1 *Biophysical models reveal the relative importance of transporter proteins and*
2 *impermeant anions in chloride homeostasis*
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28 *Running Title*
29 The cellular determinants of Cl^- driving force in the nervous system.
30

31 *Impact statement*
32 Impermeant anions may influence chloride concentration in neurons, but do not appreciably change
33 chloride driving force, the parameter important for inhibitory signaling.
34

35 *Keywords*
36 Chloride, cation-chloride cotransport, impermeant anions, KCC2, Na^+/K^+ ATPase, inhibition, GABA
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Abstract

Fast synaptic inhibition in the nervous system depends on the transmembrane flux of Cl^- ions based

on the neuronal Cl^- driving force. Established theories regarding the determinants of Cl^- driving force

have recently been questioned. Here we present biophysical models of Cl^- homeostasis using the

pump-leak model. Using numerical and novel analytic solutions, we demonstrate that the Na^+/K^+-

ATPase, ion conductances, impermeant anions, electrodiffusion, water fluxes and cation-chloride

cotransporters (CCCs) play roles in setting the Cl^- driving force. Our models, together with

experimental validation, show that while impermeant anions can contribute to setting $[Cl^-]_i$ in

neurons, they have a negligible effect on the driving force for Cl^- locally and cell-wide. In contrast, we

demonstrate that CCCs are well-suited for modulating Cl^- driving force and hence inhibitory signalling

in neurons. Our findings reconcile recent experimental findings and provide a framework for

understanding the interplay of different chloride regulatory processes in neurons.

 Introduction

 Fast synaptic inhibition in the nervous system is mediated by type A γ-aminobutyric acid receptors

 (GABA$_A$Rs) and glycine receptors (GlyRs), which are primarily permeable to chloride (Cl$^-$) (Farrant and

 Kaila, 2007). Together with the neuronal membrane potential, the transmembrane gradient for Cl$^-$

 sets the driving force for Cl$^-$ flux across these receptors, controlling the properties of inhibitory

 signaling. Modification of neuronal intracellular Cl$^-$ concentration ([Cl$^-$]$_i$) has been shown to play a

 causative role in multiple neurological diseases including epilepsy, chronic pain, schizophrenia and

 autism (Rivera et al., 2004; Huberfeld et al., 2007; Price et al., 2009; Hyde et al., 2011; Tyzio et al.,

 2014). Similarly, intracellular Cl$^-$ is thought to be modulated during brain development so that

 GABAergic transmission contributes optimally to the construction of neural circuits (Ben-Ari, 2002).

 Given the importance of Cl$^-$ for brain function and dysfunction, the cellular mechanisms that control

 its transmembrane gradient and driving force are of considerable interest.

 Plasmalemmal Cl$^-$ transporters, in particular cation-chloride cotransporters (CCCs), are understood to

 be the major mechanism by which neurons regulate the driving force for Cl$^-$ permeable anion channels

 (Kaila et al., 2014). Recently, it has been suggested that in fact impermeant anions control local [Cl$^-$]$_i$

 and driving force (Glykys et al., 2014), rather than the CCCs. The majority of intracellular anions are

 impermeant to the neuronal membrane; these include ribo- and deoxynucleotides, intracellular

 proteins and metabolites (Burton, 1983). Impermeant anions induce what is known as the Donnan (or

 Gibbs-Donnan) effect (Hill, 1956; Sperelakis, 2012) – an uneven distribution of impermeant molecules

 across the membrane which is osmotically unstable. Without active ion transport to counter this

 effect, neurons would swell and burst (Kay, 2017). Animal cells, including neurons, maintain cell

 volume in the presence of impermeant anions by using the Na$^+$/K$^+$-ATPase to pump Na$^+$ out of the cell

 and K$^+$ in, along with the passive movement of water and other ions (Tosteson and Hoffman, 1960;

 Armstrong, 2003; Liang et al., 2007; Kay, 2017). This pump-leak mechanism, whilst stabilizing cell

 volume, also establishes the negative resting membrane potential and transmembrane Na$^+$ and K$^+$

gradients, which serve as energy sources for the coupled transport of other molecules, including Cl⁻

by CCCs. In the absence of active Cl⁻ transport, $[Cl^-]_i$ is set by the membrane potential; i.e. the Nernst

potential of Cl⁻ ($E_{Cl} = \frac{RT}{F} \ln \frac{Cl_i}{Cl_o}$) equals the transmembrane potential (V_m).

The transmembrane Cl⁻ gradient and the driving force ($DF=V_m-E_{Cl}$) for Cl⁻ permeable ion channels are

therefore the outcome of multiple, dynamically interacting mechanisms. This makes experimental

investigation of the determinants of Cl⁻ driving force difficult, particularly at a local level.

Computational models based on established biophysical first principles are productive means for

exploring the roles of cellular mechanisms in generating local Cl⁻ driving force. Here, we establish

numerical and novel analytic solutions for an inclusive model of Cl⁻ homeostasis to elucidate the

determinants of the neuronal driving force for Cl⁻.

We demonstrate that baseline $[Cl^-]_i$ is a product of the interaction of Na^+/K^+-ATPase activity, the

average charge of impermeant anions, ion conductances, CCCs and water permeability. Consistent

with recent experimental reports (Glykys et al., 2014), and our own experimental validation using

electroporation of anionic dextrans and optogenetic probing of E_{GABA}, we find that impermeant anions

can contribute to setting $[Cl^-]_i$. However, we find that they can only affect the Cl⁻ driving force by

modifying active transport mechanisms, and then only negligibly. Impermeant anions therefore do

not appreciably modify synaptic signaling properties, contrary to recent experimental interpretation

(Glykys et al., 2014). In contrast, we demonstrate using biophysical models and gramicidin perforated-

patch clamp recordings that CCCs selectively regulate substantial changes in the Cl⁻ driving force. This

is consistent with a meta-analysis of experimental data from the field, which shows a strong

correlation between the activity of the specific CCC, KCC2, and Cl⁻ driving force. The ability of CCCs to

specifically modulate Cl⁻ at a local level depends on the characteristics of Cl⁻ electrodiffusion in the

structure concerned, demonstrated using multicompartment modelling. Together, our models

provide a theoretical framework for understanding the interplay of chloride regulatory processes in

neurons and interpreting experimental findings.

Results

A biophysical model based on the pump-leak mechanism demonstrates the importance of the sodium-

potassium ATPase for setting transmembrane ion gradients including chloride.

To compare the effects of impermeant anions and Cl^- cotransport on Cl^- homeostasis we first

developed a single compartment model based on the pump-leak formulation (Tosteson and Hoffman,

1960; Kay, 2017) (**Fig. 1A**). This model, defined by a set of differential equations, incorporated

mathematical representations of the three major permeable ion species Cl^-, K^+, Na^+ as well as

impermeant anions (X^z) with average charge z. Permeable ions could move across the cellular

membrane via passive conductances according to each ion's respective electrochemical gradient.

Further, active transport of Na^+ and K^+ by the Na^+/K^+-ATPase and cotransport of Cl^- and K^+ by the

cation-chloride cotransporter KCC2 with a non-zero conductance g_{KCC2} of 20 $\mu S/cm^2$ unless otherwise

stated were included. Finally, our formulation accounted for the dynamics of cell volume (w),

intracellular osmolarity (Π) and transmembrane voltage (V_m) (see Materials and Methods).

Importantly, regardless of initial starting concentrations of permeant or impermeant ions, cell volume

or V_m, the model converged to stable fixed points without needing to include any means for 'sensing'

ion concentration, volume or voltage. In the case of permeable ions, initial concentrations did not

even determine the final cellular volume. For example, despite initiating the model with different

starting concentrations of Cl^- (1, 15, 40 and 60 mM respectively, **Fig. 1B**), the $[Cl^-]_i$ always converged

to the same stable concentration of 5.2 mM, a typical baseline $[Cl^-]_i$ for adult neurons, and volume

always converged to 2.0 pL, a typical volume for hippocampal neurons (Ambros-Ingerson and Holmes,

2005). The model is robust in the sense that its convergence to a stable steady state does not depend

on a narrow set of parameters and initial values. Another proposed model of the Na^+/K^+-ATPase with

experimental validation (Hamada et al., 2003) parameterized for our steady state values had similar

activity in terms of net flux rate in our range of $[Na^+]_i$ and invariance to initial permeable ionic

concentrations, (**Figure 1 – figure supplement 1A-B**).

Consistent with previous results (Xiao et al., 2002; Dierkes et al., 2006; Dijkstra et al., 2016), 'turning

off' the activity of the Na^+/K^+-ATPase in our model led to a progressive collapse of transmembrane ion

gradients, progressive membrane depolarization and continuous and unstable cell swelling. Such

effects could be reversed by reactivation of the pump (**Fig. 1C**). The relative activity of the Na^+/K^+-

ATPase sets the final stable values for ion concentrations (of both permeable and impermeant ions),

V_m and volume (**Fig. 1C**). When we increased the activity of the Na^+/K^+-ATPase in our model, the final

steady-state concentration for K^+ increased, whilst Na^+ and Cl^- dropped to levels that approximate

those observed in mature neurons (**Fig. 1D**). At the same time, the final, stable-state membrane

potential and cell volume also decreased. Interestingly, as has been observed previously (Fraser and

Huang, 2004), beyond a certain level, further increases in Na^+/K^+-ATPase activity have negligible

effects on cell volume and transmembrane voltage. For subsequent analysis, we chose a 'default'

effective pump rate for the Na^+/K^+-ATPase of approximately 1.0×10^{-2} $C/(dm^2.s)$, that is a pump rate

constant of 10^{-1} $C/(dm^2.s)$ (equation (2)), and average intracellular impermeant anion charge (z) of -

0.85 as extrapolated from reasonable cellular ionic concentrations and osmolarity (Lodish et al., 2009;

Raimondo et al., 2015). This resulted in steady-state ion concentrations and membrane potentials that

approximate those experimentally observed in mature neurons: Cl^- 5 mM; K^+ 123 mM; Na^+ 14 mM; X^z

155 mM; and a V_m of -72.6 mV (Jiang and Haddad, 1991; Diarra et al., 2001; Tyzio et al., 2008). We

were able to corroborate the numerical solutions for final steady-state values by developing a

parametric-analytic solution (**Supplementary file 1**). We observed exact correspondence between the

numerical and analytic solutions within our model (**Fig. 1D**). In subsequent analyses, this novel analytic

solution allowed us to explore rapidly a large parameter space to determine how various cellular

attributes might affect Cl^- homeostasis.

