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## The Folding Pathway of an Ig Domain is Conserved On and Off the Ribosome

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Classification: Biological Sciences, Biophysics and Computational Biology

29 **Abstract**

30

31 Proteins that fold cotranslationally may do so in a restricted configurational space, due to the  
32 volume occupied by the ribosome. How does this environment, coupled with the close  
33 proximity of the ribosome, affect the folding pathway of a protein? Previous studies have  
34 shown that the cotranslational folding process for many proteins, including small, single  
35 domains, is directly affected by the ribosome. Here, we investigate the cotranslational folding  
36 of an all- $\beta$  immunoglobulin domain, titin I27. Using an arrest peptide-based assay and  
37 structural studies by cryo-EM, we show that I27 folds in the mouth of the ribosome exit  
38 tunnel. Simulations that use a kinetic model for the force-dependence of escape from arrest,  
39 accurately predict the fraction of folded protein as a function of length. We used these  
40 simulations to probe the folding pathway on and off the ribosome. Our simulations - which  
41 also reproduce experiments on mutant forms of I27 - show that I27 folds, while still  
42 sequestered in the mouth of the ribosome exit tunnel, by essentially the same pathway as free  
43 I27, with only subtle shifts of critical contacts from the C to the N terminus.

44

45

46 **Significance Statement**

47 Most proteins need to fold into a specific three-dimensional structure in order to function.

48 The mechanism by which isolated proteins fold has been thoroughly studied by experiment  
49 and theory. However, in the cell proteins do not fold in isolation, but are synthesized as linear  
50 chains by the ribosome during translation. It is therefore natural to ask at which point during  
51 synthesis proteins fold, and whether this differs from the folding of isolated protein  
52 molecules. By studying folding of a well characterized protein domain, titin I27, stalled at  
53 different points during translation, we show that it already folds in the mouth of the ribosome  
54 exit tunnel, and that the mechanism is almost identical to that of the isolated protein.

55

## 56 **Introduction**

57 To what extent is the cotranslational folding pathway of a protein influenced by the presence  
58 of the ribosome and by the vectorial emergence of the polypeptide chain during translation?  
59 Recent studies have shown that small proteins can fold inside the ribosome exit tunnel (e.g.,  
60 the small zinc finger domain ADR1a) (1), while other proteins can fold at the mouth of the  
61 tunnel (e.g., the three-helix bundle spectrin domains) (2); however some proteins may be  
62 simply too large to fold within the confines of the ribosome (e.g., DHFR) (3). The nature of  
63 cotranslational protein folding is determined by a number of biophysical factors, including  
64 the folding properties of the isolated protein (4-9), together with the effects the ribosome  
65 itself may have on the folding process (10-16). Due to the spatial constraints imposed upon  
66 the nascent chain by the confines of the tunnel, and effects due to the close proximity of the  
67 ribosome itself, the ribosome has been shown to influence directly the cotranslational folding  
68 of small proteins and single domains: The stability of folded or partly folded states may be  
69 reduced when folding occurs close to, or within the confines of, the ribosome (17, 18); the  
70 folding kinetics are expected to be correspondingly altered, with the rate of folding likely to  
71 be decreased and the unfolding rate increased, in close proximity to the ribosome(18).  
72 Interactions of the folded state or nascent polypeptide with the ribosome may also be either  
73 stabilising or destabilising (19, 20). Since translation is vectorial in nature, it is possible that  
74 when proteins fold cotranslationally they fold via different pathways than those used when  
75 proteins fold outside the ribosome, or when isolated proteins fold *in vitro* (2, 11, 21-24).  
76 However, addressing these issues is challenging, because standard protein folding methods  
77 are not directly applicable to cotranslational folding.

78 The folding of the protein close to the ribosome generates a pulling force on the  
79 nascent chain. This force has been probed by single molecule (25) as well as arrest peptide  
80 (AP) experiments (1-3). In this work, we use such arrest peptide-based cotranslational force-  
81 measurement experiments, simulations, and structural studies to investigate how the  
82 ribosome affects the folding of titin I27, a small all- $\beta$  immunoglobulin domain with a  
83 complex greek-key fold; the stability, kinetics and folding pathway of I27 has been  
84 extensively characterized in previous studies of the isolated domain (26, 27). In this study we  
85 investigate whether I27 can begin to fold in the confines of the ribosome, and if the folding  
86 pathway observed in the isolated domain is conserved during cotranslational folding. Results  
87 from all three techniques show that I27 folds in the mouth of the ribosome exit tunnel; our  
88 simulations correctly capture the onset of folding in I27 and three mutant variants, allowing

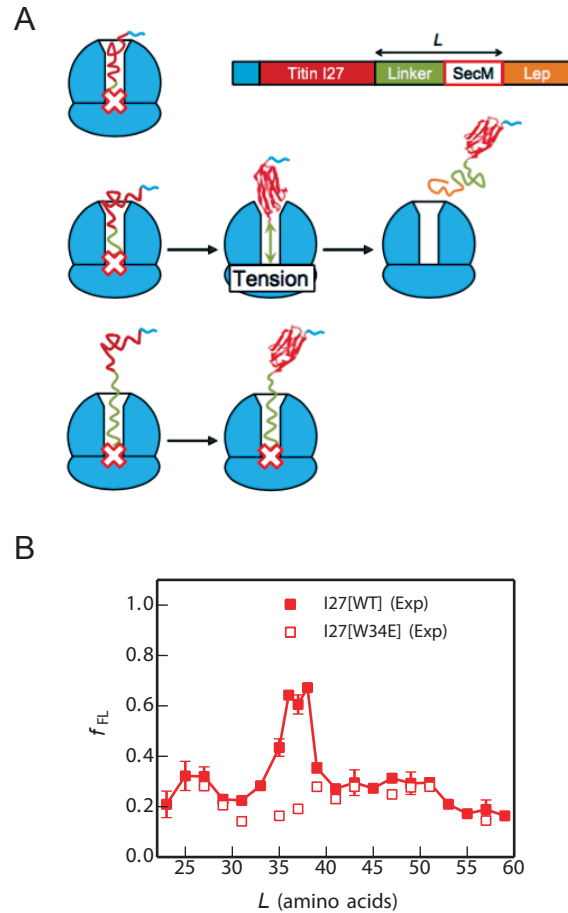
89 us to predict how destabilisation of regions that fold early and late in the isolated domain  
90 affect folding on the ribosome. Our simulations further show that the folding pathway of I27  
91 is largely unaffected by the presence of the ribosome, except for small but significant changes  
92 observed for contacts near the N and C termini.

## 93 **Results**

### 94 **I27 folds close to the ribosome**

95 In order to gain insight into when I27 can commence folding on the ribosome, we employed  
96 an arrest peptide force-measurement assay (28) carried out using the PURE *in vitro*  
97 translation system, as described in (1-3). In these experiments, the *E. coli* SecM arrest peptide  
98 (AP) is used to stall the nascent protein chain temporarily during translation. The yield of  
99 full-length protein which escapes stalling in a defined time interval ( $f_{FL}$ ), determined from  
100 SDS-PAGE gels, provides a proxy for the pulling force exerted on the nascent chain by the  
101 protein as it folds (1-3) (Figure 1A). By measuring  $f_{FL}$  for a set of constructs where the length  
102  $L$  of the linker between the target protein and the SecM AP is systematically varied, a force  
103 profile can be recorded that reflects the points during translation where the folding process  
104 starts and ends. Previous work has shown that the location of the main peak in a force profile  
105 correlates with the acquisition of protease resistance in an on-ribosome pulse-proteolysis  
106 assay (17, 29) and that the amplitude of the main force peak correlates with the  
107 thermodynamic stability of the protein (29, 30), indicating that the main peak represents a  
108 *bona fide* folding event rather than, *e.g.*, the formation of a molten globule. The sharp onset  
109 of the main force peak observed for most proteins analysed thus far (29) is also as expected  
110 for a cooperative folding event.

111 The force profile for wild-type I27 (Figure 1B) has a distinct peak at  $L = 35-38$   
112 residues (see Methods for sequences of the constructs). This peak is absent from the force  
113 profile for the mutant I27[W34E], a non-folding variant of I27, demonstrating that the peak is  
114 due to a folding event and not, for example, to non-specific interactions of the unfolded  
115 nascent chain with the ribosome. The non-zero  $f_{FL}$  for the non-folding mutant is attributed to  
116 the spontaneous rate of escape from arrest in the absence of acceleration by forces associated  
117 with folding. Since it takes  $\sim 35$  residues in an extended conformation to span the  $\sim 100$  Å  
118 long exit tunnel (31), the critical length  $L \approx 35$  residues suggests that I27 starts folding while  
119 in mouth of the exit tunnel.



