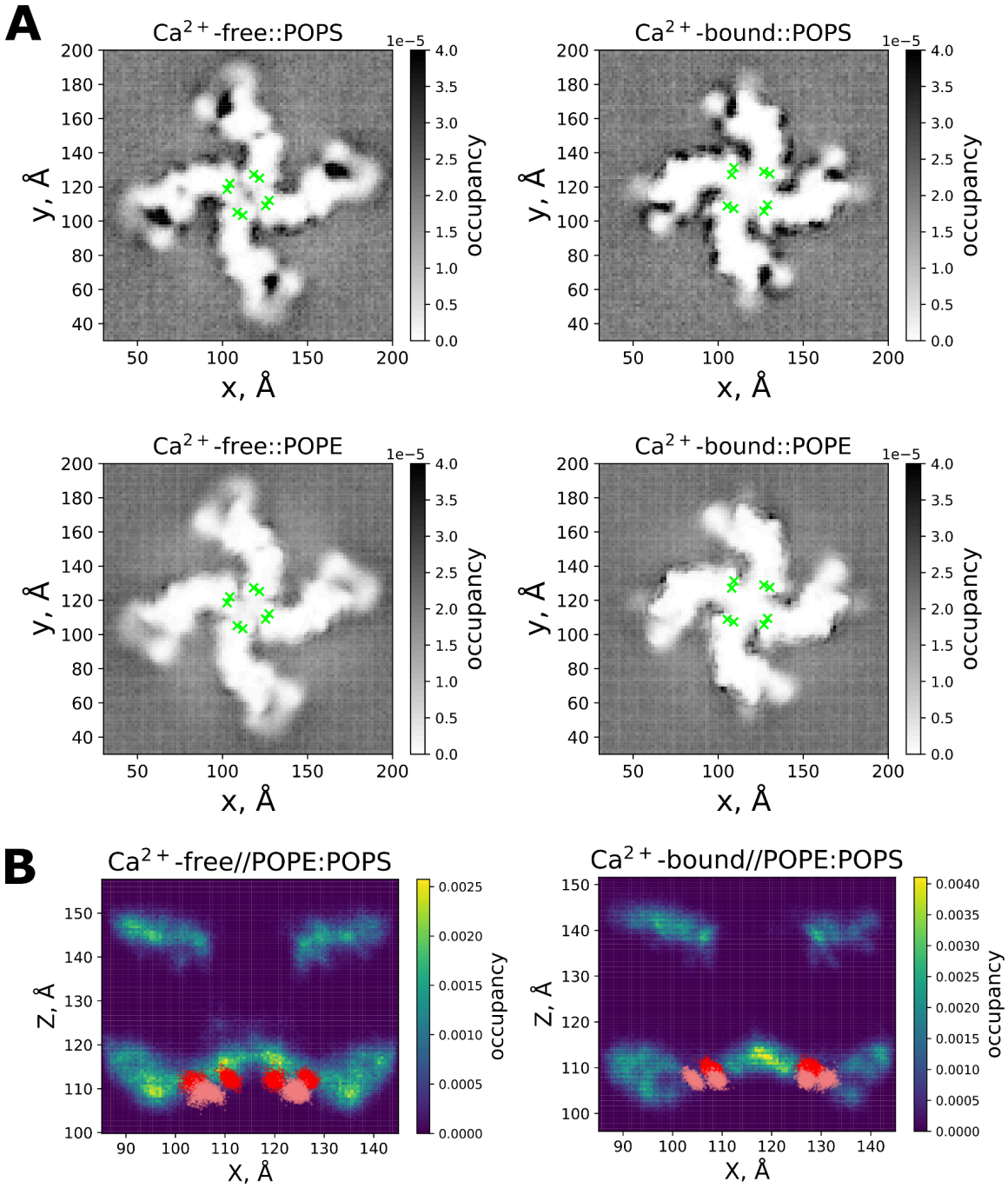
SI Fig. 4. Atomistic simulations of the Ca^{2+}-bound state at +300 mV. Time traces of (*A*) distances between E310/E313 ('S6Glu') and ^{317}KRQK^{320} of neighboring subunits, illustrating the salt bridge formation between the two charged clusters at the end of S6 helices, and (*B*) distances between CA atoms of G302 and I308 of neighboring subunits, illustrating the widening of fenestrations between S6 helices; different colors represent distances between different residue pairs (4 in total). Simulation replicas where spontaneous lipid block occurred ('Rep1' and 'Rep4') are indicated.



