Mechanisms of surface and volume light scattering from Caenorhabditis elegans revealed by angle-resolved measurements

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Abstract

The roundworm *Caenorhabditis elegans* has become a widely studied model system in biology in part due to its small size and optical accessibility. Optical imaging of *C. elegans* primarily relies on scattering of light to generate contrast between the animal and its surroundings, and/or contrast to resolve internal structures within the body. However, the mechanisms of light scattering by *C. elegans*, and how they influence image contrast, are not well understood, and optimization of contrast is typically performed in a trial-and-error fashion. To address this gap, we performed measurements of light scattering from *C. elegans* on laboratory substrates under varied illumination angles, wavelength, substrate material, bacterial seeding, and index matching conditions.   By combining these results with those of computational simulations, we find that light scattering arises via two components that exhibit distinct angular and spatial distributions: volume scattering within the body and surface scattering at the boundaries between the worm and its surroundings. We elucidated how volume and surface scattering interact with each other and collectively contribute to the overall scattering. Informed by these findings, we determined the illumination angle range, wavelength

23   range, and substrate material that optimize imaging contrast. Our findings provide practical

24   strategies for optimizing image quality across a range of imaging systems.

Introduction

The microscopic nematode *C. elegans* is one of the most widely studied model

organisms in biology due to its short life cycle, genetic accessibility, and compact anatomy [1,

2]. Its optical accessibility, fully mapped connectome [3], and rich repertoire of behaviors [4]

have made it a valuable system for *in vivo* imaging of neural activity [5-7], behavioral dynamics

[8-10], and genetic perturbations [10, 11]. Recent advances in automation [11-16] have enabled

high-throughput genetic and pharmacological screening using *C. elegans*.

One of the strengths of *C. elegans* is its amenability to optical imaging. Many assays of

behavior, morphology, aging, and development rely on obtaining images of worms with high

contrast relative to their surroundings [17-21]. Since like most biological tissues, *C. elegans*

exhibits relatively low optical absorption over optical wavelengths, most of the optical contrast is

provided by scattering, i.e. changes in direction of incident light rays [22].

In any imaging system, lenses or other optical elements are configured to focus light

collected from a certain range of angles. Two imaging modes can be differentiated by whether

the illumination angles are within the range of collection by the optical system [23, 24]. In a

bright field imaging configuration, the illumination rays are within the collection angle of the

optics. Optical scattering by the object diverts some of the light away from the collection angles,

causing the object to appear darker than its surroundings. By contrast, in a dark field imaging

configuration, the illumination light is outside the collection angle of the optical system;

scattering by the object causes some of this light to be collected, and the object appears bright

against a dark background.

The amount and angular distribution of light scattering are highly sensitive to illumination

angle, wavelength, and the optical properties of the object and background medium. In practice,

many imaging systems rely on empirically chosen parameters, without a systematic

understanding of how these variables affect image contrast.

Light scattering in biological tissues has previously been studied primarily in the context

of biomedical imaging [22, 25]. Prior theoretical descriptions of light scattering in biological

tissues include those applying Mie theory [26, 27] which describes scattering from spheres as a

function of diameter and refractive index, and the Henyey–Greenstein phase function [28, 29],

which describes the degree of anisotropy occurring during scattering. Angle-dependent

scattering profiles have been characterized for tissue sections [30-32], tissue-mimicking

phantoms [33], and single cells [32, 34]. However, to our knowledge light scattering from intact,

living, multicellular organisms such as *C. elegans* has not been examined in detail.

In this study, we developed and applied methods for angle-resolved measurements of

light scattering from *C. elegans*. Supported by experimental observations and computational

modeling, we find that light scattering from *C. elegans* arises from a combination of volume

scattering within the animal body and surface scattering due to the refractive index mismatch at

the interface between the worm and its surroundings. We probe the interaction between the

volume and surface scattering and reveal how together they shape the overall scattering. Based

on these findings, we determined optimal values for several experimental parameters including

illumination angle, illumination wavelength, and substrate material that maximize optical

contrast.

Our work provides a quantitative framework for understanding light scattering in *C.*

elegans and offers practical guidance for optimizing image quality. Moreover, the physical

principles outlined here will inform optimization strategies for imaging other organisms.

Results

Angle-resolved scattering functions

To quantitatively describe light scattering from *C. elegans*, we illuminated worms on agar

plates at specific angles using a ring of light-emitting diodes (LEDs) and captured images using

a camera and lens (Fig. 1A, Material and Methods). We imaged worms against a background

of either agar seeded with food bacteria or bare (unseeded) agar (Fig. 1B). Under illumination at

variable angles, the animals exhibit distinct spatially inhomogeneous scattering patterns (Fig.

1B). The contrast of the animal against its background also changes with θ (Fig. 1B).

