

Running headline: Spatial clustering and ecosystem resilience

Clustering and correlations: Inferring resilience from spatial patterns in ecosystems

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1 I. ABSTRACT

2 1. In diverse ecosystems, organisms cluster together in such a manner that the
3 frequency distribution of cluster sizes is a power-law function. Spatially-explicit
4 models of ecosystems suggest that *loss of such power-law clustering* may indicate
5 loss of ecosystem resilience. Hence, it is hypothesised that spatial clustering
6 properties in ecosystems - which can be readily measured using remotely-sensed
7 high-resolution data - can help infer proximity to ecosystem thresholds and may
8 even provide early warning signals of ecosystem collapse. Recent empirical and
9 simulation studies, however, don't find consistent relationships between spatial
10 clustering and ecosystem resilience. Furthermore, how spatial clustering metrics
11 relate to other well-known early warning signals of ecosystems collapse, specifi-
12 cally the phenomenon of *critical slowing down* (CSD), remains unclear.

13 2. We synthesize the literature on cluster sizes in empirical and theoretical studies
14 that show how local interactions (especially, positive feedback) among organisms
15 can cause power-law clustering. In addition, we analyse a minimal spatial model
16 of ecosystem transitions that allows us to disentangle the role of environmental
17 stressor and positive feedback on spatial patterns and ecosystem resilience.

18 3. Our literature synthesis reveals that empirically observed power-law cluster-
19 ing in ecosystems is parsimoniously explained by local positive feedback. Our
20 synthesis together with model analysis demonstrates that, depending on the
21 strength of positive feedback, emergence of power-law clustering can occur at
22 any distance from the critical threshold of ecosystem collapse. In fact, we find
23 that for systems with strong positive feedbacks, which are most likely to exhibit
24 abrupt transitions, loss of power-law clustering may not even occur prior to
25 ecosystem thresholds. We also argue that cluster-size distributions are unrelated
26 to the phenomenon of CSD.

27 4. We demonstrate that, due to CSD, a power-law feature does occur near critical
28 thresholds but in a different quantity; specifically, a power-law decay of spatial
29 correlations of ecosystem state.

30 5. We conclude that loss of power-law clustering cannot be used as a reliable
31 indicator of ecosystem resilience. Our synthesis and model analyses highlights
32 links between local positive feedback, emergent spatial properties and how they
33 may be used to interpret ecosystem resilience.

34 II. INTRODUCTION

35 Desertification of semi-arid ecosystems ([van de Koppel et al. 2002](#)), eutrophica-
36 tion of lakes ([Carpenter et al. 1999](#)), spread of diseases ([Chaves et al. 2012](#)), inva-
37 sion ([Hansen et al. 2013](#)) and community shifts in coral reefs ([Knowlton 2004](#)) are
38 some examples of state transitions or regime shifts in ecological systems. Some
39 of these transitions can be abrupt and irreversible, leading to catastrophic loss
40 of wildlife, habitats, and ecosystem services. Such transitions are also known as
41 critical transitions in the ecology literature. They happen when a system crosses
42 a certain threshold, called critical threshold, of environmental conditions. Over
43 the last decade, several studies have devised and validated methods to detect
44 the vulnerability of ecosystems to transitions ([Carpenter et al. 2011](#), [Dakos et al.](#)
45 [2012](#), [2011](#), [Eby et al. 2017](#), [Guttal and Jayaprakash 2009](#), [Kéfi et al. 2014](#), [Kéfi et al.](#)
46 [2007](#), [Scheffer et al. 2009](#)). One such method is based on the idea that patterns of
47 self-organisation in ecosystems can offer signatures of resilience ([Kéfi et al. 2014](#),
48 [Kéfi et al. 2007](#), [Rietkerk et al. 2004](#), [von Hardenberg et al. 2001](#)). Self-organised
49 patterns themselves often result from an interplay of facilitative and competitive
50 interactions among organisms ([Manor and Shnerb 2008](#), [Scanlon et al. 2007](#), [von](#)
51 [Hardenberg et al. 2010](#)). Therefore, a comprehensive understanding of how lo-
52 cal interactions between organisms scale to their spatial distribution and affect
53 ecosystem resilience, is of broad ecological interest.

54 Of many varieties of self-organization found in nature ([D’Odorico et al. 2012](#),
55 [Kéfi et al. 2007](#), [Rietkerk and van de Koppel 2008](#), [Scanlon et al. 2007](#)), we focus
56 on spatial patterns where organisms exhibit clustering of irregular size and shape
57 (see Glossary); these are found in many ecosystems such as semi-arid ecosystems,
58 mussel beds or seagrass (Fig 1). Here, the frequency distributions of these cluster-
59 sizes may follow a power-law function (henceforth referred to as *power-law clus-*
60 *tering*). These are interesting because they may imply that systems lack charac-
61 teristic size/shape (see Box 1 for a summary of properties of power-laws). Some
62 simulation and empirical studies suggest that when ecosystems are stressed, clus-
63 ters fragment leading to loss of large patches ([Kéfi et al. 2014](#), [Kéfi et al. 2007](#)). This
64 results in a qualitative change in the properties of cluster sizes, from a power-law
65 to an exponential distribution. The progressive truncation of the tail of the power-
66 law clustering has, therefore, been hypothesised to represent loss of resilience in

67 ecosystems ([Fernández and Fort 2009](#), [Kéfi et al. 2014](#), [Kéfi et al. 2007](#), [Kéfi et al.](#)
68 [2011](#), [Lin et al. 2010](#), [Weerman et al. 2012](#)).

69 Empirical evidence for this hypothesis, however, is ambiguous ([Maestre and](#)
70 [Escudero 2009](#), [Meloni et al. 2017b](#), [Weerman et al. 2012](#)). Additionally, simu-
71 lation studies in more complex models suggest that details of systems matter,
72 thus questioning the generality of these cluster based indicators ([Génin et al.](#)
73 [2018b](#), [Schneider and Kéfi 2016](#)). Nevertheless, the possibility of inferring ecosys-
74 tem resilience from a single snapshot and the increasing availability of low-cost
75 remotely-sensed spatial datasets, where these methods can be applied, is attrac-
76 tive. Therefore, an evaluation of the generality and robustness of clustering prop-
77 erties as a signature of ecosystem resilience is needed.

78 To understand issues of generality, we must discuss another class of power-
79 law behaviours that are considered universal features near/at *critical points* of
80 phase transitions. Here, we emphasize that the theoretical underpinnings of
81 ecosystem dynamics and indicators of stability are based on principles derived
82 from the theory of phase transitions and bifurcations ([Scheffer et al. 2009](#), [Stro-](#)
83 [gatz et al. 1994](#)). This theory predicts that as a system nears a critical point of
84 phase transitions, it takes increasingly longer to recover from perturbations. This
85 phenomenon of slowed recovery is called *critical slowing down* (CSD) in the con-
86 text of continuous phase transitions in the physics literature. However, a similar
87 effect of slowed recovery appears even in ecological models that show abrupt
88 transitions ([Scheffer et al. 2009](#), [Strogatz et al. 1994](#), [Wissel 1984](#)). Consequently,
89 CSD has been widely used to devise methods to detect the approach of critical
90 thresholds in ecosystems ([Scheffer et al. 2009](#), [Wissel 1984](#)). An aspect of CSD
91 that is much less known in the ecology literature is that close to, and at the crit-
92 ical point, the strength of perturbation decays as a power-law function of time
93 - indicating a very slow recovery ([Ma 2000](#), [Sethna 2006](#), [Stanley 1999](#)); this is
94 in contrast to systems far away from thresholds where perturbations decay ex-
95 ponentially fast. In fact, many power-law behaviours arise near/at continuous
96 phase transitions ([Ma 2000](#), [Sethna 2006](#), [Stanley 1999](#)).

97 We highlight an interesting contrast between the two power-law relationships
98 we have discussed thus far: While the power-laws associated with CSD are ex-
99 pected to *emerge near/at critical points* of phase transitions ([Ma 2000](#), [Sethna 2006](#)),

100 the power-laws in clustering are hypothesised to be *lost near/at critical thresh-*
101 *olds* of ecosystem collapse (Kéfi et al. 2007). It is now fairly well established
102 that many mechanisms cause emergence of power-laws even away from critical
103 thresholds (Newman 2005, Pascual and Guichard 2005, Roy et al. 2003). However,
104 the theoretical basis for why a loss of power-law clustering can indicate approach
105 to a critical threshold in ecosystem models is unclear. Furthermore, elucidating
106 relationships (if any) between the dynamical phenomenon of CSD and cluster
107 size properties, has not gained attention in the literature. Such an exercise will
108 not only prove helpful in evaluating the generality of ecosystem resilience in-
109 dicators but also reveal the crucial role of local positive feedback in ecosystem
110 patterning.