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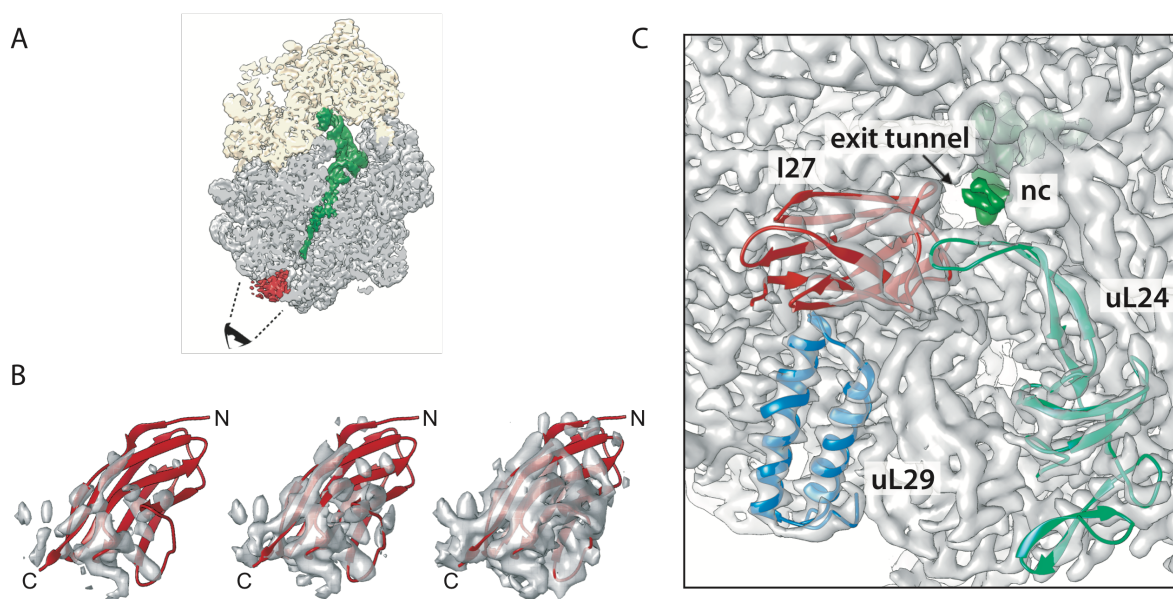
121 **Figure 1.** Cotranslational folding of the titin I27 domain by force-profile analysis. (A) The  
122 force-measurement assay (modified from (2)). I27, preceded by a His-tag, is placed  $L$   
123 residues away from the last amino acid of the SecM AP, which in turn is followed by a 23-  
124 residue C-terminal tail derived from *E. coli* LepB. Constructs are translated for 15 min. in the  
125 PURE *in vitro* translation system, and the relative amounts of arrested and full-length peptide  
126 chains produced are determined by SDS-PAGE. The fraction full-length protein,  $f_{FL}$ , reflects  
127 the force exerted on the AP by the folding of I27 at linker length  $L$ . At short linker lengths  
128 (top), there is not enough room in the exit tunnel for I27 to fold, little force is exerted on the  
129 AP, and the ribosome stalls efficiently on the AP ( $f_{FL} \approx 0$ ). At intermediate linker lengths  
130 (middle), there is enough room for I27 to fold but only if the linker segment is stretched,  
131 force is exerted on the AP, and stalling is reduced ( $f_{FL} > 0$ ). At long linker lengths (bottom),  
132 I27 has already folded when the ribosome reaches the last codon in the AP, and again little  
133 force is exerted on the AP ( $f_{FL} \approx 0$ ). (B) Force profiles for the I27 domain (solid squares) and  
134 the non-folding (nf) mutant I27[W34E] (open squares). The standard error of  $f_{FL}$  is calculated  
135 for values of  $L$  where three or more experiments were performed.

136

### 137 **Cryo-EM shows that I27 folds in the mouth of the exit tunnel**

138 To confirm that the peak in the force profile corresponds to the formation of a folded I27  
139 domain, we replaced the SecM AP with the stronger TnaC AP (32-34) and purified stalled  
140 ribosome-nascent chain complexes (RNCs) carrying an N-terminally His-tagged I27[L=35]

141 construct (see Methods). The construct was expressed in *E. coli*, RNCs were purified using  
142 the N-terminal His-tag, and an RNC structure with an average resolution of 3.2 Å (SI  
143 Appendix, Fig. S1) was obtained by cryo-EM. In addition to the density corresponding to the  
144 TnaC AP, a well-defined globular density (~4.5-9 Å resolution) was visible protruding from  
145 the exit tunnel (Figure 2A). Given the flat ellipsoidal shapes of the protruding density and of  
146 the I27 structure, there is only one way to fit the NMR structure of I27 (PDB 1TIT (35)) that  
147 gives a good Fourier-shell correlation between the isolated I27 density and the map generated  
148 from the I27 PDB model (SI Appendix, Fig. S2). In the fitted model, the C-terminal end of  
149 I27 extends into the exit tunnel and a  $\beta$ -hairpin loop on ribosomal protein uL24 is lodged in a  
150 cavity in I27 (Figure 2B and Supporting Video S1). The I27 domain further packs against  
151 ribosomal protein uL29 and ribosomal 23S RNA (Figure 2C), as if it is being pulled tight  
152 against the ribosome by the nascent chain. We conclude that the peak at  $L = 35$ -38 residues in  
153 the force profile indeed represents the cotranslational folding of the I27 domain at the tunnel  
154 exit.



155

156 **Figure 2.** Cryo-EM structure of I27[ $L=35$ ] RNCs. (A) Cryo-EM reconstruction of the I27–  
157 TnaC[ $L = 35$ ] RNC. The ribosomal small subunit is shown in yellow, the large subunit in  
158 grey, the peptidyl-tRNA with the nascent chain in green, and an additional density  
159 corresponding to I27 at the ribosome tunnel exit in red. The black cartoon eye and dash lines  
160 indicate the angle of view in panel (C). The density contour level for feature visualization is  
161 at 1.7 times root-mean-square deviation (1.7 RMSD). (B) Rigid-body fit of the I27 domain  
162 (PDB 1TIT) to the cryo-EM density map displaying from high (left) to low (right) contour  
163 levels at 2.6, 2.0 and 1.4 RMSD, respectively. N and C represent the N and C termini of the  
164 I27 domain, respectively. (C) View looking into the exit tunnel (arrow) with density for the  
165 nascent chain (nc) in dark green. Ribosomal proteins uL29 (blue; PDB 4UY8), uL24 (light  
166 green; the  $\beta$  hairpin close to I27 domain was re-modeled based on PDB 5NWY) and the

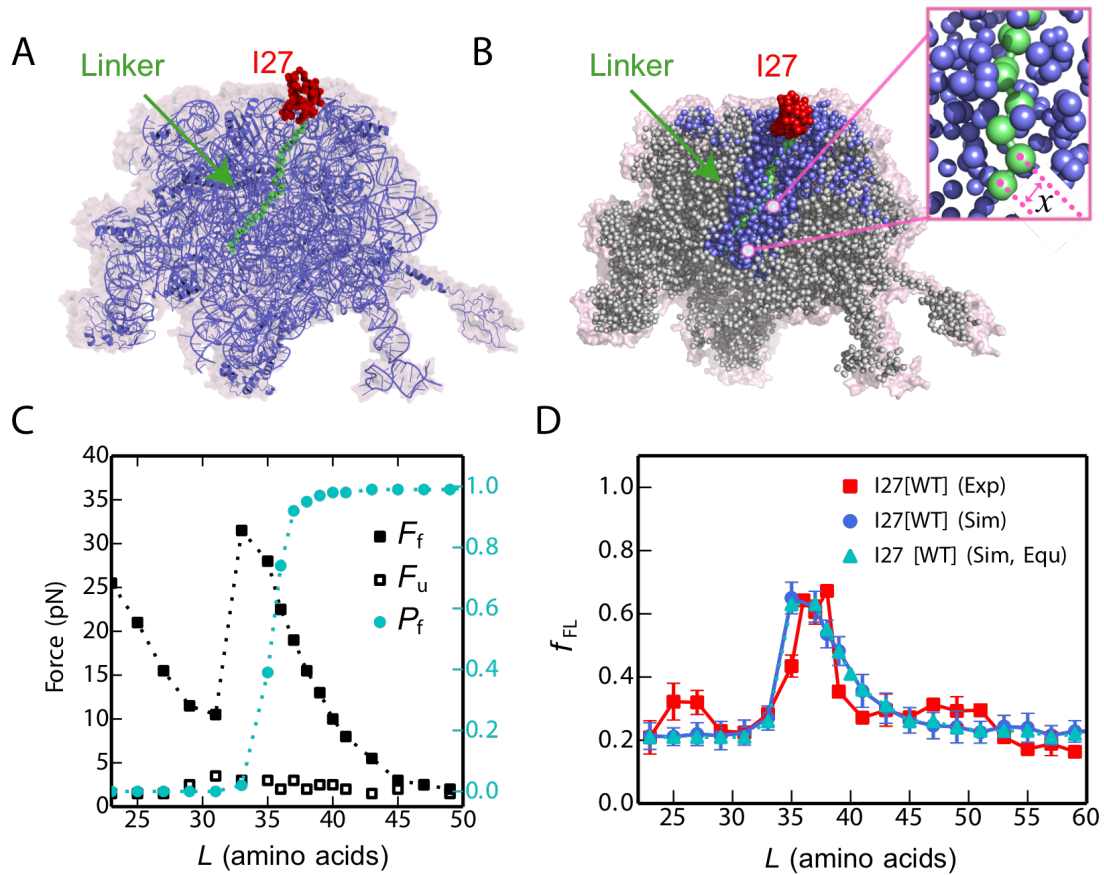
167 fitted I27 domain (red) are shown in cartoon mode; 23S RNA and proteins not contacting I27  
168 are shown as density only. The density contour level is at 5 RMSD excluding tRNA, nascent  
169 chain and I27 domain, which are displayed at 1.7 RMSD.