We aimed to quantitatively describe the observed angle-dependent optical scattering and

image contrast. Using the images recorded with different illumination angles, we calculated the

scattering function p(θ), which has units of sr^{-1} (inverse steradian) and describes the probability

per solid angle of a photon scattering by a polar angle θ relative to its original trajectory (Fig. 1A)

[35, 36]. The value of the scattering function at θ = 0º represents the probability of a ray to be

transmitted without scattering.

We measured scattering functions under red light illumination (peak wavelength 629 nm) for

C. elegans, agar, and bacteria-coated (seeded) agar (Fig. 1C). For both *C. elegans* and agar

backgrounds the scattering function is highest at θ = 0º, decreases from 0º to 60º, is roughly

constant between 60º and 140º and increases modestly between 140º and 180º.

To determine the effect of scattering angle on image contrast, we defined a contrast index

equal to the difference in irradiance between worm and background, both measured at the

camera, divided by the sum of irradiance of the worm and background (see Material and

Methods). We calculated the contrast index as a function of θ (Fig. 1D). For $\theta = 0^\circ$, the contrast

indices are negative (approximately -0.85) for worms on both seeded and unseeded agar

plates. This is because at $\theta = 0^\circ$, the system is in the bright field regime, in which the worm

appears darker than its surroundings.

For $\theta > 0^\circ$, in the dark field imaging regime, the scattering function for the background

decays with angle at a faster rate than that of the worm (Fig. 1C), causing the contrast indices to

increase sharply with θ and reach maximum of about 0.65 for θ around 15° (Fig. 1D). For $\theta >$

15°, the contrast indices decrease with θ and reach minimum between 60° and 90°. For θ close

to 180°, the image contrast slightly rises for worms on unseeded plates, and we observe a

negative contrast index for worms on seeded plates (Fig. 1D).

Fig. 1. Angle-resolved, wavelength-dependent scattering functions and image contrast for *C. elegans* on agar substrate. (A) Experimental setup. (B) Representative images of worms and background at different illumination angles (FOV: 1.11 × 0.56 mm). (C) Scattering functions for *C. elegans* on seeded and unseeded agar substrate under red light illumination.

108   (D) Image contrast as a function of illumination angle θ for *C. elegans* on seeded and unseeded

109   agar plate under red light illumination. (E) Scattering functions for *C. elegans* on seeded agar

110   plates under red, green, and blue light illumination. (F) Image contrast index for *C. elegans* on

111   seeded agar as a function of illumination angle and wavelength.

Wavelength dependence of scattering functions

Optical scattering from biological tissue typically has a strong wavelength dependence

[37, 38]. We asked how scattering from *C. elegans* varies with wavelength. We measured the

angle-resolved scattering functions for *C. elegans* on seeded agar plates under red (peak

wavelength 629 nm), green (516 nm), and blue (447 nm) LED illumination (Fig. 1E). The

scattering functions for green and blue light show similar overall shapes as that of red light. We

found that light scattering with longer-wavelength illumination is higher for worms but lower for

seeded agar, which results in a higher image contrast for red illumination than for green or blue

(Fig. 1F). The maximum contrast index for red illumination (approximately 0.7) is about twice

than of blue (approximately 0.35). For θ between 90º to 180º, scattering functions with longer

wavelengths are lower for both worms and seeded agar and the contrast indices for red are

higher than green and blue. For all three wavelengths, we observe negative contrast when θ is

near 180º. These results suggest that red illumination is better than green or blue in terms of

generating contrast in images.

Our results show that the scattering functions near θ = 0º are higher for longer

wavelength for both *C. elegans* and seeded agar. It indicates that the proportion of the light that

is forward transmitted without scattering increases with wavelength, consistent with results from

prior studies that the optical reduced scattering coefficient (μ_s') generally decreases with

wavelength [37].

Empirical analysis of volume and surface scattering

We sought to understand the mechanisms contributing to the observed angle-dependent

scattering functions from *C. elegans*. We reasoned that light scattering from *C. elegans* could

arise from two sources: (1) volume scattering within the animal body due to refractive index

variations and (2) surface scattering from refractive index mismatch at the boundaries between

the worm and the air and/or the substrate (Fig. 2A).

To help distinguish between these two types of scattering, we first isolated the volume

scattering by minimizing the surface scattering via approximate refractive index matching. Most

biological tissues have an index of refraction slightly higher than that of water [39]. The *C.*

elegans cuticle has been measured to have a refractive index about n = 1.35-1.37, only slightly

higher than that of water [40]. To achieve approximate index matching, we placed a gelatin pad

(n ≈ 1.36) on the agar plate in order to eliminate the worm-to-air interface and surround the

animals with media with similar refractive index to that of the worm (Fig. 2B). After index

matching, the observed scattering function should be almost entirely due to volume scattering.

