

1 **Cell adhesion and fluid flow jointly initiate genotype spatial distribution in**
2 **biofilms**

3
4 Short title: Cell adhesion and fluid flow in biofilm early development

5
6 Ricardo Martínez-García¹, Carey D. Nadell^{2,3*}, Raimo Hartmann², Knut Drescher², Juan A.
7 Bonachela^{4,5*}

8
9
10
11 1. Department of Ecology and Evolutionary Biology, Princeton University. Princeton NJ, 08544,
12 USA. 2. Max Planck Institute for Terrestrial Microbiology, 35043 Marburg, Germany. 3.
13 Department of Biological Sciences, Dartmouth College, Hanover, NH 03755, USA. 4. Marine
14 Population Modeling Group, Department of Mathematics and Statistics, University of
15 Strathclyde, Glasgow, G1 1XH, Scotland, UK. 5. Department of Ecology, Evolution, and Natural
16 Resources, Rutgers University, 14 College Farm Road, New Brunswick, NJ 08901, USA.

17
18 *Correspondence: juan.bonachela@rutgers.edu, carey.d.nadell@Dartmouth.edu

19

20 **Abstract**

21 Biofilms are microbial collectives that occupy a diverse array of surfaces. The function and
22 evolution of biofilms are strongly influenced by the spatial arrangement of different strains and
23 species within them, but how spatiotemporal distributions of different genotypes in biofilm
24 populations originate is still underexplored. Here, we study the origins of biofilm genetic structure
25 by combining model development, numerical simulations, and microfluidic experiments using the
26 human pathogen *Vibrio cholerae*. Using spatial correlation functions to quantify the differences
27 between emergent cell lineage segregation patterns, we find that strong adhesion often, but not
28 always, maximizes the size of clonal cell clusters on flat surfaces. Counterintuitively, our model
29 predicts that, under some conditions, investing in adhesion can reduce rather than increase clonal
30 group size. Our results emphasize that a complex interaction of fluid flow and cell adhesiveness
31 can underlie emergent patterns of biofilm genetic structure. This structure, in turn, has an outsize
32 influence on how biofilm-dwelling populations function and evolve.

33 **Author summary**

34 Biofilms are bacterial groups, often attached to surfaces, in which a broad variety of cooperative
35 and competitive interactions typically occur. The spatial organization of different strains and
36 species within biofilm communities strongly influences their global functioning, but little is known
37 about how such structure arises. Combining experiments on *V. cholerae* and simulations of a
38 cellular automaton, we show that the complex interaction between bacterial traits (cell adhesion)
39 and environmental factors (fluid flow intensity) strongly influences the early origins of biofilm
40 spatial structure. In most cases, we found that highly-adhesive strains form larger clusters than the
41 weakly-adhesive ones. Against intuition, however, we also found the opposite outcome: weakly-
42 adhesive tend to form larger clusters than the highly adhesive ones when flows are weak or the
43 population density of colonizing cells is high.

44

45 **Introduction**

46 In addition to living as planktonic cells in liquid environments, bacteria often form dense
47 conglomerates attached to surfaces, termed biofilms. Biofilms are one of the most widespread
48 forms of life on Earth, and they are deeply embedded into global scale processes such as
49 biogeochemical cycling [1]. They also play a central role in the interaction between bacteria and
50 multicellular organisms, including humans, as biofilm production enhances antibiotic tolerance [2]
51 and influences bacterial pathogenesis and microbiome functioning [3]. From a biotechnological
52 point of view, biofilms are used to purify wastewater and to control catalysis reactions, including
53 those involved with biofuels [4]. Biofilms are also the primary source of biological fouling in
54 industrial settings [5].

55 Within a biofilm, cells are typically embedded in a matrix of extracellular polymeric
56 substances (EPS) made of proteins, lipids, nucleic acids and polysaccharides [6]–[8]. The secretion
57 of the matrix, together with other products such as digestive enzymes, nutrient chelators, and
58 adhesins, provides biofilm-dwelling bacteria with increased metabolic versatility, tolerance to
59 exogenous stress and resistance to fluid shear [9]–[15]. The functioning and evolutionary stability
60 of behaviors that alter the local environment – including secretion phenotypes, which usually affect
61 nearest-neighbors the most strongly – in turn depend on the spatial arrangement of secreting versus
62 non-secreting strains and species (i.e., different genotypes) in a biofilm community [16]. For
63 example, intra-strain cooperative behaviors are more likely to be evolutionarily stable when
64 different cell lineages are segregated in space, with typical interaction distances between cells
65 being strongly influenced by the diffusivity of secreted products, biofilm architecture, and
66 environmental flow conditions [16]–[19]. Spatially constrained interaction is well known to be
67 important in ecology broadly, and there are numerous examples of spatial structure influencing
68 evolution in biofilm communities [20]–[22]. Thus, spatial structure in biofilms, once it arises, has
69 a large impact on their form and function. The means by which biofilm strain and species structure
70 originates in the first place, however, are less well understood.

71 At the early stages of biofilm formation, planktonic cells encounter and transiently adhere
72 to surfaces. Bacteria possess sophisticated mechanisms for deciding whether to remain in place,
73 depending on substratum properties and environmental quality [23]–[26]. Having committed to
74 biofilm formation, surface-residing cells secrete additional and diverse adhesion factors, including
75 extracellular matrix material. These secretions, in combination with growth, death, and steric
76 interactions between cells, strongly impact biofilm spatial organization [16], [27]–[30].
77 Environmental features, such as surface chemistry and fluid flow, are also key to biofilm
78 development. In cases where flow influences cell surface motility, flow regime and environmental
79 geometry can exert a dramatic effect on the spatial spread of surface-bound bacteria [31], [32].
80 Fluid flow is also likely to play a key role in the deposition and spatial arrangement of different
81 strains and species within biofilms [15], [33]–[35]. In spite of its putative importance, we have a
82 limited understanding of how flow, surface colonization processes, and cell adhesion interact to
83 influence the spatial strain structure of nascent bacterial communities. Targeting this knowledge
84 gap is the primary goal of the present study.

85 We performed experiments with matrix-producing or non-producing strains of the model
86 biofilm-forming bacterium *Vibrio cholerae*, the causative agent of pandemic cholera in humans.
87 We aimed to use a simplified, ecologically neutral scenario, in which mixed strains are genetically
88 identical except for fluorescent labels, to provide a first step towards understanding how key
89 environmental features interact with cell adhesion and population density to control the initial
90 distribution of cell lineages on a surface [36], [37]. Based on these experiments, we developed a

91 cellular automaton, with which we considered different scenarios that included varying flow
92 strengths, densities of founder cells, and variable cell adhesiveness. Our study of surface
93 occupation patterns motivated the use of spatial correlation functions as a quantitative method to
94 characterize the contribution of adhesiveness and flow regime on the origins of clonal clustering
95 spatial structure. The results, although obtained for *V. cholerae*, will more generally improve our
96 understanding of the patterns with which microbes colonize abiotic and biotic surfaces. These
97 initial patterns of surface occupation are key to the longer-term biofilm architectures that endure
98 to impact bacterial ecology, evolution, and pathogenesis.

99

100 **Results**

101 **Surface colonization experiments**

102 To isolate the influences of adhesiveness, flow, and population density on surface colonization
103 regimes, we used strains of *V. cholerae* without flagella that either produce extracellular matrix
104 constitutively, or not at all [38]. As *V. cholerae* does not use gliding or twitching motility to roam
105 on glass surfaces after attachment [39], differences in surface occupation by our strains could be
106 specifically attributed to their difference in matrix production. Individual cells of *Vibrio cholerae*
107 are capable of attaching to surfaces in the absence of extracellular matrix secretion, but matrix
108 production augments surface and cell-cell adhesion, and is essential for producing three-
109 dimensional biofilm structures. The direct contribution of matrix production to biomass
110 accumulation in biofilms, relative to loss of cells into the passing flow, has been demonstrated in
111 our previous work [38], [40]. Cell motility in the planktonic phase, which influences surface
112 exploration [39], [41]–[45], is not included here and will be the focus of future work.

113 For each strain (matrix-producing, and non-producing), a red- and blue-fluorescent version
114 was constructed by engineering fluorescent protein expression constructs on the chromosome.
115 Founder cells were inoculated in polydimethylsiloxane (PDMS) microfluidic chambers as 1:1 co-
116 cultures of the blue and red variants of the matrix-producing strain, or, in separate experiments,
117 blue and red variants of the matrix non-producing strain. Flow rate was maintained at 0.1 $\mu\text{L}/\text{min}$
118 through chambers measuring 500 μm wide, 100 μm tall, and 7000 μm long. Bacteria within a
119 growth chamber were thus identical regarding the production of matrix, differing only in their
120 color. Nutrients were continuously provided in the inflow. We focused on the early stages of
121 biofilm growth before large 3D structures could form; thus, growth was limited by the availability
122 of space on the surface, and not by the access to nutrients in the influent medium. It is important
123 to note that, even in these early phases of biofilm growth (once cell clusters reach 4-8 bacteria),
124 cells capable of producing matrix will have begun to do so [46].

125 Experiments were stopped when the biofilm population fully occupied the basal surface,
126 as judged by eye. The data generated by our experiments consisted of 2D surface occupation
127 patterns composed by clusters of different lineages that express either the blue or the red
128 fluorescent tag (Fig 1). Surface occupation was captured by fluorescence microscopy. Images were
129 acquired in the largest viewing fields allowed by our microscope constraints, measuring 60 μm x
130 60 μm (923 x 923 pixels), with 60 such viewing fields comprising an entire chamber. Note that,
131 since the snapshots analyzed in the experiments correspond to tiles within a larger total area in the
132 growth channel, there can be exchange of cells across tiles through detachment and re-attachment
133 of individuals. See Materials and Methods for a more detailed description of our experimental
134 approaches and strain engineering, and S1 Fig.a for a schematic representation of the experimental
135 procedure.

136

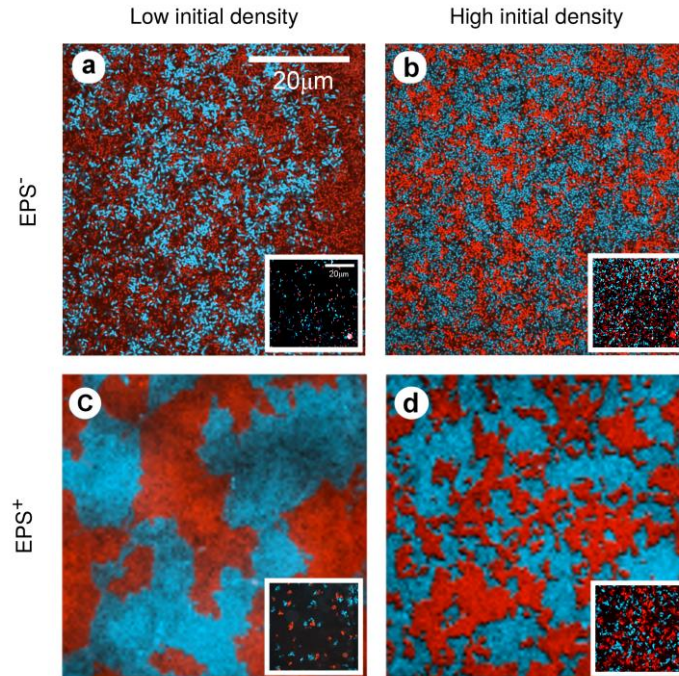


Fig. 1. Experimental colonization patterns. Snapshots of one field of view at confluence for both matrix-producing (bottom) and non-producing (top) strains at low (left) and high (right) initial cell densities. The inset of each panel shows the initial distribution of founder cells. Initial densities: a) 0.01 cells/ μm^2 , b) 0.162 cells/ μm^2 , c) 0.012 cells/ μm^2 , d) 0.113 cells/ μm^2 .

Modelling framework

To explore the mechanisms underlying the experimental results, and to extend our predictions to a broader set of environmental flow conditions and cell adhesion strengths, we developed a probabilistic cellular automaton capturing the essential features of the experimental system. In our model, we consider two strains with identical non-dimensional cell adhesiveness, σ , and initial density of colonizing cells, $\rho_0/2$, that compete for the occupation of empty space on a discrete two-dimensional lattice. The density of founder cells is defined by the fraction of initially occupied lattice squares. In the absence of extensive surface motility, adhesiveness varies inversely with the probability that a cell detaches from the surface. This may occur either because of shoving between cells or because of flow, which detaches cells and relocates them downstream. We will use here a real number in $[0,1]$ to represent adhesiveness, with $\sigma = 1$ indicating strong adhesion and $\sigma = 0$ weak adhesion. The only difference between strains within a given experiment is therefore a binary variable for the cell color, c , which is later used to analyze the arrangement of different cell lineages.

The dynamics of the model has two main ingredients: (i) birth and (ii) flow-induced cell detachment and relocation. We assume that these two processes are stochastic and independent (S1 Fig b,c). Time is discretized in short intervals of fixed length dt ; within each time step, a random cell reproduces (*i.e.* divides) with probability p_b , and another random cell may be detached and eventually relocated with probability p_d . The detachment probability depends on cell adhesiveness and flow intensity, whereas cell transport, both in the direction of the flow and transversely to it, is entirely determined by flow intensity, f , which we define using a normalized non-dimensional parameter that takes values in $[0,1]$. The flow structure in our microfluidic

165 devices is laminar, so we assume that flow intensity fixes the maximum distance that cells may be
166 transported downstream. $f = 1$ represents intense flows under which cells can be transported a
167 maximum distance equal to lattice length, and $f = 0$ represents no flow and therefore no cell
168 detachment and transport. Cell transport in the direction transverse to the flow is bounded by the
169 distance traveled downstream (see Materials and Methods). Since surface colonization occurs over
170 short time scales and resources are continuously supplied by the inflowing nutrient medium, we
171 do not include cell death in the model. In our experiments cells can in principle detach from one
172 viewing field and re-attach in another viewing field downstream; we implement this possibility in
173 our simulations using periodic boundary conditions. Cells that exit the system through one of the
174 borders due to long-range relocation re-enter through the opposite side, which is equivalent to cell
175 relocations originating upstream and balances out the anisotropic effects introduced by the
176 presence of a directional flow.

