Title:

A minimal "push-pull" bistability model explains oscillations between quiescent

and proliferative cell states.

Sandeep Krishna[a,1] and Sunil Laxman[b,1]

[a]Simons Centre for the study of Living Machines, National Centre for Biological

Sciences, Tata Institute of Fundamental Research.

[b]Institute for Stem Cell Biology and Regenerative Medicine.

GKVK post, Bellary Road

Bangalore 560065

[1]Correspondence: sandeep@ncbs.res.in and sunil@instem.res.in

Abstract

A minimal model for oscillating between quiescent and growth/proliferation states,

dependent on the availability of a central metabolic resource, is presented. From the

yeast metabolic cycles (YMCs), metabolic oscillations in oxygen consumption are

represented as transitions between quiescent and growth states. We consider metabolic

resource availability, growth rates, and switching rates (between states) to model a

relaxation oscillator explaining transitions between these states. This frustrated

bistability model reveals a required communication between the metabolic resource that

determines oscillations, and the quiescent and growth state cells. Cells in each state

reflect memory, or hysteresis of their current state, and "push-pull" cells from the other

state. Finally, a parsimonious argument is made for a specific central metabolite as the

controller of switching between quiescence and growth states. We discuss how an

oscillator built around the availability of such a metabolic resource is sufficient to

generally regulate oscillations between growth and quiescence, through committed

transitions.

Keywords: quiescence, growth, metabolic oscillation, bistability, acetyl-CoA, NADPH

Introduction

While all cells can exist in a variety of states, two opposite ends of the spectrum are the "growth" state (leading to mitotic division and proliferation), and a non-proliferative "quiescent" state. The quiescent state, operationally defined here as a reversibly non-dividing state, is the predominant state of all living cells (Lewis and Gattie, 1991; Gray *et al.*, 2004). Understanding how cells reversibly transition from a quiescent state, to a growth state coupled with cell division and proliferation (henceforth called "growth" in this manuscript) is therefore a fundamental biological question. Current explanations for how cells commit to growth and cell division account for metabolic regulation, biomolecule synthesis, and regulated progression through the cell cycle, presenting multiple, integrated mechanisms of information transfer within a cell that lead to the eventual growth outcome.

However, when a population of genetically identical cells are present in a uniform environment, how can individual cells within such a population decide to switch between a quiescent (effective "G0") state and a growth/proliferation state? Indeed, such heterogeneity of cell states within populations is widely observed and acknowledged. Numerous examples exist in nearly all systems studied, from simple eukaryotes like the budding yeast, to complex mammalian systems (Cooper, 1998, 2003; Coller *et al.*, 2006; Daignan-Fornier and Sagot, 2011; Klosinska *et al.*, 2011; De Virgilio, 2012; Dhawan and Laxman, 2015), with multiple molecular events correlating with transitions between growth and quiescence. For any population transitioning into either of these states, experimentalists have asked: (i) what hallmarks allow discrimination between actively proliferating and G0 cells? (ii) how do cells transit back and forth between these two states? And (iii) how are different signals processed and integrated into an appropriate cellular response? The regulation of the final cellular outcome occurs at multiple levels, including differential gene expression programs, and signaling responses to growth factors, which can be different depending upon the type of cell or organism studied. At its very core, however, this transition between quiescent and growth states is a metabolic problem; cells must be in a metabolic state capable of

committing to growth/proliferation, and must sense this state, which the pushes cells

towards growth. Indeed, several lines of evidence now reiterate a primary metabolic

determinant for cells committing to a growth state (exiting quiescence), or remaining in a

quiescent state (Futcher, 2006; Daignan-Fornier and Sagot, 2011; Laporte *et al.*, 2011;

Cai and Tu, 2012; De Virgilio, 2012; Lee and Finkel, 2013; Dhawan and Laxman, 2015;

Kalucka *et al.*, 2015; Kaplon *et al.*, 2015). While multiple factors can regulate the

transition between quiescence and growth, all such studies suggest that without this

core metabolic transformation, switching states is impossible. Given this absolute

metabolic requirement to switch to growth, if there is an isogenic ("identical") population

of cells present in a uniform environment, how can there be a two-state outcome where

some cells undergo growth/proliferation, while the rest remain quiescent?

Surprisingly, there are few rigorous theoretical, mathematical models that attempt to

provide a conceptual framework sufficient to explain this, and suggest experimentally

testable predictions. This is in contrast to the extensive, elegant, and often prescient

models that have been built to explain progress through the classical cell division cycle

(CDC), by incorporating existing experimental data of phase specific cell-cycle

activators and inhibitors (Tyson and Novak, 2001; Tyson *et al.*, 2003; Ferrell *et al.*,

2009; Tyson and Novák, 2015). Such modeling of the CDC has a long history

(examples include (Goldbeter, 1991; J, 1991; Norel and Agur, 1991; Novak and Tyson,

1993; Ferrell *et al.*, 2009; Tyson and Novák, 2015)), and these types of theoretical

studies have revealed biological possibilities that were experimentally determined only

much later (such as (Cross *et al.*, 2002; Pomerening *et al.*, 2003; Wei *et al.*, 2003;

Mirchenko and Uhlmann, 2010)). Given this, there is considerable value in building

coarse-grained but rigorous theoretical models to understand switching between

quiescence and growth states. In such a model, the switching between quiescence and

growth states could be treated as a biological oscillation (Tyson *et al.*, 2003; Novák and

Tyson, 2008; Tsai *et al.*, 2008; Ferrell *et al.*, 2009), while considering a dependence on

a metabolic "resource" as a driver of the oscillator. For building such a model, we

therefore require extensive experimental data from biological systems where metabolic

oscillations are demonstrably closely coupled with exiting quiescence/entering the CDC.

Such data are readily available from the budding yeast, S. cerevisiae. Yeast have been

the instrumental cellular model in revealing processes that define both the CDC

(Hartwell, 1974), and the quiescence cycle (Gray *et al.*, 2004; Daignan-Fornier and

Sagot, 2011; Daignan-Fornier B and Sagot I, 2011; De Virgilio, 2012; Dhawan and

Laxman, 2015). The classical CDC involves progression through the G1, S, and G2/M

phases. In contrast, during a quiescence (or effective "G0") cycle, cells remain non-

dividing, but can exit quiescence and enter the G1 phase of the cell cycle to

subsequently complete the CDC.

Experimentally dissecting specific processes driving entry into, and exit from,

quiescence (into the CDC) is challenging in asynchronous, heterogeneous cultures of

cells. However, synchronized yeast populations in well-mixed cultures (as manifest by

oscillations in oxygen consumption) have long been observed and studied using batch

and chemostat conditions limited for a carbon source (glucose or ethanol), which are

subsequently fed continuously with limited concentrations of glucose or ethanol

(Chance *et al.*, 1964; Hommes, 1964; Hess and Boiteux, 1971; Satroutdinov *et al.*,

1992; Keulers *et al.*, 1996; Jules *et al.*, 2005; Lloyd and Murray, 2005). Gene

expression studies from such glucose-limited yeast metabolic cycles or oscillations (we

will utilize the term YMC henceforth in this manuscript for consistency) showed that a

majority of the genome is expressed highly periodically, further revealing a molecular

organization of growth and quiescent states (Klevecz *et al.*, 2004; Tu *et al.*, 2005;

Futcher, 2006; Mellor, 2016). In general, both the shorter (Klevecz *et al.*, 2004; Murray

et al., 2007), and the longer oxygen consumption oscillations in yeast (Tu *et al.*, 2005)

showed this general pattern. Notably, genes associated with biosynthesis and growth

(comprehensively further described in (Brauer *et al.*, 2008)) typically peak during a high

oxygen consumption phase in the YMC (Tu *et al.*, 2005; Rowicka *et al.*, 2007; Slavov

and Botstein, 2011, 2013), while genes that mark autophagy, vacuolar function and a