170

## 171 **Coarse-grained molecular dynamics simulations recapitulate I27 folding on the** 172 **ribosome**

173 The yield of folded protein in arrest peptide experiments has been used as a proxy for the  
174 pulling forces that are exerted on the nascent chain at different points during translation in all  
175 studies to date (1, 2, 29). Here, to further elucidate the molecular origins of these forces and  
176 provide a quantitative interpretation of the observed folding yield of I27, we have calculated  
177 force profiles based on coarse-grained MD simulations (see Methods). Briefly, in the MD  
178 model, the 50S subunit of the *E. coli* ribosome (36) (PDB 3OFR) and the nascent chain are  
179 explicitly represented using one bead at the position of the C $\alpha$  atom per amino acid, and three  
180 beads (for P, C4', N3) per RNA base (Figure 3A). The interactions within the protein were  
181 given by a standard structure-based model (37-39), which allowed it to fold and unfold.  
182 Interactions between the protein and ribosome beads were purely repulsive (40) and the  
183 ribosome beads were fixed in space, as in previous simulation studies (18). I27 was  
184 covalently attached to unstructured linkers having the same sequences as those used in the  
185 force-profile experiments (Figure 3B) and the C terminus of the linker was tethered to the last  
186 P atom in the A-site tRNA (41) with a harmonic potential, allowing the force exerted by the  
187 folding protein to be directly measured. The potential chosen was stiff enough that  
188 displacements caused by typical pulling forces were smaller than 1 Å. For each linker length  
189  $L$ , we used umbrella sampling to determine the average force exerted on the AP by the  
190 protein in the folded and unfolded states while arrested, as well as the populations of those  
191 two states (Figure 3C). We also estimated the folding and unfolding rates directly from  
192 folding/unfolding simulations.

193



194

195 **Figure 3.** MD simulations of cotranslational folding of I27. (A) 50S subunit of the *E. coli*  
 196 ribosome (PDB 3OFR) with I27[L=35] attached via an unstructured linker. (B) Coarse-  
 197 grained model for I27 (red) and linker (green), with surrounding ribosomal pseudo-atoms in  
 198 blue. Pseudo-atoms with grey colour are not used in the simulations. The instantaneous force  
 199 exerted on the AP is calculated from the variation in the distance  $x$  between the C-terminal  
 200 Pro pseudo-atom and the next pseudo-atom in the linker (see inset). (C) Average forces  
 201 exerted on the AP by the unfolded state ( $F_u$ , empty symbols) and folded state ( $F_f$ , filled  
 202 symbols) of I27 at different linker lengths  $L$ . The average fraction folded I27 for different  
 203  $L$ ,  $P_f$ , is shown in cyan on the right axis. Free energy profiles at each linker length are shown  
 204 in SI Appendix, Fig. S4. (D) Experimental (red square) force profiles for cotranslational  
 205 folding of I27. Force profiles calculated from simulations using full kinetic scheme or pre-  
 206 equilibrium model are shown in blue circle and cyan triangle respectively. The RMSD of the  
 207  $f_{FL}$  between experiment and simulation is 0.08.

208

209

210 Given the experimentally-determined force-dependence of the escape rate  $k(F)$  (25), here  
 211 approximated by a Bell-like model (42), we can calculate the expected escape rate while the  
 212 protein is in the unfolded or folded state, from which the fraction full-length protein obtained  
 213 with a given linker length and incubation time can be determined from a kinetic model, as  
 214 described in Methods. The calculated  $f_{FL}$  profile for I27 is shown in Figure 3D (see also SI

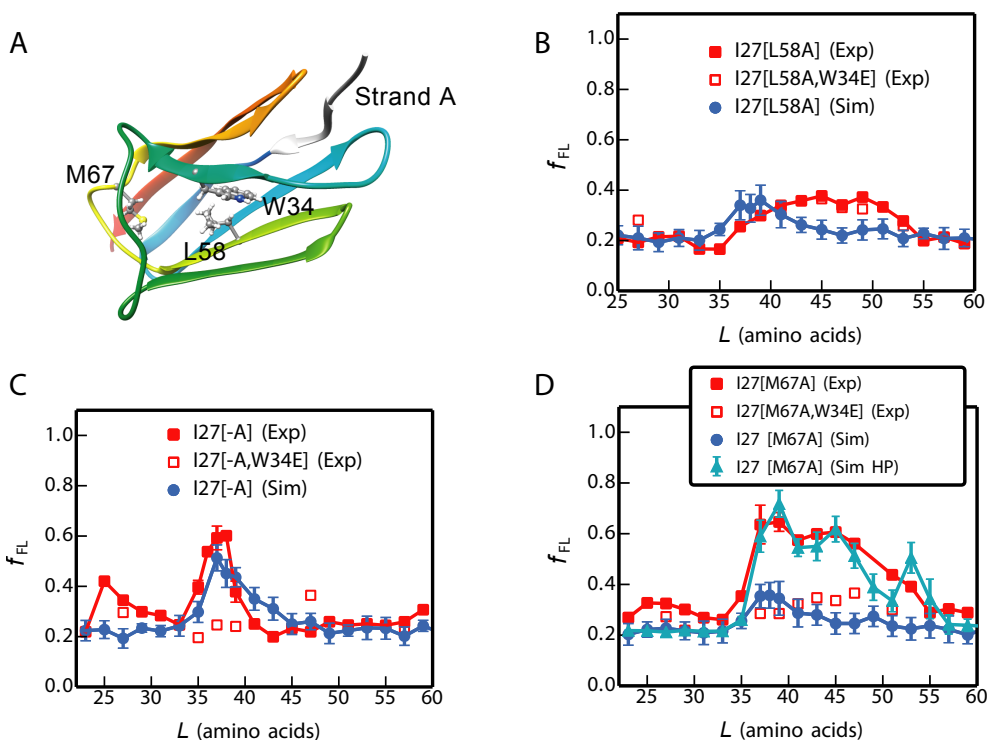


215 Appendix, Fig. S5) for the full solution of the kinetic model, as well as for an approximation  
216 in which the folding and unfolding rates are assumed to be faster than the escape rate (“pre-  
217 equilibrium”). Both results are very consistent with each other, as well as with the  
218 experimental profile. The peak in the folding yield arises as consequence of two opposing  
219 effects, the force exerted by the folding protein and population of the folded state, which  
220 respectively decrease and increase as the linker length increases. In the simulations with the  
221 I27[L=35] construct, the folded I27 domain is seen to occupy positions that largely overlap  
222 with the cryo-EM structure (Supporting Video S2). Overall, these results suggest that the MD  
223 model provides a good representation of the folding behaviour of the I27 domain in the  
224 ribosome exit tunnel. To show that the simulation model is not specific to I27, we have also  
225 applied it to another two proteins with different topologies for which experimental force  
226 profiles have been recorded, Spectrin R16 (all- $\alpha$  fold) and S6 ( $\alpha/\beta$  fold) (2, 29). In these  
227 cases, we also obtain force profiles similar to experiment (SI Appendix, Fig. S6 and S7).

228

### 229 **Force profiles of I27 variants probe the folding pathway**

230 To test whether the cotranslational folding pathway is the same as that observed for the  
231 isolated I27 domain *in vitro*, we investigated three destabilised variants of I27, both by  
232 simulation and experiment. One mutation in the core, Leu 58 to Ala (L58A), located in  $\beta$ -  
233 strand E (Figure 4A) destabilizes the protein by 3.2 kcal mol<sup>-1</sup>, and removes interactions that  
234 form early during folding of the isolated domain, playing a key role in formation of the  
235 folding nucleus ( $\phi$ -value = 0.8) (26). Two further mutations, M67A and deletion of the N-  
236 terminal A-strand, remove interactions that form late in the folding of I27 (*i.e.*, both mutants  
237 have low  $\phi$ -values (26, 27)). The A-strand is the first part of I27 to emerge from the  
238 ribosome, while M67 is located in a part of I27 that is shown by cryo-EM to be located in  
239 very close proximity to a  $\beta$  hairpin loop of ribosomal protein uL24 in I27-TnaC[L=35] RNCs  
240 (SI Appendix Fig. S3A). The interaction with the I27 domain shifts the tip of this uL24  
241 hairpin by about 6 Å compared to its location in other RNC structures (SI Appendix Fig.  
242 S3B).



243

244 **Figure 4.** Simulations capture the experimental force profiles for mutant I27 domains. (A)  
 245 Mutated residues in I27 (sticks). (B-D) Experimental (red) force profiles and calculated ones  
 246 from full kinetic scheme (blue) for (B) I27[L58A], (C) A-strand deletion mutant I27[-A], (D)  
 247 I27[M67A]. I27[M67A] (Sim HP) represents a simulation in which hydrophobic interactions  
 248 between I27[M67A] and ribosome proteins uL23/uL29 are included. Experimental force  
 249 profiles for non-folding mutants that contain an additional W34E mutation are shown as red  
 250 open squares. The RMSD of the  $f_{FL}$  between experiment and simulation for I27[L58A], I27[-  
 251 A] are 0.07 and 0.08 respectively. For I27[M67A], the  $f_{FL}$  RMSD is 0.07 between  
 252 experiment and simulation (Sim HP).