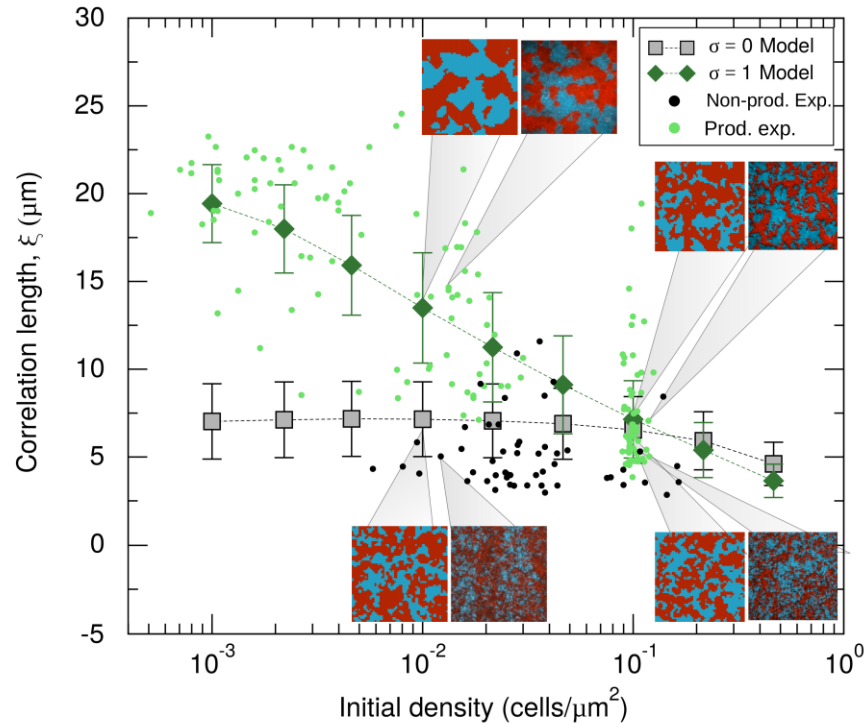
177 Finally, each run of the model was stopped when 95% of the positions of the lattice were
178 occupied, which avoids the high number of shoves that occur when surface coverage is nearly
179 complete and which have a negligible effect on the final coverage pattern. This condition is similar
180 to that used in terminating the experimental runs, which were stopped when the bottom surface of
181 the chamber was nearly completely covered by cells. See Materials and Methods for further details
182 on the modeling approach.

183

184 **Experimental output and model validation.**

185 To characterize the patterns of bacterial surface occupation obtained experimentally (Fig 1), we
186 measured their clonal correlation lengths, ζ , and studied their dependence on the initial population
187 density. The correlation length is obtained from the spatial autocorrelation function, $C(r)$, which
188 provides a measure of the order in spatially-extended systems by quantifying how its spatial
189 elements co-vary with one another on average, as a function of spatial separation distance r . For a
190 given separation distance r , the autocorrelation is positive if individuals separated by r tend to be
191 of the same type, negative if they tend to be of different types, and zero if there is no consistent
192 relationship between them. The correlation length is, thus, the shortest distance for which two
193 spatial elements of the patterns are statistically independent. See Materials and Methods for a
194 detailed definition of the correlation function. Because this distance is related to the typical cluster
195 size within the field-of-view, from an ecological perspective, the mean correlation length
196 quantifies the expected lineage segregation in the surface occupation pattern (see Materials and
197 Methods). When two matrix-secreting strains colonize the chamber, the correlation length of the
198 confluence pattern increases as the total initial density of cells decreases (green dots in Fig 2).
199 However, if the two strains are matrix non-secreting (and therefore only very weakly adhesive),
200 the correlation length does not show strong dependence on the initial density of cells in the
201 chamber (black dots in Fig 2). Note that the lowest initial coverage densities for the two cases are
202 different; matrix-secretors could be initiated at very low densities for which non-secreting strains
203 did not give viable results. This limitation on initial population density was most likely due to the
204 relative ease with which non-secreting strains are removed by flow.

205



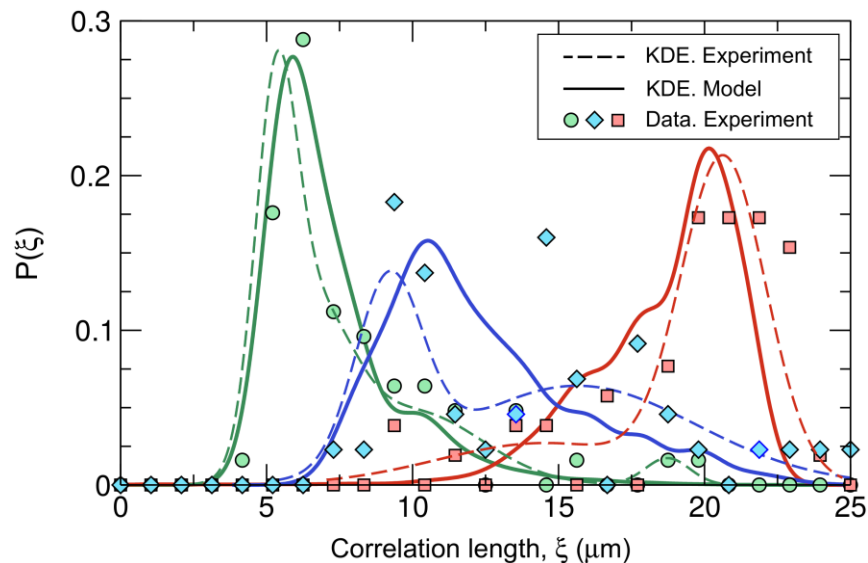
206
207 **Fig 2. Model validation: correlation length comparison.** Experimental correlation lengths
208 measured in the matrix-producing (pale green dots) and non-producing (black dots) strain, and
209 their model equivalent $\sigma = 1$ (dark-green diamonds), respectively $\sigma = 0$ (gray squares).
210 Numerical results are shown for flow intensity $f = 1$, which gives the best agreement with the
211 experiments, averages taken over 2×10^6 independent realizations. Error bars represent the
212 standard deviation. The insets show snapshots of colonization patterns obtained in the
213 experiments (right) and the model (left) at initial colonization densities indicated by the gray
214 pointers.

215
216 To compare our model and experiments, we used the simulation framework to study the
217 behavior of the clonal correlation length as a function of flow intensity and system size. To keep
218 our analysis as close as possible to the experiments, we initialized each simulation with a density
219 of cells ρ_0 and assigned to each cell either the blue or red color with probability 0.5. In this manner,
220 we constructed, on average, a 1:1 (blue:red) mix of cells randomly located within the lattice. Since
221 bacteria in our experiments either produce matrix constitutively or not at all, we assumed that these
222 strains correspond in our model to the $\sigma = 1$ (highly-adhesive) and $\sigma = 0$ (weakly-adhesive) cases,
223 respectively. In addition, we parametrized the spatial scale of the model to mimic the experimental
224 device. We used a square lattice of lateral length $L = 60$ sites, which represents each of the ($60 \mu\text{m}$
225 $\times 60 \mu\text{m}$) field-of-view tiles of the experimental system (i.e., corresponding to a lattice mesh size
226 $dx = dy = 1 \mu\text{m}$), and assuming an approximate cellular cross section of $1 \mu\text{m}^2$ [29], we limited the
227 maximum occupancy of each position of the lattice to only one cell. Finally, since we are interested
228 in the final occupancy patterns, regardless of the temporal scale at which colonization takes place,
229 we fixed the birth rate to minimize the computational time. This parametrization leaves flow
230 intensity, f , as the only parameter that is free in the model but fixed in the experiments. Since flow
231 intensity is defined in terms of a non-dimensional quantity in the model, we established a

232 connection between its value in the experiments and the model by finding the best quantitative
233 agreement between model-produced and experimental patterns. For a broad range of flow
234 intensities (S2 Fig), the theoretical results confirm the qualitative trend observed with our
235 experiments: clonal correlation length and total initial density are negatively correlated for highly-
236 adhesive cells, but nearly uncorrelated for weakly-adhesive cells. However, we found the best
237 quantitative agreement for the mean correlation length between experiments and simulations in
238 the strong flow limit $f = 1$, for which the simulation results are shown together with the
239 experimental data in Fig 2. The correlation length is also quantitatively, but not qualitatively,
240 affected by the simulated “field of view” (or tile size); spatial segregation increases for larger
241 systems, but the trends of the $\sigma = 0$ and $\sigma = 1$ curves are independent of system size for $f = 1$. A
242 more detailed analysis of the effect of system size in our simulations is provided in S1 Text.

243 As a last part of the model validation effort, we obtained the simulation (highly-adhesive,
244 $\sigma = 1$) and experimental (matrix-producer) distributions resulting from the correlation lengths
245 obtained with independent replicates, and compared one versus the other for different initial
246 densities (Fig 3). To compute the distributions, we divided the experimental measures in three
247 ranges of initial densities (low, intermediate and high according to the clusters of experimental
248 data observed in Fig 2) and used fast adaptive kernel density estimation in which the bandwidth
249 of the kernel varies across the dataset. These algorithms are particularly useful to estimate
250 asymmetric distributions with a fat tail in one extreme and a thinner tail on the other [47]. The
251 model and experimental distributions agree, especially in the high and low-density limits at which
252 more experimental replicates were gathered. Note that, in both extremes, the estimated
253 distributions are skewed (S3 Fig), suggesting that the median is a better measure of the central
254 tendency of the distribution than the mean. However, because both measures do not seem to differ
255 significantly (see S2 Fig and S4 Fig) whereas the mean provides less noisy results, we will focus
256 hereafter on the mean and the standard deviation as indicators of central tendency and dispersion,
257 respectively.

258



259

260 **Fig 3. Model validation: distribution of correlation lengths for matrix-producing strains.**

261 Estimated theoretical (full line) and experimental (dashed line) correlation length distributions.

262 The symbols represent the experimental distribution prior to smoothing estimations. Each color

263 represents a range of colonizing cell densities: green, 10^{-1} cells/ μm^2 for the model and high

264 density experimental data (cluster of data around 10^{-1} cells/ μm^2 in Fig 2); blue, 10^{-2} and 2.15×10^{-2}
265 cells/ μm^2 for the model and intermediate density experimental data ($7 \times 10^{-3} < \rho_0 < 3 \times 10^{-2}$
266 cells/ μm^2); and red, 10^{-3} and 2.15×10^{-3} cells/ μm^2 in the model and low density experimental data
267 ($\rho_0 < 5 \times 10^{-3}$ cells/ μm^2).

268

269 **Model predictions. Interaction between bacterial traits and flow intensity.**

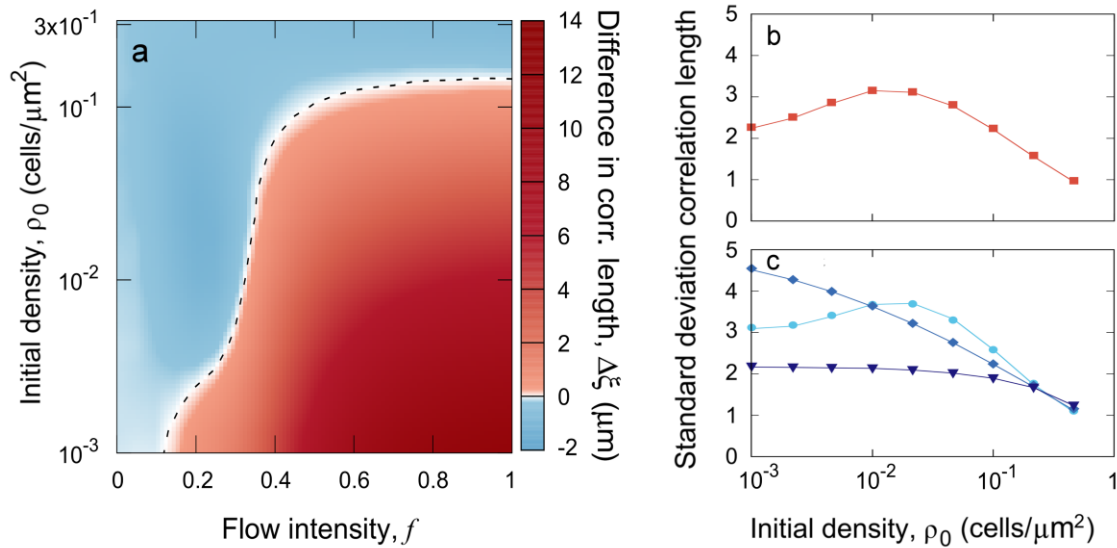
270 As discussed above, we consider founder density and adhesiveness as the traits of interest for our
271 strains in this study. Both of these traits are influenced by genetically encoded factors, such as
272 matrix secretion, as well as by environmental factors, such as habitat turnover and surface
273 chemistry [48]. To the extent that adhesion and surface colonization density are under bacterial
274 control, we consider these traits here to be part of a general strategy set for influencing surface
275 occupation [49]. We explore the effects of the flow on colonization strategies by studying how
276 diverse combinations of flow strength, adhesiveness, and initial population density influence final
277 patterns of surface occupation. As described above, numerical simulations were initiated with a
278 1:1 mixture of red and blue strains that have the same adhesiveness.

279 As shown in S2 Fig, the mean correlation length decreases as the initial density increases
280 for any flow intensity and any cell adhesiveness. This trend is applicable also for weakly adhesive
281 cells ($\sigma = 0$), although for the highest flow intensities the trend is only evident for very high initial
282 densities. The results are more convoluted when looking at a range of adhesiveness for a fixed
283 initial density (S5 Fig). Lower ρ_0 (cells/ μm^2) conditions show null or positive association between
284 adhesiveness and correlation length, whereas higher initial densities show a null or slightly
285 negative interdependence.

286 In order to assess how the different colonization strategies would be influenced by the
287 flow, we quantified the difference between the correlation length reached by highly-adhesive
288 strains ($\sigma = 1$) and weakly-adhesive strains ($\sigma = 0$) as a function of flow intensity and initial
289 population density. Intuitively, one might expect that populations of highly-adhesive cells
290 universally obtain larger clonal clusters, and indeed, this outcome does occur broadly, especially
291 with increasing flow speed. When flow is strong, less-adhesive cells are frequently removed from
292 the surface, exposing new area for attachment and growth and generally causing population
293 admixture. However, there is a considerable region of the parameter space in which populations
294 of weakly-adhesive cells show the larger clonal clusters (higher correlation length) at confluence,
295 especially when flow is weak, or when initial population density is high (Fig 4a). The difference
296 in correlation length between highly and weakly-adhesive cells becomes more pronounced in
297 larger systems. Weakly-adhesive strains form larger clusters than the highly-adhesive ones for a
298 larger set of flow intensities, and this difference in cluster size can be quantitatively of similar
299 magnitude to the one gained by highly-adhesive strains in the strong flow limit (further details on
300 the effects of the system size are provided in S1 Text).