Fig. 2C depicts the scattering functions for *C. elegans* and bacteria-seeded agar in the

index-matching experiment. For θ = 0º, scattering functions for the index-matched case are

higher than for the non-index-matched case for both worm and background, implying that index

matching increases the forward transmission. For θ > 0º, index matching caused the scattering

function to decrease. For the background, scattering function for the index-matched case is

lower than for non-index matching for θ near 0º and 180º, and slightly higher elsewhere.

We calculated the contrast indices due to volume scattering from the worm and

background (Fig. 2D). For θ = 0º, in the bright field imaging regime, the image contrast for index

matching is similar to non-index matching. For θ > 0º, in the dark field imaging regime, the

image contrast reaches maximum around ~10º for index matching and decreases with θ. In

general, index matching reduces image contrast (Fig. 2D), implying that the surface scattering

helps improving the image contrast.

As in the non-index matched case, we observed that volume scattering in *C. elegans* is

dependent on illumination wavelength and red light generates higher image contrast than green

or blue (Fig. S1). Our results show that optimal image contrast can be obtained by applying red

illumination with an angle of 10º-15º relative to the optical axis. These findings inform useful

guidelines for optimizing image quality for *C. elegans* in liquid-based assays [19, 41, 42] or other

index-matched environments.

166   Mechanisms of volume and surface scattering

The experimental observations of optical scattering in the non-index matched and index

matched scenarios pose several questions. What mechanisms contribute to the angular

dependence of the optical scattering from *C. elegans* and how does it shape the nonlinear

change of the image contrast (Fig. 2C and D)? What is the interplay between the volume and

the surface scattering and how do they contribute to the overall scattering? Why does the

optical scattering exhibit spatially inhomogeneous patterns across the worm's body, sometimes

appearing to highlight the outline of the animal (Fig. 1B)?

To explore the mechanisms underlying the angle-dependent scattering, we created a

simple computational model of the worm and its surroundings (Fig. 2A). We modeled the adult

C. elegans as a cylinder (diameter = 50 µm, length = 1000 µm, n = 1.33) on the surface of a

rectangular cuboid (n = 1.33) representing the substrate and a concave shaped film (n = 1.33)

formed at the either side (Fig. 2A and Fig. S2) representing the capillary water film that forms

near the points of contact between the worm and the agar substrate [43]. We set the volume

scattering within the cylinder to follow the Henyey-Greenstein phase function [28], widely used

to model light scattering within biological tissues [29], with anisotropy parameter g equal to 0.83

and scattering mean free path 12 µm (optimization of free parameters is described in Materials

and Methods). To simulate the index matched case, we eliminated the worm-to-air interface by

adding a medium (n = 1.33) representing the index matching pad (Fig. 2B).

Using these models, we simulated scattering functions for both the index matched and

non-index matched cases. We observed that the simulated scattering functions show good

consistency with the experimental observations for both the index matched and non-index

matched cases (Fig. 2E). Our computational models successfully capture important features of

experimental data. For θ = 0º, scattering functions for index matching is higher than non-index

matching. For θ > 0º, the scattering functions decrease with θ, and scattering for index matching

is lower than non-index matching. These findings suggest that the presence of surface

scattering increases the overall scattering (θ > 0º). Both scattering functions increase for large

θ, between 140º and 180º. The back scattering detected at θ = 180º is due to Fresnel reflection

(optical reflection occurring at interfaces between media of different refractive indices) (Fig. 2A

and B) [44]. The flat horizontal upper surface of the index-matching medium (Fig. 2B) creates a

Fresnel reflection for the index matched scenario that is much stronger than that of the non-

index matched case.

In addition to calculating the scattering due to volume scattering only (via index

matching), we calculated the scattering due to surface scattering alone (by removing volume

scattering), the sum of the volume only and surface scattering only models, and finally the

scattering in the complete model incorporating both volume and surface scattering (Fig. 2F).

In the surface scattering only model, the worm and water film act as an approximately

cylindrical lens. Due to the resulting lensing effect, the surface scattering is strong for θ

between 0º and 50º and reaches a minimum around 90º. Next, it increases between 90º and

180º, due to Fresnel reflections occurring at the worm-to-air interface. In comparison, the

volume scattering alone exhibits a more uniform angular distribution, gradually decreasing with

scattering angle except for peaks near at 0º and 180º due to transmission and Fresnel

reflections, respectively.