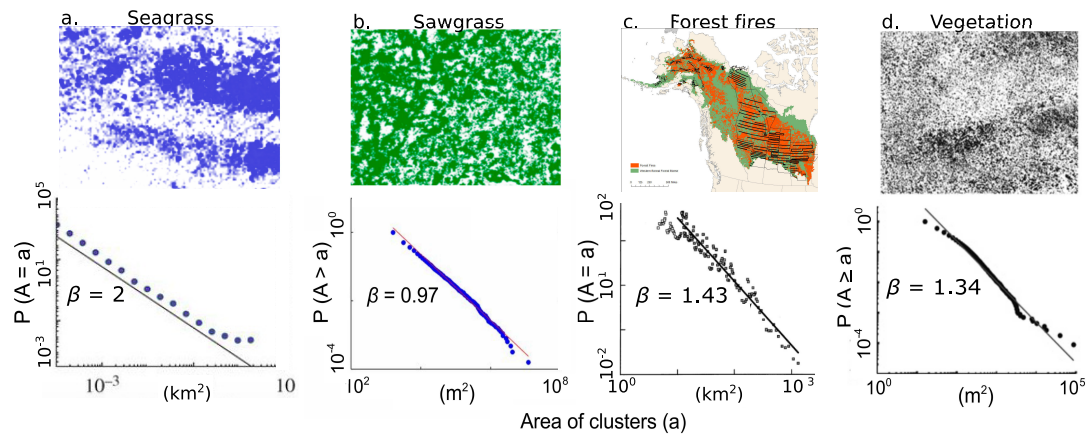


Figure 1: Power-law cluster size distributions in different ecosystems (top panel) and representative snapshots which are not necessarily from the same study area or time period (bottom panel). (a) West broad ledges seagrass near the isle of scilly (Irvine et al. 2016) (b) Saw-grass in everglades wetlands, USA (Foti et al. 2013) (c) Forest fires in Alaskan boreal forests, USA (1990-91) (Malamud et al. 1998) and (d) Vegetation in Kalahari, Namibia (Scanlon et al. 2007). Top row image credits: (a) modified from (Duffy et al. 2017), (b) <https://doi.org/10.1016/j.ecss.2017.11.001> [CC BY] <http://creativecommons.org/licenses/by/4.0/>, (Foti et al. 2013), (c) U.S. Geological Survey, Department of the Interior/USGS U.S. Geological Survey Map created by Tyler Lewis/USGS. (d) Data (1995-2014) (Scanlon et al. 2007).

111 In this article, we review and synthesize the literature on how local ecological
112 processes lead to the formation and dynamics of clusters, and how the resulting
113 spatial patterns relate to ecosystem stability. Owing to the interdisciplinary na-
114 ture of the study, we introduce important terms and concepts via summaries in
115 Boxes 1, 2 and 3 and a glossary in Table 1.

116 We summarize our main findings here: First, our synthesis reveals the im-

117 portance of local positive feedback in the emergence of power-law clustering in
118 various ecosystems. To probe the relationship between positive feedback, clus-
119 tering and resilience, we use a spatially-explicit model which, unlike previous
120 relatively complex models, decouples the effects of positive feedback and envi-
121 ronmental stress. Together with synthesis of previous studies, our analyses en-
122 ables us to demonstrate that power-law clustering (or loss thereof) is unrelated
123 to resilience. We then demonstrate how CSD - a universal feature of dynamical
124 systems near thresholds - manifests as a power-law decay of spatial correlations.
125 We discuss the important role of positive feedback in shaping clustering proper-
126 ties and suggest future directions of research to quantify patterns/dynamics of
127 clustering and to infer ecological interactions.

128 Glossary

- 129 1. Regime shifts: Changes in qualitative nature of ecosystem states. These changes can be
130 abrupt or gradual functions of the underlying drivers.
- 131 2. Critical threshold: The value of an environmental condition (such as rainfall) and/or state
132 variable (e.g. woody cover) at which a system undergoes an abrupt regime shift. In some
133 ecology papers, it is used interchangeably with critical point but here we avoid doing so.
- 134 3. Resilience: The amount of change a system can withstand without transitioning to an al-
135 ternative state. In the model described in Box 2, we interpret resilience as the distance to
136 the threshold driver (or density).
- 137 4. Stability: The rate at which a system recovers to its original equilibrium from small pertur-
138 bations.
- 139 5. Critical point: In the physics literature, this term refers to the value of driver at which the
140 system typically undergoes a *continuous* phase transition from one state to the other.
- 141 6. Critical slowing down: The phenomenon in which systems near threshold of transitions
142 are slow to recover from perturbations.
- 143 7. Positive feedback: Interactions between individuals that results in enhanced reproduction
144 and/or reduced death rates of both individuals.
- 145 8. Cluster: A set of individuals who are within a minimum distance (typically the nearest
146 neighbor distance) of at least one member of the same set.
- 147 9. Scale-free: A quantity having infinite average value, thus lacking a characteristic scale.
148 Also see Box 1.
- 149 10. Percolation: In the physics literature, percolation is the movement/spread of an agent
150 through the entire extent of the system via a connected path of sites.
- 151 11. Percolation density: The lowest density of occupied sites at which a fully connected path in
152 the system is possible. At the same density, we observe a scale-free distribution of cluster
153 sizes in the landscape.
- 154 12. Spatial autocovariance function: Covariance between states at two locations as a function
155 of the distance between them. Also see Box 3.

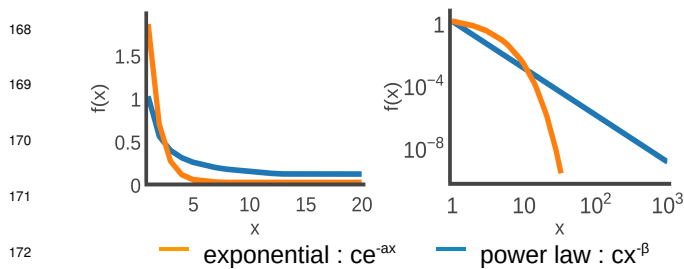
156 13. Power spectrum/Spectral density function: Strength of fluctuations as a function of fre-
 157 quency; it is the Fourier transform of the autocovariance function. Also see Box 3.

158 **BOX 1: POWER-LAW AND SCALE-FREE BEHAVIOURS**

159 Biology is replete with examples of self-organised spatial clustering (Guichard
 160 et al. 2003, Rietkerk and van de Koppel 2008, von Hardenberg et al. 2001). In
 161 some cases, clumps have a wide range of sizes such that the frequency of occur-
 162 rence of clumps of a particular size (denoted by x) decays as a power function of
 163 the size i.e. $f(x) = cx^{-\beta}$ (defined for all clusters above a size $x > x_{min}$ with c and
 164 β being constants). Below we describe two interesting properties of this function.

165
 166

167 **Heavy-tailedness**



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 173 **Figure 2:** The plot on the left shows that the power-law
 174 function has a heavier tail, i.e. higher frequency ($f(x)$) of
 175 occurrence of large events, than in an exponential func-
 176 tion. The plot on the right shows that power-law func-
 177 tion is a straight line on log-log axes; the heavier tail of
 power-law is evident here too.

The power-law frequency distribution has much higher occurrences of extreme events than predicted by commonly used distributions such as Gaussian or exponential distributions (Fig 2); this feature of the power-law distribution is also called *heavy-tailedness*.

Scale-free power-laws

Power-laws with an exponent $\beta \leq 2$ mathematically describe features *that lack a characteristic size/length scale*. To see this, we observe that when $\beta \leq 2$ the mean of this distribution is infinite. Exact expressions for the mean (\bar{x}) and variance (σ_x^2) of the (normalised) power-law probability density function, denoted by $p(x)$, are given by

$$\bar{x} = \int_{x_{min}}^{\infty} xp(x)dx = \begin{cases} x_{min}^{\frac{\beta-1}{\beta-2}} & \text{if } \beta > 2 \\ \infty & \text{if } 1 < \beta \leq 2 \end{cases}$$

$$\sigma_x^2 = \int_{x_{min}}^{\infty} (x - \bar{x})^2 p(x) dx = \begin{cases} x_{min}^2 \frac{\beta-1}{(\beta-3)(\beta-2)^2} & \text{if } \beta > 3 \\ \infty & \text{if } 1 < \beta \leq 3 \end{cases}$$

178 Thus, there is no characteristic size or typical length scale in this distribution,
179 when $\beta \leq 2$, and therefore the distribution is called *scale-free*. Power-law distri-
180 butions of biological quantities with exponents $\beta \leq 2$ are therefore intriguing.
181 Such distributions, however, are not uncommon and have been documented in
182 various ecosystems (Fig 1).

183 **BOX 2: SPATIAL MODEL WITH POSITIVE FEEDBACK**

184 Several spatial models in ecology try to explain power-law clustering but due
185 to their relative complexity, it is difficult to clearly elucidate the role of positive
186 feedback on clustering and resilience (Guichard et al. 2003, Kéfi et al. 2007, Man-
187 rubia and Solé 1997, Scanlon et al. 2007). To address this problem, we employ
188 a simple spatially-explicit model with only two parameters. In this model, we
189 consider a discrete two-dimensional space where each grid cell is updated prob-
190 abilistically depending on states of cells in its neighbourhood. The simplicity
191 of this model allows us to independently tune, and thus study effects of, envi-
192 ronmental driver and positive feedback on spatial patterns via two parameters p
193 and q , respectively. See Fig 3 for a schematic of the update rules; detailed model
194 description is available in Appendix A and was first described in the physics lit-
195 erature in Lübeck (2006) and has been recently adopted in the context of regime
196 shifts (Eby et al. 2017). Using this model, we study the effect of positive feedback
197 (q) on steady-state density (defined as proportion of occupied sites) and spatial
198 patterns (quantified via cluster size distributions and spatial power-spectrum) as
199 a function of the environmental driver p . We add that reducing p in this model
200 can also be interpreted as increasing environmental stress.