301

302



303

304 **Fig 4. Model output: mean cluster size and variability.** a) Difference in correlation length
 305 resulting from investing in cell adhesion for different flow intensities and initial colonization
 306 densities. The dashed line indicates the values of f and ρ_0 at which this difference is equal to
 307 zero. Averages are taken over 5×10^4 independent realizations. b, c) The standard deviation of the
 308 correlation length is a proxy for lineage segregation variability in highly-adhesive strains, (b; $\sigma =$
 309 1) and weakly-adhesive cells (c; $\sigma = 0$). Averages are taken over 2×10^6 independent realizations
 310 of the model.

311

312 Finally, the correlation length of clonal clusters is highly variable in our experiments with
 313 constitutively matrix-secreting cells, especially for intermediate colonizing population density (Fig
 314 3). In light of this observation, we used the model to investigate how flow intensity influences
 315 variability in the correlation length for highly- and weakly-adhesive cells and continuously varying
 316 initial surface density. For low flows, the variability in clonal cluster size follows the same trend
 317 for highly-adhesive and weakly-adhesive strains, reaching its maximum values at intermediate
 318 initial densities (Fig 4b, 4c, S6 Fig). Differences between strains emerge as flow intensity
 319 increases. On the one hand, highly-adhesive cells cannot be detached or shoved, and thus their
 320 cluster size variability is not influenced by flow speed (S6 Fig). Such variability in the correlation
 321 length is also quantitatively influenced by system size, although the curve maintains its concavity
 322 as a function of the initial population density (S1 Text). On the other hand, as flow speed increases,
 323 the dispersion in the weakly-adhesive strain correlation length transitions from a convex form to a
 324 uniformly decreasing function of initial population density (Fig 4c). This pattern holds for strains
 325 with intermediate adhesiveness, although the influence of flow intensity on correlation length
 326 variability decreases as adhesiveness increases (S6 Fig).

327

328 Discussion

329 Combining experiments in microfluidic devices with numerical simulations of a cellular
 330 automaton, we have developed a framework for quantifying strain mixture versus segregation in
 331 the coverage patterns that emerge from bacterial expansion competition on 2-D flat surfaces. We
 332 used experimental data to validate the core assumptions of the model framework, which permitted

333 us to make predictions for a broad set of ecological scenarios defined by the intensity of
334 environmental flow, surface colonization density, cell adhesion properties, and the extension of
335 the colonized surface.

336 Microbes occupy a vast variety of surfaces, often subject to a wide range of fluid flow
337 intensities. A common example of surface attachment stressed by laminar flow-induced shear
338 forces is chitin colonization in marine environments [9], an important feature of the natural
339 ecology of many *Vibrio* species. Typical surface colonization densities are also likely to vary
340 widely depending on the species, environmental conditions, and local demographics of bacterial
341 communities. Among the mechanisms that control seeding density, some are under bacterial
342 control, and others are not. For example, chemotaxis toward surfaces and the active production of
343 adhesins/extracellular matrix can modulate cell surface occupation, but so too will ambient
344 population density conditions in the planktonic phase, local flow patterns, and the chemistry of the
345 surface bacteria attempt to colonize [25], [48], [50], [51]. Decreasing the initial colonization
346 density increases the typical distance between founder cells and thus the territory that can be
347 potentially occupied by each of them and its descendants [52]. In our experiments and simulations
348 with highly-adhesive strains subject to strong flows, this translates into a negative correlation
349 between cell lineage cluster size and initial cell density, consistent with previous reports in other
350 species [53]. In populations of weakly-adhesive cells, however, flow encourages spatial mixing of
351 genetic lineages by detaching cells and transporting them to other positions in the local
352 environment, which reduces the sensitivity of the final pattern to the initial conditions. As a result,
353 when flows are strong and colonization densities are moderate to low, investment in cell-cell and
354 cell-surface adhesion results in stronger clonal clustering of cell lineages.

355 It follows from intuition that populations of highly-adhesive cells might generate coherent
356 clonal clusters more easily than less adhesive cells. And indeed, this result was observed in our
357 experiments and for many model conditions. However, there was a broad region of the model
358 parameter space in which the opposite behavior was predicted. This exception occurred at low
359 flow strengths and, independently of flow strength, when the initial population density of
360 colonizing cells was very high. In each of these two cases, we found that a different mechanism
361 underlies such counterintuitive result. For the former case, if flows are weak cell relocations occur
362 over short distances, which alleviates local competition for space within large clusters instead of
363 mixing the population. Weakly adhesive strains thus form larger clusters than highly-adhesive
364 strains *via* limited dispersal. For the latter case, when surfaces are almost fully occupied during
365 the colonization phase, populations of highly adherent cells (which resist removal by flow) fix the
366 initial state of the system into one of randomly distributed cell lineages. In populations of weakly
367 adhesive cells, however, the vast majority of cells that detach cannot re-attach to the surface
368 elsewhere and are lost to the flow output. The positions from which detached cells were removed
369 are then occupied by descendants of neighboring cells that had managed to remain in place. If the
370 detached cell was originally surrounded by cells of its same lineage, then the empty space is filled
371 by a new cell within the same lineage and the update has no effect; in a mixed region, however,
372 the growth will tend to reduce mixing and thus to increase the clonal correlation length of the
373 system. Therefore, populations of highly-adhesive cells are not universally expected to show
374 stronger spatial genetic structure than populations of less adherent cells; the structure depends on
375 the ecological conditions and bacterial traits controlling surface colonization density, as well as
376 the environmental flow regime.

377 Complex surface attributes, such as its topology and chemical properties, are not explored
378 here but are expected to influence cluster sizes in some natural environments by increasing the

379 complexity of fluid flow patterns, inducing short-range cell relocation and modifying the long-
380 range relocation mechanism. Furthermore, in our simulations and experiments, surfaces are
381 unoccupied prior to cell inoculation. In *V. cholerae* and other biofilm-forming organisms, matrix
382 production is known to prevent planktonic cells from entering the biofilm, thus providing a
383 competitive advantage to resident cells during surface colonization processes [54]. The tendency
384 of cells to adhere to one another and form large clusters is likely to fall under selection based on
385 the size of resource patches in a given environment. Resources matching has been intensively
386 addressed in animal ecology, both from the perspective of optimizing the search process [55]–[58],
387 and including its demographic implications [59]–[61]. Given our model results, we speculate a
388 relationship between adhesion, surface attraction, and the variance of cell lineage cluster size that
389 could determine the ability of microbes to cope with variability in nutrient patchiness. Exploring
390 the role of these three parameters is a future line of research expanding upon this study.

391 The emergent spatial structure of cell lineages during biofilm growth is important to
392 numerous other facets of microbial ecology, especially for the evolutionary trajectories of social
393 phenotypes [16]. Many phenotypes associated with biofilm formation and the pathogenesis of
394 bacterial infections, for example, are secreted factors such as digestive enzymes and nutrient-
395 chelating molecules [66]. In many cases, these secreted compounds may enable a biofilm, as a
396 collective, to degrade complex polymers – including host tissues – that otherwise would be
397 inaccessible [9], [12]. Since secreted enzymes can be costly to produce and may benefit all cells
398 in the immediate surroundings, their evolutionary stability often relies on population structure,
399 which can promote preferential interaction among cells of a single strain. If cells are mostly
400 surrounded by neighbors of the same lineage, cooperative cells are more likely to interact with
401 clonemates, which are also cooperative, promoting the evolutionary stability of the cooperative
402 phenotype in question [21], [67]. Other forms of cell-cell interaction, on the other hand, are only
403 effective in mixed population structures; these include, for example, cross-feeding mutualisms in
404 which different cell types depend on close proximity to benefit each other [17]–[19]. Antagonistic
405 phenotypes, such as toxin secretion (e.g. Type VI-mediated attack), also depend on mixed
406 population structure to be optimally effective [68]–[73].

407 Given the relationship between spatial structure and the evolutionary stability of different
408 secretion phenotypes, we might expect surface colonization and adhesion strategies to coevolve
409 with the ability to produce extracellular public goods, as well as toxins. This would be consistent
410 with the coevolution of cooperation and dispersal more generally, either via movement in motile
411 organisms or passive transport in sessile species, which has been well-explored [74]–[80]. Varying
412 surface colonization and adhesion is just one of several means through which spatial structure can
413 be altered by microbes in the process of biofilm growth [16]. Previous reports have shown that
414 some organisms, such as the social amoeba *Dictyostelium discoideum*, preferentially adhere to
415 clonemates and promote aggregation of genotypes during collective movement [81]. For many
416 microbes, the expansion of growing cell groups toward a source of limiting nutrients tends to
417 promote the spontaneous segregation of different strains due to genetic drift along the advancing
418 group front [36]. After colonizing a surface, matrix-guided motility heavily influences early
419 biofilm structure in some strains of *Pseudomonas aeruginosa* [82]. During cell group growth,
420 phenotypes like toxin secretion also promote strain segregation by enforcing positive feedback on
421 the local frequency of each self-immune toxin-secreting strain of *V. cholerae* [72]. Combined with
422 constraints due to surface properties, this array of biological forces can yield complex dynamics
423 of spatial organization in microbial communities that we are just beginning to understand [22].

424 Other factors will also impact the evolution of adhesion phenotypes, including the relative
425 advantage of highly-adhesive cells against less-adhesive cells in direct competition [54], and the
426 trade-off between competitive surface adhesion and the ability to disperse to new habitats for later
427 growth [38], [40], [83]. This is only one of many dimensions of surface-associated microbial
428 behavior, which can include sophisticated mechanisms of surface departure and re-attachment, as
429 well as various forms of individual and collective surface motility [26]. Disentangling the impacts
430 of these different adhesion and detachment principles is an important area for future study.

431

432 **Materials and Methods**

433 ***V. cholerae* strain engineering.**

434 We conducted surface colonization experiments using *V. cholerae*, a model organism for biofilm
435 formation on a broad range of surfaces. In order to control the several genes that are regulated by
436 the flagellum activity and by quorum sensing, we first deleted *flaA*, which encodes the flagelling
437 core protein, and *hapR*, which encodes the quorum sensing master regulator. This results in a
438 double mutant $\Delta flaA \Delta hapR$ that produces EPS and therefore termed EPS⁺. Second, we produced a
439 triple mutant strain by deleting *vpsL*, a gene required for EPS production. The resulting
440 $\Delta flaA \Delta hapR \Delta vpsL$ strain never produces EPS and we thus call it EPS⁻. Finally, we derived two
441 versions of the EPS⁺ and the EPS⁻ strain: one that expresses the teal fluorescent protein *mTFPI*
442 and one that expresses the ref fluorescent protein *mKate*. This difference in the fluorescence
443 protein is the only difference between otherwise genetically identical strains in our mixes, and it
444 will allow us to distinguish different lineages in the surface colonization pattern.

445

446 **Experimental protocol.**

447 We performed bacterial surface occupation experiments using microfluidic culture methods.
448 Chambers were 500 μm wide, 100 μm high and 7 mm long, and were constructed from
449 poly(dimethylsiloxane) bonded to glass coverslips. Overnight cultures of the EPS⁺ and EPS⁻
450 strains were normalized to an optical density at 600 nm of 1.0, mixed to create a 1:1 culture of red
451 and blue cells, and back-diluted either 1:100, 1:1000, or 1:10000 prior to being introduced into the
452 chambers. The cultures were then incubated at room temperature for one hour to allow cells to
453 attach to the glass coverslip. Varying the planktonic culture density in this manner allowed us to
454 vary the initial population density on the glass surface. Following this attachment period, sterile
455 M9 medium with 0.5% glucose was introduced to the chamber at 0.1 $\mu\text{L}/\text{min}$, using a high
456 precision syringe pump (Harvard Apparatus). The chambers were fixed to the stage of an inverted
457 spinning disk confocal microscope (Nikon, Andor), which was used to capture images of the cell
458 populations residing on the coverslip glass. The entire surface of each chamber was imaged once
459 per hour until surface coverage was complete as judged by eye.

460

461 **Model details.**

462 The two main ingredients of our model are:

463 (i) *Reproduction*. Bacteria reproduce at a given rate μ : every time step a cell division takes
464 place with probability $p_b = \mu dt$, where the length of the time step dt (i.e. temporal resolution of
465 the simulations) is fixed such that $p_b < 1$. Since fluorescent protein constructs have no fitness effect
466 [38], we set the same reproduction rate for both strains. In addition, since in each experiment both
467 strains equally invest in adhesion, we ignore the potential cost of adhesiveness here. Finally, since
468 we are interested in the final spatial cell distributions, our results are independent of the specific
469 value used for μ , which is fixed to minimize computational time. Newborn cells occupy a randomly

470 chosen site among the available places within the Moore neighborhood of the parental cell (eight
471 lattice positions surrounding the parental cell). If there is no empty position, the new cell will try
472 to shove one of the resident (i.e. existing) cells in the neighborhood and occupy its position. The
473 outcome of a shoving attempt is determined by a displacement probability, p_s , defined in terms of
474 non-dimensional adhesiveness as:

$$p_s = \frac{1 - \sigma}{2}. \quad (1)$$

475
476 With this definition, highly-adhesive cells ($\sigma = 1$) are never displaced by newborns,
477 whereas weakly-adhesive residents and newborns will have the same probability to be shoved due
478 to low cell-surface adhesion (i.e. $p_s=0.5$). From each shoving event, two possibilities ensue: (i) the
479 resident cell remains in its position, and the newborn is displaced to one of the empty neighboring
480 sites of the resident, or (ii) the newborn cell takes the position of the resident, which is displaced
481 to one of its empty neighboring sites. In both scenarios, if the complete neighborhood of the
482 resident cell is occupied, the losing cell is removed from the system with the outflow. Note that
483 this formulation truncates a cascade of shoving events that might take place for weakly adhesive
484 cells as a cluster of bacteria expands from its center. In this manner, we are assuming that shoving
485 events can only occur on short spatial scales before one cell must be released into the passing flow
486 to relieve the pressure of increasing local population density.