By comparing the sum of the volume only and surface scattering only models with the

complete model incorporating both volume and surface scattering, we can assess the degree of

interaction between volume and surface scattering. If volume and surface scattering occur

independently, we would expect the sum of the volume only and surface scattering only to equal

the complete model. This is indeed the case for large angles, from about 60º to 150º (Fig. 2F)

because Fresnel reflection from the upper worm surface is independent of volume scattering

occurring inside the animal. However, for angles less than 60º the complete scattering is much

smaller than the sum of the surface and volume scattering models. This is because rays

coming from below the worm must first go through volume scattering in the worm's body and

then surface scattering before being collected by the detector. In this case, the incident ray

directions are largely randomized by the volume scattering, causing most of them to scatter into

larger angles. On the other hand, the lensing effect of surface scattering in the complete model

helps refracting the rays to reach the detector, causing the complete scattering to be somewhat

higher than that of pure volume scattering (Fig. 2F).

Fig. 2. Analysis of volume and surface scattering from *C. elegans*. (A) Schematic diagram

of light scattering from *C. elegans* on a moist substrate. SS: surface scattering; VS: volume

scattering; FR: Fresnel reflection. (B) Schematic diagram of light scattering from *C. elegans* on

the substrate underneath an index-matching media. (C) Angle-resolved scattering functions for

C. elegans and seeded agar for non-index matched (non-IM) and index matched (IM) cases. (D)

Angle-resolved image contrast of *C. elegans* on seeded agar for non-IM and IM cases. (E)

Simulated angle-resolved scattering functions obtained from the computational model for non-

IM and IM cases, overlaid with the experimental results. (F) The simulated scattering functions

for the volume scattering only, the surface scattering only, the sum of the volume and surface

scattering, and the complete scattering.

Spatial distribution of light scattering

We asked how the volume and surface scattering contribute to the spatially

inhomogeneous patterns observed in the worm images as a function of angle (Fig. 1B). First,

we measured the cross-sectional scattering profile for the animal at various illumination angles

(Fig. 3A-D). We noticed that in some images the worm appeared brightest along its edges,

whereas in others it appeared brightest at its center. To quantify how much the light scattering

is spatially located around the edges of the worm body compared to the center, we defined the

edge-to-middle contrast index (EMCI) (see Materials and methods). The EMCI has a range from

-1 to 1 with positive EMCI indicating the scattering is brighter around the edges than in the

middle. We measured EMCI as a function of θ for the index matched and non-index matched

cases (Fig. 3E and F).

In the non-index matched case, for $\theta < 10°$, stronger scattering is experimentally

observed around the edges and weaker in the middle, yielding positive EMCI (Fig. 3E). For $\theta >$

$10°$, scattering is weaker around the edges and stronger in the middle, yielding negative EMCI.

When θ approaches $180°$, the EMCI turns positive again. In contrast, the EMCI for index-

matching is negative for all the experimentally tested θ except for $0°$ (Fig. 3F). These results

indicate that high scattering observed around the edges for θ close to 0º and 180º for the non-

index matched scenario are due to the surface scattering. In comparison, the volume scattering

consistently concentrates around the middle of the animal body with different θ.

Next, we ask whether the observed spatial distribution of scattering can be reproduced

by our computational model depicted in Fig. 2A and B. The profiles of scattering across the

simulated worm for the representative θ in the scenarios of non-index matching and index

matching are shown in Fig. 3B and D, respectively. The simulated EMCIs as functions of θ for

both scenarios are respectively shown in Fig. 3E and F. The results from our simulation show

high consistency with the experimental data, further supporting the validity of our computational

model.

Combining the results from Fig. 2F, we can describe the mechanisms underlying the

angle-dependent spatial distribution of the scattering. For θ = 0º, in the bright field imaging

regime, the light incident to the animal is scattered out of the collection angle of the camera (Fig.

1A), causing the animal to appear darker than the background (Fig. 1B, Fig. 3A-D), yielding

negative image contrast (Fig. 1D and Fig. 2D) and positive EMCI (Fig. 3E and F). In the dark

field imaging regime, for small θ > 0º, the surface scattering is dominant and contributes to the

brightened boundaries of the animal due to the lensing effect, yielding positive EMCI (Fig. 3B

and E). For larger scattering angles, the volume scattering is dominant, and the scattering

spatially distributes around the center, yielding negative EMCI (Fig. 3E and F). The negative

contrast index and the positive EMCI observed near θ = 180º for worm on bacteria-seeded agar

(Fig. 1B, D, F, Fig. 3A, B, and E) is due to the Fresnel reflection from the substrate.

These findings hold important implications for choosing optimal illumination angles for

different purposes. For some applications where identification and tracking the animals are

more important, smaller angles would be preferable since the surface scattering highlights the

object boundaries and overall contrast is highest. For cases in which resolving the internal

structures of *C. elegans* is more important, the illumination should be set to higher angles where

the volume scattering is dominant.