201 **Positive feedback and abrupt regime shifts:** Stronger positive feedback
202 in ecosystems are known to cause non-linear and even abrupt responses to
203 stress (Kéfi et al. 2010, 2016, Xu et al. 2015b). In our model too, when positive

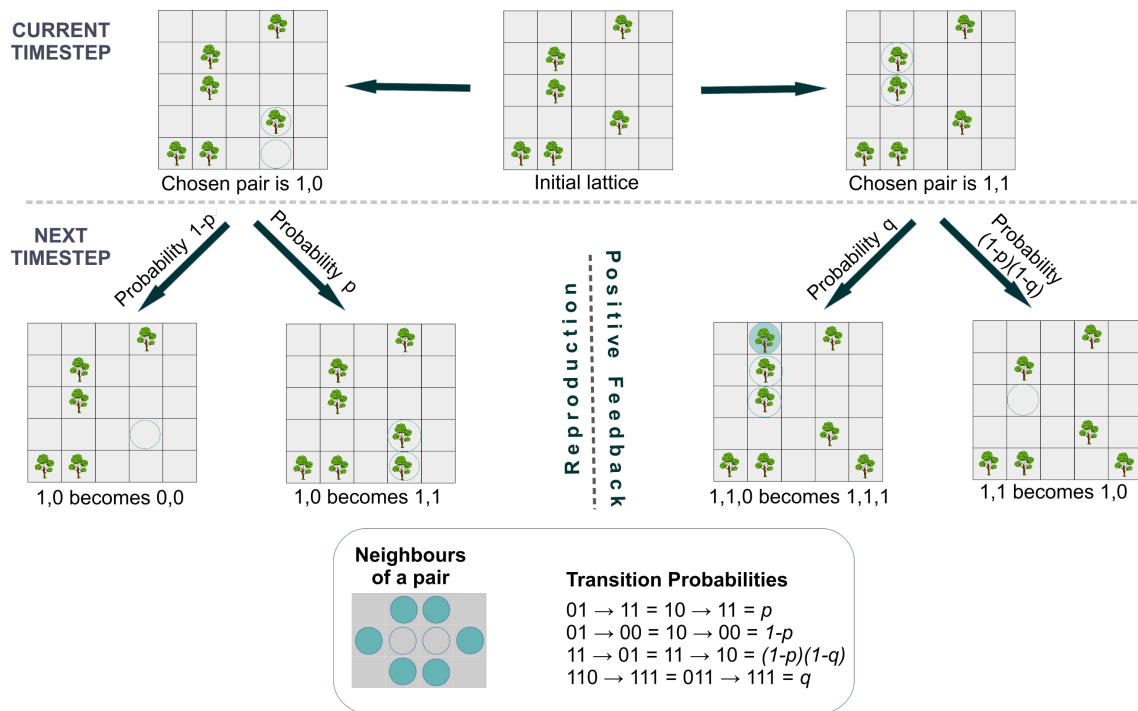


Figure 3: Schematic representation of the model and simulation procedure, for a given 'Initial lattice' shown at the centre of the top row. The parameter p represents baseline birth rate whereas q represents the strength of local positive feedback; reducing p in this model can be interpreted as increasing environmental stress. Light blue circles represent (randomly) chosen cells to update. Depending on the states of chosen cells, the update scheme results in baseline birth or death (left part of second row), or increased birth or reduced death due to positive feedback (right part of the second row). The box at the bottom shows (i) neighbours of a focal pair of cells and (ii) model update rules captured via transition probabilities.

204 feedback is weak, the system undergoes a continuous transition from an occu-
 205 pped to a bare state as we increase environmental stress (Fig 4a). As positive
 206 feedback strength increases, the system can maintain a high density state even
 207 for higher levels of stress; but the system also exhibits an abrupt transition to a
 208 bare state when the stressor crosses the critical threshold. Henceforth, we refer
 209 to the point of transition (defined by either driver value (p) or density (ρ)) from
 210 an occupied to a bare state as *threshold*. When we specifically refer to a continu-
 211 ous transition, we call it a *critical point* whereas the corresponding term for the
 212 discontinuous transition is *critical threshold* (also see Glossary).

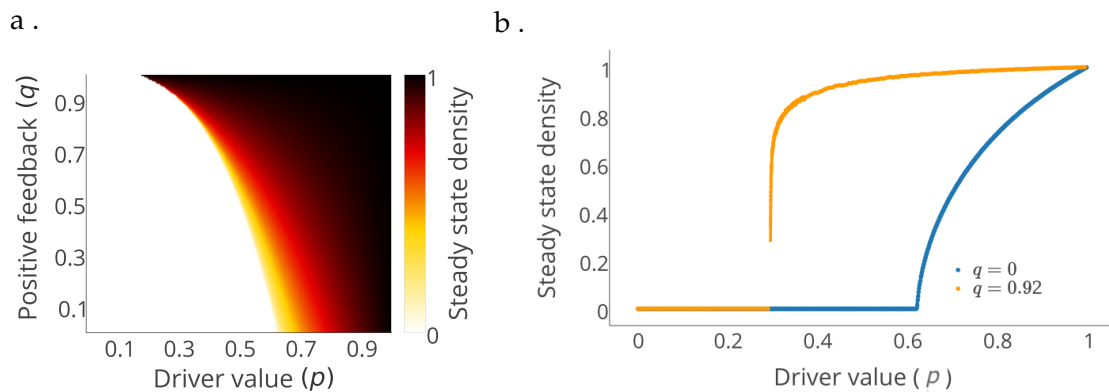


Figure 4: Positive feedback increases the non-linearity and cause an abrupt collapse in response to stress. (a) Steady-state density as a function of driver value (ρ) and positive feedback (q). (b) shows a closer look for two values of positive feedback, $q = 0$ (continuous transition) and 0.92 (discontinuous transition). Lattice size 1024×1024 .

213 III. POSITIVE FEEDBACK PROMOTES POWER-LAW CLUSTERING AT RELA- 214 TIVELY LOW DENSITIES

215 Power-law cluster size distributions are seen in diverse ecosystems, includ-
216 ing drylands, mussel beds, seagrass beds, sawgrass and forest fires (Fig 1). Many
217 studies propose that these systems are likely shaped by positive feedback (Aguiar
218 and Sala 1994, Boada et al. 2017, Dell et al. 2016, Foti et al. 2013, Guichard et al.
219 2003, Kéfi et al. 2007, Knowlton 2004, Maestre et al. 2003, Scanlon et al. 2007).
220 For example, in semi-arid landscapes, seed germination and seedling survival
221 probabilities are higher in the neighbourhood of other plants than out in the
222 open (Aguiar and Sala 1994, Maestre et al. 2003). This results from reduced light
223 and heat stress as well as increased water availability to young saplings in the
224 vicinity of adult plants. Similarly, in mussel-beds, steadfast attachment of mussel
225 to the substrate is directly dependent on the attachment of neighbours (Guichard
226 et al. 2003). In macroalgal beds, recruitment and survival of macroalgal fronds
227 shows density dependence due to the protection offered by neighbours from her-
228 bivory by sea urchin and fish (Boada et al. 2017, Dell et al. 2016).

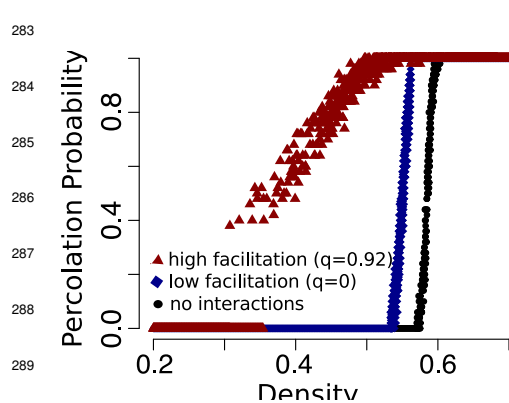
229 To understand how positive feedback promotes such clustering, it is insight-
230 ful to first discuss how power-law cluster-size distributions are also realised in
231 ‘null models’ that are devoid of any interactions among organisms. In spatial
232 null models, individuals are initially assigned to random locations on a two-

233 dimensional discrete lattice. They then either die or give birth to an offspring
234 at a rate that does not depend on presence/absence of any individual on the
235 landscape (Kéfi et al. 2011). Consequently, the proportion of occupied sites in the
236 landscape (henceforth called density) changes from nonzero values to zero (a bare
237 state) as a gradual function of decreasing birth (or increasing death) rates (Grim-
238 mettt 1999, Kéfi et al. 2011). These null models correspond to classic models in
239 the physics literature in the context of a phenomenon called percolation (Stauffer
240 1979). The lowest density at which there is a non-zero probability of emergence
241 of a fully connected path in the system is called the *percolation density*; at the per-
242 colation density (denoted by ρ_p), the system also shows a scale-free clustering. In
243 other words, despite the lack of positive feedback in these null models, a power-
244 law cluster size distribution with $\beta < 2$, and hence scale-free clustering, occurs
245 at the percolation density (ρ_p). The value of percolation density depends on the
246 geometry of the landscape. For ecological contexts, a relevant geometry is that of
247 two dimensional square lattice where the percolation density is 0.59 (Grimmett
248 1999, Stauffer 1979).