"quiescence" state peak during a steady, low oxygen consumption phase. Strikingly, in

these continuous YMC cultures, cell division is tightly gated to a temporal window. Cells

divide synchronously only once during each metabolic cycle (Küenzi and Fiechter,

1969; Tu *et al.*, 2005; Robertson *et al.*, 2008; Laxman *et al.*, 2010) and remain in a non-

dividing state during the rest of the cycle. The non-dividing population in the low oxygen
consumption phase exhibits typical hallmarks of quiescent cells (Tu *et al.*, 2005, 2007;
Shi *et al.*, 2010; Cai *et al.*, 2011; Shi and Tu, 2013; Dhawan and Laxman, 2015).
Furthermore, in each YMC, during the tight temporal window when cells do divide, the
culture has two, visibly distinct sub-populations: dividing and nondividing (Tu *et al.*,
2005; Robertson *et al.*, 2008; Laxman *et al.*, 2010). These data have suggested a close
coupling between the metabolic and the cell division cycles. Importantly, the YMC itself
is metabolite/nutrient regulated, and controlled by the amount of available glucose. The
distinct phases of the YMC correspondingly show a separation of metabolic processes
(Tu *et al.*, 2005, 2007; Murray *et al.*, 2007; Machné and Murray, 2012), and several
lines of evidence suggest that key metabolite amounts are critical for entering or exiting
a proliferative or a non-proliferative state (Murray *et al.*, 2003, 2007; Tu *et al.*, 2007; Shi
et al., 2010; Cai *et al.*, 2011; Machné and Murray, 2012; Mellor, 2016). These studies
collectively indicate the following: (i) a separation of two states (proliferative, and
effectively G0) in cell populations, dependent on metabolic states, and (ii) a loose
metabolic framework within which it may be possible to study transitions between
quiescence and growth transitions. Thus, these studies provide extensive experimental
data using which a theoretical, mathematical model can be built to sufficiently explain
oscillations between a "quiescent" state and a "growth" state.
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Here, we use existing data from these YMCs to build a robust, general model for
oscillations between a quiescent and a growth state. Importantly, the model
necessitates the requirement of a tripartite communication - between the metabolic
resource, the quiescent cells, and the cells exiting quiescence and entering growth - in
order for the cells to sustain oscillation between these two states. The model oscillations
depend on an underlying bistability, suggesting that cells in either state exhibit
hysteresis, or memory, of their states. Finally, using this model, we show how two
central metabolites, thought to be critical for entry into a growth state, satisfy the
required criteria for the currency that controls oscillations between these two cell

states. Collectively, we provide a coarse-grained, sufficiency model to explain general

principles of how cells can oscillate between a quiescent and growth state, depending

upon amounts and utilization of an internal metabolic currency.

Results

Apparent bistable states during yeast metabolic cycles

Yeast cells grown to a high cell density (in batch culture mode) in a chemostat, and

when subsequently fed limited amounts of glucose medium, spontaneously undergo

robust oscillations in oxygen consumption (YMCs) (Figure 1A) and (Klevecz *et al.*, 2004;

Tu *et al.*, 2005; Murray *et al.*, 2007; Silverman *et al.*, 2010; Burnetti *et al.*, 2015), with

the period of each oscillation ranging from ~2.5-5 hours (Figure 1A). For these

oscillations to occur, the batch culture typically needs to first be starved for a few hours

(Figure 1A), during which time all glucose is depleted, and all cells enter a non-dividing

state (although the extended starvation is not an absolute requirement, as observed

historically in breweries). After starvation, when cells are continuously provided limited

glucose in the medium, the oscillations in oxygen consumption spontaneously start and

continue indefinitely (Figure 1A). Comprehensive gene expression analysis across

these longer-period oscillations (1.5-4.5 hr cycles) has revealed highly periodic

transcript expression (Tu *et al.*, 2005; Rowicka *et al.*, 2007), and proteins encoded by

these transcripts can be binned into three general classes (Figure 1B, 1C). These

represent "growth genes" during the high oxygen consumption phase, followed by the

rapid decrease in oxygen consumption coupled with "cell division" (Figure 1B, 1C)

(Kudlicki *et al.*, 2007; Rowicka *et al.*, 2007). The cells exhibiting the "growth" signature

during the high oxygen consumption phase all go on to enter and complete the CDC

(42). Finally, the YMC enters a state of ~stable oxygen consumption, where the gene

expression profile revealed a "quiescent"-like state (Figure 1B, 1C). Mitotic cell division

is tightly gated only to a narrow window (Figure 1B, 1C). Interestingly, in this phase,

only a fixed fraction of the cells (~35%) (and not all cells) divide during each cycle

(Figure 1D). During the stable oxygen consumption phase, there are almost no budding

cells observed (Figure 1D). Note: given that this is a controlled chemostat system, the

overall cell number/density is constant throughout these oscillations (Klevecz *et al.*,

2004; Tu *et al.*, 2005), which becomes important for our mathematical model.

Defining the two states and apparent bistability

If these data are more grossly binned into groups, there appears to be ~2 effective

equilibrium states in this system. If binned based on the gene/metabolic patterns, there

is the oxidative phase (high oxygen consumption) closely coupled to growth,

immediately followed by the reductive mitotic phase, which depends upon (and follows

directly from), the oxidative phase. Indeed, experimental data suggest that these two

steps, the growth and proliferation steps, are irreversibly coupled (Laxman *et al.*, 2010).

This can therefore be conceived as one bin, representing a "growth" state. The

extended, low oxygen consumption phase where there is a long, steady build-up of

resources, can be viewed as a second bin. Both these states or bins appear to be

somewhat stable, contained systems, with what appears to be a transition or inflection

point leading to a committed switch to the other state. Thus, there appears to be an

apparent cellular state bistability occurring during these oscillations in oxygen

consumption. The stable, low oxygen consumption phase can therefore be practically

envisioned as representing the non-dividing, "quiescent" state (Q), while the rapid

increase in oxygen consumption followed by the reduction in oxygen consumption

phase represents the "growth" state (G) (Figure 1E). Considering this, our objective was

to build a mathematical model that conceptualized the oscillations in oxygen

consumption as oscillations between these two (Q and G) states.

For this, we first needed to define what plausible, broad scenarios this YMC system

might fit into. We therefore considered the currently accepted explanations for

commonly observed cellular heterogeneity within clonal populations. Many microbial

cells at high cell densities put out "quorum/alarmone" molecules that affect the entire

population, and lead to collective behavior along with heterogeneity (Miller and Bassler,

2001; Schauder *et al.*, 2001; Whitehead *et al.*, 2001; Zhu *et al.*, 2003; Chen *et al.*, 2004;

Farewell *et al.*, 2005; Srivatsan and Wang, 2008). Other possibilities emerge from

metabolic resource sharing, seen widely in systems ranging from microbial populations

to cancer cells (Veening *et al.*, 2008; Cairns *et al.*, 2011; Campbell *et al.*, 2015, 2016).

This extends to regulation at the levels of metabolic specialization and stochastic gene

expression resulting in phenotypic heterogeneity (Avery, 2006; Ibanez *et al.*, 2013;

Holland *et al.*, 2014; Ackermann, 2015; Sumner and Avery, 2017). From within this

range of possibilities, we envisaged three general scenarios that could result in the type

of oscillations (Q <-> G) seen in the YMC, and could make biological sense (Figure 1F):

(i) there could be the production and secretion of a resource by a sub-population of cells

("feeders"), which is taken up by other cells that will go on to divide; (ii) there could be

the secretion and accumulation of a metabolite that is sensed and taken up by only

some cells (but is not consumed); (iii) there is a build up of a metabolite, which is

consumed by the cells at some threshold concentration (Figure 1F). Starting from these

scenarios, we built simple models to test which one could create an oscillatory system

between the two states, which can come from an apparent bistability in the system.

Figure 1: Apparent two-state bistability during Yeast Metabolic Cycles.

A) A representative YMC, indicating stable oscillations in oxygen consumption (based on dissolved

oxygen dO2) in yeast cultures, reflecting the yeast metabolic cycle. Note that the stable oscillations are

driven by restricted feeding.

B) A more detailed illustration of each oscillation cycle, also indicating the phases of the

YMC.

C) Functional outputs based on gene expression studies (from (33)), which clearly define the oxygen

consumption phases of the YMC into a general "growth/proliferation" phase, and a "quiescence" phase.