487 (ii) *Cell detachment and relocation.* At every time step, we also check for potential cell
488 relocations that occur due to fluid flow passing above the cell monolayer. Since flow enters the
489 experimental chambers from one direction only (left to right), we assume that cells can only be
490 removed by flow if their neighboring position on the left is empty. This simplification implements
491 a drafting effect that is supported by basic fluid mechanics calculations reported by [84]: cells are
492 protected from drag by neighbors that sit on the surface immediately upstream. Therefore, the
493 detachment probability is zero when the directly adjacent up-stream site is occupied, and p_d
494 otherwise. We define p_d using a combination of the non-dimensional flow strength, f , and cell
495 adhesiveness:

$$p_d = f(1 - \sigma), \quad (2)$$

496 where, for simplicity, we assume that f is normalized and therefore can take any value between 0
497 and 1. According to Eq. (2), highly-adhesive cells cannot be detached, whereas weakly-adhesive
498 cells will be dislocated with a probability given only by the strength of the flow. Because it is not
499 possible experimentally to track detached and re-attached individual cells over the full length of
500 the microfluidic growth chambers to inform our model, we hypothesized a mechanism for long-
501 range surface re-attachment. We could thus make predictions of the spatial structure of the
502 population at confluence and directly check them against experimental results. In our simulations,
503 once a cell has been detached, a landing position is calculated using the following rules that
504 account for flow directionality. The distance traveled in the direction of the flow, Δx , is determined
505 by a random integer uniformly distributed between 0 and fL , whereas the distance traveled in the
506 transversal direction, Δy , is obtained as a random integer uniformly distributed between $-\Delta x$ and
507 Δx . If the sorted position was already occupied, then the detached cell is removed from the system,
508 which accounts for bacterial loss with the outflow. With these rules, cells can only relocate to
509 positions downstream of the flow orientation, unless they pass through the system boundaries due
510 to periodic boundary conditions, which recovers the isotropy in the surface-occupation patterns.
511 On the other hand, detached cells can freely drift perpendicular to the flow. A summary of the
512 model parameters and their numerical values is provided in S1 Table.

513
514
515
516
517

Characterization of surface occupation patterns: the correlation length.

We characterize bacterial surface occupation patterns using the spatial autocorrelation function, $C(r)$, which can be mathematically defined as,

$$C(r) = \frac{\langle c(R)c(R+r) \rangle - \langle c(R) \rangle \langle c(R+r) \rangle}{\langle c^2(R) \rangle - \langle c(R) \rangle^2} \quad (3)$$

518
519
520
521
522
523
524

where c is the binary variable that represents the lineage color (and thus takes value 1 or 2 depending on whether the lattice cell is occupied by a blue or red cell), and $\langle . \rangle$ represents an average over all the elementary spatial units of the system, which are labeled by the index R . Given the use of periodic boundary conditions in our cellular automaton and cell mixing across adjacent tiles in the experimental device, surface occupation patterns are isotropic and the average over the angular variable can be done.

525 In microscopy images, the elementary unit is the pixel (0.065 μm), whereas in the
526 simulations, it is the lattice position (1 μm). Note that the normalization factor ensures that the
527 correlation function reaches 1 when two positions have a perfect correlation. In addition, the
528 uncorrelated average product, $\langle c(R) \rangle \langle c(R+r) \rangle$, force the correlation function to be zero when
529 two locations are completely independent from each other. The correlation length is thus given by
530 the first zero of the correlation function (S7 Fig). The spatial autocorrelation function given in Eq.
531 (3) is related to the radial distribution function, often used to describe how density varies as a
532 function of distance from a reference particle in a system of multiple particles.

533
534

Acknowledgments

536 The authors would like to thank J. Dushoff for early discussions related to the project, to E. Bendo
537 for assistance in data collection during early stages of the experiments, and to S.A. Levin for early
538 discussions and a critical reading of the manuscript. They are also grateful to IFISC (CSIC-UIB)
539 for the use of its computational infrastructure. This work is funded by the Gordon and Betty Moore
540 Foundation through Grant GBMF2550.06 to R.M-G; J.A.B. is supported by the Marine Alliance
541 for Science and Technology for Scotland (MASTS) pooling initiative, funded by the Scottish
542 Funding Council (HR09011) and contributing institutions. C.D.N. received support from the
543 Alexander von Humboldt Foundation and the Cystic Fibrosis Foundation (STANTO15RO). K.D.
544 received support for this work from the Max Planck Society, Behrens Weise Foundation, European
545 Research Council (StG-716734), Human Frontier Science Program (CDA00084/2015-C), and
546 Deutsche Forschungsgemeinschaft (SFB987).

547

References

- 549 [1] H. W. Paerl and J. L. Pinckney, "A mini-review of microbial consortia: Their roles in
550 aquatic production and biogeochemical cycling," *Microb. Ecol.*, vol. 31, no. 3, pp. 225–
551 247, 1996.
- 552 [2] C. A. Fux, J. W. Costerton, P. S. Stewart, and P. Stoodley, "Survival strategies of infectious
553 biofilms," *Trends Microbiol.*, vol. 13, no. 1, pp. 34–40, 2005.
- 554 [3] S. Macfarlane, B. Bahrami, and G. T. Macfarlane, "Mucosal biofilm communities in the
555 human intestinal tract," *Adv. Appl. Microbiol.*, vol. 75, pp. 111–143, 2011.
- 556 [4] B. Halan, K. Buehler, and A. Schmid, "Biofilms as living catalysts in continuous chemical

- 557 syntheses,” *Trends Biotechnol.*, vol. 30, no. 9, pp. 453–465, 2012.
- 558 [5] J. Jass and J. Walker, “Biofilms and biofouling,” in *Industrial Biofouling: detection,*
559 *prevention and control*, J. Walker, S. Surman, and J. Jass, Eds. New York: John Wiley &
560 Sons, 2000, pp. 1–12.
- 561 [6] H. Flemming and J. Wingender, “The biofilm matrix.,” *Nat. Rev. Microbiol.*, vol. 8, no. 9,
562 pp. 623–33, 2010.
- 563 [7] H.-C. Flemming, J. Wingender, U. Szewzyk, P. Steinberg, S. A. Rice, and S. Kjelleberg,
564 “Biofilms: an emergent form of bacterial life,” *Nat. Rev. Microbiol.*, vol. 14, no. 9, pp. 563–
565 575, 2016.
- 566 [8] I. W. Sutherland, “The biofilm matrix - An immobilized but dynamic microbial
567 environment,” *Trends Microbiol.*, vol. 9, no. 5, pp. 222–227, 2001.
- 568 [9] K. Drescher, C. D. Nadell, H. A. Stone, N. S. Wingreen, and B. L. Bassler, “Solutions to
569 the public goods dilemma in bacterial biofilms,” *Curr. Biol.*, vol. 24, no. 1, pp. 50–55, 2014.
- 570 [10] C. Absalon, K. van Dellen, and P. I. Watnick, “A communal bacterial adhesin anchors
571 biofilm and bystander cells to surfaces,” *PLoS Pathog.*, vol. 7, no. 8, 2011.
- 572 [11] K. P. Rumbaugh, S. P. Diggle, C. M. Watters, A. Ross-Gillespie, A. S. Griffin, S. A. West,
573 M. Auer, A. N. Hamood, K. P. Rumbaugh, A. F. Read, and E. Al., “Quorum sensing and
574 the social evolution of bacterial virulence.,” *Curr. Biol.*, vol. 19, no. 4, pp. 341–5, Feb. 2009.
- 575 [12] H. J. Gilbert and G. P. Hazlewood, “Bacterial cellulases and xylanases,” *J. Gen. Microbiol.*,
576 vol. 139, pp. 187–194, 1993.
- 577 [13] A. Monmeyran, A.-F. Bitbol, P. Thomen, C. Douarche, and N. Henry, “Bacterial biofilm
578 under flow : First a physical struggle to stay , then a matter of breathing,” *PLoS One*, vol.
579 12, no. 4, p. e0175197, 2017.
- 580 [14] P. S. Stewart, “Mini-review: Convection around biofilms,” *Biofouling*, vol. 28, no. 2, pp.
581 187–198, 2012.
- 582 [15] C. J. Rupp, C. A. Fux, and P. Stoodley, “Viscoelasticity of *Staphylococcus aureus* Biofilms
583 in Response to Fluid Shear Allows Resistance to Detachment and Facilitates Rolling
584 Migration Cory,” *Appl. Environ. Microbiol.*, vol. 71, no. 4, pp. 2175–2178, 2005.
- 585 [16] C. D. Nadell, K. Drescher, and K. R. Foster, “Spatial structure, cooperation, and competition
586 in biofilms,” *Nat. Rev. Microbiol.*, vol. 14, no. 9, pp. 589–600, 2016.
- 587 [17] S. Estrela and S. P. Brown, “Metabolic and Demographic Feedbacks Shape the Emergent
588 Spatial Structure and Function of Microbial Communities,” *PLoS Comput. Biol.*, vol. 9, no.
589 12, 2013.
- 590 [18] S. Mitri, J. B. Xavier, and K. R. Foster, “Social evolution in multispecies biofilms.,” *Proc.*
591 *Natl. Acad. Sci. U. S. A.*, vol. 108 Suppl, no. Supplement_2, pp. 10839–46, 2011.
- 592 [19] B. Momeni, A. J. Waite, and W. Shou, “Spatial self-organization favors heterotypic
593 cooperation over cheating,” *Elife*, vol. 2013, no. 2, pp. 1–18, 2013.
- 594 [20] J. B. Xavier and K. R. Foster, “Cooperation and conflict in microbial biofilms,” *Proc. Natl.*
595 *Acad. Sci. USA*, vol. 104, no. 3, pp. 876–881, 2007.
- 596 [21] W. W. Driscoll and J. W. Pepper, “Theory for the evolution of diffusible external goods,”
597 *Evolution (N. Y.)*, vol. 64, no. 9, pp. 2682–2687, 2010.
- 598 [22] A. Stacy, L. McNally, S. E. Darch, S. P. Brown, and M. Whiteley, “The biogeography of
599 polymicrobial infection,” *Nat. Rev. Microbiol.*, vol. 14, no. 2, pp. 93–105, 2016.
- 600 [23] G. A. O’Toole and R. Kolter, “Initiation of biofilm formation in *Pseudomonas fluorescens*
601 WCS365 proceeds via multiple, convergent signalling pathways: A genetic analysis,” *Mol.*
602 *Microbiol.*, vol. 28, no. 3, pp. 449–461, 1998.

- 603 [24] A. P. Tomaras, C. W. Dorsey, R. E. Edelman, and L. A. Actis, “Attachment to and biofilm
604 formation on abiotic surfaces by *Acinetobacter baumannii*: Involvement of a novel
605 chaperone-usher pili assembly system,” *Microbiology*, vol. 149, no. 12, pp. 3473–3484,
606 2003.
- 607 [25] P. M. Merritt, T. Danhorn, and C. Fuqua, “Motility and chemotaxis in *Agrobacterium*
608 *tumefaciens* surface attachment and biofilm formation,” *J. Bacteriol.*, vol. 189, no. 22, pp.
609 8005–8014, 2007.
- 610 [26] G. A. O’Toole and G. C. L. Wong, “Sensational biofilms: Surface sensing in bacteria,”
611 *Curr. Opin. Microbiol.*, vol. 30, no. 1999, pp. 139–146, 2016.
- 612 [27] W. P. J. Smith, Y. Davit, J. M. Osborne, W. Kim, K. R. Foster, and J. M. Pitt-Francis, “Cell
613 morphology drives spatial patterning in microbial communities,” *Proc. Natl. Acad. Sci.*, p.
614 201613007, 2016.
- 615 [28] Y. A. Millet, D. Alvarez, S. Ringgaard, U. H. von Andrian, B. M. Davis, and M. K. Waldor,
616 “Insights into *Vibrio cholerae* Intestinal Colonization from Monitoring Fluorescently
617 Labeled Bacteria,” *PLoS Pathog.*, vol. 10, no. 10, 2014.
- 618 [29] K. Drescher, J. Dunkel, C. D. Nadell, S. van Teeffelen, I. Grnja, N. S. Wingreen, H. A.
619 Stone, and B. L. Bassler, “Architectural transitions in *Vibrio cholerae* biofilms at single-
620 cell resolution,” *Proc. Natl. Acad. Sci.*, vol. 113, no. 14, pp. E2066–E2072, 2016.
- 621 [30] J. Yan, A. G. Sharo, H. A. Stone, N. S. Wingreen, and B. L. Bassler, “*Vibrio cholerae*
622 biofilm growth program and architecture revealed by single-cell live imaging,” *Proc. Natl.*
623 *Acad. Sci.*, pp. 1–7, 2016.
- 624 [31] A. Siryaporn, M. K. Kim, Y. Shen, H. A. Stone, and Z. Gitai, “Colonization, competition,
625 and dispersal of pathogens in fluid flow networks,” *Curr. Biol.*, vol. 25, no. 9, pp. 1201–
626 1207, 2015.
- 627 [32] A. Persat, H. a Stone, and Z. Gitai, “The curved shape of *Caulobacter crescentus* enhances
628 surface colonization in flow,” *Nat. Commun.*, vol. 5, no. May, p. 3824, 2014.
- 629 [33] Y. Shen, A. Siryaporn, S. Lecuyer, Z. Gitai, and H. A. Stone, “Flow directs surface-attached
630 bacteria to twitch upstream,” *Biophys. J.*, vol. 103, no. 1, pp. 146–151, 2012.
- 631 [34] R. Rusconi, J. S. Guasto, and R. Stocker, “Bacterial transport suppressed by fluid shear,”
632 *Nat. Phys.*, vol. 10, no. 3, pp. 212–217, 2014.
- 633 [35] A. Persat, C. D. Nadell, M. K. Kim, F. Ingremeau, A. Siryaporn, K. Drescher, N. S.
634 Wingreen, B. L. Bassler, Z. Gitai, and H. A. Stone, “The Mechanical World of Bacteria,”
635 *Cell*, vol. 161, no. 5, pp. 988–997, 2015.
- 636 [36] O. Hallatschek, P. Hersen, S. Ramanathan, and D. R. Nelson, “Genetic drift at expanding
637 frontiers promotes gene segregation.,” *Proc. Natl. Acad. Sci. U. S. A.*, vol. 104, no. 50, pp.
638 19926–30, 2007.
- 639 [37] A. J. Dumbrell, M. Nelson, T. Helgason, C. Dytham, and A. H. Fitter, “Relative roles of
640 niche and neutral processes in structuring a soil microbial community,” *Int. Soc. Microb.*
641 *Ecol. Journa*, vol. 4, no. 3, pp. 337–345, 2010.
- 642 [38] C. D. Nadell and B. L. Bassler, “A fitness trade-off between local competition and dispersal
643 in *Vibrio cholerae* biofilms.,” *Proc. Natl. Acad. Sci. U. S. A.*, vol. 108, no. 34, pp. 14181–
644 14185, 2011.
- 645 [39] A. S. Utada, R. R. Bennett, J. C. N. Fong, M. L. Gibiansky, F. H. Yildiz, R. Golestanian,
646 and G. C. L. Wong, “*Vibrio cholerae* use pili and flagella synergistically to effect motility
647 switching and conditional surface attachment,” *Nat. Commun.*, vol. 5, pp. 1–8, 2014.
- 648 [40] J. Yan, C. D. Nadell, and B. L. Bassler, “Environmental fluctuation governs selection for