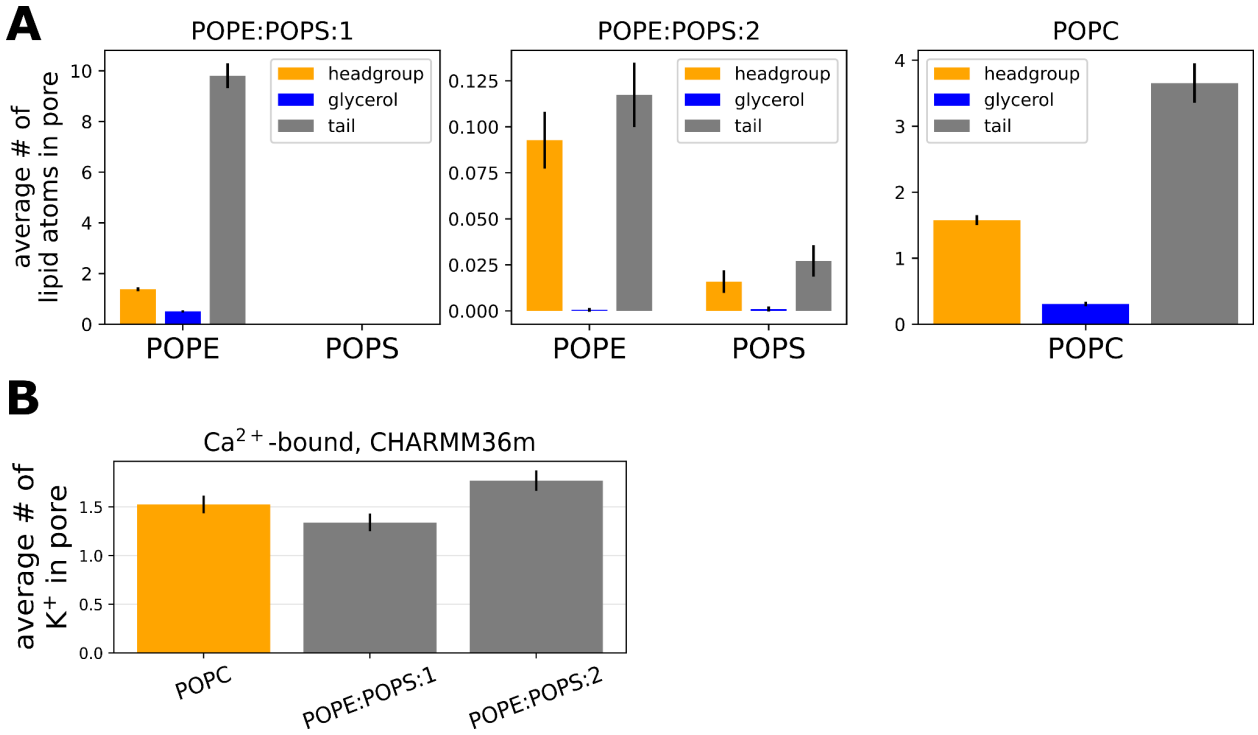
SI Fig. 5. Distribution of water beads in the pore (*A*) without, and (*B*) with restraints on water preventing it from entering the pore in Martini3 simulations.



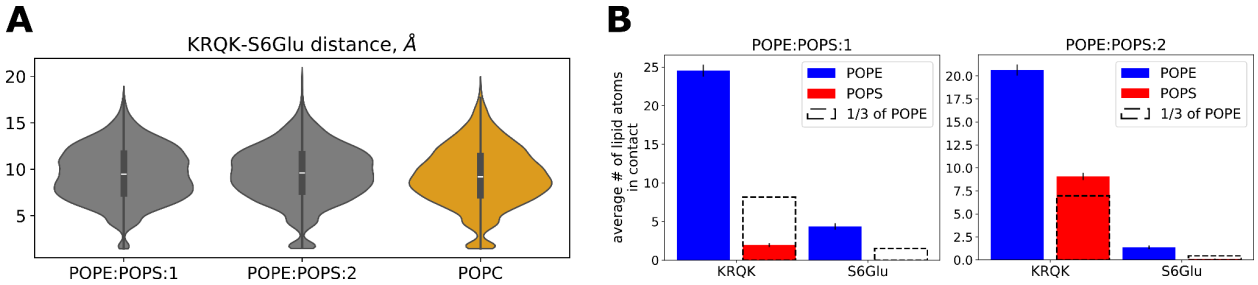
SI Fig. 6. Relative frequencies of lipid beads occupying the pore in (*A*) Ca^{2+}-free (hydrated pore) and (*B*) Ca^{2+}-bound states in Martini3 simulations; colored bars represent the headgroup beads (excluding the glycerol beads), blue corresponds to positively charged beads, red to negative and green to non-charged. (*C*) Distribution of the number of individual lipid residues simultaneously occupying the pore in Martini3 simulations.