253

254 The simulated force profile for the L58A variant predicts a much lower force peak than for  
 255 wild-type I27; likewise, the experimental force peak is lower and broader than for wild-type,  
 256 extending from  $L = 37$ -53 residues (Figure 4B). The  $f_{FL}$  values are very similar to those  
 257 obtained for I27[L58A,W34E], a non-folding variant of I27[L58A]. Therefore, the weak  
 258 forces seen at  $L \approx 40$ -50 residues are not due to a folding event, indicating that I27[L58A]  
 259 does not exert an appreciable force due to folding near the ribosome.

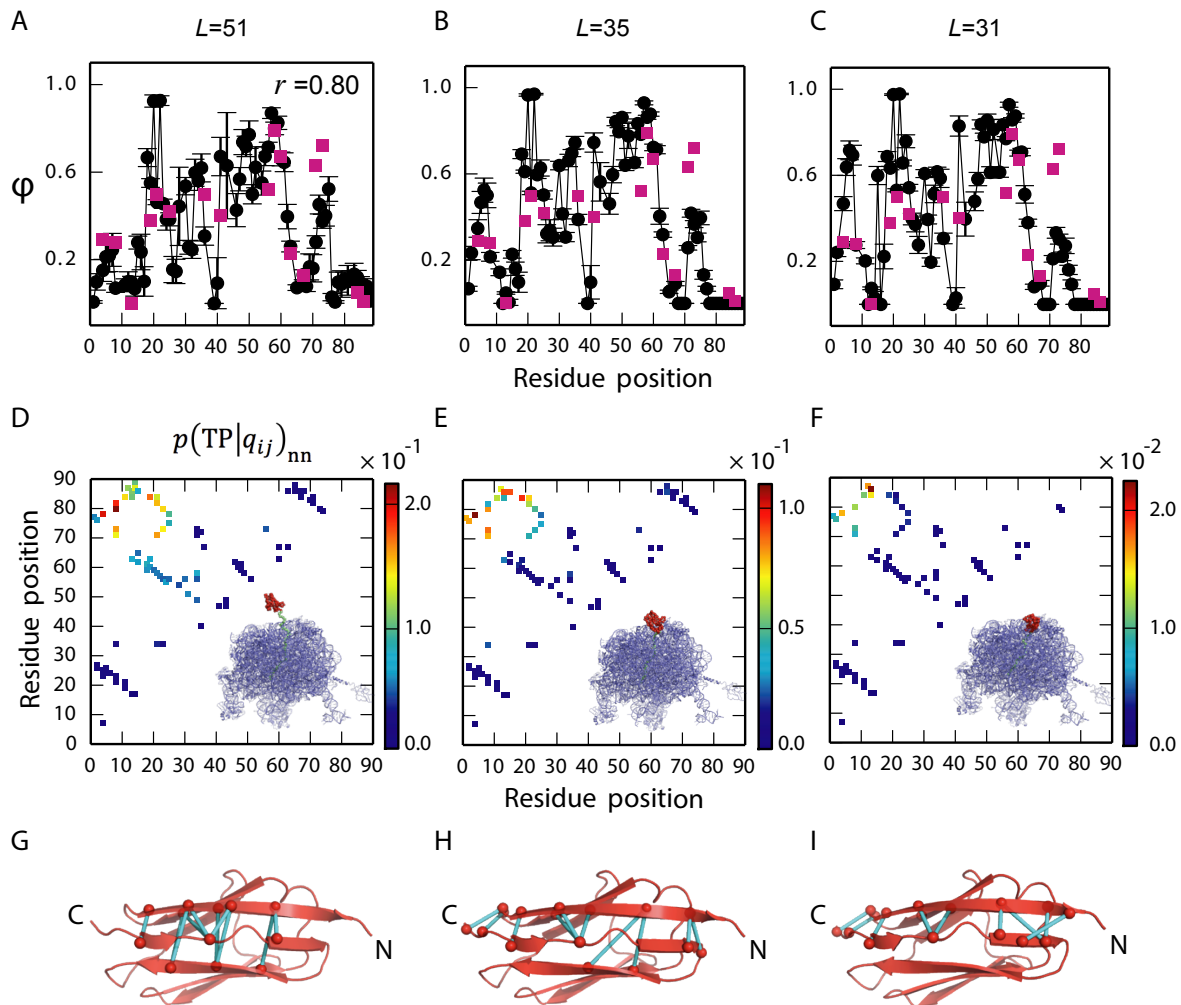
260 The A-strand comprises the first seven residues of I27 and removal of this strand,  
 261 I27[-A], results in a destabilisation of  $2.78 \text{ kcal mol}^{-1}$ ; however, both the simulated and  
 262 experimental force profiles for I27[-A] are very similar to those for wild-type I27 (Figure  
 263 4C). Residue M67 is located in the E-F loop, and mutation to alanine results in a  
 264 destabilisation of  $2.75 \text{ kcal mol}^{-1}$ ; for this variant, folding commences at  $L \approx 35$  residues as  
 265 for wild-type I27, but the peak is much broader (Figure 4D). Non-folding control experiments

266 for variants I27[-A,W34E], and I27[M67A,W34E] (Figure 4C and D) show that the peaks in  
267 the force profiles for these variants are due to a folding event. These results show that  
268 deletion of the A-strand and destabilisation of the E-F loop do not affect the onset of  
269 cotranslational folding of I27, but that the M67A mutation increases the width of the folding  
270 transition. The simulation model used for the other mutants does not predict such a broad  
271 peak, suggesting that it may be necessary to include additional factors to reproduce the data  
272 for M67A. One possibility which may explain the result would be favourable interactions of  
273 the folded M67A mutant with the ribosome surface. The ribosomal surface proteins uL23 and  
274 uL29 have been suggested to form a potential interaction site for nascent proteins such as  
275 trigger factor (43), signal recognition particle (44) and SecYE (45). Here we have explored  
276 the hypothesis that the broad force peak of mutant M67A might due to interactions between  
277 an exposed hydrophobic cavity on I27[M67A] resulting from the mutation, and hydrophobic  
278 surface residues of ribosomal proteins uL23 and uL29. By introducing such interactions into  
279 the model, we are able to obtain a broad peak in the force profile very similar to that seen in  
280 experiment (Figure 4D).

### 281 **The folding pathway is only subtly affected by the presence of the ribosome.**

282 To compare the folding pathways when the protein is folding near the tunnel exit or outside  
283 the ribosome, we estimated  $\phi$ -values based on the transition paths of I27 folding on the  
284 ribosome from our coarse-grained simulations, using a method introduced previously(46).  
285 The transition paths are those regions of the trajectory where the protein crosses the folding  
286 barrier, here defined as crossing between  $Q = 0.3$  and  $Q = 0.7$ . For each linker length, 30  
287 transition paths were collected from MD simulations. To reduce the uncertainty in the  
288 experimental reference data, we only compared with experimental  $\phi$ -values if the change in  
289 folding stability between the mutant and the wild type is sufficiently large ( $|\Delta\Delta G| > 7$  kJ/mol)  
290 (47). As seen in Figure 5A, when the linker is long ( $L = 51$  residues) and I27 is allowed to  
291 fold outside the ribosome, the calculated  $\phi$ -values are consistent (Spearman correlation  
292  $r=0.80$ ) with the experimental values obtained for the folding of isolated I27 *in vitro* (26). For  
293 shorter linker lengths ( $L = 31$  and 35 residues), calculated  $\phi$ -values remain largely unchanged  
294 except for a slight increase near the N terminus (around residues 3-6) and a slight decrease  
295 near the C terminus (around residues 72-74) (Figure 5B and C).

296



297

298 **Figure 5.** Simulated folding pathways for ribosome-tethered I27. LH column,  $L=51$ ; middle  
 299 column,  $L=35$ ; RH column,  $L=31$ . Top panels: Simulated  $\phi$ -values for I27 (black).  $\phi$ -values  
 300 determined by *in vitro* folding of purified I27 are shown as red squares. At  $L=51$  the  
 301 simulated  $\phi$ -values match well with experiment (spearman correlation  $r=0.80$ ). At  $L=35$  and  
 302  $L=31$  the simulated  $\phi$ -values are higher at the N terminus and lower at the C terminus, than  
 303 the experimental values, reflecting a change in importance of these regions when I27 folds in  
 304 the confines of the ribosome. Middle row: Relative probability that if a particular contact is  
 305 formed then the protein is on a folding trajectory,  $p(\text{TP}|q_{ij})_{\text{nn}}$ . When the protein is  
 306 constrained the limiting factor is formation of a few key contacts. A cartoon of the ribosome  
 307 with I27 in red is shown on each panel. Bottom row: The top ten most important contacts are  
 308 coloured in cyan on the native structure.