D) Observed cell division during the YMC. Cell division is tightly gated to a narrow window of the YMC.

Note that only a fraction of cells, and not all cells, divide during this window of each cycle.

E) Reducing the oxygen consumption (dO2) oscillation into a two-state (Q state and G state) system. The

apparent bistability is also illustrated.

F) Plausible biological scenarios that could result in an oscillation between Q and G states, based on

observed phenomena. These scenarios are considered for building the model.

A "push-pull" model, requiring communication between the Q state, G state and

the resource, produces oscillatory behavior

Model framework for a two-state yeast population

In order to model such a two-state population of cells, the variables to consider would

be the following: (a) The number of cells in the quiescent state and in the growth state,

(b) some indicator of resource availability (dependent on the accumulation and

consumption of the resource) which could modulate the switching rate between Q and

G states, and the growth rate.

Thus, using this framework, we build the following equations that can describe the

dynamics of a two-state population of yeast cells in a well-mixed system:

"Change in Q population over time":

$dQ/dt = \nu_{GQ}G - \nu_{QG}Q - \phi Q,$ (1)

"Change in G population over time":

$dG/dt = \gamma G - \nu_{GQ}G + \nu_{QG}Q - \phi G$, (2)

where $Q(t)$ is the number of cells in the quiescent state at time t, $G(t)$ the number of

cells in the growing/dividing state, each ν represents a switching rate, $\phi(t)$ is the

chemostat outflux rate (which could vary with time), and γ is the growth rate of cells in

the growing/dividing state. If we further assume that the chemostat is working in a mode

that maintains the total population (or density) of cells at some constant level, i.e., the

outflux from the chemostat balances the growth of cells at all times, this means $\phi(t) =$

$\gamma G/(G + Q)$. In this case, the population dynamics can be described by a single

equation:

$dq/dt = \nu_{GQ}(1 - q) - \nu_{QG}q - \gamma(1 - q)q$, (3)

where $q \equiv Q/(G + Q)$ is the fraction of cells in the quiescent state.

Next, we assume that the cells contain some 'resource' that they require for growth,

without making any further assumptions about the resource. Let $a(t)$ denote the

concentration per cell of this resource at time t, and let σ denote the rate at which

additional amounts of this resource enter each cell from the surroundings (where the

resource is replenished due to the influx of fresh medium into the chemostat). a is

depleted both by dilution due to the outflux (at a rate $\gamma(1-q)$ as explained above), as well

as by consumption by growing cells (this rate is also proportional to $\gamma(1-q)$, which is the

net rate of production of new cells). The dynamics of this resource can thus be

described by the equation:

"Change in resource over time":

$da/dt = \sigma - \mu\gamma(1 - q)a - \gamma(1 - q)a$,   (4)

where μ is a proportionality constant that sets just how much resource is consumed by a

growing cell, compared to the amount that is depleted by dilution.

In writing equations 3 and 4, we have assumed that all cells have the same amount of

this internal resource a. A less restrictive assumption that still gives the same equation

is to assume that a represents the average concentration of the resource across the

population of cells, but that the distribution of resource levels is similar for Q and G

cells. Further, the same equations also model the case where the resource is not an

intracellular one, but an extracellular one - σ then is just reinterpreted as the rate at

which the resource is added to the extracellular medium either by an external feed or by

secretion of the resource by the cells themselves (e.g., by making σ dependent on q).

By choosing which of the parameters in the above equations are zero or non-zero, and

how they depend on q and/or a, this framework can be used to model a variety of

scenarios, which subsume the broad, biological scenarios illustrated in Figure 1E.

These mathematically distinct scenarios are described below (and illustrated in Figures

2A and 2B):

1. A sub-population of feeder cells (in the Q state) secrete a resource that is sensed by

other cells that can grow and divide (G state); resource accumulation σ increases with

q.

Such a scenario can be modelled with the G cells either consuming the resource ($\mu \neq$

0), or only sensing but not consuming the resource ($\mu = 0$) in the processes of

growing/dividing. The growth rate in the G state may be a constant, or may depend on

the level of the resource (e.g., γ proportional to a). There are three sub-scenarios for

how cells may switch between the two states:

a. There is no switching between Q and G states (ν_{QG} and ν_{GQ} both zero).

b. There is random switching between Q and G states (ν_{QG} and/or ν_{GQ} are non-zero

constants).

c. Switching between Q and G states is dependent on cell density and/or the resource

level (ν_{QG} and ν_{GQ} both functions of q and/or a).

2. All cells produce and secrete a resource that is sensed only by a sub-population of

(G) cells that can grow and divide, i.e., σ is a constant. As in scenario 1, the G cells may

or may not consume the resource, the growth rate in the G state may or may not

depend on the level of the resource, and there are three sub-scenarios for how cells

may switch between the two states: no switching, random switching or density/resource

dependent switching.

3. There is a build up of a resource, which is directly supplied from outside into the

chemostat medium (σ is a constant). This metabolite is sensed or consumed by the G

cells when they grow/divide. Again, the growth rate in the G state may or may not

depend on the level of the resource and switching may work in one of three ways: none,

random or density/resource dependent switching.

While scenarios (2) and (3) may appear mechanistically very different, they are in fact

mathematically no different from each other; both result in a constant production of the

resource (Figure 2B). Hence, we need not distinguish between these two. Testing all

the scenarios above, using equations 3 and 4, we show in the next section that

oscillations are not possible in the absence of switching, or even with random switching,

when there is no substantial time delay between resource utilization and division events

(as assumed in writing equations 3 and 4). Thus, scenarios 1c, 2c and 3c are the only

possibilities left that give oscillations (Figure 2C). This means that the switching

between Q and G states is a stochastic event, but with a probability that depends on the

resource level, and/or the density of cells in the Q or G state, implying some form of

communication between the resource, the cells in the Q state and the cells in the G

state.

Figure 2: A "push-pull" model for oscillations arising from an underlying bistability between Q and G

states.

A) A range of biologically plausible scenarios from Figure 1F, now broken down into precise categories,

where parameters affecting the rates of proliferation (g), switching

between Q and G states (n), as well as consumption (m) and supply (s) of the resource are included. The

variations in these parameters are used to build and test our model.

B) Schematic illustration of Figure 1A, indicating feedback loops and parameters considered, to test for

possible oscillations between Q and G states. For clarity, potential feedback loops caused by the

parameters being dependent on the resource a are not shown, but are included in our models.

C) A hysteretic oscillator, based on switching between Q and G states, a required communication

between Q, G and the resource, and the oscillation of the amounts of resource itself that controls the Q<-

>G transitions (see the Methods for the parameter values that that produce this dynamics). In the left

panel: The thin black curve shows the path traced by the oscillation in the q-a plane, the thick dashed line

is the curve along which production of resource exactly balances consumption/dilution, and the solid

black dots trace the high and low branches of the steady state q levels when the resource level is held

constant (the grey rectangle indicates the region of bistability). In the right panel: blue and green curves

show, respectively, the fraction of quiescent cells and the resource level as a function of time.

Some necessary conditions for oscillations

Within the framework of our model we can show that a density-dependent switching rate

is necessary to get oscillations.

(i) *No oscillations in the absence of switching:*

When both ν_{QG} and ν_{GQ} are zero, then equation 3 above becomes:

$dq/dt = -\gamma(1 - q)$. (5)

As long as γ is always positive, irrespective of its dependence on *a*, this has only one

stable steady-state solution, *q* = 0, because the rate of change of *q* is always negative.

And this is globally stable, i.e., every initial value of *q* (except *q* = 1) will flow to *q* = 0.

The *q* = 1 state is an unstable steady state, i.e., any fluctuations away from it, however

small, will result in the system moving to *q* = 0. Thus, there can be no oscillations in the

absence of switching.

(ii) *No oscillations with constant parameters:*

When all the parameters in equations 3 and 4 are constants, independent of q and a,

then no oscillations are possible because eq. 3 becomes independent of eq. 4, and

therefore, being a one dimensional ordinary differential equation without explicit time-

dependence, cannot show oscillations (an oscillation in q requires that dq/dt take both

positive and negative values for the same value of q, for at least within some range of q,

and this is not possible for a 1D ordinary differential equation).