249 In many ecosystems, densities that correspond to power-law clustering are
250 typically lower than the above mentioned percolation density of the null model.
251 For example, regions in the Kalahari show power-law cluster-size distributions
252 of vegetation for densities ranging from 0.14 to 0.54 (Scanlon et al. 2007); a bulk
253 of these areas also exhibit power-laws with exponent $\beta < 2$, and are thus scale-
254 free (see Box 1 and Glossary). Power-law cluster-size distributions observed in
255 several other ecosystems also show exponents within the scale-free range (e.g.
256 Fig 1). To explain such power-law clustering, many spatial ecological models of
257 ecosystems have been developed (Grassberger 1993, Guichard et al. 2003, Kéfi
258 et al. 2007, Manrubia and Solé 1997, Scanlon et al. 2007). These models often in-
259 corporate ecosystem-specific processes and are consequently complex, involving
260 many parameters. Nevertheless, they have commonalities. For example, they all
261 assume local positive feedback in some form that causes increased birth (or re-
262 duced death) rates of individuals who are surrounded by others (Box 2). Below,
263 we explain how local positive feedback can lower the percolation density.

264 The emergence of power-law clustering depends on how local interactions be-
265 tween individuals scale to cluster dynamics. Even in spatial null models that are

266 devoid of any positive interactions, clusters form entirely due to random filling
267 of the lattice; furthermore, larger clusters are more likely to merge with other
268 clusters and therefore have higher growth rates. Theory predicts that power-law
269 clustering emerges whenever clusters grow in proportion to their size, a phe-
270 nomenon known as proportionate growth (Grimmett 1999, Manor and Shnerb
271 2008, Stauffer 1979). Such growth occurs at the percolation density of 0.59 for the
272 spatial null models for a square lattice. In models with local positive interactions,
273 empty sites near an existing cluster of occupied states are more likely to become
274 occupied. This not only expands the original cluster but also increases chances
275 of merger of this cluster with a nearby cluster. This dynamic of clusters is con-
276 trast to the spatial null model where expansion as well as merger of clusters are
277 driven entirely by the random filling of the landscape. Therefore, in models with
278 positive feedback, proportionate growth and scale-free cluster size distributions
279 (i.e. a power-law with $1 < \beta < 2$) occur at densities lower than the percolation
280 density of the null model (Scanlon et al. 2007). This may offer a potential expla-
281 nation for the observed low densities at which power-law clustering is seen in
282 many ecosystems (Fig 1).



290 Figure 5: Percolation probability
291 changes from zero to nonzero at density
292 0.59 for the spatial null model, 0.53 for
293 low positive feedback ($q = 0$) and 0.31
294 for high positive feedback ($q = 0.92$).
295 For each of these cases, these transitions
296 in percolation probability occur exactly
297 at the density where power-law cluster
298 size distributions are observed; see Fig 6.
Lattice size 256×256 .

We support this argument by showing how percolation density changes with positive feedback in our model (Box 2). To do so, we use the concept of *percolation probability* which is defined as the probability of occurrence of a fully-connected path of occupied cells in the landscape. In Fig 5, we display the percolation probability as a function of density for two different values of positive feedback and the spatial null model. We then identify percolation density as the lowest density at which this probability is non-zero. We find that the percolation density is lower for the system with higher positive feedback, consistent with our synthesis of previous theoretical and empirical stud-

299 ies discussed above. Interestingly, we also observe that weak positive feedback
300 leads to continuous change in percolation probability whereas strong positive
301 feedback, owing to stronger nonlinear response of the system, makes it discon-
302 tinuous (Fig 5).

303 Putting empirical studies together with spatially-explicit models of clustering,
304 we conjecture that strong positive feedback is likely to be the key interaction low-
305 ering percolation density in many ecosystems.

306 **IV. CLUSTER-SIZE DISTRIBUTIONS ARE NOT INDICATORS OF ECOSYSTEM** 307 **RESILIENCE**

308 Let us now use the above link we established between positive feedback and
309 cluster-size distributions to address the larger question: how general is the rela-
310 tionship between cluster-size distributions and ecosystem resilience?

311 Studies over the last decade have suggested that fragmentation of large clus-
312 ters leads to a thinning of the tail of the cluster-size distribution. Consequently,
313 this causes loss of a power-law clustering, which can be used as an indicator of a
314 stressed and less resilient ecosystem ([Génin et al. 2018a](#), [Kéfi et al. 2014](#), [Kéfi et al.](#)
315 [2007](#)). A corollary to this hypothesis is that ecosystems with power-law cluster-
316 ing are relatively farther from critical thresholds, and hence are likely to be more
317 resilient. The evidence for this hypothesis in both models and data, however,
318 has been ambiguous ([Maestre and Escudero 2009](#), [Meloni et al. 2017b](#), [Moreno-](#)
319 [de las Heras et al. 2011](#), [Schneider and Kéfi 2016](#)). Therefore, the generality of the
320 relationship between cluster sizes and resilience remains unknown.

321 To resolve this, let us first consider how positive feedback affects both spatial
322 clusters and resilience. As we argued in section III, positive feedback lowers the
323 percolation density. Additionally, positive feedback promotes abrupt transitions
324 and increases the threshold value of density from which the regime shift happens
325 (Fig 4). Therefore, we hypothesize that, depending on the strength of positive
326 feedback, power-law clustering can occur at any distance from the threshold of
327 regime shift. We make predictions for two scenarios: we predict that in systems

328 with *weak positive feedback*, the distance between percolation density and thresh-
329 old of regime shift will be relatively large. Thus, increasing stress and an ap-
330 proach to threshold follows the previously expected pattern of loss of power-law
331 clustering (Kéfi et al. 2014). In contrast, for systems with *strong positive feedback*,
332 which are most likely to exhibit abrupt transitions, the distance between percola-
333 tion density and the critical threshold of collapse will be negligible or even zero.
334 Hence, power-law clustering may occur at the critical threshold itself and the loss
335 of power-law clustering cannot be used as a resilience indicator.

336 To buttress our arguments, we analyse the model presented in Box 2. Indeed,
337 our model analysis confirms our expectations: A weak positive feedback scenario
338 shows that percolation density (ρ_p) is relatively far from the threshold of transi-
339 tion (ρ_p) (Fig 6a and inset); moreover, we find that loss of power-law clustering
340 and appearance of thin-tailed (exponential) cluster-size distribution precedes the
341 transition (see Appendix C). Our model reveals that this distance between the
342 density of threshold of transition and percolation density reduces as a function
343 of positive feedback and becomes even zero for large values of positive feedback
344 (Fig 6c, d). Consequently, the qualitative features of cluster size distribution (e.g.
345 being a power-law, truncated power-law or exponential) do not follow a general
346 trend as a function of ecosystem stress see Appendix C. In Fig 6b and inset, we
347 show a case where a strong positive feedback scenario shows a power-law clus-
348 tering occurring very near, even possibly at, the critical threshold of collapse. Put
349 together, our model analyses suggests that the relationship between cluster-sizes
350 and ecosystem resilience heavily depends on the strength of positive feedback
351 in the ecosystem. We recall that systems with strong positive feedback are most
352 likely to exhibit abrupt shifts; it is precisely in these systems that the expected
353 trend of cluster-size distributions, of loss of power-law as the system approaches
354 thresholds, is least likely to be true. This questions the generality as well as po-
355 tential utility of cluster-size distributions as indicators of ecosystem resilience.

356 We discuss above results in light of theoretical studies which too have found
357 the association of cluster-size distributions with resilience to be tenuous (Génin
358 et al. 2018b, Schneider and Kéfi 2016). These studies investigate spatially-explicit
359 models of dryland vegetation and forest gap dynamics. They include, for exam-
360 ple, a lowered grazing-induced mortality for individuals with more neighbours,