Figure 1: A biophysical model of ion dynamics based on the pump-leak mechanism demonstrates the

importance of the sodium-potassium ATPase for setting transmembrane ion gradients including

chloride. (A) A single cell compartment was modelled as a cylinder with volume changes equivalent to

changes in cylindrical radius. Dynamics of membrane permeable potassium (purple, K^+), sodium (pink,

Na^+) and chloride (green, Cl^-) ions were included. Impermeant anions (orange, X^z) had an average

intracellular charge z. The KCC2 transporter moved Cl^- and K^+ in equal parts according to the

transmembrane gradient for the two ions. The sodium-potassium ATPase transported 3 Na^+ ions out

for 2 K^+ ions moved into the cell. (B) Regardless of intracellular starting concentrations of the

permeable ions, the model converged to identical steady state values for all parameters without

needing to include any means for 'sensing' ion concentration, volume or voltage. We show the result

for Cl⁻ as a time series of $[Cl^-]_i$ (top panel) and volume (bottom panel). (C) The ATPase plays a key role

in maintaining steady state ion concentration, membrane voltage (V_m) and volume. Switching off the

ATPase results in a continuous increase in cell volume (bottom panel), membrane depolarization

(middle panel) and ion concentration dysregulation (top panel with colours per ion as in 'A'. All cellular

parameters recovered when the ATPase was reactivated. (D) The model's analytic solution showed

exact correspondence with steady state values generated by numerical, time series runs (dots) for

varying ATPase pump rates. Steady state values for the concentrations of the ions with colours as in A

(top panel); V_m (middle panel) and volume (bottom panel). The dashed line indicates the default ATPase

pump rate used in all simulations unless specified otherwise. The result in (B) is replicated with Hamada

et al.'s experimentally validated model of the ATPase in Figure 1 – figure supplement 1B, and the

models' respective pump fluxes are compared in figure supplement 1A.

Membrane chloride conductance affects steady-state intracellular chloride concentration only in the
presence of cation-chloride cotransport.

Using the analytic solution, we investigated how changes in baseline ion conductance for the major

ions in our model (g_K, g_{Na} and g_{Cl}) affected Cl⁻ homeostasis. We calculated the steady-state values for

the Cl⁻ reversal potential (E_{Cl}) and potassium reversal potential (E_K), resting membrane potential (V_m)

and volume (w) whilst independently manipulating the conductance for each ion (**Fig. 2**). Increasing

the baseline potassium conductance (g_K) resulted in E_{Cl}, E_K and V_m converging to similar steady-state

values (**Fig. 2A**) without significantly affecting cell volume. We were also able to replicate the classic

dependence of membrane potential on $log([K^+]_o)$ (**Figure 2 – figure supplement 1**). In contrast,

increasing the baseline sodium conductance (g_{Na}) beyond 20 $\mu S/cm^2$ resulted in a steady

depolarization of E_{Cl}, E_K and V_m with accompanying cell swelling (**Fig. 2B**).

The effect of manipulating Cl⁻ conductance (g_{Cl}) depended on the activity of concurrent cation-chloride

cotransport by KCC2 (**Fig. 2C**). In the presence of active KCC2 at very low values of g_{Cl}, the steady state

$[Cl^-]_i$ is such that E_{Cl} approaches E_K. This follows because in the absence of alternative Cl⁻ fluxes, KCC2

utilizes the transmembrane potassium gradient to transport Cl⁻ until E_{Cl} equals E_K. With increasing g_{Cl}

however, E_{Cl} increases, moving away from E_K towards V_m, and at very high Cl⁻ conductances E_{Cl} and V_m

approached similar values in our model. Without the activity of KCC2, g_{Cl} had no effect on steady state

E_{Cl}, E_K, V_m or volume, with g_{cl} able to be very low but greater than zero (**Fig. 2D**). In this instance E_{Cl}

always equals V_m as the movement of Cl⁻ across the membrane is purely passive. Without the activity

of KCC2, there can be no driving force for Cl⁻ flux at steady state (V_m-E_{Cl} = 0). Our model therefore

behaved in a manner consistent with established theoretical predictions (Kaila et al., 2014).

Figure 2: Membrane chloride conductance affects steady-state intracellular chloride concentration only in the presence of cation-chloride cotransport. Steady state values for different ionic conductance were calculated using the model's analytic solution. (A) Steady state E_{Cl} (green), E_K (purple), V_m (black) and volume w (bottom panel) were calculated at different K^+ conductances (g_K). Increasing g_K resulted in a convergence of steady state E_{Cl}, E_K and V_m. (B) Increasing Na^+ conductance g_{Na} resulted in a progressive increase in steady state E_{Cl}, E_K and V_m with accompanying cell swelling. (C) In the presence of active cation-chloride cotransport (g_{KCC2} = 20 µS/cm²), increasing Cl^- conductance shifted steady state E_{Cl} from E_K toward V_m. (D) In the absence of KCC2 activity, g_{Cl} had no effect on steady state parameters. E_{Cl} equals V_m in all instances. Dashed lines indicate the default values for g_K, g_{Na} and g_{Cl}.

In Figure 2 – figure supplement 1 the classic dependence of the membrane potential on log([K$^+$]$_o$) is

shown.

<u>Cation-chloride cotransport sets the chloride reversal and driving force for transmembrane chloride</u>

<u>flux.</u>

Next, we used our single-cell unified model to explore how the activity of cation-chloride cotransport

affects Cl$^-$ homeostasis. In our model, the activity of KCC2 is set by the conductance of KCC2 (g_{KCC2}).

Using the numerical formulation with the default values described in Figure 1 and 2, we steadily

increased g_{KCC2} from 20 µS/cm^2 to 370 µS/cm^2 and tracked changes to E$_{Cl}$, E$_K$, V$_m$ and volume. Increasing

KCC2 activity over time caused a steady decrease in [Cl$^-$]$_i$ reflected by a hyperpolarization of E$_{Cl}$ (**Fig.**

3A). V$_m$ decreased only modestly, resulting in an increase in the driving force for Cl$^-$ flux that tracks the

increase in g_{KCC2}. This effect saturates as E$_K$ constitutes a lower bound on E$_{Cl}$. Importantly, increases in

g_{KCC2} resulted in persistent changes to E$_{Cl}$ and the driving force for Cl$^-$. Employing alternate models for

KCC2 (Fraser and Huang, 2004; Lewin et al., 2012; Raimondo et al., 2012) and the ATPase (Hamada et

al., 2003) did not change this result when compensation for parameterization was given, although

different KCC2 models result in different kinetic rates for Cl$^-$ and K$^+$ transport (**Figure 3 – figure**

supplement 1 and 2). Using the analytic solution to our model we calculated how KCC2 activity affects

steady state values of E$_{Cl}$, E$_K$, V$_m$, w and Cl$^-$ driving force (**Fig. 3B**). In confirmation of our findings in

Figure 2D, with no KCC2 activity (g_{KCC2} = 0), E$_{Cl}$ equaled V$_m$ and the Cl$^-$ driving force was zero. As we

increased g_{KCC2}, steady state E$_{Cl}$ pulled away from V$_m$ and approached E$_K$. This resulted in an increase

in Cl$^-$ driving force (V$_m$-E$_{Cl}$) with steady state values of 11.3 mV at our chosen default value of g_{KCC2}. The

results obtained with our model are therefore fully consistent with the view that CCCs, in this case

KCC2, establish the driving force for Cl$^-$.

To test this theoretical finding, we performed gramicidin perforated patch-clamp recordings from CA3

hippocampal neurons in rat organotypic brain slices whilst activating Cl$^-$ permeable GABA$_A$ receptors

with muscimol (10 µM), in order to measure the $GABA_A$R driving force, which approximates Cl^- driving

force (**Fig. 3C**). We then tested our model predications by applying the CCC blocker, furosemide (1

mM) (**Fig. 3D, E**). We noted that after furosemide was introduced the E_{GABA} became significantly more

depolarized (baseline: -78.81 ± 2.78 mV vs furosemide: -71.59 ± 3.04 mV, *n* = 10, *p* = 0.01, *paired t-*

test) whilst there was no significant difference in V_m (-69.92 ± 1.77 mV vs -71.51 ± 2.61 mV, *n* = 10, *p*

= 0.36, *paired t-test*) (**Fig. 3E**). This reflects a significant change in the $GABA_A$R driving force (8.89 ±

3.44 mV vs 0.08 ± 4.57 mV, *n* = 10, *p* = 0.04, *paired t-test*). Furthermore, we then noted that this

change in driving force persisted for at least 15 minutes post application of furosemide (**Fig. 3F**). These

results are consistent with our model predictions and demonstrate how the application of a CCC

blocker reduces the $GABA_A$R driving force (and hence Cl^- driving force) by selectively depolarizing E_{GABA}

with negligible effects on V_m.

In addition, we sought experimental data from the literature to determine whether changes in KCC2

activity correlate with alterations to steady-state $[Cl^-]_i$. We focused on changes in KCC2 expression

level, as this is likely to be a strong predictor of changes in KCC2 activity. Indeed, in a meta-analysis of

7 studies and 8 experiments from our review of 26 studies, weighted for methodological biases and

data quality, we observed a significant correlation (R^2 = 0.796, p < 0.001) between the change in KCC2

expression and Cl^- driving force (**Fig. 3G**). Absolute changes in V_m were less than 2 mV in all but one

study, meaning that the change in driving force could be ascribed to significant shifts in E_{GABA} (R^2 =

0.045, p < 0.001). The outlier data point (showing a 8.45 mV change) was from a study into the effects

of acute stress, where other factors could have transiently influenced V_m (MacKenzie and Maguire,

2015). The meta-analysis supports the prediction that cation-chloride cotransport by KCC2 is an

important determinant of $[Cl^-]_i$ (R^2 = 0.83, p < 0.001, 9 studies) and driving force (see **Supplementary**

file 2 Table S2-1 for raw data, and the scoring table for weighting in **Table S2-2**).