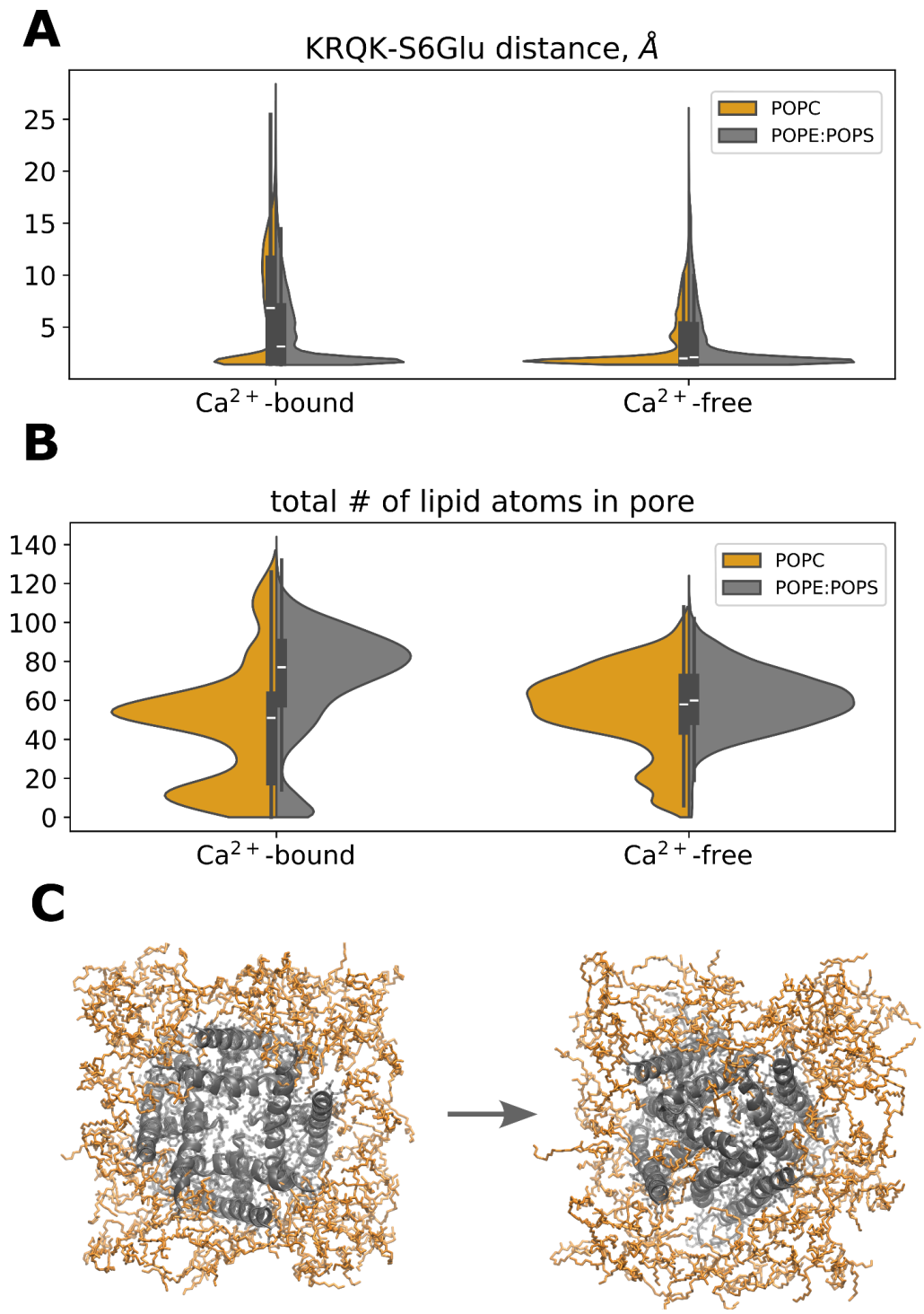
SI Fig. 7. (*A*) Lipid distribution around the channel in Martini3 simulations; E310/E313 residues in S6 helices are shown as reference (green markers). (*B*) Occupancies of PO4 beads (corresponding to phosphate groups; only beads within 3 nm of E310/E313 in the membrane plane were counted), shown in a plane perpendicular to the membrane. Note that the densities are integrated along the Y axis, and the density in the pore region of the channel (X = 110..120 Å) does not necessarily represent PO4 beads occupying the pore.