- 649 plasticity in biofilm production,” *ISME J.*, Mar. 2017.
- 650 [41] S. Kojima, K. Yamamoto, I. Kawagishi, and M. Homma, “The Polar Flagellar Motor of
651 *Vibrio cholerae* Is Driven by an Na⁺ Motive Force,” *J. Bacteriol.*, vol. 181, no. 6, pp. 1927–
652 1930, 1999.
- 653 [42] P. I. Watnick and R. Kolter, “Steps in the development of a *Vibrio cholerae* El Tor biofilm,”
654 *Mol. Microbiol.*, vol. 34, no. 3, pp. 586–595, 1999.
- 655 [43] E. Lauga, W. R. DiLuzio, G. M. Whitesides, and H. A. Stone, “Swimming in circles: motion
656 of bacteria near solid boundaries.,” *Biophys. J.*, vol. 90, no. 2, pp. 400–12, 2006.
- 657 [44] A. P. Berke, L. Turner, H. C. Berg, and E. Lauga, “Hydrodynamic attraction of swimming
658 microorganisms by surfaces,” *Phys. Rev. Lett.*, vol. 101, no. 3, pp. 1–4, 2008.
- 659 [45] M. L. Gibiansky, J. C. Conrad, F. Jin, V. D. Gordon, D. A. Motto, M. A. Mathewson, W.
660 G. Stopka, D. C. Zelasko, J. D. Shrout, and G. C. L. Wong, “Bacteria use type IV pili to
661 walk upright and detach from surfaces.,” *Science (80-.)*, vol. 330, no. 6001, p. 197, 2010.
- 662 [46] V. Berk, J. C. N. Fong, G. T. Dempsey, O. N. Develioglu, X. Zhuang, J. Liphardt, F. H.
663 Yildiz, and S. Chu, “Molecular architecture and assembly principles of *Vibrio cholerae*
664 biofilms.,” *Science*, vol. 337, no. 6091, pp. 236–9, 2012.
- 665 [47] Z. I. Botev, J. F. Grotowski, and D. P. Kroese, “Kernel density estimation via diffusion,”
666 *Ann. Stat.*, vol. 38, no. 5, pp. 2916–2957, 2010.
- 667 [48] L. K. . Ista, J. A. . Callow, Maureen E.; Finlay, S. E. . Coleman, A. C. . Nolasco, R. H. .
668 Simons, J. A. . Callow, and G. P. Lopez, “Effect of Substratum Surface Chemistry and
669 Surface Energy on Attachment of Marine Bacteria and Algal Spores,” *Appl. Environ.*
670 *Microbiol.*, vol. 70, no. 7, pp. 4151–4157, 2004.
- 671 [49] J. van Gestel and M. A. Nowak, “Phenotypic Heterogeneity and the Evolution of Bacterial
672 Life Cycles,” *PLoS Comput. Biol.*, vol. 12, no. 2, pp. 1–23, 2016.
- 673 [50] R. Dillon and L. Fauci, “A microscale model of bacterial and biofilm dynamics in porous
674 media,” *Biotechnol. Bioeng.*, vol. 68, no. 5, pp. 536–547, Jun. 2000.
- 675 [51] P. D. Steinberg, R. D. E. Nys, and S. Kjelleberg, “Chemical Cues for Surface Colonization,”
676 *J. Chem. Ecol.*, vol. 28, no. 10, pp. 1935–1951, 2002.
- 677 [52] D. P. Lloyd and R. Allen, “Competition for space during bacterial colonization of a surface,”
678 *J. R. Soc. Interface*, vol. 12, no. 110, p. 20150608, 2015.
- 679 [53] J. van Gestel, F. J. Weissing, O. P. Kuipers, and A. T. Kovács, “Density of founder cells
680 affects spatial pattern formation and cooperation in *Bacillus subtilis* biofilms.,” *ISME J.*,
681 vol. 8, no. 10, pp. 2069–79, 2014.
- 682 [54] J. Schluter, C. D. Nadell, B. L. Bassler, and K. R. Foster, “Adhesion as a weapon in
683 microbial competition.,” *ISME J.*, vol. 9, no. 1, pp. 139–49, 2015.
- 684 [55] E. L. Charnov, “Optimal foraging theory: the marginal value theorem,” *Theor. Popul. Biol.*,
685 vol. 9, pp. 129–136, 1976.
- 686 [56] S. Matsumura, R. Arlinghaus, and U. Dieckmann, “Foraging on spatially distributed
687 resources with sub-optimal movement, imperfect information, and travelling costs:
688 Departures from the ideal free distribution,” *Oikos*, vol. 119, no. 9, pp. 1469–1483, 2010.
- 689 [57] E. Ranta, P. Lundberg, and V. Kaitala, “Resource Matching with Limited Knowledge,”
690 *Oikos*, vol. 86, no. 2, pp. 383–385, 1999.
- 691 [58] W. F. Fagan, E. Gurarie, S. Bewick, A. Howard, R. S. Cantrell, and C. Cosner, “Perceptual
692 Ranges, Information Gathering, and Foraging Success in Dynamic Landscapes,” *Am. Nat.*,
693 vol. 189, no. 5, pp. 000–000, 2017.
- 694 [59] A. Hastings, “Can spatial variation alone lead to selection for dispersal?,” *Theor. Popul.*

- 695 *Biol.*, vol. 24, no. 3, pp. 244–251, 1983.
- 696 [60] S. A. Levin, D. Cohen, and A. Hastings, “Dispersal strategies in patchy environments,”
- 697 *Theor. Popul. Biol.*, vol. 26, no. 2, pp. 165–191, 1984.
- 698 [61] R. S. Cantrell, C. Cosner, and Y. Lou, “Evolution of dispersal and the ideal free
- 699 distribution,” *Math. Biosci. Eng.*, vol. 7, no. 1, pp. 17–36, Jan. 2010.
- 700 [62] A. J. Grimbergen, J. Siebring, A. Solopova, and O. P. Kuipers, “Microbial bet-hedging: the
- 701 power of being different,” *Curr. Opin. Microbiol.*, vol. 25, pp. 67–72, 2015.
- 702 [63] A. Solopova, J. van Gestel, F. J. Weissing, H. Bachmann, B. Teusink, J. Kok, and O. P.
- 703 Kuipers, “Bet-hedging during bacterial diauxic shift,” *Proc. Natl. Acad. Sci.*, vol. 111, no.
- 704 20, pp. 7427–7432, 2014.
- 705 [64] H. J. E. Beaumont, J. Gallie, C. Kost, G. C. Ferguson, and P. B. Rainey, “Experimental
- 706 evolution of bet hedging,” *Nature*, vol. 462, no. 7269, pp. 90–93, 2009.
- 707 [65] J.-W. Veening, W. K. Smits, and O. P. Kuipers, “Bistability, Epigenetics, and Bet-Hedging
- 708 in Bacteria,” *Annu. Rev. Microbiol.*, vol. 62, no. 1, pp. 193–210, 2008.
- 709 [66] F. Harrison, “Bacterial cooperation in the wild and in the clinic: Are pathogen social
- 710 behaviours relevant outside the laboratory?,” *BioEssays*, vol. 35, no. 2, pp. 108–112, 2013.
- 711 [67] M. A. Nowak, “Five rules for the evolution of cooperation,” *Science*, vol. 314, no. 5805,
- 712 pp. 1560–1563, 2006.
- 713 [68] S. A. West, “Cooperation and Competition Between Relatives,” *Science*, vol. 296, no. 5565,
- 714 pp. 72–75, 2002.
- 715 [69] R. D. Alexander, “The evolution of social behavior,” *Annu. Rev. Ecol. Syst.*, vol. 5, pp. 325–
- 716 383, 1974.
- 717 [70] A. Gardner, S. A. West, and A. Buckling, “Bacteriocins, spite and virulence,” *Proc. Biol.*
- 718 *Sci.*, vol. 271, no. 1547, pp. 1529–35, 2004.
- 719 [71] A. Gardner and S. A. West, “Spite and the scale of competition,” *J. Evol. Biol.*, vol. 17, no.
- 720 6, pp. 1195–1203, 2004.
- 721 [72] L. McNally, E. Bernardy, J. Thomas, A. Kalzigi, J. T. Pentz, S. Brown, B. Hammer, P.
- 722 Yunker, and W. Ratcliff, “Killing by Type VI secretion drives clonal phase separation and
- 723 the evolution of cooperation,” *Nat. Commun.*, vol. 8, p. 14371, 2017.
- 724 [73] V. Bucci, C. D. Nadell, and J. B. Xavier, “The Evolution of Bacteriocin Production in
- 725 Bacterial Biofilms,” *Am. Nat.*, vol. 178, no. 6, pp. E162–E173, 2011.
- 726 [74] A. Gardner and S. A. West, “Demography, altruism, and the benefits of budding,” *J. Evol.*
- 727 *Biol.*, vol. 19, no. 5, pp. 1707–1716, 2006.
- 728 [75] R. A. Johnstone and M. A. Cant, “Sex differences in dispersal and the evolution of helping
- 729 and harming,” *Am. Nat.*, vol. 172, no. 3, pp. 318–330, 2008.
- 730 [76] R. Kümmerli, A. Gardner, S. A. West, and A. S. Griffin, “Limited dispersal, budding
- 731 dispersal, and cooperation: An experimental study,” *Evolution*, vol. 63, no. 4, pp. 939–949,
- 732 2009.
- 733 [77] S. T. Powers, A. S. Penn, and R. A. Watson, “The concurrent evolution of cooperation and
- 734 the population structures that support it,” *Evolution (N. Y.)*, vol. 65, no. 6, pp. 1527–1543,
- 735 2011.
- 736 [78] T. B. Taylor, A. M. M. Rodrigues, A. Gardner, and A. Buckling, “The social evolution of
- 737 dispersal with public goods cooperation,” *J. Evol. Biol.*, vol. 26, no. 12, pp. 2644–2653,
- 738 2013.
- 739 [79] C. A. Aktipis, “Know when to walk away: contingent movement and the evolution of
- 740 cooperation,” *J. Theor. Biol.*, vol. 231, no. 2, pp. 249–260, 2004.

- 741 [80] M. E. Hochberg, D. J. Rankin, and M. Taborsky, “The coevolution of cooperation and
742 dispersal in social groups and its implications for the emergence of multicellularity,” *BMC*
743 *Evol. Biol.*, vol. 8, p. 238, 2008.
- 744 [81] E. A. Ostrowski, M. Katoh, G. Shaulsky, D. C. Queller, and J. E. Strassmann, “Kin
745 discrimination increases with genetic distance in a social amoeba,” *PLoS Biol.*, vol. 6, no.
746 11, pp. 2376–2382, 2008.
- 747 [82] K. Zhao, B. S. Tseng, B. Beckerman, F. Jin, M. L. Gibiansky, J. J. Harrison, E. Luijten, M.
748 R. Parsek, and G. C. L. Wong, “Psl trails guide exploration and microcolony formation in
749 *Pseudomonas aeruginosa* biofilms,” *Nature*, vol. 497, no. 7449, pp. 388–91, 2013.
- 750 [83] Y. Yawata, O. X. Cordero, F. Menolascina, J.-H. Hehemann, M. F. Polz, and R. Stocker,
751 “Competition-dispersal tradeoff ecologically differentiates recently speciated marine
752 bacterioplankton populations,” *Proc. Natl. Acad. Sci. U. S. A.*, vol. 111, no. 15, pp. 5622–
753 7, 2014.
- 754 [84] G. B. Chapman and G. R. Cokelet, “Flow resistance and drag forces due to multiple adherent
755 leukocytes in postcapillary vessels,” *Biophys. J.*, vol. 74, no. 6, pp. 3292–3301, 1998.
756
757