Since Cl$^-$ cotransport by KCC2 depends on the transmembrane gradient for K$^+$, which in turn is

established by the Na$^+$/K$^+$-ATPase, we explored how the Na$^+$/K$^+$-ATPase pump rate affected E_{Cl} in the

presence of KCC2 activity (g_{KCC2} = 20 μS/cm^2). As depicted in **Fig. 3H**, the Na$^+$/K$^+$-ATPase pump is a

critical determinant of intracellular Cl$^-$ homeostasis. Of interest, above a pump rate constant of 10^{-3}

C/(dm^2.s), where cell volume is stable, even many fold changes in pump rate have minimal effects on

E_{Cl} and V_m. It therefore seems unlikely that neurons would adjust the Na$^+$/K$^+$-ATPase to modulate Cl$^-$

driving force.

Figure 3: Cation-chloride cotransport sets the chloride reversal and driving force for transmembrane

chloride flux. (A) Increasing KCC2 activity in our model by increasing g_{KCC2} from 20 μS/cm^2 to 370 μS

/cm^2 resulted in a persistent decrease in E_{Cl} (green), a minimal decrease in E_K (purple), V_m (black) and

volume w (bottom panel). This resulted in a permanent increase in the DF for Cl$^-$ via a change in E_{Cl}

from -83.9 mV to -93.2 mV (red). (B) The steady state values for E_{Cl}, E_K, V_m (top panel), volume (middle

panel) and Cl$^-$ driving force (DF, bottom panel, red) at different KCC2 conductances. Increasing KCC2

activity resulted in a decrease in steady state E_{Cl} and DF. E_K represented a lower bound on E_{Cl} at high

KCC2 conductances. Similar results were noted for other kinetic models of KCC2 (figure supplement 1

and 2). (C) Schematic showing experimental setup. Gramicidin perforated patch-clamp recordings

were performed on CA3 pyramidal cells from rat hippocampal organotypic brain slices. (D) Insets depict

$GABA_AR$ currents elicited by somatic application (20 ms) of muscimol (10 μM) at different voltages.

Calibration: 500 ms, 500 pA. Current-voltage (IV) plots with holding current (reflecting membrane

current, grey) and total current (reflecting membrane current plus the muscimol-evoked current, black

or pink) were drawn to calculate changes in V_m, E_{GABA} and DF before (left, black) and after (right, pink)

furosemide was applied. Voltages were corrected for series resistance error. (E) Population data

showing significant changes in E_{GABA} and DF but not V_m five minutes after furosemide application. (F)

Changes in DF over time show a significant decrease from baseline once furosemide was introduced.

(G) Meta-analysis of experimental studies demonstrates a correlation between KCC2 activity (%

change) and Cl$^-$ DF (mV, top plot, red) but not membrane potential (mV, bottom plot, grey), confirming

the role of KCC2 for establishing the neuronal Cl$^-$ gradient in adult tissue. The data and scoring system

used to generate the regression can be found in Supplementary file 2 (Tables S2-1 and S2-2). (H)

Sufficient Na$^+$/K$^+$-ATPase activity is critical for steady state ionic gradients, V_m (top panel, colours as in

'A') and volume (bottom panel), but these variables are relatively stable near the default pump rate

*(dashed line). 'ns', non-significant; *p <0.05; ** p<0.01.*

 <u>Altering the concentration of intracellular or extracellular impermeant anions, without changing the</u>

 <u>average charge of impermeant anions, does not affect the steady state gradient or driving force for</u>

 <u>chloride.</u>

To determine the effect of impermeant anions on Cl^- homeostasis we first explored whether adjusting

the concentration of impermeant anions ($[X]_i$), while maintaining a constant average impermeant

anion charge (z), had any impact on E_{Cl}, E_K, V_m or volume. We initiated the full single-compartment

model with different starting concentrations of impermeant anions (but the same average charge, z =

-0.85), and observed that regardless of the initial concentration of impermeant anions, over a period

of minutes, the cell adjusted its volume to give an identical steady-state impermeant anion

concentration (**Fig. 4A**, $[A]_i$ = 155 mM). Analytically, it can be shown that the number of moles of X

determines completely the volume of the compartment, while the permeant ions alone cannot be

used to predict steady state volume (Kay, 2017). Similarly, all initial impermeant anion concentrations

resulted in identical steady state values of E_{Cl} (-83.8 mV), E_K (-95.1 mV) and V_m (-72.6 mV) (**Fig. 4B**).

This shows that simply adjusting the amount of impermeant anions within a cell has no persistent

effect on $[Cl^-]_i$.

We then tested the effect of dynamically adding impermeant anions with the default average charge

either intracellularly (**Fig. 4C**) or extracellularly (**Fig. 4D**). While impermeant anions are being added to

the cell, the membrane potential hyperpolarizes and E_{Cl} decreases. However, following the cessation

of impermeant anion influx, E_{Cl}, E_K, V_m and $[X]_i$ return to steady state values due to compensatory

changes to cell volume (**Fig. 4C**). There are transient transmembrane fluxes of all ions while anions are

added into the cell, and in particular the inward flux of the cations Na^+ and K^+, such that the sum $[X]_i$

+ $[Cl^-]_i$ is not necessarily kept constant during impermeant anion addition (**Figure 4 – figure**

supplement 1). Impermeant ions (with default average charge) were added to the extracellular space,

which is effectively an infinite bath in the model, while proportional decreases in $[Cl^-]_o$ were applied

to correct for the changes to charge and osmotic balance. Additions in the extracellular space,

similarly, resulted in a temporary depolarization of E_{Cl} and V_m, but no persistent shift in these

parameters (**Fig. 4D**). The addition of extracellular impermeant anions did however result in a small

compensatory decrease in cell volume secondary to the large shifts in $[Cl^-]_i$ required to maintain the

proportion of $[Cl^-]_i$ to $[Cl^-]_o$ according to the Nernst potential. In summary, there is no lasting effect on

the reversal potential or driving force for Cl^- if only the concentration of a neuron's intracellular or

extracellular impermeant anions is altered.

 Figure 4: Adding intracellular or extracellular impermeant anions, without changing the average

charge of impermeant anions, does not affect the steady state gradient or driving force for chloride.

(A) Initiating the model with different starting concentrations of intracellular impermeant anions ($[X]_i$)

with the same average charge z (orange, top panel), led to compensatory volume changes (bottom

panel) which resulted in identical steady state concentrations. (B) Steady state E_{Cl} (green), E_K (purple)

and V_m (black) were identical regardless of initial $[X]_i$. Final volume however was a linear function of

initial $[X]_i$ (bottom panel). (C) Addition of impermeant anions of the same average charge (z) caused

transients shifts in E_{Cl} (green, top panel), E_K (purple), V_m (black) as well as $[X]_i$ (orange, bottom panel)

for the duration of the influx, and sustained increases in volume (black, middle panel). No persistent

changes in E_{Cl}, E_K or V_m were observed. The full time-dependent ionic and water fluxes for the

experiment are shown in Figure 4 – figure supplement 1, which shows that the inward flux of

impermeant anions causes fluxes of all other ions. (D) Similarly, the addition of extracellular

impermeant anions in an osmoneutral manner causes transient shifts in the permeable ion gradients

(top panel, colours as in 'C'), and sustained changes in cellular volume (black, middle panel) as well as

the extracellular X and extracellular and intracellular Cl^- concentrations.

Changing the average charge of impermeant anions can drive substantial shifts in the reversal

potential for chloride, but has negligible effects on chloride driving force.

We next sought to determine how changes in the average charge of the impermeant ions (z) might

influence the driving force on Cl^-. Such changes in average z could be associated with various cellular

processes, including post-translational modifications of proteins that decrease their charge without

changing the absolute number of protein molecules. To investigate this parameter, we modified the

average z of intracellular impermeant anions from -0.85 to -1 whilst measuring accompanying changes

in E_{Cl}, E_K, V_m and cell (**Fig. 5A**). We found that this shift to a more negative average z resulted in both

a transient and persistent decrease in E_{Cl}, E_K and V. Importantly, the shifts in E_{Cl} were accompanied by

broadly matching shifts in E_K and V, which resulted in a small change in the driving force for Cl^- of < 0.2

mV. Both numerical and analytic calculation of steady state values for E_{Cl}, E_K and V_m in our model

showed that changing the average charge of impermeant anions, while substantially affecting E_{Cl}, had

very small effects on the driving force for Cl$^-$ (**Fig. 5B**). By shifting z within reasonable ranges for

mammalian neurons (Lodish et al., 2009; Raimondo et al., 2015), and assuming osmo- and electro-

neutrality, only shifts of < 1 mV could be generated. In addition, although the absolute number of

impermeant anions (moles) remained constant throughout the process of modifying z, cell volume

shifted, and as a consequence modest alterations to the concentration of impermeant anions

occurred as well.

Next, instead of adjusting the charge of existing anions in intracellular impermeant anions as described

above, we directly added impermeant anions of differing valence to the cell (**Fig. 5C and D**). This had

the effect of both increasing the absolute quantity of impermeant anions and adjusting the average

charge of impermeant anions. The 'addition' of impermeant anions in this way models the *de novo*

synthesis of impermeant anion species, or their active transport into the cell. This process also

resulted in both transient and persistent changes to E_{Cl}, E_K and V, which was dependent on the extent

that average z was altered. Again, whilst the large additions of impermeant anions could substantially

alter the Cl$^-$ reversal potential, this had a negligible effect on the driving force for Cl$^-$ due to matching

shifts in V_m. Driving shifts in E_{Cl} in this manner also resulted in changes to cell volume (**Figure 5 – figure**

supplement 1).