309

310 To obtain a more detailed picture regarding the relative importance of different native  
 311 contacts in the folding mechanism, we computed the conditional probability of being on a  
 312 transition path (TP), given the formation of a contact  $q_{ij}$  between residues  $i$  and  $j$ ,  
 313  $p(\text{TP}|q_{ij})_{\text{nn}}$  (48). This quantity indicates which native contacts are most important for

314 determining a successful folding event.  $p(\text{TP}|q_{ij})_{\text{nn}}$  is closely related to the frequency of the  
315 contact  $q_{ij}$  on transition paths  $p(q_{ij}|\text{TP})$ , but is effectively normalized by the probability that  
316 the contact is formed in non-native states  $p(q_{ij})_{\text{nn}}$ , and can be expressed as:

$$317 \quad p(\text{TP}|q_{ij})_{\text{nn}} = \frac{p(q_{ij}|\text{TP})p(\text{TP})_{\text{nn}}}{p(q_{ij})_{\text{nn}}} \quad [1]$$

318 where  $p(\text{TP})_{\text{nn}}$  is the fraction of non-native states which are on transition paths at  
319 equilibrium. The subscript nn means that only the non-native segments of a trajectory are  
320 included, *i.e.*, unfolded states and transition paths; the native, folded state is not included in  
321 the calculation since native contacts are always formed in this state. The simulations suggest  
322 that formation of native contacts between the N and C termini is somewhat more important  
323 when folding takes place in the mouth of the exit tunnel ( $L = 31$  residues) than far outside the  
324 ribosome ( $L = 51$  residues) (Figure 5D-F, upper left-hand corner in the panels). This is likely  
325 due to the greater difficulty of forming these contacts (examples are shown in Figure 5G-I)  
326 under ribosomal confinement; therefore, forming them becomes more critical in enabling the  
327 protein to fold.

328

## 329 **Discussion**

330 Using a combination of MD simulation, force-profile measurements and cryo-EM, we have  
331 investigated the cotranslational folding pathway of the 89-residue titin I27 domain. I27 has  
332 been extensively characterised in previous *in vitro* folding studies (26, 27, 49-60). Results  
333 from all three techniques show that wild-type I27 folds in the mouth of the ribosome exit  
334 tunnel; in the cryo-EM structure of I27-TnaC[ $L=35$ ] RNCs, I27 packs against ribosomal  
335 proteins uL24, uL29, and ribosomal 23S RNA. This is in apparent contrast to a previous  
336 NMR study on another Ig-like protein, in which the domain was shown to acquire its native  
337 fold (as reflected in the NMR spectrum) only when fully outside the ribosome tunnel, at  $L =$   
338 42-47 residues linker length (20).

339 In order to determine the molecular origin of the measured force profile, we  
340 performed molecular dynamics simulations of I27 folding on the ribosome, varying the  
341 length of the linker sequence between the arrest peptide and the I27 domain. We calculated  
342 the pulling force directly from the simulations and translated this into yield of folded protein  
343 using a kinetic model parameterized based on known release kinetics of the SecM AP. This  
344 enabled us to recapitulate the experimental arrest peptide force measurement profile, and

345 therefore relate  $f_{FL}$  directly to the force exerted on the arrest peptide. Our simulations  
346 demonstrate the direct effect that the restoring force of the nascent chain can have on  
347 determining when the protein folds on the ribosome. We show that  $f_{FL}$  depends upon a  
348 combination of the force exerted by the folded protein and the fraction of folded protein at  
349 the given linker length  $L$ .

350 In order to relate how destabilization of regions that fold early and late in the isolated  
351 domain affects folding on the ribosome, we used simulations to predict the onset of folding in  
352 three mutant variants of I27. A previous  $\phi$ -value analysis of I27 (26) showed that early  
353 packing of the structurally central  $\beta$ -strands drives the folding of this domain, while  
354 peripheral strands and loop regions pack later in the folding process. Mutations in the folding  
355 core (such as L58A) slow folding, whereas mutations in the periphery have no effect on  
356 folding rates (26). L58 is a key residue in the critical folding nucleus and almost fully packed  
357 in the transition state, in isolated domain studies. The simulated and experimental force  
358 profiles of I27 [L58A] show that this variant does not fold in or near the exit tunnel; hence,  
359 destabilisation of the central folding core prevents folding close to the ribosome. Since  
360 isolated I27[L58A] is fully folded, it is likely that this variant can only fold cotranslationally  
361 at longer linker lengths, when it is no longer in close proximity to the ribosome and exerts  
362 little force on the nascent chain.

363 Our experiments show that I27 variants destabilized in regions of the protein that are  
364 unstructured, or only partially structured, in the transition state, are still able to commence  
365 folding close to the ribosome. The force profiles reveal that the onset of folding of mutants  
366 with the A-strand deleted, or with the Met 67 to Ala mutation in the E-F loop, is the same as  
367 for wild-type although these have a similar destabilisation as L58A (Figure 4). The broader  
368 peak observed experimentally for M67A is harder to interpret. A plausible explanation is that  
369 the mutation introduces non-specific interactions of the folded domain with the ribosome  
370 surface, and we have shown that incorporating such interactions into the simulations could  
371 reproduce the results. An additional factor may be that the mutation is in a region that  
372 interacts closely with ribosomal protein uL24 in the wild-type cryo-EM structure (SI  
373 Appendix Fig. S3).

374 Our simulations reproduce the onset of folding in the three mutant variants of I27 (Figure 4),  
375 and so give us the confidence to investigate how confinement within the ribosome affects the  
376 folding pathway of I27. We used simulations to investigate the folding of I27 arrested on the  
377 ribosome at various linker lengths, using a Bayesian method for testing the importance of

378 specific contacts on the folding pathway, as well as by computing  $\phi$ -values (Figure 5).  
379 Overall, we find that the mechanism and pathway of folding are robust towards variation in  
380 linker length and relatively insensitive to the presence of the ribosome; small but significant  
381 changes are observed only for contacts near the N and C termini. These shifts are consistent  
382 with the greater importance of forming N-terminal contacts when the C terminus is  
383 sequestered within the exit tunnel, possibly to compensate for loss of contacts at the C  
384 terminus.

385 In our kinetic modelling, we found that we obtained similar results with or without the  
386 assumption that folding and unfolding are fast relative to the escape rate, suggesting that this  
387 “pre-equilibrium” assumption is justified, at least for this protein. The reason for its validity  
388 in the case of I27 can be seen by comparing the folding and unfolding rates with the force-  
389 dependent escape rate of  $\sim 2.4 \times 10^{-3} \text{ s}^{-1}$  obtained at the highest forces of  $\sim 20 \text{ pN}$  (c.f., Fig.  
390 3C). Folding and unfolding rates at different linker lengths can be obtained by combining the  
391 linker-length dependence of the rates from simulation with the known folding/unfolding rates  
392 for isolated I27 from experiment (SI Appendix, Fig. S9). The presence of the ribosome  
393 increases the unfolding rate at shorter linker lengths so that it is faster than the maximum  
394 escape rate, while not slowing the folding rate sufficiently for it to drop below the escape  
395 rate. Note that the unfolding rate does drop below the maximum escape rate at larger linker  
396 lengths, but by that point the folded population is already almost 100%, so the pre-  
397 equilibrium assumption still gives accurate results. Although this assumption appears to be  
398 justified in the case of I27, it is probably not true in general, and it will be interesting to  
399 investigate for slower-folding proteins in future.

400 The arrest peptide experiments, in which a protein exerts a force due to folding in  
401 some ways resemble atomic force microscopy or optical tweezer experiments in which an  
402 external force is applied to the protein termini. It is important to note, however, that the  
403 nature and effect of the forces exerted on the folding protein by tethering to the ribosome are  
404 very different than is the case for pulling on both termini by an external force. For example,  
405 forces of the magnitude seen in this work (up to  $\sim 20 \text{ pN}$ ) tend to have very little effect on the  
406 unfolding rate when applied to the termini of I27, due to the similarity in extension of the  
407 folded and transition states (61); by contrast folding rates are dramatically slowed, even by  
408 very small forces, due to the large difference in extension of between unfolded and transition  
409 states (56). The forces arising from tethering to the ribosome are due to the folding of the  
410 protein itself rather than an external device. They arise from the constriction of available

411 configuration space, particularly for folded and partially folded states, as well as from any  
412 attractive interactions between the protein and the ribosome. Our simulations suggest that for  
413 I27, reducing the linker length speeds up unfolding and slows folding rates by similar factors.  
414 Thus, it seems that comparisons to the effects of forces exerted by AFM and optical tweezer  
415 experiments need to be performed with care.

416 We have previously shown that  $\alpha$ -helical proteins can fold co-translationally (2),  
417 perhaps unsurprising since helical structures are dominated by short-range interactions and  
418 helices can form within the ribosome tunnel itself (62, 63). Here, our equilibrium arrest-  
419 peptide assay and structural studies reveal that an all- $\beta$  protein, titin I27, is able to fold within  
420 the mouth of the ribosome exit tunnel, despite its folding being dominated by long-range  
421 interactions. Molecular simulations, accounting for the effect of the entropic restoring force  
422 on protein stability, reproduce the yield of protein from experiments remarkably well. These  
423 simulations reveal that I27 folds on the ribosome by the same pathway as when the protein  
424 folds away from the confines of the ribosome. We note that a similar conclusion has been  
425 reached by Guinn et al. (64) for another small protein, src SH3, using a completely different  
426 experimental approach which combines optical tweezer experiments and chemical denaturant  
427 to characterize the folding pathway of src SH3. Thus, the evidence so far suggests that single-  
428 domain proteins, both  $\alpha$ -helical and  $\beta$ -sheet, can fold close to the ribosome. On the other  
429 hand, while all- $\beta$  proteins appear to fold by a similar pathway with or without the ribosome  
430 present, there is evidence for  $\alpha$ -helical proteins forming partially structured cotranslational  
431 intermediates (11, 65) or folding by different pathways on the ribosome (2). This mechanistic  
432 difference may relate partly to the small contact order of helical proteins, allowing partially  
433 folded states to be more stable than for all- $\beta$  proteins. The situation for multidomain proteins  
434 is likely to be still more complicated, as some studies have already indicated (11, 23, 66, 67).