758

759 **Supporting information**

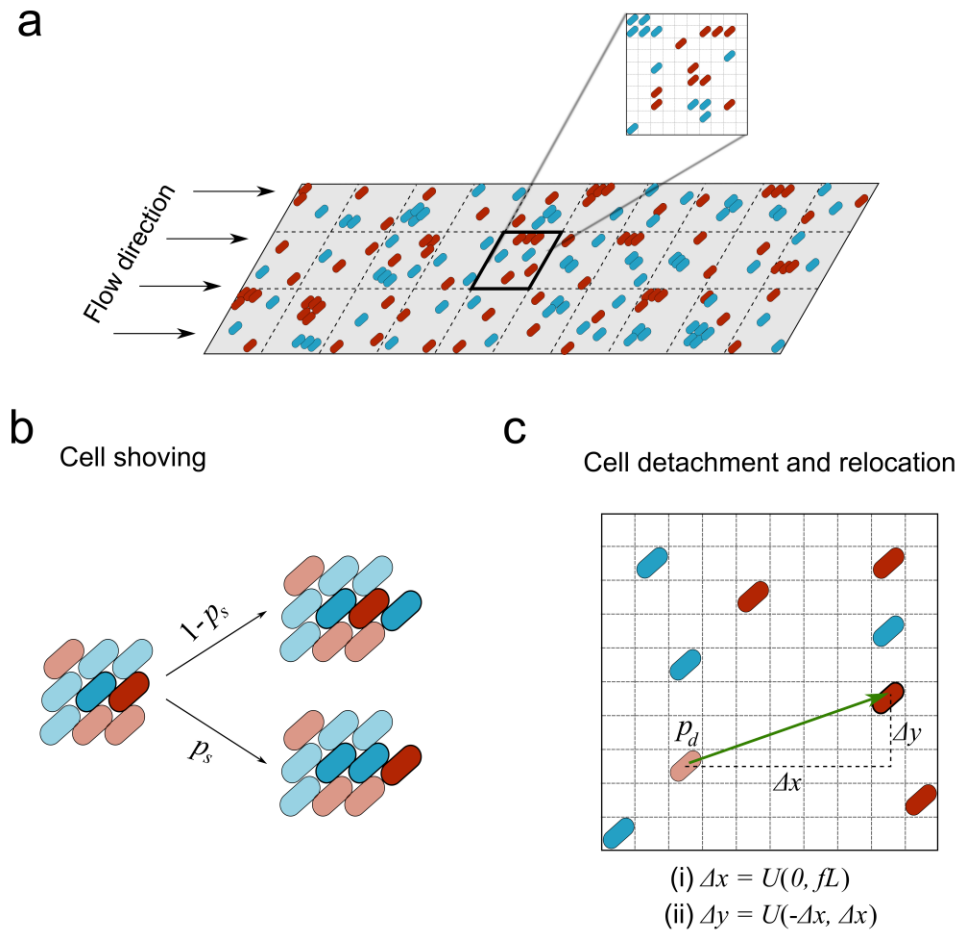
Symbol	Name	Cause	Value
σ	Cell adhesiveness	BT	Free parameter in [0,1]
f	Flow intensity	EF	Free parameter in [0,1]
ρ_0	Founder cell density	BT by EF interaction	Free parameter in $[10^{-3}, 0.5]$ <i>cells/μm^2</i>
μ	Reproduction rate	BT	Fixed parameter, 0.57 a.u.
L	Lateral lattice size	EF	Fixed parameter, 60 μm
dt	Time step	--	Fixed parameter, $1/L^2$ a.u.
dx	Lattice mesh	--	Fixed parameter, 1 μm
p_d	Detachment probability	BT by EF interaction	$f(1 - \sigma)$
p_s	Shoving probability	BT by BT interaction	$\frac{1 - \sigma}{2}$
Δx	x -distance traveled	EF	Random, in $[0, fL]$
Δy	y -distance traveled	EF	Random, in $[-\Delta x, \Delta x]$

760

761

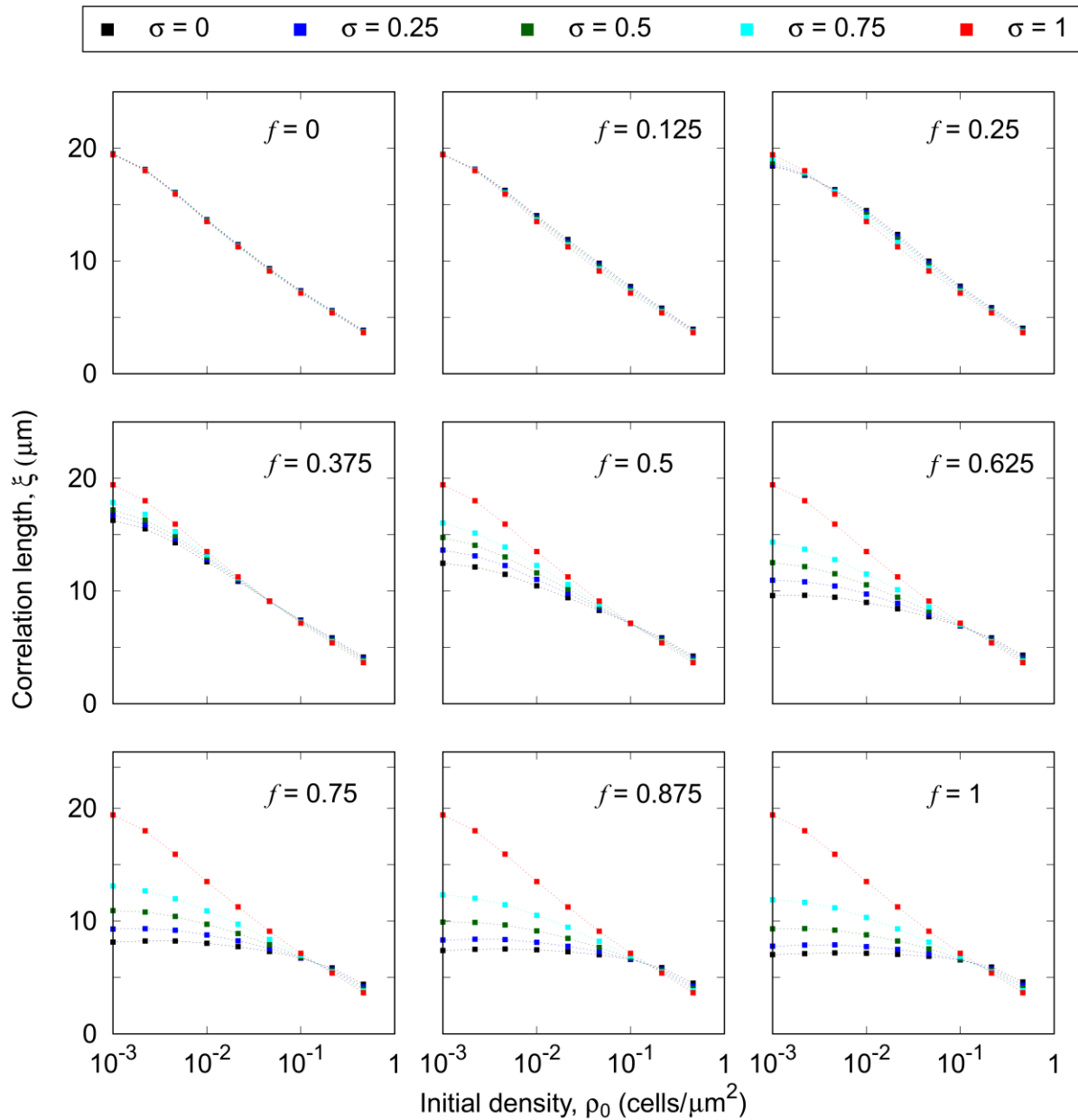
762

S1 Table. List of parameters used in the model, including whether it represents an environmental factor (EF), a bacterial trait (BT) or an interaction between them.



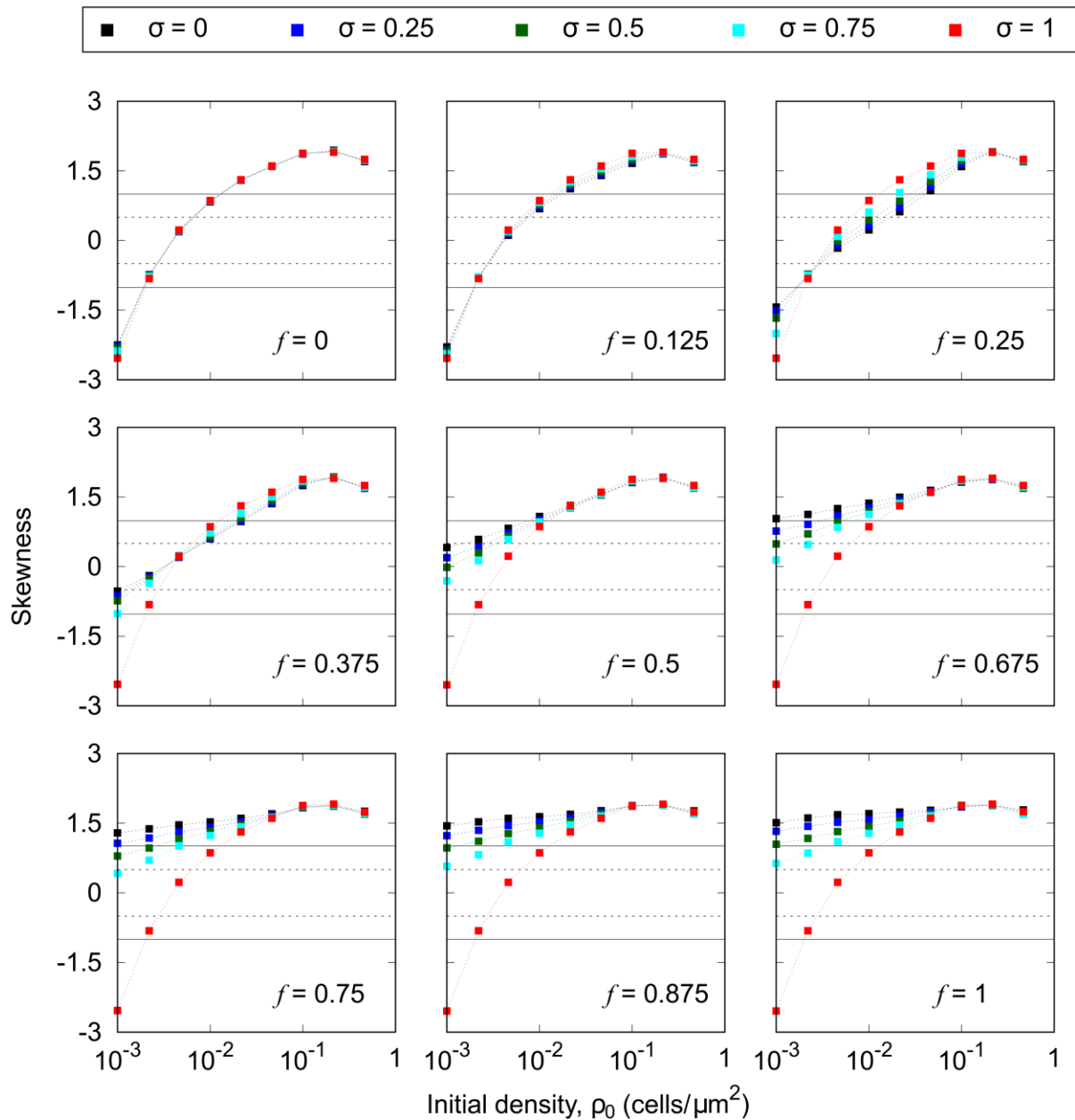
763
764
765
766
767
768
769
770
771
772
773

S1 Fig. Schematic of the experimental setup and the model updating rules. a) Schematic of the division of the experimental chamber in tiles and model representation of one of the tiles, as a 2D lattice with one cell at each lattice box. b) Cell displacement due to shoving following cell division occurs with probability p_s . With complementary probability $1-p_s$, the resident cell keeps its position and the newborn jumps to one of the adjacent empty position. c) Cells may be detached from the surface of the chamber with probability p_d and transported to a new emplacement following the relocation rules explained in the text with periodic boundary conditions (Materials and Methods). $U(a,b)$ indicates a uniformly distributed random variable between a and b . f is the flow intensity and L the lattice lateral length.



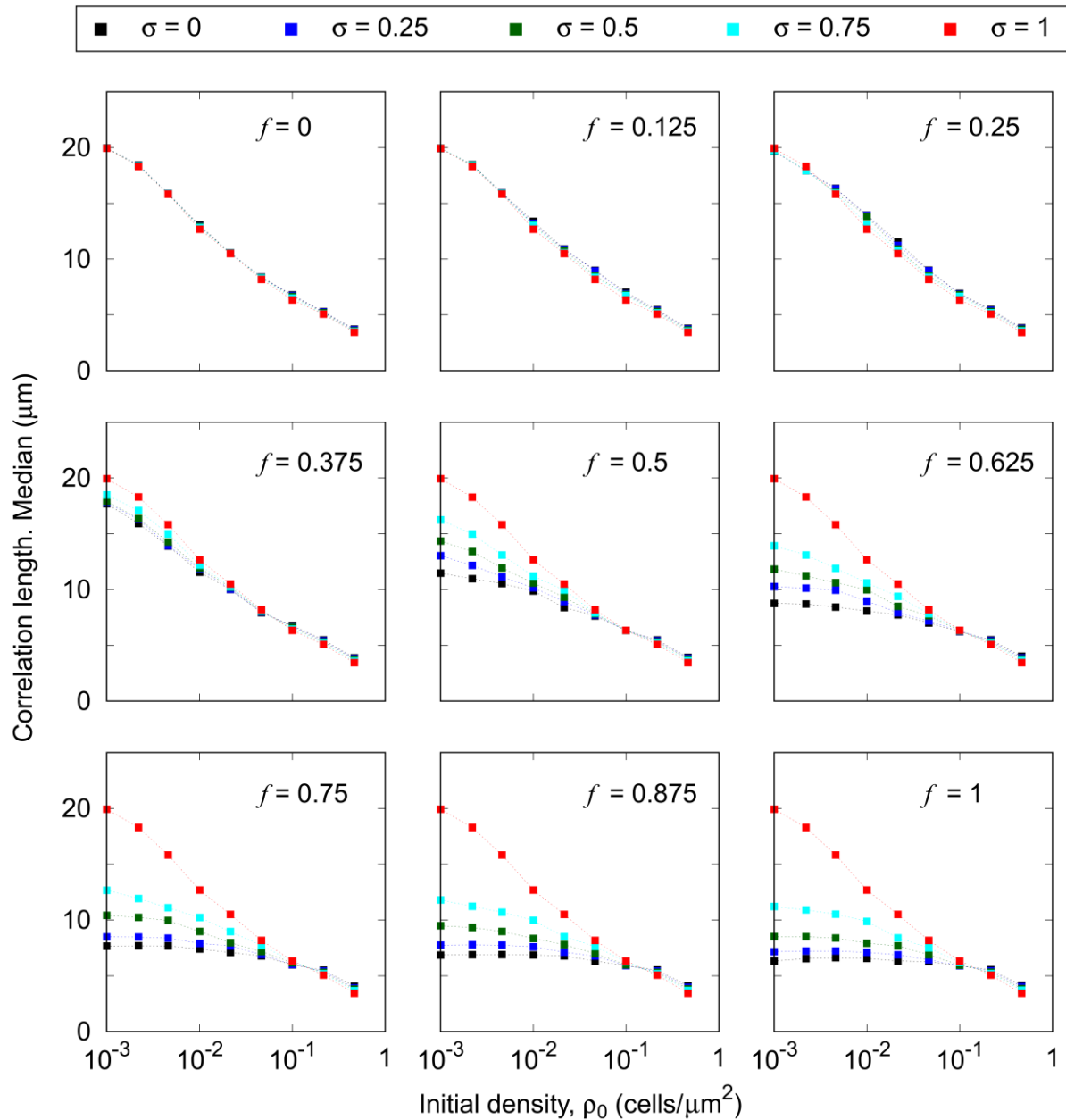
774
775
776
777
778

S2 Fig. Correlation length versus initial density. Mean correlation length, ξ , for different colonization strategies (σ, ρ_0) in several ecological conditions given by the flow intensity f . Each curve represents a cell adhesiveness σ . The color code is maintained in all the panels. Averages are taken over 2×10^6 independent model realizations.