To experimentally test our biophysical modelling predictions, we used photo-activation of ChR2

expressing GABAergic interneurons and whole-cell patch clamp recordings of mouse organotypic CA3

hippocampal pyramidal neurons to measure V_m, E_{GABA}, and GABA driving force (which approximates

Cl$^-$ driving force) before and after addition of impermeant anions (**Fig. 5D**). To add impermeant anions

to the recorded cell, we used single cell electroporation of fluorescently tagged anionic dextrans

(Alexa Flour 488, see Materials and Methods). Successful addition of impermeant anions could be

confirmed visually by observing strong and stable fluorescence of the anionic dextran restricted to the

recorded cell (**Fig. 5D, E**). To drive impermeant anions into the cell, negative voltage steps (20 ms, 0.5

– 1 V) were applied to the electroporation pipette necessarily resulting in direct membrane

depolarization, which recovered over a period of 1-5 minutes (**Fig. 5D**). Once this acute perturbation

had settled, as predicted by our model, changing the average charge of impermeant anions, following

the addition of highly negatively charged dextrans to the cell, resulted in a stable, mean negative shift

in V_m from a baseline of -66.95 ± 3.96 mV to -73.96 ± 3.06 mV (n = 6, *P* = 0.03, *Wilcoxon test,* **Fig. 5E,**

F). Again in line with our predictions, addition of impermeant anions also resulted in a significant

reduction of resting E_{GABA} (a proxy for E_{Cl}) from baseline values of -72.45 ± 1.95 mV to -77.25 ± 1.35

mV (*P* = 0.03, *Wilcoxon test*). Importantly, however, similar shifts in V_m and E_{GABA} resulted in an

undetectable shift in GABA and hence Cl^- driving force (5.49 ± 4.43 mV vs 3.33 ± 3.46 mV, n = 6, *P* =

0.22, *Wilcoxon test,* **Fig. 5E, F**).

Our single compartment model of Cl^- homeostasis, in conjunction with experimental validation,

demonstrates that whilst the adjustment of average impermeant anion charge can significantly affect

E_{Cl}, this results in negligible changes to the driving force for Cl^-. This contrasts with the results shown

earlier, where adjusting the activity of cation-chloride cotransport modulates both E_{Cl} and the driving

force for Cl^- substantially.

Figure 5: Adjusting the average charge of impermeant anions shifts the chloride reversal potential with

negligible effects on the driving force for chloride. (A) Decreasing the average charge of impermeant

anions from -0.85 (the default value) to -1 (orange, bottom panel), without changing the absolute

number of intracellular impermeant anions caused a persistent decrease in E_{Cl} (green, top panel), E_K

(purple) and V_m with moderate increases in volume (middle panel) in our default single compartment

model. Negligible changes in Cl$^-$ driving force (ΔDF = 0.16 mV, red) were observed. (B) Analytic solution

(solid lines) for different impermeant anion average charge (z) exactly matches the steady state values

from numerical, time series runs (dots) based on adjusting average z as in 'A'. Steady state E_{Cl}, E_K, V_m

(top panel) and $[X]_i$ (middle) increased with increasing average z, while changes in average z resulted

in very small changes in Cl$^-$ DF (bottom, red). The vertical dashed line indicates the values at the chosen

default average z of -0.85. (C) Influx of a species of impermeant anions with average charge -1.5, that

decreased average z (bottom panel) from -0.85 to -1 and increased the number of impermeant anions

also caused persistent decreases of E_{Cl}, E_K and V_m as in 'A', but with larger increases in cell volume

(middle panel). Again, very small persistent changes in Cl$^-$ DF were observed. The volume changes for

different flux amounts and charges are illustrated in Figure 5 – figure supplement 1. (D) Top, schematic

of the experimental setup where whole-cell recordings were made from CA3 pyramidal cells in mouse

organotypic brain slices. Impermeant anions (orange) were delivered via electroporation of the

negatively charged fluorescent dextran Alexa Flour 488 via a pipette positioned near the soma of the

recorded cell. GABA$_A$R currents were elicited via photo-activation (100 ms, 470 nm LED via objective)

of ChR2-expressing GAD2+ interneurons (green cells) in the presence of 5 μM CGP-35348 to block

GABA$_B$Rs. Lower trace, current clamp recording showing V_m changes during electroporation of anionic

dextran.   Confocal image demonstrating cell-localised fluorescence of the anionic dextran

electroporated in 'E'. (E) Top, widefield images with electroporation pipette (orange dashed lines) and

the recording pipette (white dashed lines). Note increased fluorescence in the soma after

electroporation. Below, insets show GABA$_A$R currents evoked by photo-activation of GAD2+

interneurons at different holding potentials. Calibration: 1 s, 100 pA. IV plots were used to calculate

V_m, E_{GABA} and DF before (left) and after (right) electroporation. (F) Top, population data showing

significant decreases in mean V_m and E_{GABA} but not DF five minutes after electroporation. Below:

changes in DF over time. Point at which electroporation occurred marked with orange arrow. 'ns',

non-significant; *p<0.05.

<u>Impermeant anions drive small shifts in chloride driving force by modifying the Na^+/K^+-ATPase pump</u>

<u>rate under conditions of active chloride cotransport.</u>

We next set out to determine how, and under what conditions, the modification of impermeant

anions could potentially generate the very small persistent shifts in Cl⁻ driving force we observed in

our models. Due to their small size (< 1 mV) these were not detectable during the experimental

validation. First, we repeated the simulation performed in Figure 5A by changing the average charge

of impermeant anions in the cell, but under conditions where the Na^+/K^+-ATPase effective pump rate

(J_p) was either a cubic function of the transmembrane Na^+ gradient (default condition) or was fixed at

a constant value (**Fig. 6A**). In the case where the pump rate was fixed, adjusting the average charge of

impermeant anions generated no persistent change in Cl⁻ driving force (**Fig. 6A**). Modifying

impermeant anions caused a significant change in steady-state intracellular Na^+ concentration when

J_p was kept constant. However, small shifts in Cl⁻ driving force occurred only when the effective pump

rate was variable, in which minor changes to $[Na^+]_i$ caused significant changes to J_p, which in turn

resulted in a small shift in Cl⁻ driving force. There is a direct relationship between the average charge

of impermeant anions (z), $[Na^+]_i$, the effective Na^+/K^+-ATPase pump rate and Cl⁻ driving force. This

relationship was abolished when the effective Na^+/K^+-ATPase pump rate was held constant by

removing its dependence on Na^+ (**Fig. 6B**). In addition, even large variations in effective pump rate

near the default value caused negligible shifts in Cl⁻ driving force of < 1 mV. These results were similar

when using the experimentally-matched ATPase model by Hamada et al. (2003), with slight

differences in final values (**Figure 6 – figure supplement 1A**). These small, impermeant anion driven,

Na^+/K^+-ATPase pump dependent shifts in Cl^- driving force are completely dependent on the presence

of cation-chloride cotransport. In the absence of KCC2, there is no Cl^- driving force as $E_{Cl} = V_m$ (**Fig. 6F**).

We then tested whether relaxing the condition of transmembrane osmoneutrality might also alter

impermeant anion induced effects on Cl^- driving force. We modelled a situation where increases in

cell surface area beyond a certain 'resting' surface area generated a hydrostatic pressure (membrane

tension), which could balance an osmotic pressure difference of 10 mM between the intra and

extracellular compartments (see **Fig. 6C**, schematic in **Fig. 6E** and Materials and Methods). In this case,

adding impermeant anions of default charge z resulted in constrained increases in cell volume, which

were accompanied by persistent transmembrane differences in osmolarity and intracellular Na^+

concentration. This was sufficient to generate small differences in driving force for Cl^- of < 0.2 mV for

reasonable increases in cell surface area (Nichol and Hutter, 1996; Dai et al., 1998). Again, this was

entirely due to Na^+ driven shifts in the Na^+/K^+-ATPase effective pump rate. By removing the

dependence of Na^+/K^+-ATPase activity on Na^+ concentration, addition of impermeant anions no longer

generated persistent shifts in Cl^- driving force (**Fig. 6C**). Because the model is also directly dependent

on $[Na^+]_i$, the results were similar when using the experimentally-matched ATPase model by Hamada

et al. (2003) (**Figure 6 – figure supplement 1B**). We observed a direct relationship between

transmembrane osmotic gradient, $[Na^+]_i$, the effective Na^+/K^+-ATPase pump rate and Cl^- driving force.

This relationship was removed when the effective Na^+/K^+-ATPase pump rate was held constant, with

no changes in Cl^- driving force seen despite the generation of the same shift in the transmembrane

osmotic gradient (**Fig. 6D**).

In summary, changes in Cl^- driving force generated by changing the ionic contributions to cellular

charge (by altering the average charge of impermeant anions) or osmoneutrality (by increasing the

contribution of hydrostatic pressure) are due to the alteration of the dynamics of active ion transport

mechanisms in the cell. However, these effects are negligible in magnitude and cannot contribute

 significantly to setting physiologically observed Cl⁻ driving forces. It is worth reiterating that any non-

 zero Cl⁻ driving force is entirely dependent on the presence of active Cl⁻ cotransport. In our model, in

 the absence of KCC2, neither the Na⁺/K⁺-ATPase nor impermeable anions can shift Cl⁻ out of

 equilibrium (**Fig. 6F**).

<u>Figure 6: Impermeant anions drive small shifts in chloride driving force by modifying the Na^+/K^+-</u>

<u>ATPase pump rate under conditions of active chloride cotransport.</u> *(A) E_{Cl} (green), E_K (purple) and V_m*

(black) (top panel), $[Na^+]_i$ (pink, middle panel) and average impermeant anion charge z (orange, lower

panel) over time in the default single compartment model. Changing the average z from -0.85 to -1

generated small, persistent Cl^- driving force (DF) shifts (arrows, red) only when the effective ATPase

pump rate (J_p) was variable (solid line) and not when J_p was kept constant (dashed line). (B) Solving

analytically across different values of z with either a variable J_p (solid lines) or a constant J_p (dashed

line), demonstrates the direct relationship between Na^+ (pink), effective pump rate (J_p) (top panel) and

DF (lower panel, red). (C) E_{Cl} (green), E_K (purple) and V_m (black) (top panel), $[Na^+]_i$ (pink, middle panel)

and cell volume (black, lower panel) over time in the default single compartment model. Impermeant

anions of the same average charge of the cell were added. A volume constraint was incorporated by

adding a hydrostatic force dependent on membrane tension (dashed lines), which resulted in an

impermeable anion induced transmembrane osmotic differential. This caused a small change in DF

when J_p was variable (dashed line), but not when J_p was held constant (dotted lines). (D) Solving

analytically across osmolarity differences demonstrates the direct relationship between Na^+ (pink),

effective pump rate (J_p) (top panel) and DF (lower panel). Note, the small changes in DF. (E) Schematic

explaining the mechanism through which impermeant anion induced cell swelling in the presence of

volume constraints (i.e. membrane tension) result in steady states with equal but non-zero osmotic

($osmo_P$) and hydrostatic pressures (H_P), causing transmembrane osmotic differences (t_1). This causes

small changes in Na^+, and hence J_p. (F) All Na^+/K^+-ATPase pump rate-related shifts in the DF require

KCC2 activity; in the absence of activity (dashes), no shifts in driving force can occur. In Figure 6 –

supplement figure 1, we show that the results in (A) and (C) are similar when an experimentally-

matched model of the Na^+/K^+-ATPase is used.