Fig. 3. Spatial distribution of light scattering from *C. elegans*. (A) Images recorded at

representative angles for *C. elegans* on seeded agar, without an index matching media (FOV =

1.11 × 1.11 mm). The red dashed line highlights the pixels used for calculating the cross-

sectional scattering profiles at the midpoint of the animal. (B) Scattering profiles across the

animal body at the representative angles for the non-index matched case. E1, M, and E2

represent the first edge portion, the middle portion, and the second edge portion of the profile

used for calculating EMCI. (C) Images recorded at representative angles for *C. elegans* on

seeded agar underneath an index matching media. (D) Scattering profiles across the animal

body at the representative angles for the index matched case. (E) Experimental and simulated

EMCIs as functions of θ for the non-index matched (non-IM) scenario. (F) Experimental and

simulated EMCIs as functions of θ for the index matched (IM) scenario.

Optimal illumination for microscopy

To illustrate the consequences of our findings about mechanisms of volume and surface

scattering for microscopy of *C. elegans*, we acquired images of worms using imaging systems

with higher magnification and numerical aperture (Materials and methods), under various angles

of illumination (Fig. 4).

In bright field microscopy (Fig. 4A), the angular range of illumination θ is smaller than

that collected by the objective lens. Some of the light incident on the worm scatters out of the

collection angles, causing the worm to appear darker than its background. When θ is small, the

light collected from the animal is dominated by low-angle scattering (Fig. 4A), which is largely

contributed by surface scattering (Fig. 2F). The surface scattering helps improve the contrast of

the edges with a tradeoff of lower resolution for the internal structures (Fig. 4B), consistent with

the results from our spatial analysis (Fig. 3A, B, and E). To better visualize internal features, the

angular range of illumination should be larger so that more light scattered to high angles due to

volume scattering can be collected (Fig. 4C). Compared to an image using narrow angle

illumination (Fig. 4B), wider-angle illumination better resolves internal structures within the body,

such as pharynx, intestine, and embryos (Fig. 4C).

In dark field microscopy (Fig. 4D), the illumination is outside the acceptance angle of the

objective lens. Optical scattering diverts some of the light into the acceptance angle, and the

worm appears bright against a dark background (Fig. 4E-F). When the illumination angle is low,

more low-angle scattering (surface scattering) is collected, contributing to bright edges and high

contrast against the background (Fig. 3A, B, E, and Fig. 4E). In contrast, larger θ causes more

high-angle scattering (volume scattering) to be collected, which facilitates the visualization of

internal structures (Fig. 4F). Compared to Fig. 4E, higher-angle illumination helps resolving the

internal organs (Fig. 4F and G).

For both bright field and dark field imaging, it is important for the illumination angles to

be either inside or outside of the lens' acceptance range, but not both simultaneously.

Illuminating with a range of angles both inside or outside of the lens' acceptance range results a

combination of bright field and dark field images, resulting in low image contrast (Fig. 4H).

These results show that the illumination for microscopy needs to be optimized according

to the goals of the study. Low-angle illumination highlights the boundaries and improves the

image contrast, which is beneficial for animal tracking and the identification of small

independent objects, such as embryos and larvae. By comparison, high-angle illumination helps

resolve internal structures, if needed.

Fig. 4. Effects of illumination angle on bright field and dark field microscopy. (A)

Schematic diagram of bright field microscopy. LAS: low-angle scattering; HAS: high-angle

scattering. (B) A bright field image of *C. elegans* under narrow-angle illumination (θ < 1.7°) FOV:

1.0 x 0.5 mm. (C) A bright field image of *C. elegans* under wider-angle illumination (θ < 8.7°). (D)

Schematic diagram of dark field microscopy. (E) A dark field image of *C. elegans* under low-

angle illumination (θ = 20°). (F) A dark field image of *C. elegans* under higher-angle illumination

(θ = 57°). (G) Computationally inverted version of F yields an image similar to bright field

illumination. (H) A combination of bright field and dark field images of *C. elegans* under

illumination applied both inside and outside of the collection angle range of the objective lens (θ

< 20°).

Substrate-dependent image contrast

Image contrast is often limited by the amount of scattering in the substrate background.

While *C. elegans* is commonly cultured on agar media, other gelling agents may also be used

instead of agar [9, 45, 46]. We asked how the composition of the substrate might affect image

contrast. We prepared three substrates including agar, gelatin, and gellan gum (Materials and

methods), at concentrations that yielded similar gel strengths, and measured scattering

functions and image contrast for animals on these substrates (Fig. 5).

Our results show that the scattering functions for all three substrates decrease with θ but

show different magnitudes (Fig. 5A). Overall scattering from gelatin and gellan gum is smaller

than from agar. This results in the highest image contrast for animals on gelatin and gellan gum,

and lowest on agar (Fig. 5B).