(iii) *No oscillations for random (density-independent) switching:*

A less restrictive assumption is that ν_{GQ} and ν_{QG} are constants (which includes zero -

we've already examined the case where both are zero above), but γ and σ may be

functions of q and/or a. In the scenarios we examine, γ may be an increasing function of

a (all scenarios), while σ may be an increasing function of q (scenario 1). In this

situation, the dependence of each variable on the other is `monotonic' (dq/dt is a

decreasing function of a, while da/dt is an increasing function of q). Equations with such

monotonic dependencies have been studied mathematically in detail (Pigolotti *et al.*,

2007; Tiana *et al.*, 2007), which show explicitly that when such a coupled set of

equations has only two variables (here, q and a), then sustained oscillations are not

possible. Intuitively, there is not enough time delay in such a small two-leg feedback

loop to destabilise the overall negative feedback that pulls the variables into a single

stable steady-state value.

Hysteretic oscillator based on the two-state model

Apart from there being broadly two states, a second crucial observation from the

experiments is that there is a distinct separation of timescales. The transitions from a

situation where almost 100% of cells are in the Q state to one where 30-40% are in the

G state, and vice versa, are very rapid. Whereas, between these two transitions the

dynamics proceeds on much slower timescales. A simple way to obtain such a two-

timescale oscillator from this two-state model uses the strategy of `frustrated bistability'

previously suggested by (Krishna *et al.*, 2009). It requires three ingredients: (1) a

negative feedback loop between q and a, (2) bistability in q in the absence of the

feedback, and (3) the assumption that changes in q happen on a relatively fast

timescale compared to changes in a. While the first can be achieved in several ways,

the two simplest, biologically plausible, scenarios are where growing cells consume a

resource a and: (i) the growth rate γ is proportional to the resource a; or, (ii) one or both

switching rates depend on a such that the net switching rate from G to Q decreases with

a. However, the third requirement of separation of timescales means that the switching

rates must be at least several-fold higher than γ and σ. This means that the term $\gamma q(1-q)$

in equation 3 is practically negligible and hence the dependence of γ on a, or lack of it,

would have little effect. We therefore concentrate on the case (ii) where the switching

rates depend on a to implement the negative feedback, and for simplicity keep γ

independent of a.

Bistability in q in the absence of the feedback implies that when a is kept fixed, for some

range of a values, equation 3 should allow two stable steady state levels of q, one lower

and one higher. This is shown in Figure 2C left panel, where one can see the high and

low 'branches' traced by the solid black circles - every point on these branches is a

stable steady state q can attain for the corresponding a value, using a version of

equation 3 derived from scenario 3 in Figure 2B (see Methods for the full equation).

When the resource a is sufficiently small, then there is only one high steady-state level

possible for q. Similarly, when a is sufficiently large, there is only one low steady-state

possible. However, for intermediate values of a, the system exhibits bistability and both

low and high steady-state levels co-exist. In this bistable region, which steady-state

level q attains depends on where it started (i.e., its 'initial condition'). Importantly, in

these oscillations, the system exhibits a 'memory' (or a 'hysteresis') - the steady-state

level that q eventually settles into depends on the history of the system.

When there exists such bistability, then one can get oscillations from the system

described by equations 3 and 4, provided the switching rates are a few-fold faster than

the rates of consumption and accumulation of the resource (Krishna et $al.$, 2009), as

follows: when q is high, a increases due to lack of consumption, so the system creeps

along the high branch in Figure 2C left panel (see the trajectory traced by the thin black

line) until it hits the edge of the bistable region. At that point, cells start switching to the

G state, which happens relatively rapidly due to the separation of timescales. Thus, the

trajectory "falls off" the edge down to the low branch. On the low branch, with more G

cells, the now increased consumption of the resource causes a to start decreasing,

leading to the system creeping down along the low branch. When the system reaches

the left edge of the bistability, the trajectory jumps up to the high branch as cells rapidly

switch to the Q state. For a range of parameter values, this settles into a stable

oscillation, as shown in Figure 2C right panel, which shows how q and a vary with time

as one follows the black trajectory in Figure 2C left panel.

For this kind of oscillation, as we have demonstrated in the previous section, ν_{QG} and/or

ν_{GQ} must necessarily be functions of q, not constants independent of q. This can be

interpreted as a form of '*quorum/cell number sensing*' - implying some form of cell-cell

communication (or a cell density dependent phenomenon). More specifically, we find

that choosing either ν_{QG} to be a decreasing step-function of q (as in Figure 2C), or ν_{GQ}

to be an increasing step-function of q (see Supplemental Figure S1) is sufficient to

produce frustrated bistability. Other shapes that we have not explored may also produce

bistability, and hence oscillations. However, our purpose here is not to find the 'best-fit'

model, but rather to demonstrate the basic ingredients which are sufficient to produce

hysteretic oscillations that are similar to the experimental observations. The requirement

for ν_{QG} to be a decreasing step-function of q, or ν_{GQ} to be an increasing step-function of

q, is basically a requirement for a "push-pull" mechanism - the more the Q cells, the

more other Q cells get pulled to remain in that state, and the more G cells get pushed to

switch away from their state, and vice versa. Irrespective of the precise molecular

means by which this is achieved, cell-cell communication is a necessary ingredient for

implementing such a push-pull mechanism.

Possible variations in the shape of the oscillations

From our gross model explained in Figure 2, we obtain predictable oscillations with a

specific pattern. The model oscillations exhibit a fast drop in q when exiting the

predominantly quiescent phase, followed by a slow(er) drop, and then a rapid rise back

to a high q level. Experimentally however, a few variations within the general oscillation

patterns are known to occur, depending upon the strain background (Burnetti *et al.*,

2015). In the CEN.PK strain (our major reference system, from where the gene and

metabolite oscillation datasets were obtained (Tu *et al.*, 2005, 2007; Mohler *et al.*,

2008)) dO2 levels (which we equate with q) show a fast drop, a slow further drop, and

then a rapid rise (Figure 3A scenario (i)). However, as comprehensively described in

(Burnetti *et al.*, 2015), three other variations have been extensively documented.

Following a fast drop in dO2 levels, some strains then show a slower drop followed by a

more extended low dO2 phase (bump), and a fast rise in dO2 (Figure 3A, scenario (ii)).

Other strains show an overall fast drop in dO2, an extended low dO2 phase and bump,

and a fast rise (Figure 3A, scenario (iii)), or a fast drop in dO2 (increased oxygen

consumption), followed by a slower, extended rise in dO2 (Figure 3A, scenario (iv)).

Can our model explain this small diversity of shapes seen during the overall drop and

rise in oxygen concentrations?

In the model, the shape observed depends on the shapes of the two branches of q

steady-states (solid black circles in Figure 2C, left panel). Because the lower branch

starts at a q value of around 0.5 and then increases as a increases, therefore there is a

slow drop in q after the fast drop. To produce the experimental dO2 oscillations in other

yeast strains (as shown in Figure 3A), the lower branch must have a different shape.

For example, for strains which show a slow increase after the first rapid decrease of q,

the low q branch must *decrease* as a increases. Similarly, the other waveforms would

involve other shapes of the lower or higher branches. In Figure 3C we show that simple

changes in the dependence of the switching rate ν_{QG} on a produce different waveforms

for the oscillations. Here we've shown how to get different shapes of the low-q phase of

the oscillation by manipulating the lower branch of the bistability - changes to the high-q

phase could similarly be easily made by manipulating the upper branch. The main point

is that the shape of the waveform is primarily determined by the shape of the bistability

branches, which in turn are determined by how ν_{QG} and ν_{GQ} depend on q and a. Thus,

our model predicts that these switching rates are what must vary between strains that

show different oscillation waveforms.

Figure 3

Figure 3: Diverse waveforms in the oscillations: experimentally observed and model predictions.

A) Experimentally observed patterns of oscillations in dissolved oxygen/ oxygen consumption, which is

dependent on yeast strain backgrounds and chemostat growth conditions.