<u>Changes in cation-chloride cotransport activity generate local differences in chloride reversal and</u>

<u>driving force, which depend on cytoplasmic diffusion rates.</u>

An important functional question is how Cl$^-$ driving force might be modified at a local level within a

neuron. We considered local persistent changes of Cl$^-$ driving force for the case of active

transmembrane Cl$^-$ fluxes (**Fig. 7**) and impermeant anions (**Fig. 8**) by extending the single compartment

model described above into a multi-compartment model or 'virtual dendrite.' This dendrite was 100

μm in length and consisted of ten compartments, each of 10 μm length and 1 μm diameter. The

compartments contained the same mechanisms and default parameterization as the single

compartment model described above. Compartmental volume was changed by altering the radius,

while holding the length constant. In addition, all ions, except impermeant anions, could move

between compartments by electrodiffusion (**Fig. 7A** and Materials and Methods).

To explore the local effects of CCC activity, we increased g_{KCC2} from our default value of 20 μS/cm^2 to

600 μS/cm^2 in the second last compartment of the virtual dendrite exclusively. This resulted in a

persistent decrease in E_{Cl}, concurrent with a modest decrease in V_m, resulting in a permanent increase

in Cl$^-$ driving force and minimal change in compartment volume (**Fig. 7B**). The spatial precision of this

alteration depended strongly on the diffusion constant for Cl$^-$. With a Cl$^-$ diffusion constant of

2.03x10^{-7} dm^2.s^{-1}, these alterations spread widely through the virtual dendrite. For example, the

change in Cl$^-$ driving force was 4.8 mV in the furthermost compartment (90 μm apart) as compared to

5.9 mV in the compartment manipulated. When we decreased the Cl$^-$ diffusion constant by one order

of magnitude, the change in Cl$^-$ driving force was 7.3 mV in the compartment in which KCC2 was

adjusted, but only 1.8 mV in the furthermost compartment from the site of manipulation (**Fig. 7C**).

These findings suggest that local differences in cation-chloride cotransport activity can drive spatially

restricted differences in Cl$^-$ driving force under conditions of constrained Cl$^-$ diffusion, however under

conditions of typical ionic diffusion the effect of Cl$^-$ transport by KCC2 is relatively widespread.

Figure 7: Local changes in KCC2 activity generate local differences in chloride reversal and driving force only under conditions of constrained chloride diffusion. (A) Schematic depicting the multi-compartment model representing a virtual dendrite of length 100 μm. The virtual dendrite consists of 10 compartments of length 10 μm and initial radius 0.5 μm. Each compartment contains the same mechanisms and default parameterization as the single compartment model. All ions, except impermeant anions, could move between compartments by electrodiffusion. (B) Top panel, E_{Cl} (green), E_K (purple), V_m (black) and calculated DF (arrows, red) from the 2nd from top compartment (indicated with a white triangle) where the conductance of KCC2 was increased. The insets depict the diameter, and absolute change from baseline of E_{Cl}, V_m and DF for all compartments before (i), during (ii) and after (iii) the activity of KCC2 was selectively increased. This resulted in E_{Cl} decreasing in all compartments with minimal changes to E_K and V. Consequently, the Cl- DF (red) increased. In this case the diffusion constant for Cl- (D_{Cl} = 2.03 x 10^{-7} dm^2.s^{-1}) resulted in E_{Cl} and DF changes being widespread across the virtual dendrite. (C) Reducing the Cl- diffusion constant to 0.2 x 10^{-7} dm^2.s^{-1} resulted in a localized effect of compartment specific KCC2 activity increases on E_{Cl} and DF.

Local impermeant anions do not appreciably affect the local driving force for chloride.

Following the last result, we considered whether changing impermeant anions in part of a dendrite could create a local area with a different Cl- driving force compared to the rest of the cell. We first

added impermeant anions of average charge (z) exclusively to the second-last compartment of the

virtual dendrite while measuring the Cl$^-$ reversal, V_m and Cl$^-$ driving force in all compartments. During

addition of the impermeant anions, E_{Cl} and V_m decreased with an accompanying decrease in Cl$^-$ driving

force. However, following cessation of impermeant anion influx, all parameters returned to baseline

levels, except for the volume of that specific compartment, which showed a modest increase (**Fig. 8B**).

This suggests that local addition of impermeant anions of average charge has no local effect on Cl$^-$

homeostasis but can affect the volume of the compartment concerned.

Next, we again added impermeant anions to the second-last compartment of the virtual dendrite, but

this time we added impermeant anions with a more negative charge than average (z = -1). This

resulted in the average charge of impermeant anions in that specific compartment becoming more

negative (**Fig. 8C**). During the addition of the impermeant anions, E_{Cl} and V_m decreased across the

dendrite, but with small accompanying shifts in Cl$^-$ driving force. Following cessation of local

impermeant anion influx, a persistent shift in E_{Cl} and V_m was observed specifically in the compartment

manipulated. However, this generated a negligible, persistent change in Cl$^-$ driving force (< 0.01 mV

change in driving force for a compartment specific change in z from -0.85 to -0.93), and only within

that specific compartment of the virtual dendrite. Again, impermeant anion addition resulted in a

permanent increase in the volume of the compartment concerned. This finding suggests that local

impermeant anions can adjust the Cl$^-$ reversal potential locally, but are not well-placed to cause

significant, permanent shifts in the driving force for Cl$^-$. Indeed, electrodiffusion may further limit the

degree to which changes in impermeant ion charge can modify driving forces through alterations in

active ionic transport: the resulting permanent Cl$^-$ driving force changes in the multi-compartment

model are many times smaller than the shifts in the single compartment version (as compared to **Fig.**

6B).

Figure 8: Local changes in impermeant anions do not establish the local driving force for chloride. (A)

Multi-compartment model of a 10 compartment 100 μm virtual dendrite as in Figure 7. (B) Top panel,

E_{Cl} (green), E_K (purple), V_m (black) and calculated DF (arrows, red) within the compartment where

additional impermeant anions were exclusively added (indicated with a white triangle). The insets

depict the diameter, and absolute change from baseline for E_{Cl}, V_m and DF for each compartment of

the virtual dendrite before (i), during (ii) and after (iii) impermeant anions were added. The selective

addition of impermeant anions of average z to the 2nd from top compartment resulted in transient but

non-permanent shifts in E_{Cl}, V_m and the Cl^- DF in all compartments. The volume of the compartment

where impermeant anion addition occurred increased permanently. (C) Traces and insets as in 'B'

showing the addition of impermeant anions of different charge in order to decrease average z in the

2nd from top compartment specifically. Note that during addition of impermeant anions, E_{Cl}, E_K, V_m

changed. We also observed persistent decreases in E_{Cl} and V_m in the compartment manipulated, with

a negligible change in DF (0.01 mV). Again, impermeant anion addition resulted in an increase in the

volume of the specific dendritic compartment manipulated.

Discussion

The driving force for Cl⁻ is a fundamental parameter affecting the excitability of neuronal networks

(Raimondo et al., 2017). Recently, impermeant anions, rather than CCCs, have been suggested as the

primary determinants of the neuronal driving force for Cl⁻ (Glykys et al., 2014). Here we have explored

the determinants of the Cl⁻ driving force in neurons by deriving theoretical models based on

biophysical first principles. We show that the Na^+/K^+-ATPase, baseline K^+, Na^+ and Cl⁻ conductances,

average charge on impermeant anions, water permeability and CCCs, likely all play roles in setting

neuronal $[Cl^-]_i$. However, our findings suggest that while impermeant anions can contribute to setting

the $[Cl^-]_i$ in neurons, they can only affect Cl⁻ driving force by modifying the activity of active transport

mechanisms (i.e. the Na^+/K^+-ATPase). Our modelling and experimental data demonstrate that under

physiologically relevant conditions, impermeant anions do not alter the Cl⁻ driving force significantly.

In contrast, CCCs are well placed to modulate Cl⁻ driving force and hence inhibitory signaling.

Previous theoretical models, which account for the dynamics of Cl⁻ ions, have been useful in

determining how changes to the driving force for Cl⁻ are critical for controlling the effect of synaptic

inhibition in the brain (Qian and Sejnowski, 1990; Staley and Proctor, 1999; Doyon et al., 2011; Jedlicka

et al., 2011; Lewin et al., 2012; Mohapatra et al., 2016). Whilst these models have included the Na^+/K^+-

ATPase, the interacting dynamics of several ion species, CCCs (Doyon et al., 2011; Krishnan and

Bazhenov, 2011), electrodiffusion (Qian and Sejnowski, 1989) and impermeant anions and volume

regulation (Dijkstra et al., 2016), none have combined all of these mechanisms to explore how their

combination determines the local driving force for Cl⁻. Our theoretical approach is based on the pump-

leak formulation (Tosteson and Hoffman, 1960). It suggests that mammalian cells maintain their

volume under osmotic stress generated by impermeant anions and the Donnan effect by employing

active transport of Na^+ and K^+ using the Na^+/K^+-ATPase (Armstrong, 2003; Kay, 2017). A Donnan

equilibrium, a true thermodynamic equilibrium requiring no energy to maintain it, is not possible in

cells with pliant membranes like neurons (Sperelakis, 2012).

Our model conforms to the pump-leak formulation: abolishing the activity of the Na^+/K^+-ATPase leads

to cell swelling, progressive membrane depolarization and rundown of ionic gradients, including that

of Cl^-. Therefore, the Na^+/K^+ ATPase is a fundamental cellular parameter that stabilizes cell volume and

determines all ionic gradients including that of Cl^- and hence must be considered in any attempt to

model ion homeostasis. Interestingly however, we demonstrate that above a certain level of Na^+/K^+-

ATPase activity, even many fold changes in pump rate have minimal effects on volume, E_{Cl} and V_m.