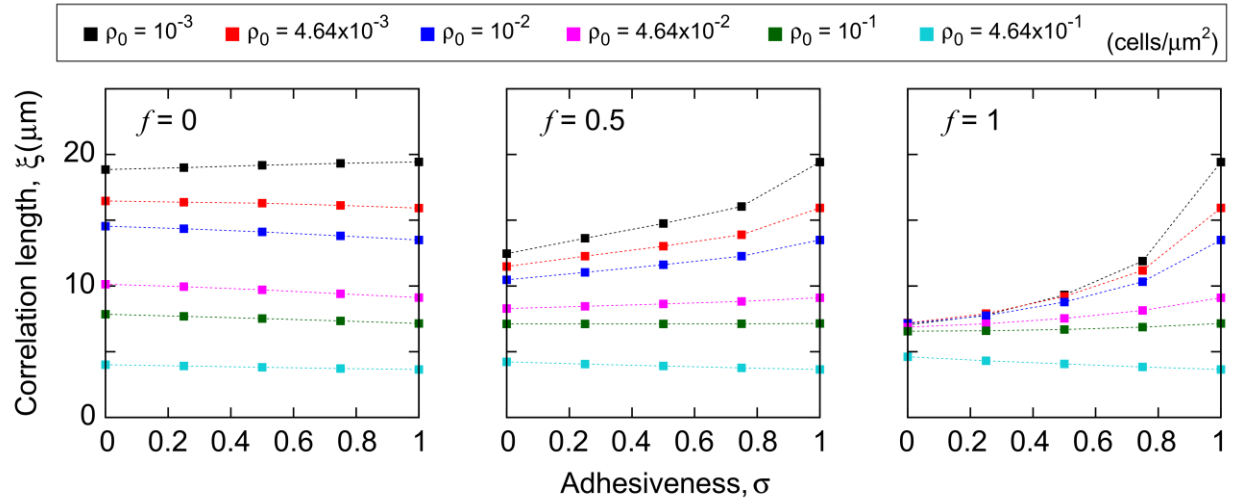
779
780
781
782
783
784
785
786

S3 Fig. Skewness of the correlation length. Skewness of the distribution of correlation lengths for different colonization strategies (σ , ρ_0) and ecological conditions, given by the flow intensity f . Each curve represents a value of the adhesiveness σ , whose color code is maintained in all the panels. The skewness is obtained from 2×10^6 independent realizations of the model. Horizontal dashed lines in each panel indicate the values ± 0.5 and the full lines, ± 1 . Skewness in the interval $[0.5, 1]$ in absolute value indicate that the data are moderately skewed, and if the skewness greater 1 in absolute value, then the distribution is highly skewed.



787
788
789
790
791

S4 Fig. Median correlation length. Median of the correlation length distribution for different colonization strategies (σ , ρ_0) and ecological conditions given by the flow intensity f . Each curve represents a value of the adhesiveness σ . The color code is maintained in all the panels. The median is obtained from a set of 2×10^6 independent model realizations.



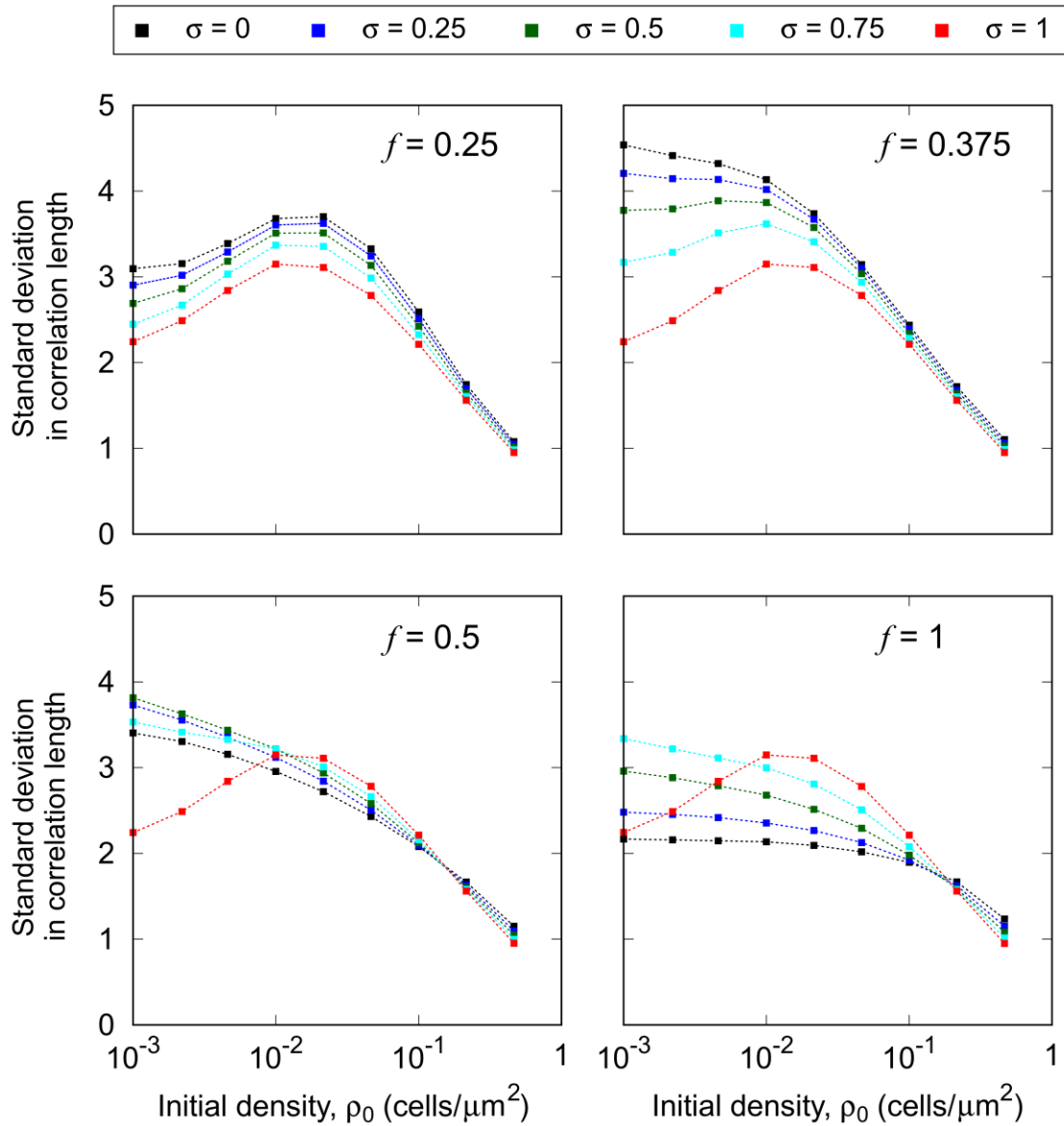
792

793

794 **S5 Fig. Correlation length versus cell adhesiveness.** Mean correlation length, ξ , for different
795 colonization strategies (σ , ρ_0) in several ecological conditions given by the flow intensity f . Each
796 curve represents a value of the initial density, ρ_0 . The color code is maintained in all the panels.

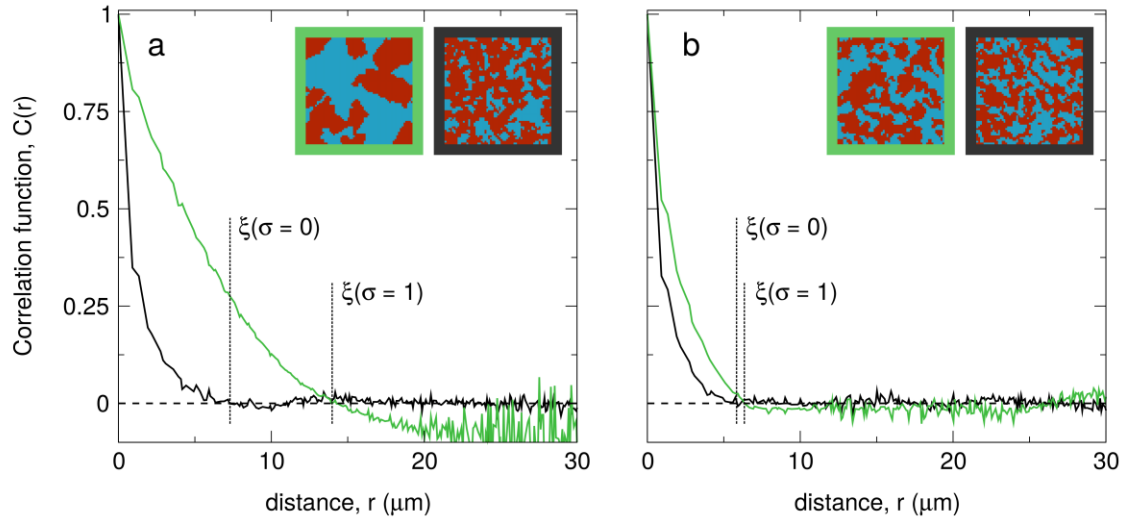
796

Averages are taken over 2×10^6 independent model realization.



797
798
799
800

S6 Fig. Cluster size variability. a) $f = 0.25$, b) $f = 0.375$, c) $f = 0.5$, d) $f = 1$. Each curve represents the standard deviation in ζ for a given adhesiveness, σ . Color code is maintained in all the panels. Averages are taken over 2×10^6 independent model realizations.



801
802
803
804
805
806
807
808
809
810

S7 Fig. Correlation function of individual model realizations. Correlation functions obtained for single realizations of the model at low (panel a; $\rho_0 = 10^{-3}$ cells/ μm^2) and high (panel b; $\rho_0 = 10^{-1}$ cells/ μm^2) initial density of cells. Correlation functions are obtained for the patterns shown in the snapshots. The color code indicates whether the pattern corresponds to $\sigma=1$ (green) or $\sigma=0$ (black) strains. The dashed lines point the value of the correlation length in each case, defined as the first zero of the correlation function.

811
812

813 **S1 Text. Size effect analysis.**

814

815 Our experimental results have been obtained using a square observation window of lateral
816 length $L = 60\mu m$ embedded within a much larger microfluidic device. Using the simulation
817 framework, we investigated whether the spatial measures in the occupation patterns are influenced
818 by the size of the focal system.

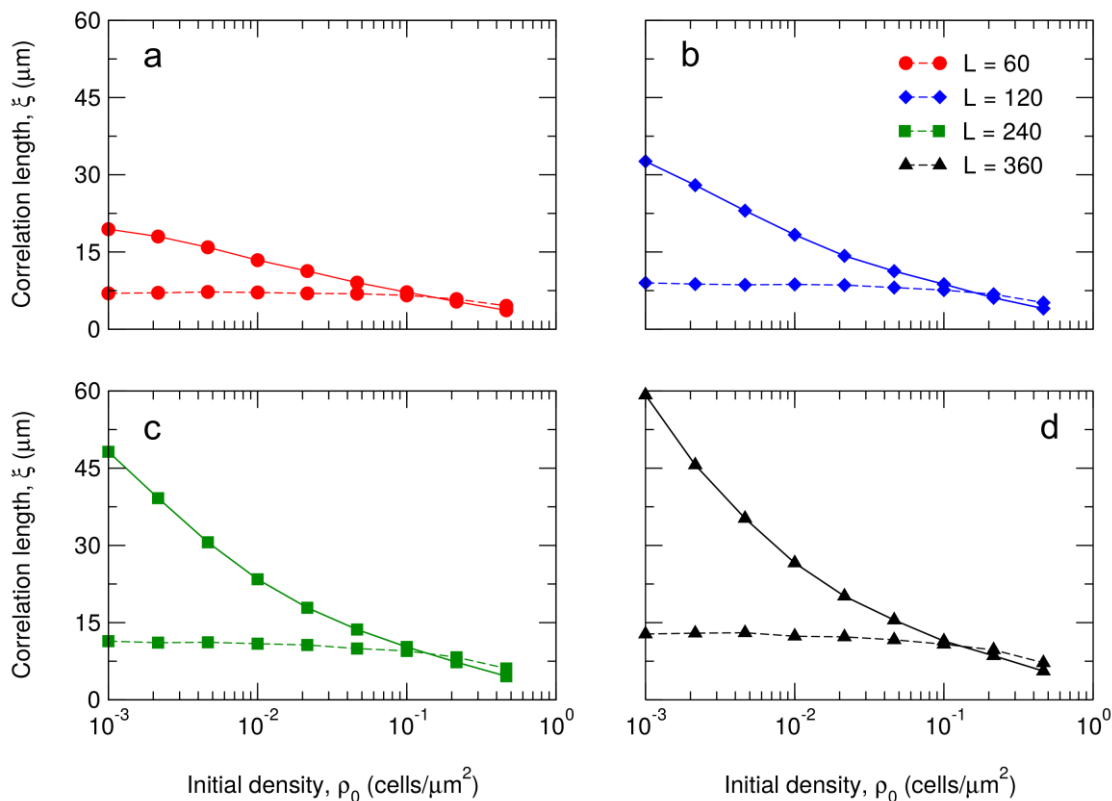
819 First, we focused on the correlation length, for both highly-adhesive and weakly-adhesive
820 strains, in the intense flow limit ($f = 1$). As shown in the main text, in this regime the model
821 accurately reproduces experimental results if the same focal area is used in both approaches.
822 Numerical simulations on larger systems confirm that both strains maintain the same qualitative
823 trends across simulated areas, and although the curves are quantitatively affected by the simulated
824 area, they intersect at the same value of the initial population density (Fig A1). The sublinear
825 scaling of the correlation length with system size, suggests a saturation of the correlation length in
826 the limit in which $\xi \ll L$ for any initial density and cell adhesiveness (Fig A2). Next, we prepared
827 a simulation setup in which we divided a system of lateral length $L = 120\mu m$ in four tiles of lateral
828 size $60\mu m$, and simultaneously measured the correlation length in the total system and in each of
829 the tiles. To ensure that the initial population density was constant for the whole system and each
830 tile, we initialized every tile with a total population density ρ_0 ($\rho_0/2$ of each strain on average).
831 Focusing on the intense flow limit ($f = 1$), the distance traveled by relocated cells in the direction
832 of the flow is a random number between 0 and L , so for a given focal area, the population mixing
833 depends on whether the system is isolated or embedded in a bigger one. However, the use of
834 periodic boundary conditions, as discussed in the main text, minimizes differences in the
835 correlation length for strong flows (Fig A3). The residual difference in the correlation length is
836 due to the fact that, in small isolated systems, the periodic boundary conditions can introduce small
837 additional correlations, since detached cells that exit the system through one of the borders and re-
838 enter through the opposite may be relocated close to their original position. These events are
839 equivalent to limited dispersal and hence tend to increase clonal cluster size. However, as it is
840 shown in Fig A3, their effect is negligible, reinforcing the validity of our periodic boundary
841 conditions.