We observed that the scattering functions for all the substrates are dependent on

illumination wavelength (Fig. S3) and longer-wavelength illumination generates higher image

contrast (Fig. 5B, Fig. S3).

We also notice that the imprint due to animal's movement causes significant scattering

on gellan gum for $\theta < 10°$ (Fig. S3F), which leads to reduced image contrast (Fig. 5B, Fig. S3F).

These imprints may not be favorable for animal identification or tracking but could be helpful for

studying locomotion trajectories.

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356 **Fig. 5. Scattering functions and image contrast for different substrate polymers.** (A)

357 Angle-resolved scattering functions for *C. elegans* on unseeded media made with agar, gelatin,

358 and gellan gum, under red illumination. (B) Image contrast for *C. elegans* on unseeded agar,

359 gelatin, and gellan gum under red, green, and blue illumination. Each dot represents a

360 measurement at an individual illumination angle.

Discussion

In this work, we investigated the mechanisms of light scattering from the nematode *C.*

elegans using angle-resolved measurements and computational modeling. Our findings reveal

that the light scattering from *C. elegans* arises from two primary sources: volume scattering

within the animal body and surface scattering at the interfaces between the animal, air, and the

moist substrate. These two types of scattering exhibit distinct angular and spatial profiles and

contribute differently to the overall scattering. Surface scattering is more prominent at small and

large angles, while volume scattering dominates at intermediate angles (Fig. 2). Volume

scattering is highest at the central region of the animal (Fig. 3) and helps resolve internal

anatomical structures (Fig. 4). On the other hand, surface scattering enhances the animal's

boundaries (Fig. 3), which may aid animal detection and tracking but obscure internal features

(Fig. 4). The interplay between the volume and surface scattering suggests that the optimal

imaging configuration should be tailored to the specific applications, for example, behavioral

tracking versus anatomical visualization.

We found that the image contrast is highly sensitive to multiple parameters, including the

illumination angle, wavelength, and the type of imaging substrates. These findings inform a set

of imaging parameters that enhance image contrast. The simple strategy we provided here does

not require specialized optics and labels, and therefore it is readily adaptable to many systems,

for example, microscopy [17, 47-50], high-throughput screening setups [10, 18, 19, 51, 52],

microfluidic chambers [41, 53, 54], and automation systems [12, 52, 55, 56].

Together, this study provides a biophysical understanding for the light–tissue interactions

in multiple-scattering organisms and offers physically informed strategies for optimizing the

image quality in a wide range of imaging systems. Our findings hold immediate value for

researchers using *C. elegans* as a model system, and broader relevance for biological imaging

across disciplines, such as genetics, neuroscience, and pharmacology.

It is important to acknowledge the limitations of our work. First, some sources of error

may affect the accuracy of the scattering function measurements. For example, the imperfect

alignment of the components to the optical axis would also contribute to the discrepancy in the

measurements of the angle. The LEDs, the FOV of the camera, and the iris of the lens have

non-zero sizes, such that the scattered lights we collected are from a finite range of angles.

Therefore, in order to minimize the angular uncertainty, we selected lights for particular angles

by intentionally setting the numerical aperture of the lens to be low (F-number f/16, NA ≈ 0.031),

with an acceptable cost of lower spatial resolution of the images (Fig. 1B, Fig. 3A and C).

There are a few factors that may affect the agreement of our computational model to the

experimental data. First, the animal is modeled as a simple cylinder, where the geometry of

head and tail and the bending of the body are not captured. Second, optical absorption from the

animal is not modeled. Third, we did not model the scattering due to the substrate or bacteria,

causing the background noise underestimated in the detected signal. Next, the optical

simulation is based on geometric ray tracing where the wave properties of the light, such as

diffraction, interference, and polarization, are not modeled.

This study focuses on a single developmental stage (day-1 adult) of a specific *C.*

elegans strain (N2). Variations in body morphology, internal structures, or cuticle properties

across life stages or genetic backgrounds may alter light scattering patterns [17]. Future work

could explore how scattering profiles change across developmental stages and in genetic

mutants with altered anatomy.

The calculation of the scattering function and image contrast is on an individual-animal

basis, which may neglect the variability between different organs and tissues. For example,

pharyngeal muscles are less scattering compared to embryos or intestine (Fig. 4F). Future work

will investigate the scattering profiles for different anatomical structures.

We expect the mechanisms of light scattering we identified here to have better

explanatory power for transparent, absorption-minimal biological samples, such as other

nematodes [57], fruit fly larvae [58], zebrafish larvae [59], and human organoids [60].