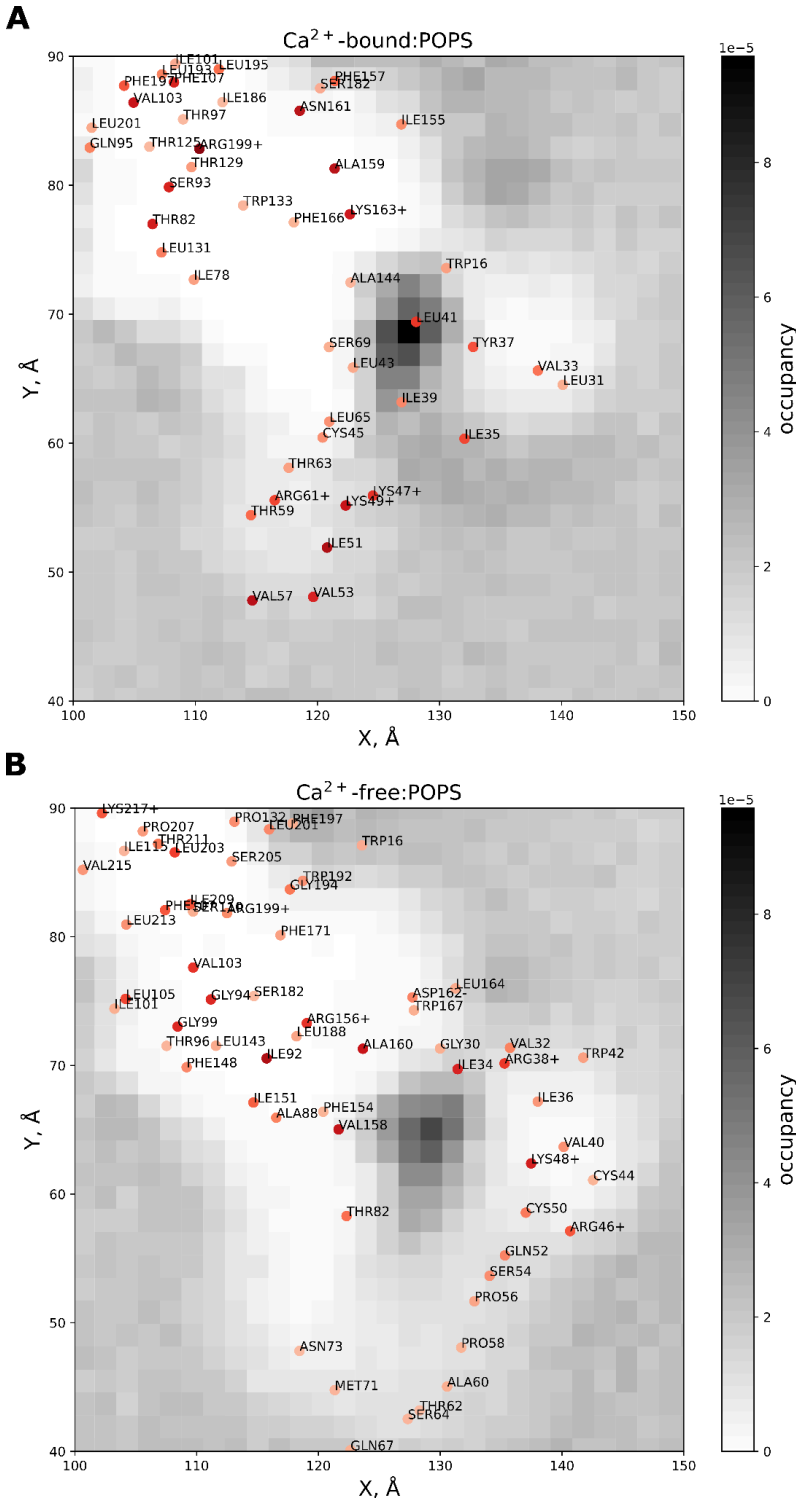
SI Fig. 8. (*A*) Occupancy of the pore by atoms belonging to either the headgroup (excluding glycerol), glycerol or the tail of a given lipid, and (*B*) average number of K^+ ions in the pore in atomistic simulations (Ca^{2+}-bound state, 0 mV).