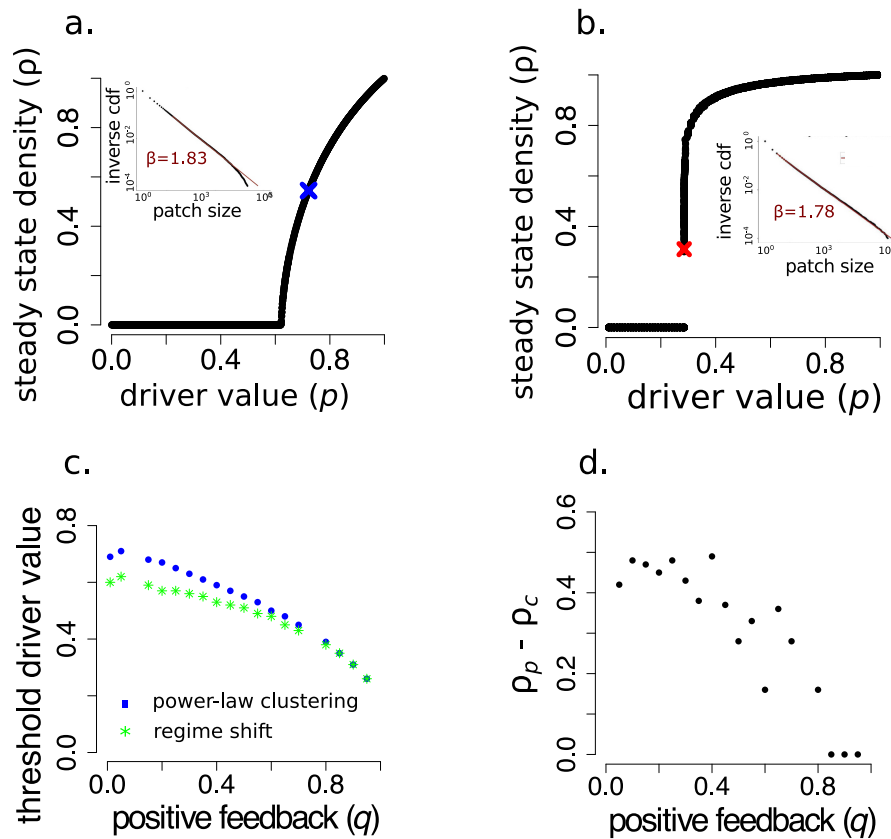


Figure 6: Strength of positive feedback, rather than distance to thresholds, determines the density at which power-law (scale-free) clustering occurs. The values of driver and density at which we find a power-law distribution are shown as crosses in the phase-diagrams (a) and (b), with their insets showing the corresponding inverse cumulative distribution function (CDF) of the patch-sizes. (a) When positive feedback is weak ($q = 0$), power-law clustering occurs far from ecosystem transition, consistent with previous hypotheses. (b) When positive feedback is strong ($q = 0.92$), power-law clustering can occur close to (or even at) the critical threshold of collapse. For the fitted function $kx^{-\beta}$ wherein $k = (\beta - 1)x_{min}^{\beta-1}$, $x_{min} = 17$ in (a) and 3 in (b); lattice size used: $= 1024 \times 1024$. (c) shows the driver values at which power-law clustering moves closer to the threshold of transition as positive-feedback (q) increases. (d) shows that $\rho_p - \rho_c$, i.e. the difference between the density at which patches follow a power-law distribution (ρ_p) and the density of the transition threshold (ρ_c), reduces as positive feedback (q) increases. For (c) and (d), lattice size of 256×256 was chosen to reduce computational time. See Appendix C for cluster size distributions at other values of p and q .

361 a process termed associative protection. When the associative protection is high,
 362 they find power-law clustering at/near the critical thresholds of collapse. These
 363 results are consistent with our synthesis because associative protection in their
 364 model (i.e., reduced mortality for plants with neighboring plants) is analogous to
 365 increased positive feedback in our model (which causes reduced death rates for
 366 individuals with neighbors).

367 Synthesizing our results together with these recent studies, we argue that
368 cluster-size distribution primarily depends on the strength of the positive feed-
369 back and that it cannot be employed as an indicator of ecosystem resilience.
370 Furthermore, since cluster-size distributions do not primarily depend on prox-
371 imity to critical threshold in these stochastic and spatial ecological models, we
372 conclude that it is also unrelated to critical slowing down (CSD); we recall that
373 CSD is a generic dynamical feature of systems near critical thresholds. See next
374 section on how CSD influences spatial properties and causes power-law features
375 in them.

376 V. SCALE-FREE SPATIAL CORRELATIONS MAY ARISE AT CRITICAL THRESH- 377 OLDS OF ECOSYSTEM COLLAPSE

378 So far, we have demonstrated that cluster-size distribution do not represent
379 resilience and hence cannot reliably indicate imminent regime shifts. However,
380 the theory of phase transitions posits the emergence of scale-free features near/at
381 critical points. Here, using our spatially-explicit ecological model, we illustrate
382 how critical slowing down - a canonical features of dynamical systems near
383 thresholds - causes scale-free behaviour in the spatial autocovariance function
384 (Fig 7; Box 3;).

385 As an ecosystem approaches a critical threshold, its return to equilibrium
386 state, when perturbed, becomes increasingly slower. This phenomenon of crit-
387 ical slowing down (Ma 2000, Scheffer et al. 2009, Wissel 1984) has two impli-
388 cations - increased spatial correlations (Dakos et al. 2010) and increased spatial
389 variance (Guttal and Jayaprakash 2009). To understand this, consider how a per-
390 turbation from the equilibrium state at any location in the ecosystem will spread
391 in space. First, owing to slowed dynamics, the perturbation lives longer and,
392 via spatial connectedness in the system, propagates to larger distances in the
393 system (Ma 2000, Sethna 2006). Consequently, a measure of spread of perturba-
394 tion, the correlation length, increases (Dakos et al. 2010, Ma 2000, Sethna 2006).
395 Second, as the perturbations persist for longer duration, further disturbances
396 enhance amplitudes of the fluctuations. This manifests as increasing spatial vari-

397 ance in the ecosystem as it moves towards the threshold (Guttal and Jayaprakash
398 2009). Here, we consider the spatial autocovariance function, defined as covari-
399 ance of local densities at two locations separated by a distance r (Box 3). This
400 function captures both spatial variance and correlations.

401 Before we illustrate computations of spatial autocovariance for our model, we
402 make a couple of technical remarks. First, physicists often refer to the autoco-
403 variance function as the ‘correlation function’; some theoretical papers in ecology
404 also do the same (Roy et al. 2003). Here, we have adopted the standard termi-
405 nology that is used in quantitative ecology literature (Eq 2 in Box 3). Second, the
406 theory of phase transition predicts critical slowing down and consequent scale-
407 free behaviour at critical points of *continuous phase transitions* (Ma 2000, Sethna
408 2006). However, it has been shown that signatures of CSD are present, albeit
409 with a relatively less magnitude, even in ecological models exhibiting abrupt
410 transitions (Dakos et al. 2011, Scheffer et al. 2009). Consequently, we argue and
411 demonstrate using the simple ecological model presented in Box 2 that scale-free
412 behaviour may characterize critical thresholds of abrupt transitions as well.

413 Calculation of the spatial autocovariance function is often beset with statistical
414 and computational difficulties. Therefore, we focus on a mathematically equiva-
415 lent measure of correlations in spatial patterns via its power spectrum (Kéfi et al.
416 2014) (Box 3; Appendix D). It can be shown that the power spectrum is the Fourier
417 transform of the autocovariance function (Reif 2009). The power spectrum of a
418 spatial pattern provides a measure of the relative contribution of fluctuations at
419 different spatial frequencies in the system, to its overall pattern. It is known in the
420 ecology literature that as systems approach critical thresholds, the low frequency
421 modes begin to dominate their power spectrum (Carpenter and Brock 2010, Kéfi
422 et al. 2014). However, the full functional form of the power-spectrum is rarely
423 quantified (but see Barbier et al. (2006), Bonachela et al. (2015), Couteron (2002)
424 in the context of periodic and multi-scale patterns of dryland vegetation). Simu-
425 lations of our model shows that the power-spectrum indeed becomes scale-free
426 at critical thresholds for systems with both weak and strong positive feedback
427 (Fig 7). We explain in Box 3 that a scale-free power spectrum is indicative of a
428 scale-free autocovariance function. Thus, scale-free power-spectrum character-
429 izes the structure of spatial perturbations near/at critical thresholds of ecosystem

430 collapse.

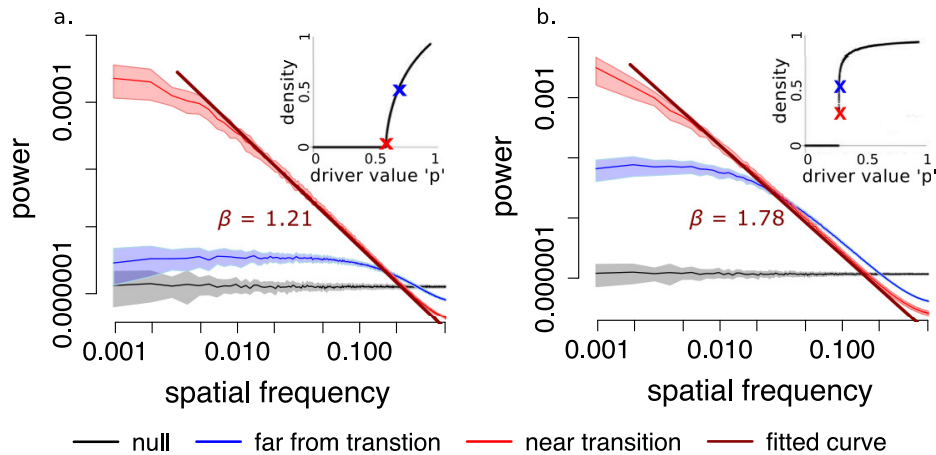


Figure 7: The power-spectrum of systems at very near (or at) the threshold of transitions decays as a power-law function of spatial frequency, (a) for $q = 0$ and (b) for $q = 0.92$. Lines represent the mean trend and bands, the SD. Insets show the location of parameter values for which power-spectrum are plotted. Blue is far from transition, Red is close/at the threshold and Grey represents the spatial null model. For the fitted function $kx^{-\beta}$, $k = 2.02 \times 10^{-7}$ when $q = 0$ and 4.65×10^{-8} when $q = 0.92$. We used $= 1024 \times 1014$ lattice.

431 BOX 3: COVARIANCE, CORRELATION AND SPECTRAL FUNCTION

432 One way to capture the spread of disturbance in a system or the length scale of
 433 spatial fluctuations, is by constructing the spatial covariance function. The *spatial*
 434 *autocovariance function* for local density ρ for a distance r is defined as

$$C(r) = \langle (\rho(\mathbf{x}) - \bar{\rho})(\rho(\mathbf{x}') - \bar{\rho}) \rangle \quad (1)$$

435 where $\bar{\rho}$ represents mean density over the entire landscape, angular brackets de-
 436 note average over all locations \mathbf{x} and \mathbf{x}' in the landscape that are separated by a
 437 distance r . Ecologists widely use the correlation function which is defined as

$$K(r) = \frac{\langle (\rho(\mathbf{x}) - \bar{\rho})(\rho(\mathbf{x}') - \bar{\rho}) \rangle}{\sigma^2} \quad (2)$$

438 where σ^2 is the spatial variance of densities in the ecosystem. Thus the covariance
 439 function is a production of the correlation function and the variance.