Materials and methods

Imaging system

Our imaging system contains a LED ring comprised of 36 RGB LEDs (Qasim, 95 mm diameter), a 3D-printed culture plate holder, a camera lens (Nikon Micro-NIKKOR, effective focal length 105 mm, set to f/16, with C-mount adapter) and a CMOS camera (Imaging Source, DMK 33GP031) (Fig. 1A). All optical elements are mounted to a 1.5" diameter, vertically oriented post (Thorlabs P14) through post mounting clamps (Thorlabs C1511) and they are aligned to be centered on a common optical axis. The distance between the lens and the plate was 400 mm. The illumination angle θ can be varied by translating the LED ring along the optical axis. The LED ring can be mounted below the agar plate to provide angles <90° or above the plate to provide angles between 90° and 180°. Some angles near 0°, 90°, and 180° are unavailable due to constraints of the optical mounts. The LED array is programmable to emit red, green, or blue light. We measured the peak wavelengths of the LEDs using a fiber optic coupled spectrometer (Ocean Optics FLAME-S-XR1) to be 629, 516, and 447 nm. The spectral widths were approximately 16, 26, and 14 nm (FWHM), respectively.

Measurement of scattering functions

The angle-resolved scattering function $p(\theta)$ describes the probability of a photon scattering into a unit solid angle oriented at an angle θ relative to its original trajectory [35, 36]. Here we calculated $p(\theta)$ via the expression (Eq. 1)

$$p(\theta) = \frac{P_{cam}(\theta)}{P_{in}(\theta) \cdot \Omega}$$

437 where $P_{cam}(\theta)$ is the optical power received by the camera over the region of interest (ROI),

438 $P_{in}(\theta)$ the total optical power incident on the object, and Ω the solid angle collected by the

439 camera lens. We measured incident irradiance using an optical power meter (Coherent

440 FieldMate Laser Power Meter with OP-2 VIS sensor). To calculate $P_{cam}(\theta)$, we used the

441 expression (Eq. 2)

$$P_{cam}(\theta) = \frac{PV(\theta) \cdot A_{roi}}{k \cdot T_{exp}}$$

443 where A_{roi} is the area of the ROI, T_{exp} the exposure time, and k is the conversion coefficient

444 from irradiance to pixel value. The measured k values for our camera in response to red, green,

445 and blue light are 3.0, 3.3, and 3.5 × 10^5 $ADU \cdot mm^2 \cdot s^{-1} \cdot \mu W^{-1}$ respectively, where ADU is an

446 analog-to-digital unit.

Calculation of the contrast index

449   We defined the contrast index (CI) for an animal against its background as (Eq. 3)

$$CI = \frac{\overline{PV_w} - \overline{PV_{bg}}}{\overline{PV_w} + \overline{PV_{bg}}}$$

451 where $\overline{PV_w}$ and $\overline{PV_{bg}}$ are the average pixel values inside the contour of the worm and the

452 background. The contours of the animal were labeled using ImageJ (version 1.54g). The ROI for

453 its background was chosen as a 1.1 mm x 1.1 mm rectangular region adjacent to the animal.

454   Since the scattering function $p(\theta)$ is proportional to the pixel values $PV(\theta)$ (Eq. 1 and 2),

455 CI can be also expressed as (Eq. 4)

$$CI(\theta) = \frac{\overline{p(\theta)_w} - \overline{p(\theta)_{bg}}}{\overline{p(\theta)_w} + \overline{p(\theta)_{bg}}}$$

Computational modeling

We conducted simulations using Ansys Zemax OpticStudio 2024 R2 optical design and

analysis software in non-sequential mode. First, we constructed 3-dimensional models

representing *C. elegans* on the culture substrate in the scenarios of non-index matching and

index matching (Fig. 2A and B). The simulated animal is illuminated by 36 collimated light

sources equally spaced in the azimuthal direction and with polar angle of θ relative to the optical

axis of the camera lens, which is modeled as a paraxial lens (f = 105 mm) with an iris (6.56 mm

diameter) positioned 400 mm from the animal. The lens forms an image at an image distance

of 142.37 mm on the other side. The image is captured by an incoherent irradiance detector.

For each angle tested we performed 10 trials of ray tracing, with each trial containing 2×10^8

analyzed rays.   We recorded an irradiance map after each trial. Next, we calculated the

scattering function $p(\theta)$ based on Eq. 1.

To model the volume scattering, we set the scattering within the simulated animal to

follow the Henyey-Greenstein phase function [28] (Fig. 2A and B). The two parameters, the

anisotropy factor (g) and mean free path (MFP), were chosen by minimizing the mean squared

error between the simulated scattering function and the experimental observations in the index-

matched scenario (Fig. 2B and E). The optimized value for the anisotropy parameter was g =

0.83 and the mean free path was MFP = 12 µm.