435

## 436 **Acknowledgements**

437 This work was supported by grants from the Knut and Alice Wallenberg Foundation, the  
438 Swedish Cancer Foundation, and the Swedish Research Council to GvH, by grants from the  
439 Deutsche Forschungsgemeinschaft (DFG) GRK 1721 and FOR1805 to RB, by a DFG  
440 fellowship through the Graduate School of Quantitative Biosciences Munich (QBM) to TS,  
441 and by the Wellcome Trust (WT095195) to JC; PT and RBB were supported by the  
442 Intramural Research Program of the National Institute of Diabetes and Digestive and Kidney



443 Diseases of the National Institutes of Health; JC is a Wellcome Trust Senior Research  
444 Fellow. The cryo-EM data were collected at the Swedish National Cryo-EM Facility funded  
445 by the Knut and Alice Wallenberg Foundation, the Family Erling Persson Foundation and the  
446 Science for Life Laboratory. This work utilized the computational resources of the NIH HPC  
447 Biowulf cluster. (<http://hpc.nih.gov>)

448

## 449 **Materials and Methods**

### 450 **Enzymes and chemicals**

451 All enzymes were obtained from Thermo Scientific. Oligonucleotides were purchased from  
452 Life Technologies. In-Fusion Cloning kits were obtained from Clontech and DNA  
453 purification kits were purchased from Qiagen. PUREfrex cell-free translation system was  
454 obtained from Eurogentec. [<sup>35</sup>S]-methionine was purchased from Perkin Elmer. Instant Blue  
455 protein stain was purchased from Expedeon.

456

### 457 **DNA manipulation**

458 Titin I27 constructs for *in vitro* translation were generated in pRSET A plasmid (Invitrogen)  
459 (previously modified to remove the sequence including the entire T7 gene 10 leader and EK  
460 recognition site up to, but not including, the *Bam*H I site and replaced with a sequence  
461 encoding residues L, V, P, R, G, S) carrying the *E. coli* SecM arrest peptide  
462 (FSTPVWISQAQGIRAGP) and a truncated *E. coli* *lepB* gene, under the control of a T7  
463 promoter. Increasing linker lengths were generated in pRSET A by PCR; linear pRSET A  
464 constructs (containing the SecM AP and truncated *lepB*, but lacking I27) were generated by  
465 PCR using primers which extended the linker from 23 aa to 63 aa (in steps of 2 aa) from the  
466 direction of the C to the N terminus. I27 flanked by GSGS linkers was amplified by PCR  
467 with overhanging homology to the plasmid containing the desired linker length. Cloning was  
468 performed using the In-Fusion system (Takara Bio USA, Inc.), according to the  
469 manufacturer's instructions. The final two C-terminal residues (EL) of the 89 aa Titin I27  
470 construct are not structured in the PDB file 1TIT, and are therefore included in the linker  
471 region. The amino acid sequence of the construct I27[L=63] is as follows (I27 in bold and  
472 SecM AP underlined):

473 MRGSHHHHHHGLVPRGSGSLIEVEKPLYGVEVFVGETAHFEIELSEPDVHGQWK

474 **LKGQPLAASPDCEIIEDGKKHILILHNCQLGMTGEVSFQAANTKSAANLKVKEL**

475 SGSGKFAYGIKDPIYQKTLVPGQQNATWIVPPGQYFMMGDWMSSFSTPVWISQAQG  
476 IRAGPGSSDKQEGEWPTGLRLSRIGGIH\*\*

477 The mutants I27[−A] (lacking β-strand A), I27[L58A] and I27[M67A] were generated for  
478 each linker length by site-directed mutagenesis. For the wild-type I27 and I27[−A] constructs  
479 with  $L = 27, 35, 37, 39, 47$  and  $57$  residues, site-directed mutagenesis was performed to  
480 generate constructs with the non-functional FSTPVWISQAQGIRAGA arrest peptide  
481 (mutated residue underlined) as full-length controls, and constructs with the crucial Pro, at  
482 the end of the AP, substituted with a stop codon as arrest controls. Site-directed mutagenesis  
483 was performed to generate W34E variants as non-folding (nf) controls at  $L = 27, 29, 31, 35,$   
484  $37, 39, 41, 43, 47, 49, 51$  and  $57$  for wild-type I27;  $L = 27, 35, 37, 39, 47$  and  $57$  residues for  
485 I27[−A];  $L = 27, 41, 45, 47, 49$  and  $53$  residues for I27[L58A];  $L = 27, 29, 37, 39, 41, 43, 45,$   
486  $47$  and  $51$  residues for I27[M67A]. All constructs were verified by DNA sequencing.

#### 487 **In vitro transcription and translation**

488 Transcription and translation were performed using the commercially available PUREfrex *in*  
489 *vitro* system (GeneFrontier Corporation), according to the manufacturer's protocol, using  $250$   
490  $\mu\text{g}$  plasmid DNA as template. Synthesis of [ $^{35}\text{S}$ ]-Met-labeled polypeptides was performed at  
491  $37\text{ }^\circ\text{C}$ ,  $500$  r.p.m. for exactly  $15$  min. The reaction was quenched by the addition of an equal  
492 volume of  $10\%$  ice-cold trichloroacetic acid (TCA). The samples were incubated on ice for  
493  $30$  min and centrifuged for  $5$  min at  $20,800 \times g$  and  $4\text{ }^\circ\text{C}$ . Pellets were dissolved in sample  
494 buffer and treated with RNase A ( $400\text{ }\mu\text{g ml}^{-1}$ ) for  $15$  min at  $37\text{ }^\circ\text{C}$  before the samples were  
495 resolved by SDS-PAGE and imaged on a Typhoon Trio or Typhoon 9000 phosphorimager  
496 (GE Healthcare). Bands were quantified using ImageJ to obtain an intensity cross section,  
497 (<http://rsb.info.nih.gov/ij/>), which was subsequently fit to a Gaussian distribution using in-  
498 house software (Kaleidagraph, Synergy Software). The fraction full-length protein,  $f_{FL}$ , was  
499 calculated as  $f_{FL} = I_{FL}/(I_{FL}+I_A)$ , where  $I_{FL}$  and  $I_A$  are the intensities of the bands representing  
500 the full-length and arrested forms of the protein. For wild-type I27 and six nf control samples  
501 ( $L = 27, 35, 37, 39, 47$  and  $57$  residues), *in vitro* transcription and translation were also  
502 performed at  $37\text{ }^\circ\text{C}$ ,  $500$  r.p.m. for exactly  $30$  min. The resultant force profile was slightly  
503 higher than that obtained at  $15$  min but has essentially the same shape (SI Fig. S5).

504 The reproducibility of force profile data has been discussed previously (2). For wild-type I27,  
505 data points  $L = 61$  and  $63$  residues are a single experiment;  $L = 33, 36, 38, 45, 53, 55$  and  $59$   
506 residues are an average of  $2$  experiments; all other values of  $L$  are an average of at least  $3$   
507 experiments. For I27[−A] strand,  $L = 23, 25, 33, 41, 43, 51, 53, 55$  residues are a single

508 experiment; all other values of  $L$  are an average of 2 experiments, except  $L = 35, 37$  and  $39$   
509 residues which are an average of at least 3 experiments. For I27[L58A], all data points are a  
510 single experiment except  $L = 27, 37, 41, 45, 47, 49$  and  $53$  residues, which are an average of  
511 2 experiments. For I27[M67A],  $L = 23, 25$  and  $51 - 63$  residues are a single experiment;  $L =$   
512  $29 - 35$  residues are an average of 2 experiments;  $L = 27, 37 - 47$  and  $51$  residues are an  
513 average of at least 3 experiments. For wild-type I27 samples incubated for 30 min, all data  
514 points are a single experiment except  $L = 27, 35, 37, 39, 47$  and  $57$  residues, which are an  
515 average of 2 experiments. For non-folding controls, all data points are a single experiment  
516 except for wild-type I27  $L = 29, 31, 39, 43$  and  $47$  residues which are an average of 2  
517 experiments.