B) Altering the communication loops between Q, G and a, to change the overall oscillation waveform.

Here g (growth rate) is constant and n_{QG} is a decreasing step-function of q. To obtain different

waveforms, we vary the way the step function parameters n_m depends on a.

C) Predicted oscillation patterns from the model (as altered described in panel B). The illustrated panels

cover the range of waveforms observed experimentally in panel A. (i) same as Fig 2C; n_m decreases with

a. (ii) n_m first increases then decreases with a. (iii) n_m first decreases and then increases with a. (iv) n_m

increases with a. Additionally, in all four cases, K increases with a and other parameters have been

chosen so that the time period of oscillations is close to 4 hours (see Methods for the full equations, with

parameter values, for each case).

Predicting oscillatory outcomes based on resource availability

We have used scenario 3c (from Figure 2) to produce oscillations in Figures 2 and 3

above. We reiterate that mathematically scenarios 2 and 3 are the same, so scenario 2c

can produce exactly the same oscillations. Further, we also find that scenario 1c (where

the resource is not supplied externally, but produced/secreted by only the Q cells) is

also capable of producing similar oscillations, based on highly constrained choices for

how the production rate of the resource (σ) depends on *q* and *a* (see Supplemental

Figure S2). Thus, while scenarios 3c and 2c are identical, all three scenarios, 1c, 2c, 3c,

with appropriate choices for how the switching rates, and production and consumption

depend on the resource and fraction of quiescent cells, are sufficient to explain the YMC

oscillations. Scenario 2c and 3c are largely indistinguishable, and both appear

biologically most plausible. Given our experimental understanding of the YMC (and the

need for a consumable resource, glucose, to control the oscillations), we think scenario

3c is most likely (and we will explore this further in a subsequent section).

Breakdown of the oscillations.

In Figures 2 and 3 above, we have chosen the particular "default" values of each of the

model parameters such that the oscillation period became approximately 4 hours, to

match the experimental observations in Figure 1. Of course, varying these parameter

values changes the time period, and for large enough variation the oscillation may also

disappear. Our model predicts how the oscillation shape and period will vary, and when

oscillations will break down, in response to experimentally tunable parameters. For

instance, Figure 4A shows how the oscillations change as the resource production rate,

σ, is varied around its default value, for the same equations that produced the

oscillations in Figures 2 and 3. When σ is decreased below the default value, the

oscillation period initially increases, with more time being spent in the high-q phase. For

low enough σ, the model exhibits damped oscillations, and then as σ is lowered further,

the model exhibits the absence of oscillations, with q settling into a high steady-state

value (see Figure 4A, and also Supplemental Figure S3 for more such plots). When σ is

increased from its default value, we again find that the period initially decreases, with

less time being spent in the high-q phase. We are able to produce oscillations having a

time period as low as ~2.5 hours (see Figure 4A(iii)). When σ is increased beyond this,

the oscillation period starts increasing again, and the low-q phase of the oscillation

starts becoming pronounced (see Supplemental Figure S3). Eventually, the oscillations

disappear, with q settling into a (relatively) low steady-state value. These predictions

largely mirror known experimental observations, where decreasing or increasing feed

rate (at these scales) control oscillations similarly.

The resource production rate σ is a parameter that can be tuned relatively easily in a

chemostat by controlling the amount of fresh glucose or ethanol being supplied per unit

time. However, another parameter that may be tunable by genetic modifications is γ, the

growth rate of cells when they are in the G state. Figure 4B shows how the oscillations

vary as γ is varied. The results are qualitatively similar but inverse to what was observed

with σ variation - an increase in γ from the default value results in an increasing period,

damped oscillations and eventually no oscillations, while a decrease first results in a

decrease of period, then a distorted shape and increasing period (see Supplemental

Figure S4 for more such plots).

The location of the dashed black lines in Figure 4 (left panels) help to understand this

behaviour. Each dashed line traces the *q* and *a* values where resource production

exactly balances resource consumption/dilution. To the right of the line the production is

less than the consumption so the resource must decrease, and vice versa to the left of

the line. The closer one is to the dashed line, the slower the rate of change of *a*. As

explained in (Krishna *et al.*, 2009), oscillations occur only when this dashed black line

passes between the upper and lower bistable branches (solid black circles) - because

then the resource keeps increasing (decreasing) when it reaches the end of the high

(low) branch making the trajectory "fall off the edge" and continue the oscillation. When

σ is decreased the dashed black line shifts leftward in the plot, coming closer to the high

q branch which causes the oscillating trajectory to spend more and more time on the

high *q* branch (because it is closer to the dashed line and so the resource accumulates

slower). Eventually, as the dashed line just touches the high *q* branch, the time period of

oscillations increases to infinity (logarithmically – see Supplemental Figure S5). For σ

values lower than this critical value there is no sustained oscillation and the system

settles into a steady-state on the high branch at the point where it crosses the dashed

line. A similar behaviour happens as σ is increased and the dashed line comes closer to

the lower branch, with the only difference being that the oscillating trajectory spends

more time at lower *q* values.

A universal feature of the YMC oscillations seen in diverse yeast strains is that the time

period of the oscillations decreases with an increase in the dilution/supply rate in the

chemostat. The time period appears to be dominated by the time spent in the high-*q*

phase, which also increases with dilution/supply rate, whereas the time spent in the low-

q phase is less and *decreases* slightly with increase in the dilution/supply rate

(described in (Burnetti *et al.*, 2015)). As described above, in our model, we find that as

we vary σ or γ, there are two regimes. In one the time period is dominated by the high-*q*

phase, and the behaviour matches the above experimental observations (see

Supplemental Figure S5). However, there is also another regime, where the time period

is dominated by the low-*q* phase. Our model therefore predicts that: (i) the observed

YMC oscillations are closer to the lower end of the σ range that produces oscillations,

so one should be able to increase σ more than decrease it before breaking the

oscillations, and (ii) if one increased σ enough while remaining in the oscillatory regime

one should observe low-*q* dominated oscillations such as those shown in Supplemental

Figure S3. These are both testable predictions of our model.

Figure 4

Figure 4: Breakdown of oscillations.

A) Varying the rate of production of resource σ. (i) σ = 0.346 hr, (ii) σ = 0.400 hr

(default parameters, same as Figure 3), (iii) σ = 0.866 hr.

B) Varying the growth rate of cells γ. (i) γ = 0.500 hr, (ii) γ = 1.665 hr (default

parameters, same as Figure 3), (iii) γ = 2.000 hr. Equations used, and other parameter values, are the

same as those that produced Figs 2C and 3C(i).

Acetyl-CoA and NADPH satisfy the requirements of the consumable resource that

controls oscillations between Q and G states

Based on our model, the metabolic resource oscillates with a unique pattern, and this

drives the oscillation between the Q and G states. From the model, some resource

builds up within the cell, and is highest at the point of commitment to the switch to the G

state (Figure 5A). It is then rapidly consumed/eliminated to fall below a certain

threshold, resetting the oscillation, after which the cycle of building up for consumption

resumes. When superimposed to the actual YMC phases (and the Q to G switch), this

build-up of the resource would necessitate its highest amounts at the beginning of the

phase where cells commit to entering high oxygen consumption (Figure 5A). We note

that these features of the resource oscillation are a very robust prediction of our model.

Across all the oscillations in Figs 2-5 we see the same behaviour, and we would see

this for any parameter choice that gives oscillations because this behaviour depends

only on our assumption that the resource is consumed by growing/dividing cells and not

by quiescent cells. Therefore, according to our model, in order for any metabolite to be

the resource that controls the oscillation between the two states, this molecule must

fully satisfy the above criteria. Furthermore, for completing this switch to the G state, the

metabolite must be able to drive all the downstream biological events for growth. So do

any central metabolites satisfy these requirements, and could therefore be the internal

resource that controls these Q-G oscillations?

Figure 5: Acetyl-CoA satisfies the requirements for the metabolic resource

controlling the Q and G oscillations.

A) Predicted pattern of oscillation of the resource, during the Q and G oscillations,

based on the model (top panel, same oscillations as Figure 2C), and experimentally

observed oscillations of acetyl-CoA and NADPH during the dO2 oscillations (bottom

panel).