This might explain recent experimental findings in which periods of Na^+/K^+ ATPase inhibition using

ouabain caused modest changes to cell volume (Glykys et al., 2014). It therefore seems unlikely that

neurons adjust the Na^+/K^+-ATPase as a means for modulating Cl^- driving force.

Baseline ion conductances are another important factor that affect Cl^- driving force. Our model is

consistent with recent experimental results that demonstrate that increased neuronal Na^+

conductance (for example by activating NMDA receptors, or preventing closure of voltage-gated Na^+

channels), leads to progressive neuronal swelling, membrane depolarization and Cl^- accumulation

(Rungta et al., 2015) – the primary pathological process in cytotoxic oedema (Liang et al., 2007). We

also show that tonic neuronal Cl^- conductance only affects baseline $[Cl^-]_i$ and driving force in the

presence of CCCs. Without active Cl^- flux, which CCCs provide, there is no driving force for passive Cl^-

flux and hence no mechanism for $[Cl^-]_i$ changes resulting from selective modification of a Cl^-

conductance. This is consistent with both classic (Misgeld et al., 1986; Thompson and Gähwiler, 1989)

and recent experimental findings (Berglund et al., 2016).

In our model, we find that elevating the activity of KCC2, the most active CCC in mature neurons (Ben-

Ari, 2002), increases the driving force for Cl^- by shifting the reversal potential for Cl^- closer to that of

K^+. Interestingly, large shifts (~7 mV) in driving force were associated with very minor (1%) changes in

volume or membrane potential. As such, modulating KCC2 represents a specific means for

manipulating the neuronal Cl^- driving force. This is consistent with traditional dogma, recent and

previous experimental results (Kaila et al., 2014; Klein et al., 2017) as well as our own experimental

validation using furosemide to block the activity of KCC2, which drove significant changes in driving

force with little effect on V_m. In further support of this, our meta-analysis of numerous experimental

studies showed a strong correlation between change in KCC2 expression and Cl^- driving force, but not

between KCC2 expression and V_m. There is an ongoing debate as to whether some cotransporters,

including CCCs, might also couple water transport to ion transport (Zeuthen, 1994; MacAulay et al.,

2002; Gagnon et al., 2004; Charron et al., 2006). Although we did not model the active movement of

water by KCC2, this scenario would not alter the central importance of CCCs for setting the Cl^- driving

force.

Using our multi-compartment model, which incorporated electrodiffusion, we found that local

modification of KCC2 activity has a specific local effect on Cl^- driving force that is dependent on the

characteristics of intracellular Cl^- diffusion. Cytoplasmic Cl^- diffusion rates had to be reduced

substantially before we observed local changes in Cl^- driving force driven by KCC2 (Qian and Sejnowski,

1989; Kuner and Augustine, 2000). Whilst differences in KCC2 activity might generate a gradient in Cl^-

driving force between large subcellular structures (i.e. dendrites versus soma), our modelling results

call into question the idea of synapse-specific regulation of Cl^- driving force within the same cellular

domain (Földy et al., 2010).

Glykys et al. (2014) used Cl^- imaging and various experimental manipulations to claim that $[X]_i$ and $[X]_o$

sets $[Cl^-]_i$ and the Cl^- driving force. From our theoretical analysis, we find that modifying the amount

of impermeant anions inside or outside neurons has no persistent effect on $[Cl^-]_i$ or Cl^- driving force,

unless we include a mechanism that allows a transmembrane osmotic pressure differential to develop

that indirectly affects active transport mechanisms. Even in this case, under transmembrane pressure

differentials that do not lyse the membrane (Nichol and Hutter, 1996), Cl^- driving force changes are

negligible (< 1 mV). Recently it has been suggested that the viscoelastic properties of the cellular

cytoskeleton could allow it to take up osmotic shifts created by impermeant anion movement like a

sponge (Sachs and Sivaselvan, 2015). This would mean that one would not see as large a volume shifts

as predicted by our models. In our model we have assumed that water can pass through the neuronal

membrane to equalize osmotic differences. Although it is thought that some neurons do not express

aquaporin channels (Andrew et al., 2007), water can permeate the phospholipid bilayer (Fettiplace

and Haydon, 1980). Therefore, whilst differences in neuronal water permeability might affect the time

taken to reach steady-state, the steady state values themselves are unlikely to be affected. We

conclude that $[Cl^-]_i$ and the Cl^- driving force are not determined by the *concentration* of impermeant

anions.

However, our theoretical findings offer a potential explanation for recent experimental observations.

We show that modifying the average *charge* upon impermeant anions (i.e. z in $[X^z]_i$), rather than their

concentration, can affect $[Cl^-]_i$ and E_{Cl^-}. Relating this to prior experimental observations, Glykys et al.

(2014) used SYTO64 staining of nucleic acids and perfusion of weak organic acids in conjunction with

Cl^- imaging to suggest that $[Cl^-]_i$ depends upon internal impermeant anions ($[X]_i$). If such a manipulation

modifies the average charge on internal impermeant anions, and not concentration *per se*, this could

account for the observed changes in $[Cl^-]_i$. Glykys et al. (2014) did not measure V_m or the Cl^- driving

force in these experiments. The clear prediction from our model is that any manipulation, which

changes the average charge of impermeant anions would not appreciably affect the Cl^- driving force

because any impermeant anion driven change on E_{Cl^-} is matched by an equivalent effect on V_m due to

accompanying shifts in cation concentrations. We have provided experimental support for this

prediction by showing that whilst E_{GABA} (and E_{Cl}) can be shifted by addition of impermeant anions using

electroporation of membrane impermeant anionic dextrans, V_m is shifted in a similar direction

resulting in an undetectable change in Cl^- driving force.

Given prior theoretical predictions (Kaila et al., 2014; Voipio et al., 2014; Savtchenko et al., 2017), it is

interesting that our model reveals that changing impermeant anions could affect the Cl^- driving force

at all. We found that the small (< 1 mV) impermeant anion-driven changes in Cl^- driving force observed

in our model were caused by indirect effects on Na^+ concentration and hence Na^+/K^+-ATPase activity.

The impermeant anion-driven changes in Cl^- driving force are even smaller in the multi-compartment

model (< 0.1 mV), in which electrodiffusion allows local changes in Na^+ to dissipate. When Na^+/K^+-

ATPase activity was decoupled from the transmembrane Na^+ gradient, we found that impermeant

anions were unable to cause persistent shifts in Cl^- driving force as predicted theoretically (Kaila et al.,

2014; Voipio et al., 2014; Savtchenko et al., 2017). It is important to note that these small, impermeant

anion-Na^+/K^+-ATPase-driven shifts in Cl^- driving force are dependent on the presence of cation-

chloride cotransport in the form of KCC2 and would entail changes in energy use by the Na^+/K^+-ATPase.

In other words, active transport mechanisms are again required to drive changes in Cl^- homeostasis.

In summary, our theoretical models, which are derived from well-established physical principles, are

consistent with our own experimental data and that of others (Glykys et al., 2014; Kaila et al., 2014;

Klein et al., 2017), and suggest that impermeant anions alone cannot shift Cl^- out of equilibrium across

the neuronal membrane. Were neurons to alter impermeant anion concentrations or charge, the

resting membrane potential would be modified with little effect on the Cl^- driving force. Our work

confirms the central importance of CCC activity in determining the effects of inhibitory synaptic

transmission in the nervous system.

Materials and methods

Single compartment model. The single compartment model consisted of a cylindrical semi-permeable

membrane separating the extracellular solution from the intracellular milieu with variable volume

(**Fig. 1A**). The extracellular ionic concentrations were assumed constant (**Table 1**). Permeable ions in

the model were K^+, Na^+ and Cl^- with their usual charges, while impermeant anions X were assumed to

be a heterogeneous group of impermeant chemical species with average intracellular charge z and

extracellular charge -1. z was chosen on the basis of known resting intracellular ion concentrations

(Lodish et al., 2009; Raimondo et al., 2015) and osmolarity (Π). Bicarbonate ions were not included in

our model as a permeant anion as they were assumed to be important for acute depolarizing effects

(via GABA$_A$Rs) rather than the chronic shifts in Cl^- driving force, which are the focus of this work (Staley

and Proctor, 1999). The model included ionic leak currents for the permeable ions, Na^+/K^+-ATPase

transporters and a CCC, in this case the K^+-Cl^- cotransporter (KCC2). KCC2 and not NKCC1 is thought to

be the most active CCC in mature neurons (Ben-Ari, 2002), therefore, to maintain conceptual simplicity

only KCC2 was modelled. Cell volume (w) change was based on osmotic water flux and incorporated

a membrane surface area scaling mechanism. An analytical solution to the model at steady state was

derived using standard techniques and can be found in Supplementary file 1. The numerical model

was initialised assuming conditions close to electroneutrality and an osmotic equilibrium between the

intracellular and extracellular compartments. A forward Euler approach was used to update variables

at each time step (dt) of 1 ms. Using a smaller dt did not influence the results in Figure 1-5. Code was

written in Python 2 and is available on GitHub (https://github.com/kiradust/model-of-neuronal-

[chloride-homeostasis](https://github.com/kiradust/model-of-neuronal-chloride-homeostasis)). The GitHub repository includes a file of Supplementary figures, in which we

display transmembrane fluxes of all ions and water for relevant simulations. An example figure

displaying ionic flux for all ions is available for Figure 4C in Figure 4 – figure supplement 1.

Membrane potential. The membrane potential V_m was based on the "Charge Difference" approach of

Fraser and Huang (2004) as follows:

$$V_m = \frac{F\,([Na^+]_i + [K^+]_i - [Cl^-]_i + z[X^z]_i)}{C_m A_m} \tag{1}$$

where F is Faraday's constant, C_m is the unit membrane capacitance and A_m is calculated as the ratio

of the surface area (of the cylinder) to cell volume. The term in brackets is the sum of all ionic charges

within the cell. This approach has the advantage that the initial voltage can be calculated without

needing to assume a steady state as is required for by the Goldman-Hodgkin-Katz (GHK) equation.