440 The *correlation length* is defined as the mean of the covariance function and can
 441 be interpreted as the average distance to which local fluctuations spread. The cor-

442 relation length becomes infinite at the critical thresholds. This means that the co-
443 variance function then follows a power-law with an exponent less than 2 (Box 1).

444 The *power spectrum*, denoted by $S(k)$, is the Fourier transform of its autocovari-
445 ance function (Baugh and Murdin 2006, Reif 2009). Therefore, it can be calculated
446 as

$$S(k) = \int C(r)e^{-ikr} dr \quad (3)$$

447 At critical thresholds, we expect the spatial covariance function to exhibit a
448 power-law relation with distance

$$C(x) = c_0x^{-\alpha} \quad (4)$$

where c_0 is a constant and α is an exponent less than two. The corresponding
spectral function for an n-dimensional system is given by

$$|S(\mathbf{k})| \sim \mathbf{k}^{-(n-\alpha)}$$

449 Therefore, evidence of a power-law spectral function is also evidence of a power-
450 law autocovariance function.

451 VI. DISCUSSION

452 In this study, we set out to investigate the generality of the conclusion that
453 loss of power-law clustering in ecosystems is indicative of reducing resilience.
454 First, our synthesis reveals that power-law clustering (or lack thereof) is unrelated
455 to ecosystem resilience. We argue that this is because power-law clustering is
456 fundamentally associated with local positive feedback rather than any generic
457 dynamics of systems near critical thresholds of ecosystem collapse. Second, when
458 ecosystems are in the vicinity of critical thresholds of collapse, critical slowing
459 causes a power-law (scale-free) behaviour but in a different metric - the spatial
460 autocovariance, or spectral function, of local densities.

461 A. Local positive feedback, clustering and resilience

462 Previous ecological models that have attempted to resolve these connections
463 include complex interactions often specific to particular ecosystems (Kéfi et al.
464 2011, 2007, Meloni et al. 2017b, Scanlon et al. 2007, Schneider and Kéfi 2016). In
465 such models, many parameters contribute to local positive feedback and environ-
466 mental stress, thus making it difficult to disentangle causal links between local
467 processes and macroscopic patterns. Here, we deliberately used a simple model
468 with only two parameters representing environmental stress and local positive
469 feedback. The simplicity of the model we employed may also be seen as a limi-
470 tation. However, it helped us conclude that loss of power-law clustering is not a
471 robust indicator of approach to ecosystem transitions. Furthermore, it allowed us
472 to disentangle the effects of environmental stress and positive feedback on clus-
473 tering and resilience. Specifically, we propose a hypothesis that distance between
474 power-law clustering (percolation threshold) and critical threshold of collapse re-
475 duces as the strength of positive feedback increases.

476 Seminal ecological models that try to explain power-law clustering observed
477 in ecosystems (Kéfi et al. 2007, Scanlon et al. 2007) assume that local births/deaths
478 of trees, in addition to being positively influenced by local density, is negatively
479 regulated by *global-scale feedback*. Mechanisms such as rapid spread of water in
480 the landscape (von Hardenberg et al. 2010) are offered as potential explanations
481 for negative regulation of local growth due to global-scale vegetation density.
482 Indeed, based on our synthesis ((Manor and Shnerb 2008, 2009) and Fig 6), we
483 demonstrate that there is no need to invoke global-scale feedback; in fact, local
484 positive feedback alone can explain the emergence of scale-free clustering in these
485 systems.

486 Several empirical studies find neither scale-free clustering (Weerman et al.
487 2012, Xu et al. 2015a) nor the expected shifts of cluster-size distributions with
488 increasing stress (Casey et al. 2016). However, they attribute this to an ab-
489 sence/disruption of global negative feedback in their systems (Casey et al. 2016,
490 von Hardenberg et al. 2010, Weerman et al. 2012)(but also see Moreno-de las
491 Heras et al. (2011)). As we argued in the previous paragraph, negative feed-

492 back isn't even a necessary condition for scale-free clustering. Further, based
493 on our synthesis that cluster-sizes do not relate to resilience, we posit that these
494 empirical results are not surprising.

495 Cluster-size distributions are nevertheless relevant in ecological contexts where
496 connectivity or porosity of the landscape is of focal interest, such as in the case of
497 forest fires or disease spread (Chaves et al. 2012, Turcotte and Malamud 2004). We
498 illustrated that power-law clustering in our ecological model is associated with
499 a transition in percolation probability, representing the emergence of a spanning
500 cluster (i.e. a fully connected path) in the system. This is also seen in spatial mod-
501 els of predator-prey interactions (Roy et al. 2003). In the context of fire or disease
502 outbreaks, presence of a cluster of vegetation (susceptible individuals) allows
503 fire (disease) to easily spread within each cluster. Consequently, scale-free clus-
504 tering, which indicates a highly connected landscape, allows the possibility of
505 catastrophic fire (or disease) outbreaks. These models represent a fundamentally
506 different class of models from what we have discussed in this paper (Dickman
507 et al. 2000, Solé et al. 1999).

508 **B. Cluster-sizes and Critical slowing down**

509 Our synthesis predicts that scale-free behaviour in spatial correlations, mea-
510 sured via autocovariance or spectral functions, can characterise critical thresh-
511 olds. This feature, we argued, arises from the critical slowing down - i.e. slowed
512 response of ecosystems near threshold points - which is a generic feature of many
513 ecological transitions. How do scale-free correlations in density (described in V)
514 and scale-free clustering (Section III) relate to each other? They both indicate
515 emergence of large spatial scales in the system. However, they capture fun-
516 damentally different properties. Scale-free correlations in density indicate that
517 *perturbations* spread to large distances in ecosystems as a consequence of critical
518 slowing down. Therefore, it captures the dynamics of perturbations and hence
519 can be used to infer stability or lack thereof. In contrast, scale-free clusters indi-
520 cate the presence of large clusters, which do not correspond to dynamics of how
521 perturbations decay. Therefore, clustering properties are unrelated to resilience,

522 as we indeed demonstrate in Fig 6.

523 How efficient is it to use scale-free features of density correlations as an early
524 warning signals (EWS) of regime shifts or critical transitions (Scheffer et al. 2009)?
525 The purpose of early warning signals is to detect signatures of approach to crit-
526 ical thresholds. In that sense, computing simpler metrics of spatial autocorrela-
527 tion between neighboring sites (Dakos et al. 2010) or spatial variance (Guttal and
528 Jayaprakash 2009) may have advantages such as ease of computation and better
529 statistical reliability in comparison to characterising the complete form of autoco-
530 variance or spectral functions. On the other hand, simpler metrics are also easily
531 affected by external factors, such as increased spatial heterogeneity or external
532 variability (Dakos et al. 2010, Kéfi et al. 2014) and hence confound interpreta-
533 tions. Further investigations can reveal the relative efficacy of different spatial
534 metrics.

535 C. Future directions

536 Our synthesis suggests some exciting directions for future research. The focus
537 of recent research, as reviewed in this paper, has been to understand how local
538 interactions produce clustering properties, and how clustering properties can be
539 used to infer resilience. However, the inverse problem of inferring ecological
540 interactions from spatial images of ecosystems remains poorly studied. For ex-
541 ample, the occurrence of power-law or such heavy-tailed distributions can itself
542 be used to infer the role of local facilitative interactions in the ecosystem. Indeed,
543 one recent study does suggest that skewness of cluster-size distributions can sug-
544 gest positive feedback in dryland-vegetation systems (Xu et al. 2015a). The bigger
545 question remains open: can we quantify the strengths and spatial scales of posi-
546 tive feedback and other ecological interactions between organisms by analysis of
547 spatial images, for example via geometrical properties of clusters such as cluster
548 sizes, fractal dimensions of clusters, and the strength of spatial correlations in the
549 system.

550 With recent advancements in remote sensing and reducing costs of spatial im-
551 ages, we can also procure extensive high-resolution spatial data over time. This

552 will enable us to quantify not only patterns, as described above, but also dynam-
553 ics of various cluster properties. Unlike static properties available from a single
554 image, dynamical properties may reflect stability of ecosystems by capturing how
555 systems respond to perturbations. Theoretical works, inspired by studies of do-
556 main growth in phase transitions, describe the dynamics of clusters in simple
557 ecological models exhibiting continuous and discontinuous transitions ([Manor
558 and Shnerb 2008](#), [Weissmann et al. 2017](#)). However, much remains to be done in
559 integrating these studies with real data. This requires extensive theoretical and
560 computational studies to identify suitable metrics of clustering properties, devel-
561 opment of statistical frameworks including appropriate spatial null models and
562 finally, empirical validations/applications based on analyses of aerial images of
563 ecosystems.

564 **D. Concluding remarks**

565 Our synthesis helps us disentangle processes that generate power-law cluster
566 sizes, scale-free correlations and how they relate to ecosystems' critical thresh-
567 olds. Real world analyses however can sometimes yield misleading patterns,
568 including power-laws and scale-free behaviours due to sampling artefacts ([Plank
569 and Codling 2009](#)) or misfitting ([Clauset et al. 2009](#), [Meloni et al. 2017a](#), [White
570 et al. 2008](#)). Various ways in which patterns are misconstrued as power-laws
571 have been discussed in detail in multiple other forums ([Breed et al. 2015](#), [Clauset
572 et al. 2009](#), [Stumpf and Porter 2012](#)). Where there is a true power-law with an
573 exponent less than two, since it is indicative of diverging quantities, there is a
574 tendency to associate such a pattern with a critical phenomenon. However, scale-
575 free patterns can also arise when underlying processes operate at multiple scales
576 and due to landscape heterogeneity ([Khaluf et al. 2017](#), [Petrovskii et al. 2011](#)).
577 Naive association of observed scale-free behaviours with either criticality or sta-
578 bility is problematic. An additional challenge in interpreting spatial patterns in
579 ecosystems is to disentangle effects of underlying spatial heterogeneity from true
580 self-organisation. With the increasing availability of high-resolution spatial datas,
581 from satellites to drone based imagery, of various ecosystems, spatial analyses are
582 likely to be widely deployed in the future. Our study highlights the importance

583 of having a clear understanding of how local-interactions drive macroscopic be-
584 haviours to infer ecological interactions and resilience of ecosystems.