842 Next, we extended our analysis to consider a $L=240\mu m$ patch with various flow intensities.
843 In this scenario, system size influences the outcome of the simulations in two directions. First, the
844 set of flow strengths for which patterns of weakly-adhesive cells have larger clonal clusters than
845 those made by adhesive strains increases considerably. Second, such regions show a larger
846 difference correlation length for bigger systems (Fig A4a). This result indicates that avoiding the
847 production of adhesion substances does not entail a residual gain (slightly larger clusters without
848 the metabolic cost of matrix production) but, for a wide range of environmental flows ($f < 0.4$),
849 such gain can be very significant, as much as that of matrix-production in strong environmental
850 flows (but, again, without the metabolic cost). If, on the other hand, we are observing a small
851 system that, instead of in isolation, is within a bigger one, the flow range for which weakly-
852 adhesive cells show larger clusters segregation is reduced to very weak intensities. This shrinkage
853 of the region results from our flow strength implementation discussed above: when the small

854 system is part of a bigger one, detached cells can travel larger distances even at weak
855 environmental flows and thus the range of limited dispersal is reduced (Fig A4b compared to
856 Figure 4). All this phenomenology indicates that, in a real system, the ratio between the typical
857 distance travelled with the flow and the system size will influence considerably the quantitative
858 (but not the qualitative) behavior of our measure for genetic segregation.

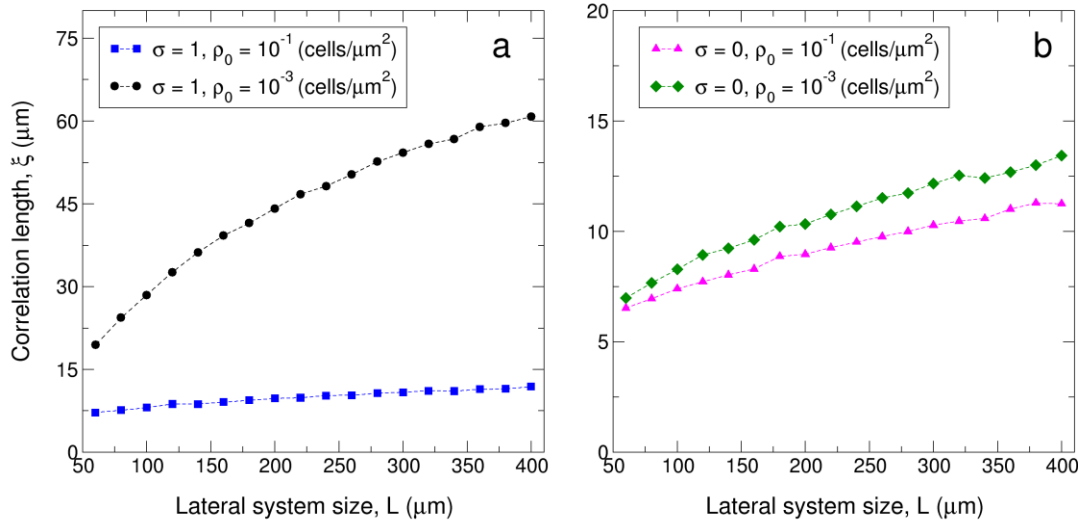
859 Finally, we analyzed the effect of the system size on cluster size variability for $\sigma = 1$ strains.
860 The standard deviation of the correlation length maintained its concavity regardless of the system
861 size, but it reached its maximum at different initial population densities (Fig A5a). Since highly-
862 adhesive cells are not relocated by the flow, the confluence pattern is strongly determined by the
863 spatial distribution of the founder population, and its correlation length variability depends on the
864 variability of the initial lineage mixing. Hence, it is the number of cells and not the density what
865 determines the position of the maximum in the standard deviation (Fig A5b). For high cell
866 numbers, it is very unlikely to randomly create a configuration with large clusters, whereas for low
867 cell numbers, the cluster size at confluence is necessarily large. In addition, for a fixed initial
868 density (or number of cells) the standard deviation increases with system size since the variability
869 in the spatial distribution of the founder population increases with system size.

870 In summary, the observation window can quantitatively affect some results of our analyses
871 as well as the regions of the parameter space in which they are expected. Therefore, not only
872 environmental forces, such as fluid flow, and bacterial traits, such as cell adhesion, are important
873 to quantify biofilm population structure. The size of the observation frames needs to be accounted
874 for as well. Importantly, however, the overall qualitative behavior of our results is not affected by
875 the size of the observation window and, therefore, any conclusion drawn for smaller surfaces can
876 be extrapolated to larger systems.



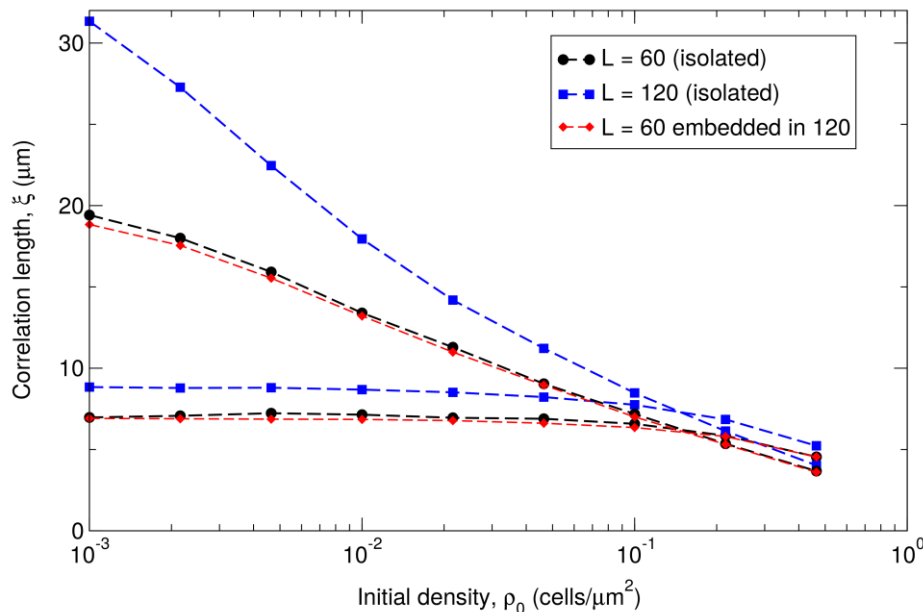
877

878 **Fig A1. Effect of surface in correlation length.** The clonal cluster size is strongly influenced by the
 879 extension of the colonized surface, although the trends of highly-adhesive and weakly-adhesive strains,
 880 and the crossing point between curves, are system size independent. Full lines correspond to $\sigma = 1$ and
 881 dashed lines to $\sigma = 0$. a) $L = 60\mu\text{m}$, b) $L = 120\mu\text{m}$, c) $L = 240\mu\text{m}$, d) $L = 360\mu\text{m}$.



882

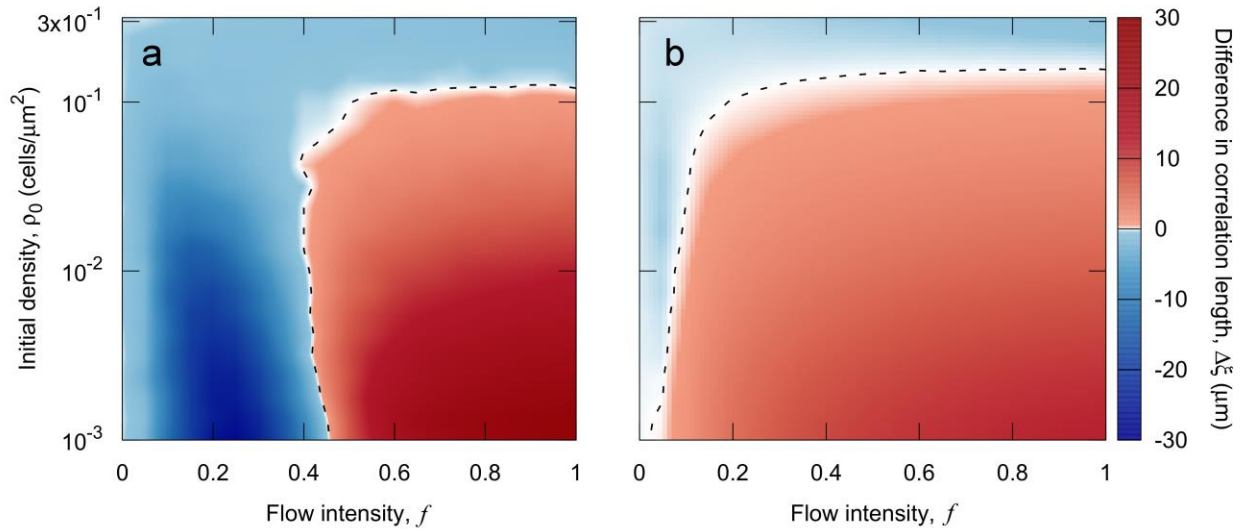
883 **Fig A2. Correlation length scaling with system size.** The correlation length scales sub-linearly, both for
 884 highly-adhesive (a) and weakly-adhesive (b) strains, with system size, which suggests a clonal cluster size
 885 saturation for large systems.



886

887 **Fig A3. Finite size effects in the correlation length.** Tiles within a larger system have the same
 888 correlation length than isolated surfaces of the same size. Simulations are run independently on systems
 889 of lateral length $L = 60\mu\text{m}$ (black circles) and $L = 120\mu\text{m}$ (blue squares). In this latter scenario, the system
 890 is divided in four tiles of lateral length $60\mu\text{m}$ and the correlation length of each of the tiles is
 891 independently obtained following the same protocol used in the $L = 60\mu\text{m}$ case.

892



893

894

895

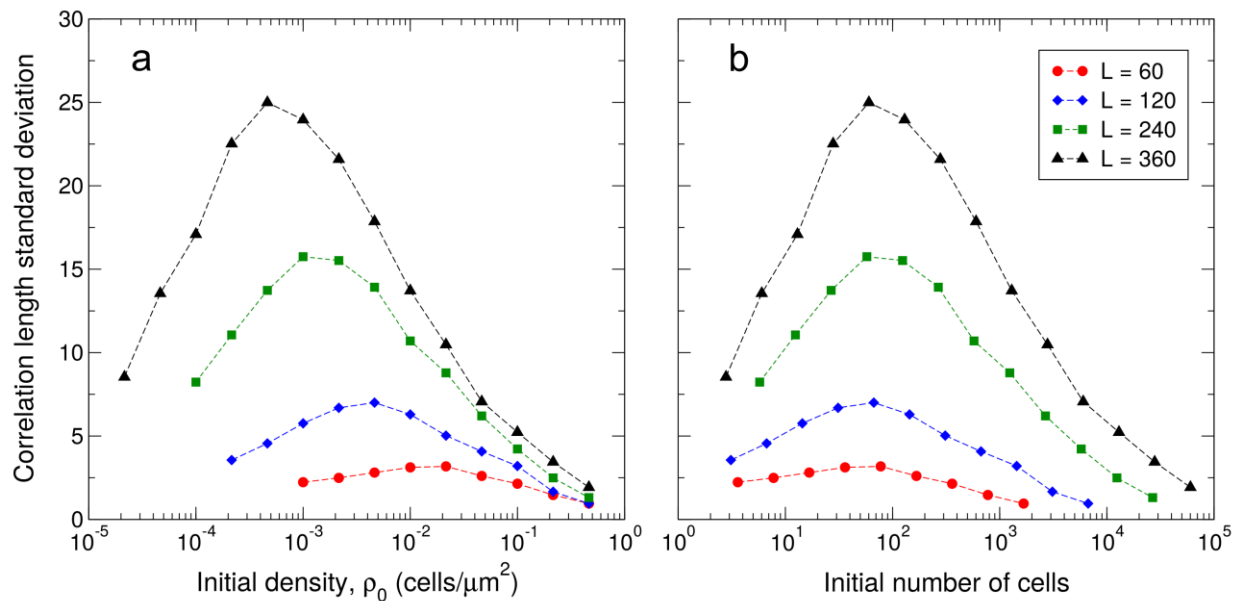
896

897

898

899

Fig A4. System size effects on the correlation length difference between highly-adhesive and weakly-adhesive strains. Correlation length differences are evaluated on a system of lateral length $L = 240\mu\text{m}$ (b) and the result is compared to what would be observed using an observation window of lateral length $L = 60\mu\text{m}$ within the system (a). In the latter case, each of the 16 observation windows is used as an independent replicate. Therefore, averages in b) are taken over 4000 replicates whereas 64000 independent realizations are gathered for the smaller system.



900

901

902

903

904

905

Fig A5. Variability in the correlation length is influenced by system size. Correlation length standard deviation versus initial population density (a), and initial number of cells (b). Color code is maintained in both panels