Permeable ion concentrations. Intracellular concentrations of the permeable ions Na$^+$, K$^+$ and Cl$^-$ were

updated individually by summing trans-membrane fluxes. Leak currents were calculated using the

standard equivalent circuit formulation $I = g(V_m - E)$. In addition, Na$^+$ and K$^+$ were transported

actively by the Na$^+$/K$^+$-ATPase, with pump rate J_p, which was approximated by a cubic function

dependent on the transmembrane sodium gradient, following Keener and Sneyd (1998):

$$J_p = P\left(\frac{[Na^+]_i}{[Na^+]_o}\right)^3, \tag{2}$$

where P is the pump rate constant. Because it is a function of the sodium gradient, J_p decreases as

[Na$^+$]$_i$ depletes. This formulation has been shown to be similar to more accurate kinetic models reliant

on both the Na$^+$ gradient and ATP concentration (Keener and Sneyd, 1998). To switch the ATPase

pump on or off (**Fig. 1C**), P was decreased/increased exponentially over 10-20 minutes, consistent

with previous reports of the dynamics of inhibition of the ATPase by ouabain and in turn the inhibition

of ouabain's effects by potassium canrenoate (Baker and Willis, 1972; Yeh and Lazzara, 1973). The

ATPase pumps 2 K$^+$ ions into the cell for every 3 Na$^+$ ions out and these constants must be multiplied

by J_p for each ion respectively. K$^+$ and Cl$^-$ were also modified by flux through the type 2 K-Cl

cotransporter (KCC2), J_{KCC2} (Doyon et al., 2016):

$$J_{KCC2} = g_{KCC2}(E_K - E_{Cl}),\qquad(3)$$

767   where g_{KCC2} is a fixed conductance and E_K and E_{Cl} are the Nernst potentials for K^+ and Cl^- respectively.

768   J_{KCC2} is 0 when $E_K = E_{Cl}$. The rate of change of the intracellular concentration of the three permeant ions

769   was given by the following equations, with the Nernst potentials for each ion given by $E_{ion} =$

770    $\frac{RT}{zF}\ln\left(\frac{[ion]_{out}}{[ion]_{in}}\right)$, w indicating the cell volume, and $\frac{dw}{dt}$ as described in equation (7) (**Fig. 1-5, 6A, B, 7 and**

771   **8**) or (9) (**Fig. 6C-E**):

$$\frac{d[Na^+]_i}{dt} = -\frac{A_m}{F}\left(g_{Na}(V_m - E_{Na}) + 3J_p\right) - \frac{1}{w}\frac{dw}{dt}[Na^+]_i \qquad(4)$$

$$\frac{d[K^+]_i}{dt} = -\frac{A_m}{F}\left(g_K(V_m - E_K) - 2J_p - J_{KCC2}\right) - \frac{1}{w}\frac{dw}{dt}[K^+]_i \qquad(5)$$

$$\frac{d[Cl^-]_i}{dt} = \frac{A_m}{F}\left(g_{Cl}(V_m - E_{Cl}) + J_{KCC2}\right) - \frac{1}{w}\frac{dw}{dt}[Cl^-]_i. \qquad(6)$$

Volume. Because the osmotic flux of water is expected to be faster than ion fluxes, the volume of the

cell was adjusted to preserve transmembrane osmotic balance at each time step. Change in

compartment volume, w, was calculated by changing the previous volume proportional to the

difference between intracellular and extracellular osmolarity, Π_i and Π_o respectively (Hernández and

Cristina, 1998), where v_w is the partial molar volume of water, p_w the osmotic permeability of a

biological membrane and SA the surface area:

$$\frac{dw}{dt} = v_w \cdot p_w \cdot SA \cdot (\Pi_i - \Pi_o). \tag{7}$$

786 For the calculations in Figure 6C-E, where we allowed transmembrane differences in osmolarity to

787 develop, we assumed that at rest the cylindrical cell had a radius of r_a and zero pressure across the

788 membrane, and that the tension (T) in the membrane followed Hooke's law such that the tension was

789 proportional to the difference between the dynamic circumference of the cell and that of the resting

790 state. From Laplace's law the hydrostatic pressure in the cell was given by:

$$H_p = \begin{cases} 4\pi k_m \left(1 - \dfrac{r_a}{r}\right) & r > r_a \\ 0 & , \ otherwise \end{cases} \tag{8}$$

793 where k_m is the spring constant of the membrane (Sachs and Sivaselvan, 2015). Equation (7) was thus

794 reformulated:

$$\frac{dw}{dt} = v_w \cdot p_w \cdot SA \cdot \left(\Pi_i - \Pi_o - \frac{H_p}{RT}\right). \tag{9}$$

In order to simulate extreme conditions of constrained volume, a larger k_m was employed than is

realistic (Dai et al., 1998). Intracellular ion concentrations were updated again after volume change at

each time step. Volume changes were manifested in the cylindrical compartment as change in the

radius. In Figure 1-6, the cell was initialised with diameter 10 μm and length 25 μm.

Anion flux. Impermeant anions were manipulated in the compartment in Figure 4-6 and Figure 8

through several mechanisms. Anions could be added to the compartment at a constant rate

dependent on A_m and have either the same average intracellular X charge z = -0.85 (**Fig. 4C, 6C and**

8B), or a different charge (**Fig. 5C, 6A and 8C**). In these cases, the number of moles of X in the

compartment was increased. Alternatively, the charge of a species of intracellular X was slowly

changed imitating a charge-carrying trans-membrane reaction (**Fig. 5A and 6A**). In this case, the

number of moles of intracellular X did not change and it was assumed charge imbalance was mopped

up by the extracellular milieu. Finally, extracellular X⁻ was changed in Figure 4D by removing as much

Cl⁻ as X⁻ was added, thus maintaining osmolarity and electroneutrality in the extracellular space.

Multi-compartment model. The single compartment dendrite model was incorporated in a multi-

compartment model by allowing electrodiffusion to occur between individual compartments

operating as described above. Compartments were initialised with a radius of 0.5 μm and length of 10

μm. Compartments were linked linearly without branching; 10 connected compartments in total were

used. The time step dt was decreased to 10^{-3} ms for simulations in multiple compartments. Code was

written in Python 3 and is available on GitHub (https://github.com/kiradust/model-of-neuronal-

chloride-homeostasis).

Electrodiffusion. The Nernst-Planck equation (NPE) was used to model one-dimensional

electrodiffusion, based on Qian and Sejnowski (1989). The NPE incorporates fluxes because of

diffusion and drift (i.e. the movement of ions driven by an electric field). It has been shown to be more

accurate than using $J_{diffusion}$ alone in small structures like dendrites (Qian and Sejnowski, 1989). The

NPE for J the flux density of ion C is calculated as:

$$J = -D\frac{zF^2}{RT}[C]\frac{dV_m}{dx} - DF\frac{d[C]}{dx}, \tag{10}$$

where D is the diffusion constant of ion C (**Table 1**), z is its charge, [C] is its concentration and x is the

distance along the longitudinal axis over which electrodiffusion occurs. The NPE was implemented

between compartments i and i+1, assuming the i→i+1 direction was positive, using a forward Euler

approach. The midpoints of the compartments were used to calculate dx, i.e. $dx = \frac{h_i + h_{i+1}}{2}$, and the

concentrations of C in each compartment were averaged to obtain J_{drift}, ensuring that $J_{i \to i+1} = J_{i+1 \to i}$,

where the fluxes had units of mol/(s.dm^2):

$$J_{i \to i+1} = -D \left(\frac{zF}{RT} \frac{([C]_i + [C]_{i+1})}{2} \frac{dV_m}{dx} + \frac{d[C]}{dx} \right). \qquad (11)$$

835 The flux was multiplied by the surface area between compartments and then divided by compartment

836 volume to determine the flux in terms of molar concentration (M/s), i.e. $\frac{\pi r^2}{\pi r^2 h_i} = \frac{1}{h_i}$, and finally

837 implemented numerically with a forward Euler approach. The implementation mirrored that in Qian

838 and Sejnowski (1989) for non-branching dendrites, but was adjusted for compartments whose

839 volumes can change:

$$C_{i \to i+1} = -\frac{dt}{h_i} D \left(\frac{zF}{RT} \frac{([C]_i + [C]_{i+1})}{2} \frac{(V_{m_i} - V_{m_{i+1}})}{dx} + \frac{([C]_i - [C]_{i+1})}{dx} \right). \qquad (12)$$

Systematic review. A literature search was performed to identify experimental studies that aimed to

correlate a change in KCC2 expression with changes in [Cl$^-$]$_i$. The MEDLINE database was used and

accessed via the PubMed online platform. Search terms included 'chloride', 'Cl', 'intracellular', 'KCC2',

'cotransporter', 'neuronal', 'GABA' using appropriate Boolean operators. All 26 studies that

demonstrated changes in KCC2 expression and E$_{GABA}$ were considered for the meta-analysis. As there

is a well described differential expression of KCC2 and NKCC1 at different stages of development, with

KCC2 expression increasing and NKCC1 expression decreasing across development, only studies that

used tissue older than postnatal day 7 were included (8 included data from younger animals). Other

exclusion criteria included: reporting a significant change in NKCC1 (5 studies); use of non-rodent

tissue (2 studies); no quantification of the change in KCC2 (2 studies). Nine experiments from eight

studies met all criteria and were included (Coull et al., 2003; Lagostena et al., 2010; Lee et al., 2011;

Ferrini et al., 2013; Campbell et al., 2015; MacKenzie and Maguire, 2015; Mahadevan et al., 2015;

Tang et al., 2015). However, one study did not report the change in V_m and hence was excluded from

the figure (Mahadevan et al., 2015). Data used in regression can be seen in Supplementary file 2 (Table

S2-1) and includes the analysis for regression against change in $[Cl^-]_i$. To accommodate varied

experimental preparations and techniques influencing data quality and biases, a 34-point scoring

system was designed to weight the studies (Table S2-2). A weighted least squares regression model

was then used to correlate the percentage (%) change in KCC2 expression versus change in Cl^- driving

force.