585 VII. DATA AND CODES

586 All simulation analyses codes, with simulation datasets corresponding to re-
587 sults presented in this paper, have been made publicly available at: [https://](https://github.com/ssumithra/PowerLawCriticalityPaper)
588 github.com/ssumithra/PowerLawCriticalityPaper. Detailed instruc-
589 tion on execution of these codes are also provided.

590 VIII. ONLINE SUPPLEMENTARY MATERIALS

591 **Appendix A:** Power-law Vs exponential functions.

592 **Appendix B:** Detailed model description.

593 **Appendix C:** Statistical fitting of cluster-size distributions.

594 **Appendix D:** Cluster-size distributions across the phase-diagram for low and
595 strong positive feedback.

596 **Appendix E:** Power-spectrum fitting.

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- 609 Aguiar, M. R. and Sala, O. E. 1994. Competition, facilitation, seed distribution and the
610 origin of patches in a patagonian steppe. *Oikos* 70:26–34.
- 611 Barbier, N., Couteron, P., Lejoly, J., Deblauwe, V., and Lejeune, O. 2006. Self-organized
612 vegetation patterning as a fingerprint of climate and human impact on semi-arid ecosys-
613 tems. *Journal of Ecology* 94:537–547.
- 614 Baugh, C. and Murdin, P. 2006. Correlation function and power spectra in cosmology.
615 *Encycl. of Astronomy and Astrophysics*. IOP Publishing, Bristol .
- 616 Boada, J., Arthur, R., Alonso, D., Pagès, J. F., Pessarrodona, A., Oliva, S., Ceccherelli, G.,
617 Piazzzi, L., Romero, J., and Alcoverro, T. 2017. Immanent conditions determine imminent
618 collapses: nutrient regimes define the resilience of macroalgal communities 284:20162814.
- 619 Bonachela, J. A., Pringle, R. M., Sheffer, E., Coverdale, T. C., Guyton, J. A., Caylor, K. K.,
620 Levin, S. A., and Tarnita, C. E. 2015. Termite mounds can increase the robustness of
621 dryland ecosystems to climatic change. *Science* 347:651–655.
- 622 Breed, G. A., Severns, P. M., and Edwards, A. M. 2015. Apparent power-law distribu-
623 tions in animal movements can arise from intraspecific interactions. *Journal of the Royal*
624 *Society Interface* 12:20140927.
- 625 Carpenter, S. and Brock, W. 2010. Early warnings of regime shifts in spatial dynamics
626 using the discrete fourier transform. *Ecosphere* 1:art10.
- 627 Carpenter, S., Cole, J., Pace, M., Batt, R., Brock, W., Cline, T., Coloso, J., Hodgson, J.,
628 Kitchell, J., Seekell, D., et al. 2011. Early warnings of regime shifts: A whole-ecosystem
629 experiment. *Science* 332:1079–1082.
- 630 Carpenter, S. R., Ludwig, D., and Brock, W. A. 1999. Management of eutrophication for
631 lakes subject to potentially irreversible change. *Ecological applications* 9:751–771.
- 632 Casey, S. T., Cohen, M. J., Acharya, S., Kaplan, D. A., and Jawitz, J. W. 2016. Hydro-
633 logic controls on aperiodic spatial organization of the ridge-slough patterned landscape.
634 *Hydrology and Earth System Sciences* 20:4457.
- 635 Chaves, L. F., Hashizume, M., Satake, A., and Minakawa, N. 2012. Regime shifts and
636 heterogeneous trends in malaria time series from western kenya highlands. *Parasitology*
637 139:14–25.

- 638 Clauset, A., Shalizi, C. R., and Newman, M. E. 2009. Power-law distributions in empirical
639 data. *SIAM review* 51:661–703.
- 640 Coutron, P. 2002. Quantifying change in patterned semi-arid vegetation by fourier anal-
641 ysis of digitized aerial photographs. *International Journal of Remote Sensing* 23:3407–
642 3425.
- 643 Dakos, V., Carpenter, S. R., Ellison, A. M., Guttal, V., Ives, A. R., Kéfi, S., Livina, V., Seekell,
644 D. A., van Nes, E. H., and Scheffer, M. 2012. Early warning signals of critical transitions:
645 Methods for time series. *PLoS One* 7:e41010.
- 646 Dakos, V., Kéfi, S., Rietkerk, M., Van Nes, E. H., and Scheffer, M. 2011. Slowing down
647 in spatially patterned ecosystems at the brink of collapse. *The American Naturalist*
648 177:E153–E166.
- 649 Dakos, V., van Nes, E., Donangelo, R., Fort, H., and Scheffer, M. 2010. Spatial correlation
650 as leading indicator of catastrophic shifts. *Theoretical Ecology* 3:163–174.
- 651 Dell, C. L., Longo, G. O., and Hay, M. E. 2016. Positive feedbacks enhance macroalgal
652 resilience on degraded coral reefs. *PloS One* 11:e0155049.
- 653 Dickman, R., Muñoz, M. A., Vespignani, A., and Zapperi, S. 2000. Paths to self-organized
654 criticality. *Brazilian Journal of Physics* 30:27–41.
- 655 D’Odorico, P., Okin, G., and Bestelmeyer, B. 2012. A synthetic review of feedbacks and
656 drivers of shrub encroachment in arid grasslands. *Ecohydrology* 5:520–530.
- 657 Duffy, J. P., Pratt, L., Anderson, K., Land, P. E., and Shutler, J. D. 2017. Spatial assessment
658 of intertidal seagrass meadows using optical imaging systems and a lightweight drone.
659 *Estuarine, Coastal and Shelf Science* 200:169–180.
- 660 Eby, S., Agrawal, A., Majumder, S., Dobson, A. P., and Guttal, V. 2017. Alternative stable
661 states and spatial indicators of critical slowing down along a spatial gradient in a savanna
662 ecosystem. *Global Ecology and Biogeography* 26:638–649.
- 663 Fernández, A. and Fort, H. 2009. Catastrophic phase transitions and early warnings
664 in a spatial ecological model. *Journal of Statistical Mechanics: Theory and Experiment*
665 2009:P09014.
- 666 Foti, R., del Jesus, M., Rinaldo, A., and Rodriguez-Iturbe, I. 2013. Signs of critical transi-
667 tion in the everglades wetlands in response to climate and anthropogenic changes. *Pro-*
668 *ceedings of the National Academy of Sciences* 110:6296–6300.

- 669 Génin, A., Majumder, S., Sankaran, S., Danet, A., Guttal, V., Schneider, F. D., and Kéfi, S.
670 2018a. Monitoring ecosystem degradation using spatial data and the r package spatial-
671 warnings. *Methods in Ecology and Evolution* 9:2067–2075.
- 672 Génin, A., Majumder, S., Sankaran, S., Schneider, F. D., Danet, A., Berdugo, M., Guttal,
673 V., and Kéfi, S. 2018b. Spatially heterogeneous stressors can alter the performance of
674 indicators of regime shifts. *Ecological Indicators* 94:520–533.
- 675 Grassberger, P. 1993. On a self-organized critical forest-fire model. *Journal of Physics A:*
676 *Mathematical and General* 26:2081.
- 677 Grimmett, G., 1999. What is percolation? Pages 1–31 *in* Percolation. Springer Berlin
678 Heidelberg.
- 679 Guichard, F., Halpin, P. M., Allison, G. W., Lubchenco, J., and Menge, B. A. 2003. Mussel
680 disturbance dynamics: signatures of oceanographic forcing from local interactions. *The*
681 *American Naturalist* 161:889–904.
- 682 Guttal, V. and Jayaprakash, C. 2009. Spatial variance and spatial skewness: leading
683 indicators of regime shifts in spatial ecological systems. *Theoretical Ecology* 2:3–12.
- 684 Hansen, G. J., Ives, A. R., Vander Zanden, M. J., and Carpenter, S. R. 2013. Are rapid
685 transitions between invasive and native species caused by alternative stable states, and
686 does it matter? *Ecology* 94:2207–2219.
- 687 Irvine, M., Jackson, E., Kenyon, E., Cook, K., Keeling, M., and Bull, J. 2016. Fractal
688 measures of spatial pattern as a heuristic for return rate in vegetative systems. *Royal*
689 *Society open science* 3:150519.
- 690 Kéfi, S., Eppinga, M. B., de Ruiter, P. C., and Rietkerk, M. 2010. Bistability and regular
691 spatial patterns in arid ecosystems. *Theoretical Ecology* 3:257–269.
- 692 Kéfi, S., Guttal, V., Brock, W. A., Carpenter, S. R., Ellison, A. M., Livina, V. N., Seekell,
693 D. A., Scheffer, M., van Nes, E. H., and Dakos, V. 2014. Early warning signals of ecological
694 transitions: methods for spatial patterns. *PloS One* 9:e92097.
- 695 Kéfi, S., Holmgren, M., and Scheffer, M. 2016. When can positive interactions cause
696 alternative stable states in ecosystems? *Functional Ecology* 30:88–97.
- 697 Kéfi, S., Rietkerk, M., Alados, C. L., Pueyo, Y., Papanastasis, V. P., ElAich, A., and
698 de Ruiter, P. C. 2007. Spatial vegetation patterns and imminent desertification in mediter-
699 ranean arid ecosystems. *Nature* 449:213–217.