After obtaining the optimized values for g and MFP, we simulated the complete

scattering from the animal using the non-index matched model (Fig. 2A). We optimized the

surface scattering by adjusting the shape of the capillary film such that the mean squared error

between the simulated scattering function and the experimental data was minimized (Fig. 2E

and Fig. S2).

Edge-to-middle contrast index (EMCI)

483   To calculate the EMCI for the scattering profile across the animal (Fig. 3A-D), we first

484 symmetrically divided the distribution into three portions (Fig. 3B), including the first edge

485 portion (42 μm length), the middle portion (84 μm length), and the second edge portion (42 μm

486 length). Next, we calculated the average value of the scattering function for both edge portions

487 $\overline{p(\theta)_e}$ and for the middle portion $\overline{p(\theta)_m}$. We then calculated EMCI as (Eq. 5)

$$EMCI(\theta) = \frac{\overline{p(\theta)_e} - \overline{p(\theta)_m}}{\overline{p(\theta)_e} + \overline{p(\theta)_m}}$$

Bright field and dark field microscopy

The bright field microscopy images of *C. elegans* (Fig. 4B and C) were obtained using

an objective lens (NA = 0.3, Leica 1× planapochromat) on a stereo microscope (Leica M165

FC). The angular range of illuminations was controlled by adjusting the aperture size of the

illumination condenser.

The dark field microscopy images of *C. elegans* (Fig. 4E-H) were obtained using an

objective lens (NA = 0.14, f = 40 mm, Mitutoyo 5X Plan Apochromat) together with a tube lens (f

= 100 mm, Thorlabs TTL100-A) in a custom-built infinity-corrected microscope. Similar to Fig.

1A, the illumination angle is controlled by positioning an LED ring concentric with the optical axis

at different distances to the animal.

Preparation of imaging substrates

To prepare different culture substrates used for comparison of image contrast (Fig. 5),

we dissolve powders of agar, gelatin, and gellan gum in nematode growth medium (NGM) liquid

at 100 ºC and wait until they solidify at room temperature. The concentration of the agar was set

to 17 g/L according to the standard protocol [61]. In order to obtain gelatin and gellan gum

substrates suitable for maintaining *C. elegans*, we adjusted their concentrations such that the

resultant substrates have comparable mechanical strengths as the standard 1.7% agar plate.

We measured their mechanical strengths by a puncture test, where the minimum forces for an

indenter (active area 19 mm^2) causing fractures on the substrate were recorded by a precision

balance (Mettler Toledo MS1602S). After these tests we set the concentration of gelatin to be

10.1% (average strength 79.9 ± 6.7 kPa, mean ± SD, N = 10 plates) and the gellan gum 1.125%

(average strength 79.9 ± 5.7 kPa, N = 10), such that their mechanical strengths are comparable

to the standard 1.7% agar plates (average strength 85.1 ± 9.3 kPa, N = 10).

Preparation of *C. elegans*

516   The N2 *C. elegans* strain used in this study was cultivated at 20°C on NGM plates with

517 OP50 bacteria using standard methods [61]. We performed timed egg laying assays within a 4-

518 hour window for synchronizing the ages of the animals. We used day 1 adults for all the

519 experiments. All experiments were performed at room temperature (approximately 20 ºC).

Data, materials, and software availability

The raw data for this paper is available in *SI Appendix*, Dataset S1. The design file for constructing the 3D models is available in *SI Appendix*, Design File S1. The source file for the optical simulation is available in *SI Appendix*, Software S1. The source code for the data analysis is available in *SI Appendix*, Software S2.

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Supplementary material

Fig. S1. (A) Angle-resolved scattering functions for *C. elegans* and seeded agar for the index

matched scenario under red, green, and blue illumination. (B) Image contrast index for *C.*

elegans on seeded agar as a function of illumination angle and wavelength for the index

matched case.

Fig. S2. Scattering functions of *C. elegans* obtained from the experiment (on seeded agar under 629-nm illumination) and the computational models with different geometries of capillary water film. Insets: the cross section of the simulated animal and capillary water film.

Fig. S3. Angle-resolved, wavelength-dependent scattering functions and image contrast for *C.*

elegans on unseeded media made with (A and B) agar, (C and D) gelatin, and (E and F) gellan

gum. Inset of (F): a representative image of *C. elegans* and its imprints on unseeded gellan gum

under a red illumination applied at 6º (FOV: 2.22 × 2.22 mm).

Dataset S1 (separate file). Raw data for the plots presented in this paper.

Design File S1 (separate file). Design files (in STEP format) for the 3D models representing *C.*

elegans and its surroundings.

Software S1 (separate file). Source files (in ZMX and ZPL format) for conducting the optical

simulations.

Software S2 (separate file). Source MATLAB scripts for the data analysis performed in this

study.