B) Predicted phase portrait of the the fraction of quiescent cells vs the resource per cell

based on the model (top panel), and experimentally observed oscillations in dO2 and

acetyl-CoA.

C) Predicted effect on the oscillation waveforms and the Q and G states, when a bolus

of the resource is added to cells in the Q state (see Methods for details), vs

experimentally observed data on oxygen consumption when a resource, acetate (the

trace is similar with for resources like ethanol, acetate, acetaldehyde, glucose) is added

to cells in the low oxygen consumption phase. Supplemental Fig S6 shows how the

response varies as the time of adding the bolus is varied.

D) Acetyl-CoA as a central regulator of a switch to the growth (G) state. The schematic

illustrates a cascade of biological processes leading to growth that acetyl-CoA amounts

regulate (coupled with coincident, required NADPH utilization). Note that all resources

that reset the oscillations, as indicated in (C) are utilized after they are converted to

acetyl-CoA.

Comprehensive datasets of 50-100 oscillating metabolites in the YMC exist (Murray *et*

al., 2007; Tu *et al.*, 2007; Mohler *et al.*, 2008). From these studies, the oscillations of

only two metabolites, acetyl-CoA and NADPH, fully fit the criteria demanded by our

model. The acetyl-CoA and NADPH oscillations as a function of the metabolic cycle,

and transitions between the Q and G state are shown in Figure 5A and 5B. The

oscillations of acetyl-CoA during the YMC almost perfectly superimposes with the

oscillation pattern of the hypothetical metabolic resource predicted by the model (Figure

5A). We plotted phase diagrams of the fraction of quiescent cells vs the amount of

resource in the cell (from the model), and also plotted phase diagrams from

experimental data for the dO2 oscillations plotted against acetyl-CoA amounts (Figure

5B). The two phase diagrams (from the model, and from experiments) strikingly

resemble each other (Figure 5B). This is despite the fact that the experimental data for

acetyl-CoA is of low resolution, with only a few sampling/time points covered, and also

only reflects overall (bulk population) measurements of acetyl-CoA, suggesting that the

actual phase diagram might be even more similar. Thus the model appears to capture

key universal features of these yeast oscillations, including the point of exit from low

oxygen consumption (Q) to high oxygen consumption and back (G), and the parameters

important in the waveform (i.e. the low and high oxygen consumption phases are

important, while the precise form of the dip and increase in dissolved oxygen may not

be so). The model also supports an inference that the acetyl-CoA oscillations are

sufficient to explain the bistability between the Q and G states, and retains the

hysteresis component.

Using our model, we next simulated what would happen if the resource was increased

to just above the threshold level, at a different time. In our model, during normal

oscillations, the amount of resource steadily increases, while the cells are in the Q state.

In our simulation, we provided a single bolus of the resource, while cells were in the Q

state (Figure 5C). We observed a predictable, sharp exit from the Q state and entry into

the G state (Figure 5C), effectively resetting the oscillation, which then continued and

restored itself to the normal, ~4 hour period in the next cycle. We compared this to

available experimental data, where oscillations have been reset by adding a bolus of an

external agent, typically glucose, acetate, acetaldehyde or ethanol (Murray *et al.*, 2003;

Klevecz *et al.*, 2004; Tu *et al.*, 2005; Cai *et al.*, 2011). All these agents show near-

identical patterns of resetting of oscillations (exit from Q and entry into G), and a

representative figure (for acetate addition) is shown in Figure 5C. Here, cells exit the

low oxygen consumption phase and enter and exit the high oxygen consumption phase,

and subsequently quickly restore normal (in this case ~4 hr) oscillations. This simulation

can be done in any part of the oscillation, and whenever most cells are in Q, adding a

bolus of the resource similarly resets the oscillation (Supplemental Figure S6). Also

notably, adding this resource when cells have switched to the G state does not alter the

oscillations much (Supplemental Figure S6), which is also something widely established

in experiments. Thus, the oscillations predicted by the model very closely recapitulates

the patterns of oscillations observed in experiments, how the central, controlling

resource might oscillate, and how the oscillations are affected upon perturbing the

resource. This strongly suggests that the threshold amounts of the resource are

sufficient to set the oscillations and switching between Q and G states.

Multiple lines of experimental data suggest that these two metabolites, acetyl-CoA and

NADPH, are key in controlling exit from quiescence and entry into growth (Tu *et al.*,

2007; Cai *et al.*, 2011; Cai and Tu, 2012; Machné and Murray, 2012; Shi and Tu, 2013,

2014; Mellor, 2016). Based on our knowledge of the metabolic prerequisites for entering

growth, and known functional endpoints or outcomes of these two molecules (Figure

5D), we can now make a strong, parsimony based argument suggesting that oscillations

in these two metabolites are sufficient to control oscillations between the Q and G state.

Particularly, several lines of study suggest that the entry into growth (from quiescence)

depends on carbon source utilization (Shi *et al.*, 2010; Cai *et al.*, 2011; Daignan-Fornier

and Sagot, 2011; Laporte *et al.*, 2011). As pointed out earlier, studies from the yeast

metabolic cycle show that the oscillations depend upon carbon sources (primarily

glucose) (Klevecz *et al.*, 2004; Tu *et al.*, 2005), and oscillations can be reset (to enter

the growth program) by adding acetate, acetaldehyde, etc. (Murray *et al.*, 2003; Tu *et*

al., 2005; Cai *et al.*, 2011). Notably, these carbon sources end up being converted

directly to acetyl-CoA, and can only then be utilized (Figure 5D). Additionally, a growth

program will require not just sufficient energy (ATP) to sustain the anabolic processes

within it, but also activate a program boosting anabolic processes that lead to cell

division, including making enough lipid moieties required for cell membranes and other

constituents of a new cell. Notably, acetyl-CoA satisfies all these requirements in the

following manner (Figure 5D): it directly enters the TCA cycle to generate ATP (Nelson,

DL; Cox, 2017); it can be utilized for the biosynthesis of numerous cellular metabolites,

including fatty acids, sterols, and amino acids (Nelson, DL; Cox, 2017); and directly

regulates cell growth and ribosome biogenesis by the acetylation of histones at "growth

promoting genes", especially histones at ribosome subunit, tRNA and ribi genes, and

activates their transcription by the SAGA complex (Cai *et al.*, 2011). The genes that

breakdown storage carbohydrates (such as glycogen and trehalose) that produce

acetyl-CoA all peak before the maximal acetyl-CoA concentration (Tu *et al.*, 2005;

Kudlicki *et al.*, 2007). Finally, the exit from quiescence requires the liquidation of these

storage carbohydrates (Shi *et al.*, 2010; Laporte *et al.*, 2011; Shi and Tu, 2013), and

conversion to acetyl-CoA (and the subsequent gene expression program) (Shi and Tu,

2013). Perturbations in the ability to sense and utilize acetyl-CoA (particularly for the

gene expression program) completely abolish oscillations (Cai *et al.*, 2011).

Physiologically, this anabolic commitment also absolutely requires the process of

reduction for anabolic biosynthesis, and this reductive capacity is supplied by NADPH

(Nelson, DL; Cox, 2017) (Figure 5D). NADPH is primarily synthesized from the pentose

phosphate pathway, which branches from this same central carbon network, and this

NADPH will fuel the required reductive biosynthesis to make molecules required for

anabolism (Figure 5D). Finally, genes encoding proteins that increase the synthesis of

NADPH are similarly coincident with those that lead to the generation of acetyl-CoA,

and disrupting NADPH production slightly results in a collapse of oscillations (Tu *et al.*,

2005, 2007). Relatedly, studies from the YMC show multiple other metabolite

oscillations coupled to or dependent upon NADPH, although any hierarchical

organization was not immediately apparent (Murray *et al.*, 2007). Without a necessary

coupling of the two molecules, the overall process of entry to growth cannot be

completed. There is substantial data, particularly from the studies of various cancers, to

show the close coupling of acetyl-CoA and NADPH for growth (Heiden *et al.*, 2009), as

well as direct evidence of acetyl-CoA promoting NADPH synthesis (Patra and Hay,

2014; Shan *et al.*, 2014). Summarizing, based on the pattern of oscillation of the

resource predicted by our model, acetyl-CoA and NADPH (based on production and

utilization) satisfy sufficiency requirements to be the molecules that control the Q-G

state oscillations. Our model thus strongly supports an argument for oscillations in

acetyl-CoA being sufficient to control Q-G state oscillations.