### 518 **Cloning and purification of ribosome-nascent chain complexes**

519 The I27 construct at  $L = 35$ , which is at the peak of  $f_{FL}$  (Figure 1B), was studied by cryo-EM.  
520 The SecM AP in these constructs was substituted with the TnaC AP (34) for more stable  
521 arrest, and the constructs were engineered to maintain a linker length of 35 amino acid  
522 residues. An N-terminal 8X His tag was introduced to enable purification. The amino acid  
523 sequence of the construct used was (I27 in bold and TnaC AP underlined):

524 MDMGHHHHHHHDYDIPTTLEVLFGQPGT**LIEVEKPLYGVEVFVGETAHFEIELS**  
525 **EPDVHGQWKLKGQPLAASPDC**EHEDGKKHILILHNCQLGMTGEVSFQAANTKS  
526 AANLKVKELSGSGSGSGGPNILHISVTSKWFNIDNKIVDHRP\*\*

527 The construct was engineered into a pBAD expression vector, under the control of an  
528 arabinose-inducible promoter. The translation-initiation region was optimized as described in  
529 (68). The plasmid was transformed into the *E. coli* KC6  $\Delta smpB \Delta ssrA$  strain. 4 colonies were  
530 picked and tested for expression of the RNCs at  $37^\circ\text{C}$  in Lysogeny broth (LB).

531 Large-scale purification of RNCs was carried out based on a protocol described in (34).  
532 Briefly, a single colony of the KC6 cells found to express the RNCs was picked and cultured  
533 in LB at  $37^\circ\text{C}$  to an  $A_{600}$  of 0.5. Expression was induced with 0.3% arabinose and was carried  
534 out for 1 hour. Thereafter, the cells were chilled on ice, harvested by centrifugation, and  
535 resuspended in Buffer A at pH 7.5 (50 mM HEPES-KOH, 250 mM KOAc, 2 mM  
536 Tryptophan, 0.1% DDM, 0.1% Complete protease inhibitor). Cell lysis was carried out by  
537 passing the cell suspension thrice through the Emulsifex (Avestin) at 8000 psi at  $4^\circ\text{C}$ . The  
538 lysate was cleared of cell debris by centrifugation at  $30,000\times g$  for 30 min in the JA25-50  
539 rotor (Beckman Coulter). The supernatant obtained was loaded on a 750 mM sucrose cushion

540 (in Buffer A) and centrifuged at 45, 000 x g for 24 hours in a Ti70 rotor (Beckman Coulter)  
541 to obtain a crude ribosomal pellet, which was resuspended in 200  $\mu$ l Buffer A by shaking  
542 gently on ice.

543 RNCs from the crude suspension were purified via their His tags by affinity purification  
544 using Talon (Clontech) beads, which was pre-incubated with 10  $\mu$ g/ml tRNA to reduce  
545 unspecific binding of ribosomes. The suspension was incubated with the beads for 1 hour at  
546 4°C and subsequently washed with 20 column volumes of Buffer B at pH 7.5 (50 mM  
547 HEPES-KOH, 10 mM Mg(OAc)<sub>2</sub>, 0.1% Complete Protease Inhibitor, 250 mM sucrose, 2  
548 mM Tryptophan). RNCs were eluted by incubating the Talon beads with Buffer C at pH 7.5  
549 (50 mM HEPES, 150 mM KOAc, 10 mM Mg(OAc)<sub>2</sub>, 0.1% Complete protease inhibitor, 150  
550 mM imidazole, 250 mM sucrose) for 15 minutes and subsequently collecting the flow-  
551 through. Elution was carried out thrice and the eluents were concentrated by centrifugation at  
552 40,000 rpm for 2.5 hours in a TLA 100.3 rotor (Beckman Coulter). The pellet obtained at the  
553 end of this step was gently suspended in a minimal volume of Buffer D at pH 7 (20 mM  
554 HEPES-KOH, 50 mM KOAc, 5 mM Mg(OAc)<sub>2</sub>, 125 mM sucrose, 2 mM Trp, 0.03% DDM).

#### 555 **Cryo-EM sample preparation, data collection, processing and accession codes**

556 Approximately 4 A<sub>260</sub>/ml units of RNCs were loaded on Quantifoil R2/2 grids coated with  
557 carbon (3 nm thick) and vitrified using the Vitrobot Mark IV (FEI-Thermo) following the  
558 manufacturer's instructions. Cryo-EM data was collected at the Cryo-EM National Facility at  
559 the Science for Life Laboratory in Stockholm, Sweden.

560 Data was acquired on a 300 keV Titan Krios microscope (FEI) equipped with a K2 camera  
561 and a direct electron detector (both from Gatan). The camera was calibrated to achieve a  
562 pixel size of 1.06 Å at the specimen level. 30 frames were acquired with an electron dose  
563 0.926 e<sup>-</sup>/Å<sup>2</sup>/frame and a total dose of 27.767 e<sup>-</sup>/Å<sup>2</sup> and defocus values between -1 to -3  $\mu$ m.  
564 The first two frames were discarded and the rest were aligned using MotionCor2 (69). Raw  
565 images were cropped into squares by RELION 2.1 beta 1 (70). Power-spectra, defocus values  
566 and estimation of resolution were determined using the Gctf software (71) and all 2,613  
567 micrographs were manually inspected in real space, in which 2,613 were retained. 468,015  
568 particles were automatically picked by Gautomatch (<http://www.mrc-lmb.cam.ac.uk/kzhang/>)  
569 using the *E. coli* 70S ribosome as a template. Single particles were processed by RELION 2.1  
570 beta 1 (70). After 80 rounds of 2D classification, 384,039 particles were subjected to 3D  
571 refinement using the *E. coli* 70S ribosome as reference structure, followed by 160 rounds of

572 3D classification without masking and 25 rounds of tRNA-focused sorting. One major class  
573 containing 301,510 particles (64% of the total) was further refined including using a 50S  
574 mask, resulting in a final reconstruction with an average resolution of 3.2 Å (0.143 FSC). The  
575 local resolution was calculated by ResMap (72). Finally, the final map was obtained by local  
576 B-factoring followed by low-pass filtering to 4.5 Å by RELION 2.1 beta 1 (70) in order to  
577 best demonstrate the I27 domain.

578 For interpretation of the cryo-EM density, the cryo-EM structure model (PDB 4YU8) of *E.*  
579 *coli* TnaC-stalled ribosome was fitted into corresponding density using UCSF Chimera (73).  
580 The NMR model (PDB 1TIT) of I27 domain was fitted into the extra density of TnaC-stalled  
581 ribosome using UCSF Chimera (73). Since the I27 domain represents a flat ellipsoid, we used  
582 all four major and minor axes covering all possible orientations of the model fitting within  
583 the density to validate the orientation of the fitted I27 model. Briefly, the model with four  
584 different orientations were converted into densities (8 Å) by UCSF Chimera, and the cross-  
585 correlation coefficients of each model map and the isolated I27 density were calculated by  
586 RELION 2.1 beta 1 (70). Finally, uL24 β hairpin was remodeled as the tip of the hairpin is  
587 shifted due to the existence of I27 domain.

588 Figures showing electron densities and atomic models were generated using UCSF Chimera  
589 (73). Electron densities are shown at multiple contour levels in Figure 2 and SI Appendix, Fig.  
590 S1. The contour levels relative to the root-mean-square deviation (RMSD) were calculated  
591 from the final map values. Final map contains the volume for the entire RNC including the  
592 I27 domain.

593 Coordinates for the cryo-EM map of the ribosome with the I27 domain density have been  
594 deposited at the EMDataBank under accession code EMD-xxxx. Coordinates of fitted *E. coli*  
595 TnaC-stalled ribosome (PDB 4UY8; uL24 remodeled) and I27 domain (PDB 1TIT) models  
596 for interpreting the cryo-EM map have been deposited at the ProteinDataBank under  
597 accession code xxxx.

### 598 **Coarse-grained molecular simulations**

599 The 50S subunit of the *E. coli* ribosome (PDB 3OFR (36)) and the nascent chain are  
600 explicitly represented using one bead at the position of the α-carbon atom of each amino  
601 acid, and three beads (for P, C4', N3) per RNA residue (Figure 2A). The interactions within  
602 the protein were given by a standard structure-based model (37-39), which allowed it to fold

603 and unfold. Interactions between the protein and ribosome beads were purely repulsive (40)  
604 and given by the same form of potential as for the structure-based model(37-39),

$$605 \quad V_{ij} = \varepsilon_{ij} \left[ \frac{A}{r_{ij}^{12}} - \frac{B}{r_{ij}^{10}} + \frac{C}{r_{ij}^6} \right] \quad [2]$$

606 where  $r_{ij}$  is the distance between two beads  $i$  and  $j$ ,  $\varepsilon_{ij}$  (=0.001 kJ/mol) sets the strength of  
607 the repulsive interactions. The amino acid, phosphate, sugar and base are assigned radii  $\sigma_i =$   
608 4.5, 3.2, 5.1 and 4.5 Å respectively, and coefficients in Eq. 2 for interactions between protein  
609 and ribosome beads  $ij$  are obtained from the mixing rules  $A = \sqrt{\sigma_i^{12}\sigma_j^{12}}$ ,  $B = \frac{2}{\sigma_i^{-10} + \sigma_j^{-10}}$   
610 and  $C = \sqrt{\sigma_i^6\sigma_j^6}$ .