- 700 Kéfi, S., Rietkerk, M., Roy, M., Franc, A., De Ruiter, P., and Pascual, M. 2011. Robust
701 scaling in ecosystems and the meltdown of patch size distributions before extinction.
702 *Ecology Letters* 14:29–35.
- 703 Kéfi, S., Rietkerk, M., van Baalen, M., and Loreau, M. 2007. Local facilitation, bistability
704 and transitions in arid ecosystems. *Theoretical Population Biology* 71:367–379.
- 705 Khaluf, Y., Ferrante, E., Simoens, P., and Huepe, C. 2017. Scale invariance in natural and
706 artificial collective systems: a review. *Journal of The Royal Society Interface* 14:20170662.
- 707 Knowlton, N. 2004. Multiple stable states and the conservation of marine ecosystems.
708 *Progress in Oceanography* 60:387–396.
- 709 Lin, Y., Han, G., Zhao, M., and Chang, S. X. 2010. Spatial vegetation patterns as early signs
710 of desertification: a case study of a desert steppe in inner mongolia, china. *Landscape
711 ecology* 25:1519–1527.
- 712 Lübeck, S. 2006. Tricritical directed percolation. *Journal of statistical physics* 123:193–221.
- 713 Ma, S.-K., 2000. *Modern theory of critical phenomena*. Da Capo Press.
- 714 Maestre, F. T., Bautista, S., and Cortina, J. 2003. Positive, negative, and net effects in
715 grass–shrub interactions in mediterranean semiarid grasslands. *Ecology* 84:3186–3197.
- 716 Maestre, F. T. and Escudero, A. 2009. Is the patch size distribution of vegetation a suitable
717 indicator of desertification processes? *Ecology* 90:1729–1735.
- 718 Malamud, B. D., Morein, G., and Turcotte, D. L. 1998. Forest fires: an example of self-
719 organized critical behavior. *Science* 281:1840–1842.
- 720 Manor, A. and Shnerb, N. M. 2008. Facilitation, competition, and vegetation patchiness:
721 from scale free distribution to patterns. *Journal of theoretical biology* 253:838–842.
- 722 Manor, A. and Shnerb, N. M. 2009. Multiplicative noise and second order phase transi-
723 tions. *Physical Review Letters* 103:030601.
- 724 Manrubia, S. C. and Solé, R. V. 1997. On forest spatial dynamics with gap formation.
725 *Journal of theoretical biology* 187:159–164.
- 726 Meloni, F., Granzotti, C. R. F., Bautista, S., and Martinez, A. S. 2017*a*. Scale dependence
727 and patch size distribution: clarifying patch patterns in mediterranean drylands. *Eco-
728 sphere* 8:e01690.
- 729 Meloni, F., Nakamura, G. M., Granzotti, C. R. F., and Martinez, A. S. 2017*b*. Empirical
730 evidence reveals the phase diagram of vegetation patterns in mediterranean drylands.
731 bioRxiv page 171835.

- 732 Moreno-de las Heras, M., Saco, P. M., Willgoose, G. R., and Tongway, D. J. 2011. Assessing
733 landscape structure and pattern fragmentation in semiarid ecosystems using patch-size
734 distributions. *Ecological Applications* 21:2793–2805.
- 735 Newman, M. E. 2005. Power laws, pareto distributions and zipf's law. *Contemporary*
736 *physics* 46:323–351.
- 737 Pascual, M. and Guichard, F. 2005. Criticality and disturbance in spatial ecological sys-
738 tems. *Trends in ecology & evolution* 20:88–95.
- 739 Petrovskii, S., Mashanova, A., and Jansen, V. A. 2011. Variation in individual walking
740 behavior creates the impression of a lévy flight. *Proceedings of the National Academy of*
741 *Sciences* 108:8704–8707.
- 742 Plank, M. J. and Codling, E. A. 2009. Sampling rate and misidentification of lévy and
743 non-lévy movement paths. *Ecology* 90:3546–3553.
- 744 Reif, F., 2009. *Fundamentals of statistical and thermal physics*. Waveland Press.
- 745 Rietkerk, M., Dekker, S. C., de Ruiter, P. C., and van de Koppel, J. 2004. Self-organized
746 patchiness and catastrophic regime shifts in ecosystems. *Science* 305:1926–1929.
- 747 Rietkerk, M. and van de Koppel, J. 2008. Regular pattern formation in real ecosystems.
748 *Trends in Ecology & Evolution* 23:169–175.
- 749 Roy, M., Pascual, M., and Franc, A. 2003. Broad scaling region in a spatial ecological
750 system. *Complexity* 8:19–27.
- 751 Scanlon, T., Caylor, K., Levin, S., and Rodriguez-Iturbe, I. 2007. Positive feedbacks pro-
752 mote power-law clustering of Kalahari vegetation. *Nature* 449:209–212.
- 753 Scheffer, M., Bascompte, J., Brock, W., Brovkin, V., Carpenter, S., Dakos, V., Held, H.,
754 Van Nes, E., Rietkerk, M., and Sugihara, G. 2009. Early-warning signals for critical
755 transitions. *Nature* 461:53–59.
- 756 Schneider, F. D. and Kéfi, S. 2016. Spatially heterogeneous pressure raises risk of catas-
757 trophic shifts. *Theoretical Ecology* 9:207–217.
- 758 Sethna, J., 2006. *Statistical mechanics: entropy, order parameters, and complexity*, vol-
759 *ume 14*. Oxford University Press.
- 760 Solé, R. V., Manrubia, S. C., Benton, M., Kauffman, S., and Bak, P. 1999. Criticality and
761 scaling in evolutionary ecology. *Trends in Ecology & Evolution* 14:156–160.
- 762 Stanley, H. E. 1999. Scaling, universality, and renormalization: Three pillars of modern
763 critical phenomena. *Reviews of modern physics* 71:S358.

- 764 Stauffer, D. 1979. Scaling theory of percolation clusters. *Physics reports* 54:1–74.
- 765 Strogatz, S., Friedman, M., Mallinckrodt, A. J., McKay, S., et al. 1994. Nonlinear dynamics
766 and chaos: With applications to physics, biology, chemistry, and engineering. *Computers
767 in Physics* 8:532–532.
- 768 Stumpf, M. P. and Porter, M. A. 2012. Critical truths about power laws. *Science* 335:665–
769 666.
- 770 Turcotte, D. L. and Malamud, B. D. 2004. Landslides, forest fires, and earthquakes:
771 examples of self-organized critical behavior. *Physica A: Statistical Mechanics and its Ap-
772 plications* 340:580–589.
- 773 van de Koppel, J., Rietkerk, M., van Langevelde, F., Kumar, L., Klausmeier, C., Hearne,
774 J. W., van Andel, J., de Ridder, N., Skidmore, A., Stroosnijder, L., and Prins, H. H. T. 2002.
775 Spatial heterogeneity and irreversible vegetation change in semiarid grazing systems.
776 *American Naturalist* 159:209–218.
- 777 von Hardenberg, J., Kletter, A. Y., Yizhaq, H., Nathan, J., and Meron, E. 2010. Periodic
778 versus scale-free patterns in dryland vegetation. *Proceedings of the Royal Society of
779 London B: Biological Sciences* 277:20092208.
- 780 von Hardenberg, J., Meron, E., Shachak, M., and Zarmi, Y. 2001. Diversity of vegetation
781 patterns and desertification. *Physical Review Letters* 87:198101:1–4.
- 782 Weerman, E., Van Belzen, J., Rietkerk, M., Temmerman, S., Kéfi, S., Herman, P., and
783 de Koppel, J. V. 2012. Changes in diatom patch-size distribution and degradation in
784 a spatially self-organized intertidal mudflat ecosystem. *Ecology* 93:608–618.
- 785 Weissmann, H., Kent, R., Michael, Y., and Shnerb, N. M. 2017. Empirical analysis of
786 vegetation dynamics and the possibility of a catastrophic desertification transition. *PloS
787 One* 12:e0189058.
- 788 White, E., Enquist, B., and J.L., G. 2008. On estimating the exponent of power-law fre-
789 quency distributions. *Ecology* 89:905–912.
- 790 Wissel, C. 1984. A universal law of the characteristic return time near thresholds. *Oe-
791 cologia* 65:101–107.
- 792 Xu, C., Holmgren, M., Van Nes, E. H., Maestre, F. T., Soliveres, S., Berdugo, M., Kéfi, S.,
793 Marquet, P. A., Abades, S., and Scheffer, M. 2015*a*. Can we infer plant facilitation from
794 remote sensing? a test across global drylands. *Ecological Applications* 25:1456–1462.

795 Xu, C., Van Nes, E. H., Holmgren, M., Kéfi, S., and Scheffer, M. 2015*b*. Local facilitation
796 may cause tipping points on a landscape level preceded by early-warning indicators. The
797 American Naturalist 186:E81–E90.