Discussion

In this study, we present a simple frustrated bistability model to explain how the

amounts of an internal metabolic resource can determine oscillations between a

quiescent and growth state. For this, we relied on extensive data coming from the YMC,

and represented the oscillations in dissolved oxygen (seen during YMCs) as a reflection

of growth and quiescent states (Figure 1). Our model incorporates factors dependent on
growth rate and amounts of the resource, as well as switching rates (between the G and
Q states). Importantly, the model emphasizes a necessary communication between the
cells in the quiescent state and the growth state, both of which interact with the
metabolic resource during such transitions (Figure 2). Quiescent cells "push" cells in the
growth state into quiescence, and "pull" other quiescent cells to remain quiescent, with
the feedback requirements imposed by the resource being distinct and opposite for the
Q and G states. Given this communication requirement between the Q and G states,
our model suggests that such oscillations will eventually breakdown when the cell
numbers are small and cells are no longer in contact with each other (something that
has been experimentally observed (Laxman *et al.*, 2010)). This model also provides
insight into understanding the "growth/division" rate of cells once committed to growth.
While healthy debates continue on the rate of growth in a cell and stages of the cell
cycle (Johnston *et al.*, 1977; Conlon and Raff, 2003; Jorgensen *et al.*, 2004; Brauer *et
al.*, 2008; Goranov *et al.*, 2013). our model shows that it is sufficient for oscillations to
have a fixed "growth rate" once the metabolic resource has crossed its threshold
concentration, and triggered a committed growth program, after which the growth and
division process is no longer dependent on available nutrients. This is also analogous to
studies of the CDC, which are built around committed, "no return" steps that proceed at
constant, predictable rates once committed to. In our model, because there is a
timescale separation between growth and switching rates, making the growth rate
dependent on the resource would make some quantitative difference to the rate of
accumulation/consumption of the resource, but would leave the Q-G oscillations largely
unchanged. Finally, using a parsimony based argument, we suggest that acetyl-CoA
(along with NADPH) satisfies all requirements for the resource that drive these
oscillations between the Q and G states (Figure 5). With acetyl-CoA as a resource, our
model, which builds oscillations on an underlying hysteresis, reproduces universal
features observed in these yeast metabolic oscillations, and provides a fairly simple
sufficiency argument for how cells transition between Q and G states. We reiterate that
our model only provides a paradigm to explain how the oscillations in an internal
metabolic resource is sufficient to control oscillations between quiescent and growth

states. This allows for (but doesn't include) other necessary elements in cells (e.g.,

unique gene transcription programs, or subsequent metabolic events that typically must

follow), that may also be required to build a more detailed model for Q-G oscillations.

The kind of oscillator we have built falls under the class of "relaxation oscillators", which

have been used to model a very wide variety of phenomena, ranging from electronic

oscillations to oscillating chemical reactions (Balthasar, 1926; Strogatz, 1994). These

are a subset of several possible types of oscillators that arise in biological systems, and

are especially relevant for the CDC (Novák and Tyson, 2008; Tsai *et al.*, 2008; Ferrell *et*

al., 2009; Ferrell, 2011). Relaxation oscillators typically involve the cyclic slow build-up

of some quantity, like charge in a capacitor, until it reaches a threshold level which then

triggers a "discharge" event, resulting in a rapid drop of the quantity. Thus, relaxation

oscillators are often characterised by processes happening on two very different

timescales, with the time period mainly determined by the slow process (Tyson *et al.*,

2003; Novák and Tyson, 2008; Tsai *et al.*, 2008; Ferrell *et al.*, 2009; Ferrell, 2011). This

is why, in contrast to linear, harmonic oscillators, they can produce non-smooth

oscillations like a square or sawtooth waveform. We note that the YMC oscillations

show a clear signature of multiple timescales - in Figure 1 it is evident that the exit from

quiescence (fast drop in dO2), as well as the re-entry into quiescence (fast rise in dO2),

happen at much faster timescales than the other phases of the oscillation. In our

relaxation oscillator model of the YMC, these differing timescales arise from the fact that

the switching rates are an order of magnitude larger than the rates of production and

consumption of the resource, and even the growth rate of the cells. The latter processes

are therefore what determine the time period of the YMC. Within the class of relaxation

oscillators, our models fall into a sub-class that depends on an underlying bistability,

which is 'frustrated' (Krishna *et al.*, 2009). The bistability, and the resultant hysteresis,

are what determine the threshold points at which the behaviour of the system rapidly

switches between accumulating or consuming the metabolic resource. Interestingly, our

model necessitates this strong hysteresis element within the Q and G state cells. The

phenomenon of hysteresis has been extremely well studied (and established)

particularly during many phases of the classical CDC, or proliferation cycle ((Pomerantz

and McCloskey, 1990; Tyson and Novak, 2001; Solomon, 2003; Wei *et al.*, 2003; Angeli

et al., 2004; Han *et al.*, 2005; Ferrell *et al.*, 2009; Ferrell, 2011; Yao *et al.*, 2011) and

many more). In contrast, a hysteresis phenomenon has not been extensively explored

when cells transition between a growth state and an effective quiescence state. Yet, in

such conditions where the transition between the two states is substantially determined

by a metabolic oscillator, as seen in the YMC and several other studies from simple

models like yeast, the hysteresis phenomenon is clearly revealed by our model. Given

this, experimental studies can be designed to dissect the nature of this hysteresis

phenomenon.

General features emerging from the model to understand oscillations between

quiescence and growth:

Although our model is relatively simple, uses data from a fairly elementary system, and

makes minimal assumptions, it does surprisingly well to constrain the possibilities for

how transitions between quiescence and growth are regulated. The model successfully

captures universally observed waveforms of oscillations, can reset the oscillations, can

predict how the oscillations of a resource can control the two states, and can predict

breakdown of oscillation fairly well, as observed in experiments. From the very large set

of metabolites known to oscillate during the YMC (Tu *et al.*, 2007; Mohler *et al.*, 2008),

our model constrains possibilities to a few, that oscillate in a way that can permit such a

bistability to exist. From this, and consistent with extensive experimental data

(discussed earlier, and in Figure 5), it is possible to make parsimonious arguments for

acetyl-CoA (coincident with NADPH) as the metabolic resources controlling transitions

from quiescence to growth, and *vice versa*. Our model helps differentiate this small set

of metabolites from other metabolites that are important to maintain oscillations, but not

initiate them (i.e. they may only allow the cell to continue in one state, or the other). For

example, sulfur metabolism is critical to maintain oscillations (Murray *et al.*, 2003, 2007;

Tu *et al.*, 2007). It is also essential for the completion of a growth program, post entry

into the high oxygen consuming phase. But this metabolite peaks *after* acetyl-CoA in the

YMC (Tu *et al.*, 2007), and can be viewed as a consequence of initiating a growth

program, and also critical to sustain/complete this growth program, but not to initiate the

oscillation. Substantiating this explanation is the fact that sulfur metabolism is highly

dependent upon the utilization of NADPH for reduction, and NADPH (as described

earlier) is coincident with acetyl-CoA. A similar argument can be made for the

sustained, high respiration seen in the YMC, which produces ATP that will be required

to maintain the growth program once committed to by the cell. Separately, other studies

have shown that "quiescent" cells can show metabolic oscillations without entry into the

CDC (Slavov et al., 2011). Here, these cells appear to show a commitment to the CDC

during these oscillations, based on gene expression patterns (Slavov et al., 2011). This

can also be viewed through our interpretation of the commitment of cells to the CDC

due to a central resource. Cells will commit to the CDC, which however may not be

completed if a subsequent metabolic resource, normally dependent upon the

central/controlling resource (predicted to be acetyl-CoA/NADPH here), becomes

limiting. In other words, for a cell, usually if this committing resource is at the correct

threshold, other resources should not be limiting unless artificially constrained in an

experimental set-up. In (Slavov et al., 2011), the limiting resource was phosphate, which

typically should be available and not limiting, and be assimilated into nucleotides in an

NADPH and acetyl-CoA dependent manner. If in a specific instance this resource

becomes limiting, the cells would commit to the growth/CDC state, but will not be able to

complete this, and will fall back into the Q state.