SI Fig. 9. Atomistic simulations of the full-length channel (Ca^{2+}-bound state) in POPC compared to POPE:POPS membranes. (*A*) Distributions of distances between ^{317}KRQK^{320} and E310/E313 ('S6Glu') residues of neighboring subunits. (*B*) Average number of contacts between headgroups belonging to a given lipid type, and ^{317}KRQK^{320} or S6Glu in atomistic simulations of Ca^{2+}-bound in POPE:POPS. POPE:POPS:1 and POPE:POPS:2 indicate simulations started from the two different initial configurations of the membrane.



SI Fig. 10. Atomistic simulations of the isolated pore domain. Distributions of (*A*) distances between ^{317}KRQK^{320} and E310/E313, and (*B*) number of lipid atoms (non-hydrogen) in the pore. (*C*) Simulation snapshots showing low structural stability of the isolated pore domain and high propensity to be blocked by lipids regardless of the state of the channel (here, Ca^{2+}-bound is shown).



SI Fig. 11. Occupancies of PO4 beads of POPS near the binding site located in the voltage sensor in (*A*) Ca^{2+}-bound and (*B*) Ca^{2+}-free states, with protein residues interacting with POPS highlighted; darker color indicates more frequent contacts.

SI Table 1. *List of simulated systems. Unless noted, listed simulations correspond to those of a full-length channel.*

Channel structure	Membrane composition	Force field	Voltage	Duration
Ca^{2+}-bound	POPC	CHARMM36m	0 mV	5x0.2..0.5 μs
			+300 mV	6x1 μs (started from of one of the 0 mV replicas)
Ca^{2+}-free	POPC	CHARMM36m	0 mV	5x0.2..0.5 μs
			+300 mV	3x0.2 μs (started from a 0 mV replica with no lipid block), 3x0.6..0.7 μs, each started from two 0 mV replicas with partial lipid block (6 replicas in total), 4x1 μs (started from a 0 mV replica with full lipid block)
Ca^{2+}-bound	POPE:POPS, 3:1	Martini3	0 mV	40 μs
Ca^{2+}-free	POPE:POPS, 3:1	Martini3	0 mV	40 μs (no restraints on water), ~20 μs (with restraints on water molecules preventing them from entering the pore)
Ca^{2+}-bound	POPC	Martini3	0 mV	40 μs
Ca^{2+}-free	POPC	Martini3	0 mV	40 μs (no restraints on water), ~20 μs (with restraints on water

				molecules preventing them from entering the pore)
Ca^{2+}-bound	POPE:POPS, 3:1	CHARMM36m	0 mV	4x0.2..0.5 μs (initial membrane configuration 1), 4x0.3..0.4 μs (initial membrane configuration 2)
Ca^{2+}-free	POPE:POPS, 3:1	CHARMM36m	0 mV	4x0.2..0.5 μs (initial membrane configuration 1), 4x0.2..0.3 μs (initial membrane configuration 2)
Ca^{2+}-bound, isolated pore domain	POPC	CHARMM36m	0 mV	10x1 μs
Ca^{2+}-free, isolated pore domain	POPC	CHARMM36m	0 mV	10x1 μs
Ca^{2+}-bound, isolated pore domain	POPE:POPS, 3:1	CHARMM36m	0 mV	4 different initial membrane configurations, 5x1 μs each
Ca^{2+}-free, isolated pore domain	POPE:POPS, 3:1	CHARMM36m	0 mV	4 different initial membrane configurations, 5x1 μs each