	Value	Description
Constants		
F	96485.33 C/mol	Faraday constant
R	8.31446 C/V	Universal gas constant
T	310.15 K	Absolute temperature (=37°C)
Parameters		
C_m	2×10^{-6} F/cm^2	Unit membrane capacitance (Qian and Sejnowski, 1989)
g_{Na}	20 µS/cm^2	Na$^+$ leak conductance (Kager et al., 2000)
g_K	70 µS/cm^2	K$^+$ leak conductance (Kager et al., 2000)
g_{Cl}	20 µS/cm^2	Cl$^-$ leak conductance
g_{KCC2}	20 µS/cm^2	KCC2 conductance (Doyon et al., 2016)
v_w	0.018 dm^3/mol	Partial molar volume of water (Hernández and Cristina, 1998)
p_w	0.0015 dm/s	Osmotic permeability (Hernández and Cristina, 1998)
k_m	25 N/dm	Variable for membrane tension (higher than reported (Dai et al., 1998): used in this paper to accentuate differences in osmolarity)
p	0.1 C/(dm^2.s)	Default pump rate constant
$[Na^+]_o$	145 mM	Extracellular Na$^+$ concentration
$[K^+]_o$	3.5 mM	Extracellular K$^+$ concentration
$[Cl^-]_o$	119 mM	Extracellular Cl$^-$ concentration

[X$^-$]$_o$	29.5 mM	Extracellular impermeant anion (X) concentration (Lodish et al., 2009; Raimondo et al., 2015)
D$_{Na}$	1.33 × 10^{-7} dm^2/s	Na$^+$ diffusion constant (Hille, 2001)
D$_K$	1.96 × 10^{-7} dm^2/s	K$^+$ diffusion constant (Hille, 2001)
D$_{Cl}$	2.03 × 10^{-7} dm^2/s	Cl$^-$ diffusion constant (Hille, 2001)
Variables (default steady state)		
V$_m$	-72.6 mV	Membrane potential
[Na$^+$]$_i$	14.0 mM	Intracellular Na$^+$ concentration
[K$^+$]$_i$	122.9 mM	Intracellular K$^+$ concentration
[Cl$^-$]$_i$	5.2 mM	Intracellular Cl$^-$ concentration
[X^z]$_i$	154.9 mM	Intracellular impermeant anion (X) concentration
z	-0.85	Average charge of intracellular X (Lodish et al., 2009; Raimondo et al., 2015)
w	2.0 pL (single compartment) 0.078 pL (multi-compartment)	Volume

Table 1: Constants, default parameters and usual steady state values for variables used in the

biophysical models.

Slice preparation and electrophysiology. For all experiments, organotypic slices were prepared from

rodent brain tissue. Wistar rats were used for the experiments testing the effects of CCC blockade.

However, to allow for optogenetic manipulation, a crossed mouse strain on a C57BL/6 background

was used. This mouse strain was a cross between mice expressing cre-recombinase in glutamic acid

decarboxylase 2 (GAD2) positive interneurons (GAD2-IRES-cre, Jax Lab strain 010802) and mice with a

loxP-flanked STOP cassette preventing transcription of the red-fluorescent protein tdTomato (a cre-

reporter strain, Ai4, Jax Lab strain 007914.). This created the GAD2-cre-tdTomato strain which resulted

in cre-recombinase and tdTomato expression in all GABAergic interneurons (Taniguchi et al., 2011).

The use of all animals was approved by either the University of Oxford (rat) or the University of Cape

Town (mouse) animal ethics committees.

Organotypic brain slices were prepared using 7 day old Wistar rats (CCC experiment) or crossed GAD2-

IRES-cre (Jax Lab strain 010802) and Ai14 tdtomato reporter (Jax Lab strain 007914) mice

(impermeant anion experiment) and followed the protocol originally described by (Stoppini et al.,

1991). Briefly, brains were extracted and swiftly placed in cold (4°C) dissection media consisting of

Gey's Balanced Salt Solution (GBSS #G9779, Sigma-Aldrich, USA) supplemented with D-glucose

(#G5767, Sigma-Aldrich, USA). The hemispheres were separated and individual hippocampi were

removed and immediately cut into 350 μm slices using a McIlwain tissue chopper (Mickle, UK). Cold

dissection media was used to rinse the slices before placing them onto Millicell-CM membranes

(#Z354988, Sigma-Aldrich, USA). Slices were maintained in culture medium consisting of 25% (vol/vol)

Earle's balanced salt solution (#E2888, Sigma-Aldrich, USA); 49% (vol/vol) minimum essential medium

(#M2279, Sigma-Aldrich, USA); 25% (vol/vol) heat-inactivated horse serum (#H1138, Sigma-Aldrich,

USA); 1% (vol/vol) B27 (#17504044, Invitrogen, Life Technologies, USA) and 6.2 g/l D-glucose. Slices

were incubated in a 5% carbon dioxide (CO_2) humidified incubator at between 35 - 37°C. Recordings

were made after 6-14 days in culture. Previous studies have shown that after 7 days in culture

(equivalent to postnatal day 14) GABAergic signaling and E_{GABA} has sufficiently developed to a level

resembling mature nervous tissue (Streit et al., 1989; Wright et al., 2017). For the impermeant anion

experiment, the mouse organotypic brain slices were injected one day after culture with adeno-

associated vector serotype 1 (AAV1) containing a double-floxed sequence for channelrhodopsin

(ChR2) linked to a yellow fluorescent protein (YFP) tag driven by the elongation factor 1 promoter

(UNC Vector Core, USA). The vector was diffusely injected into slices using a custom-built Openspritzer

pressurized ejection device (Forman et al., 2017). Slices were left for 6 days in culture to allow for

robust expression of ChR2-YFP in GAD2+ interneurons.

For recordings, slices were transferred to a submerged chamber which was perfused with artificial

cerebro-spinal fluid (aCSF) bubbled with carbogen gas (95% oxgen:5% carbon dioxide). aCSF was

composed of (in mM): NaCl (120); KCl (3); $MgCl_2$ (2); $CaCl_2$ (2); NaH_2PO_4 (1.2); $NaHCO_3$ (23); D-

Glucose (11) with pH adjusted to be between 7.35-7.40 using 0.1 mM NaOH. Neurons were visualized

using a 20X or 60x water-immersion objective (Olympus, Japan) on a BX51WI upright microscope

(Olympus, Japan). Widefield images were obtained using a Mightex CCE-B013-U CCD camera. For the

optogenetic experiments, we used epifluorescence microscopy to select slices that exhibited strong

ChR2-YFP expression throughout the hippocampal region. Micropipettes were prepared from

borosilicate glass capillaries with an outer diameter of 1.20mm and inner diameter of 0.69mm

(Warner Instruments, USA), using a horizontal puller (Sutter, USA). Recordings were made using

Axopatch 1D and Axopatch 200B amplifiers and acquired using pClamp9 (Molecular Devices) or

WinWCP software (University of Strathclyde, UK).

For the CCC blockade experiment gramicidin perforated patch recordings (Kyrozis and Reichling, 1995)

were performed using glass pipettes containing (in mM): 135 KCl, 4 Na_2ATP, 0.3 Na_3GTP, 2 $MgCl_2$, 10

HEPES and 80 µg/ml gramicidin (Calbiochem; pH 7.38; osmolarity, 290 mosmol/l). After obtaining a

cell-attached patch, the gramicidin perforation process was evaluated by continuously monitoring the

decrease in access resistance. Recordings were started once the access resistance had stabilized

between 20-80 MΩ, which usually occurred 20-40 min following gigaseal formation. Rupture of the

gramicidin patch, referred to as a break through, induces a large influx of the high Cl^- internal solution

into the cell. This causes a significant and permanent increase in the E_{GABA} at which time recordings

would be discarded. $GABA_AR$ activation was achieved by pressure application of muscimol (10 µM,

Tocris, UK), a selective $GABA_AR$ agonist, via a picrospritzer.   To calculate $E_{GABA,}$ $GABA_AR$ currents were

elicited at different command voltages. These were a series of 10 mV steps above and below -60 mV.

Reported membrane potentials were corrected for the voltage drop across the series resistance for

each neuron. Holding current (reflecting membrane current) and total current (reflecting membrane

current plus the $GABA_AR$-evoked current) were plotted against the corrected holding potential to

generate a current-voltage (I-V) curve. Using this graph, the E_{GABA} was defined as the potential where

the total current equals the holding current. Some of the data used for calculating E_{GABA} values in

Figure 3E was also used in a previous study (Wright et al., 2017). V_m was defined as the x-intercept of

the holding current, and the driving force as the difference between the two. To block CCC function,

the KCC2 blocker furosemide (1mM, Sigma, USA) was applied.

For the impermeant anion experiments, whole-cell recordings were utilized as electroporation

consistently ruptured the gramicidin patches. Pipettes were filled with internal solution composed of

(in mM): Kgluconate (120); KCl (10); Na2ATP (4); NaGTP (0.3); Na2-phosphocreatinine (10) and HEPES

(10). To test the effect of introducing impermeant anions, we electroporated the anionic 10 000 MW

Dextran Alexa-Flour 488 (Thermo Fisher, USA) using a separate pipette positioned near the soma of

the patched cell. This molecule is a hydrophilic polysaccharide, which is both membrane impermeant

and highly negatively charged. Pipettes were filled with a 5% dextran solution in phosphate buffered

saline and voltage pulses (5-10 of 20 ms duration, 0.5-1 V) were applied using a stimulus isolator.

Successful electroporation of the anionic dextran was confirmed visually by observing the cell fill with

the fluorescent dye. Electroporation also resulted in immediate membrane depolarization, which

recovered over a period of 1-5 mins. E_{GABA}, Vm and driving force were calculated during voltage steps

in voltage clamp mode (as above) with GABA$_A$Rs activated via endogenous synaptic release of GABA

using photo-activation of GAD2+ interneurons expressing ChR2-YFP with 100 ms pulses of blue light

using a 470 nm LED (Thorlabs), in the presence of 5 µM CGP-35348 (Torcis Bioscience, UK) to block

GABA$_B$R activation. E_{GABA}, Vm and driving force were calculated before and atleast 5 mins following

electroporation to allow for stabilization of Vm, E_{GABA} and driving force. The image of the dextran filled

cell was acquired using a confocal microscope (LSM510 Meta NLO, Car Zeiss, Jena, Germany). Analysis

was performed using custom written scripts in the MATLAB environment (MathWorks). Numeric

results are presented as mean +/- SEM.

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