Our model provides a foundation to build new models to resolve other aspects observed

during the YMCs. First, in each cycle of the YMC, a fraction of the cells exit quiescence

and divide. It is not fully clear if the same cell divides in each cycle, or if a cell that has

entered division in one cycle does not in the next, and so on. The decision to divide has

been viewed as a stochastic, but irreversible step (Laxman et al., 2010; Burnetti et al.,

2015). While our model as it stands cannot directly address these questions, the

dependence of the oscillations on the build-up and utilization of a specific resource,

allows the following argument to be made. First, the decision to divide in a cell would be

purely made by the amount of resource (acetyl-CoA) that has been built up in the cell.

Once acetyl-CoA reaches a certain threshold, the decision to divide is irreversible.

However, the build up of acetyl-CoA within an individual cell itself would be dependent

on small differences in overall metabolic homeostasis (compared to its neighbor), and

thus which cell reaches the threshold level first could be purely stochastic. Second, we

may speculate that if a cell has reached this threshold level and then used up its

resource during division, it is unlikely to be in a position to divide in the next round/next

cycle, whereas a cell that had not reached the threshold level in the previous cycle

would be best poised to divide instead. Our model does not take this into account, but it

provides a framework within which one could model the entire distribution of cells in

different Q/G states and with different levels of the resource. Despite the overall

stochastic aspect of Q-G transitions, such models would be able to make testable

predictions about the switching process even at the level of single cells. It is also

apparent that this level of synchrony requires high cell density in the system.

Separately, most studies have noted that upon initiating feeding in the chemostat, there

is a short period of tiny, non-robust oscillations. Based on our model, we would argue

that this is a situation where the quiescent cells are all now building up just sufficient

reserves of acetyl-CoA, within this stochastic process, and are starting to divide, but the

unusual steady-state condition in the chemostat will eventually lead to stable

oscillations.

Finally: Given the existing frameworks to describe Q-G state oscillations, our model is

necessarily coarse grained, and is intended only to build a more rigorous conceptual

framework within which to investigate the process of cells switching between

quiescence and growth states. For instance, it is straightforward to extend our models,

by adding space and diffusion processes, to account for scenarios where nutrients are

not well mixed and equally accessible, and where there is a high degree of spatial

rigidity within cell populations. It is also easy to alter other assumptions underlying our

model. For instance, our conclusions regarding acetyl-CoA being the driving resource

depend on an assumption we made in building the model that G cells consume the

resource. While this is biologically plausible, mathematically we could have assumed

the opposite, namely that the resource is consumed by cells in the Q state and not by

cells in the G state. In that case too our model could give similar oscillations - switching

rates would still need to be density/resource dependent, but the form of dependence on

resource would need to be reversed so that the high and low-q branches would be the

mirror images, with the low q branch being the only one at low a and the high q branch

being the only one at large values of a. And hence the waveform of the oscillating

resource would be flipped compared to Figures 2 and 3 - i.e., when q is high a would be

decreasing, while when q is low a would be increasing. If one could find a metabolite

that exhibited this waveform, then that metabolite would be an equally likely possibility

as a driver of the Q-G transition, except that it would have to act such that it caused a

switch from Q to G when it crossed a low threshold, or caused the opposite transition

when it crossed a high threshold. From the considerable data available, we have not

found a reasonable molecule with such a reversed waveform. Moreover, we know of no

process which consumes a metabolite in the Q state in the way described, so for now

acetyl-CoA driving the Q-G transition and being consumed during growth is the most

parsimonious explanation. Nevertheless, this shows how our framework could be easily

used in alternate scenarios.

Currently, existing experimental approaches to study such metabolically-driven Q-G

oscillations are very limited. Crude readouts, such as oxygen consumption, have very

limited resolution even to show the Q and G states, as the bistability begins to break

down. Gene expression analysis (even when done in single cells) is a late, end-point

readout which cannot explain this bistability but instead occurs after a switch. The key to

experimentally studying such bistability, therefore, will be the development of in vivo

intracellular metabolic sensors with excellent dynamic range and sensitivity, for

metabolites like acetyl-CoA or NADPH. This will allow the development of more precise

models to predict commitment steps, and identify differences within the population of

cells, that will help understand reversibility (between states), hysteresis and other

apparent phenomena.

Methods

Experimental methods and data sets:

Chemostat culture and cell division datasets: All dO2 data were obtained from YMCs

set up similar to already published data (Tu *et al.*, 2005, 2007; Kudlicki *et al.*, 2007;

Mohler *et al.*, 2008). In these studies, yeast cells were grown in chemostat cultures

using semi-defined medium, and yeast metabolic cycles were set up as described

earlier (Tu *et al.*, 2005; Tu, 2010). Data for cell division across three metabolic cycles

was obtained from earlier studies (Tu *et al.*, 2005; Laxman *et al.*, 2010). YMC gene

expression and metabolite datasets: Gene expression datasets were obtained from (Tu

et al., 2005; Kudlicki *et al.*, 2007), and metabolite oscillation datasets were obtained

from (Tu *et al.*, 2007; Mohler *et al.*, 2008; Cai *et al.*, 2011; Machné and Murray, 2012),

including acetyl-CoA oscillation datasets.

Parameter values and their q/a dependencies

Figures 2C, 3C(i), 4A(ii), 4B(ii), 5A and 5B (default choices):

To produce the oscillation shown in these figures, we make the following choices (within

scenario 3c):

$\gamma=1.665$ hr^{-1}. $\sigma=0.3996$ hr^{-1}, $\mu=1$, $\nu_{GQ}=16.65$ hr^{-1}, $\nu_{QG}=h(q)$, where $h(q)$ is the Hill function

$h(q) = \nu_m(1+\beta(q/K)^{20})/(1+(q/K)^{20})$ with $\beta=0.01$, $K=a^2/(0.75^2+a^2)$, $\nu_m=16.65\times(1.65-1.25K)$.

We use this Hill function with such a high Hill coefficient to approximate a step function

which drops rapidly from ν_m to $\beta\nu_m$ at $q=K$.

Figure 4, other panels:

The other panels of Fig 4 are made using exactly the same equations and parameter

choices as above, except for varying σ and γ as mentioned in the Fig 4 caption.

Figure 3C, other panels:

As above, except that

(ii) $\nu_m=16.65\times(1.65-1.25K)\times2.25K$ hr^{-1} and $\sigma=0.3596$ hr^{-1}.

(iii) $\nu_m=16.65\times(1.65-1.25K)+16.65\times1.85a^{10}/(200+a^{10})$ hr^{-1} and $\sigma=0.3297$ hr^{-1}.

(iv) ν_m=16.65x2.25K hr^{-1} and σ=0.3397 hr^{-1}.

σ values were varied in order to keep the time period close to 4 hours.

Figure 5C, addition of bolus:

Until time t = 11.5 hours, the simulation is the same as in Fig 2C. At t = 11.5 hrs, the

resource level is abruptly changed to 1.75 (just above its peak value in previous cycles,

which was 1.73), and then the simulation is continued with the same equations and

parameter values.

In all the above cases, the simulations were started, at t = 0 hours, with initial conditions

q=1 and a=10^{-6} (i.e., we start with all cells in a quiescent state and starved of the

resource). Simulations and figures were produced in Matlab. We used the ode45

differential equation integrator. The code is provided in Supplemental material. As extra

controls, we checked that the stiff solver ode15s also provided the same results for the

simulations in Fig 2 and 5, and a Mathematica notebook which repeats many of the

same simulations, using the default NDSolve algorithm within Mathematica, is also

provided with the Supplemental material.

1002

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