611 During the simulations, the positions of the ribosome atoms were fixed in space, as in  
612 previous studies (18). The linker between the AP and I27 was tethered by its C terminus to  
613 the last P atom of the A-site tRNA, but was otherwise free to fluctuate. The trajectory was  
614 propagated via Langevin dynamics, with a friction coefficient of 0.1 ps<sup>-1</sup> and a time step of  
615 10 fs, at 291 K in a version of the Gromacs 4.0.5 simulation code, modified to implement the  
616 potential given by Eq. 2 (74). All bonds (except the one used to measure force, below) were  
617 constrained to their equilibrium length using the LINCS algorithm (75). The attractive  
618 interactions between I27[M67A] and the hydrophobic residues (A, V, L, I, F, M, Y, W) on  
619 the surface of uL23 and uL29 are modelled as (76):

$$620 \quad V_{ij} = 4\varepsilon \left[ \left( \frac{\sigma}{r_{ij}} \right)^{12} - \left( \frac{\sigma}{r_{ij}} \right)^6 \right], \quad [3]$$

621 where  $r_{ij}$  is the distance between residues  $i$  and  $j$ ,  $\sigma$  is the range of the interaction and  $\varepsilon$   
622 represent the strength of the interaction.  $\sigma$  and  $\varepsilon$  are fixed at 6 Å and 5 kJ/mol respectively.  
623 Residues of I27[M67A] which are involved in the attractive interactions are defined as the  
624 ones whose heavy atoms are within 4.5 Å of any heavy atoms from residue 67 in the native  
625 state.

626 To calculate the pulling force exerted on the nascent chain by the folding of I27, the bond  
627 between the last and the second last amino acid of the SecM AP was modelled by a harmonic  
628 potential as a function the distance between these two atoms,  $x$  (Figure 3B):

$$629 \quad E = \frac{1}{2} k_s (x - x_0)^2 \quad [4]$$

630 where  $x_0$  is a reference distance. Here  $x_0$  is set to 3.8 Å, which is the approximate distance  
631 between adjacent C $\alpha$  atoms in protein structures and  $k_s$  is a spring constant, set to 3000  
632 kJ.mol<sup>-1</sup>.nm<sup>-2</sup>. The value of  $k_s$  was chosen so that the average displacement  $x - x_0$  remains  
633 below 1 Å for forces up to ~500 pN, which is much larger than the forces actually exerted by  
634 the folding protein. The pulling force on the nascent chain was measured by the extension of  
635 this bond as  $F = -k_s(x - x_0)$ .

636 I27 was covalently attached to unstructured linkers having the same sequences as used in the  
637 force-profile experiments (see Figure 2B). Linker amino acids are repulsive to both the  
638 ribosome and I27 beads, with interaction energy as described in Eq. 2.

639 The protein in its arrested state is subject to force  $F(t)$ , which will fluctuate, for example  
640 when the protein folds or unfolds. The rate of escape from arrest has been shown to be force-  
641 dependent (25); here we approximate the sensitivity to force using the phenomenological  
642 expression originally proposed by Bell (42)

$$643 \quad k(F) = k_0 e^{\beta F \Delta x^\ddagger}, \quad [5]$$

644 where  $k_0$  is a zero-force rupture rate,  $\Delta x^\ddagger$  is the distance from the free energy minimum to  
645 the transition state,  $\beta = 1/k_B T$  where  $k_B$  is Boltzmann's constant and  $T$  the absolute  
646 temperature. While there are functions to describe force-dependent rates with stronger  
647 theoretical basis, we use the Bell equation due to its simplicity and because its parameters  
648 have previously been estimated from optical tweezer experiments for the SecM AP. In all  
649 cases, we set  $k_0$  (Eq. 5) to  $3.4 \times 10^{-4} \text{ s}^{-1}$  and  $\Delta x^\ddagger$  to 3.2 Å, based on the values determined by  
650 Goldman *et al.* (they estimated  $k_0$  and  $\Delta x^\ddagger$  to be in the range of  $0.5 \times 10^{-4}$  to  $20 \times 10^{-4} \text{ s}^{-1}$  and  
651 1-8 Å, respectively) (25).

652 We assume the probability of remaining on the ribosome  $S(t) = 1 - f_{FL}(t)$  assuming that  
653  $\dot{S} = -k(F(t))$ , hence

$$654 \quad S(t) = \exp\left[-\int_0^t k(F(t))dt\right]. \quad [6]$$

655 The escape of I27 from the ribosome can be described using kinetic model shown in SI  
656 Appendix, Fig. S8 which explicitly takes into account the linker length-dependent  
657 folding/unfolding rates of the I27 nascent chain,  $k_u(L)$  and  $k_f(L)$ , on the ribosome, and the  
658 force-dependent rate of escape from ribosome:  $k(F_u(L))$  and  $k(F_f(L))$ . To estimate  $k_f$  at  
659 different linker lengths, we first carried out unbiased MD simulations to estimate the mean

660 first passage time for folding  $t_F^{\text{mfpt}}$ , from which the folding rate can be calculated as  
661  $k_f^{\text{wt}} = 1/t_F^{\text{mfpt}}$ . Similarly, the unfolding rate can be calculated from unfolding simulations as  
662  $k_u^{\text{wt}} = 1/t_u^{\text{mfpt}}$ . Since the rates in coarse-grained simulations are naturally much faster than  
663 in experiment, we globally scale the unfolding rates  $k_u^{\text{wt}}(L=21, 23 \dots 61)$  at different linker  
664 lengths so that  $k_u^{\text{wt}}$  at very long linker lengths ( $L=61$ ) is equal to the unfolding rate of isolated  
665 I27 ( $4.9 \times 10^{-4} \text{ s}^{-1}$ ). Similarly,  $k_f^{\text{wt}}(L=21, 23 \dots 61)$  is scaled so that the  $k_f^{\text{wt}}$  of I27 RNC[L=61]  
666 is equal to the unfolding rate of isolated I27 (SI Appendix, Fig. S9). For consistency with our  
667 pre-equilibrium solution, we further scale  $k_f^{\text{wt}}$  to match the stability of I27 RNC[L=61] in  
668 our simulation model, yielding  $k_u^{\text{wt}}$  and  $k_f^{\text{wt}}$  at  $L=61$  of  $4.9 \times 10^{-4} \text{ s}^{-1}$  and  $0.14 \text{ s}^{-1}$  (SI Appendix,  
669 Fig. S9) respectively. The same scaling method has been applied to the folding and  
670 unfolding rates of all mutants (I27[L58A], I27[M67A] and I27 [-A]) so that the  
671 folding/unfolding rates of the mutant RNC are consistent with the relative experimental  
672 values measured for the isolated mutants(26).

673

674 With the rates obtained from above, the time dependent survival probability  $S(t)$  is estimated  
675 by the kinetic Monte Carlo method (the Bortz-Kalos-Lebowitz algorithm (77)). The system is  
676 initialized at the state when the unfolded nascent chain just emerges from the ribosome tunnel  
677 (UA, SI Appendix, Fig. S8) at time  $t=0$ . At each Monte Carlo step, a uniform random number  
678  $\delta$  between 0 and 1 is chosen, and a transition from the current state  $s$  to state  $j$  will occur  
679 for the state  $j$  which satisfies  $\sum_{i=1}^{j-1} k_{si} < \delta \sum_{i=1}^N k_{si} < \sum_{i=1}^j k_{si}$ , where  $k_{si}$  represents the  
680 transition rate from state  $s$  to state  $i$ . The time is updated by  $t = t + \Delta t$ , where  $\Delta t =$   
681  $-(\delta')/\sum_{i=1}^N k_{si}$ .  $\delta'$  is a new number randomly chosen between 0 to 1.

682

683 The solution to the kinetic model can be simplified if we further assume that the escape from  
684 the ribosome is slow relative to the folding and unfolding of the protein. In this situation, we  
685 can approximate  $S(t)$  in terms of the mean forces experienced when the protein is unfolded,  
686  $F_u$ , or folded,  $F_f$ , and the unfolded and folded populations of  $P_u$  and  $P_f$  respectively,

$$687 \quad S(t) \approx \exp[-t[P_u k(F_u) + P_f k(F_f)]]. \quad [7]$$

688 The equilibrium properties of the system for each linker length were obtained from umbrella  
689 sampling using the fraction of native contacts  $Q$  as the reaction coordinate, allowing  $P_u, P_f$



690 and  $F_u$ ,  $F_f$  to be determined (Figure 3C). The details of the definition of  $Q$  have been  
691 previously described (48); in short,  $Q$  is defined as

$$692 \quad Q = \frac{1}{N} \sum_{(i,j)} \frac{1}{1 + e^{(r_{ij} - \lambda r_{ij}^0)}}, \quad [8]$$

693 where the sum runs over the  $N$  pairs of native contacts  $(i, j)$ ,  $r_{ij}$  is the distance between  $i$  and  
694  $j$  in configuration,  $r_{ij}^0$  is the distance between  $i$  and  $j$  in the native state,  $\lambda = 1.2$  which  
695 accounts for fluctuations when the contact is formed. A boundary of  $Q = 0.5$  is used to  
696 separate folded from unfolded states.

697 In order to characterize folding mechanism, we used transition paths from folding  
698 simulations for the  $L = 51$  case at 291 K. 50 independent simulations, each started from fully  
699 extended configurations, were carried out for 4 microseconds. The folding barriers for the  $L =$   
700 31 and 35 cases are very high at the same temperature, therefore the transition paths are  
701 obtained from unfolding simulations instead. Starting from native-like folded  
702 configurations, 50 unfolding simulations were carried out, with each trajectory being 4  
703 microseconds long. Transition paths were defined as those portions of the simulation  
704 trajectory from the last time it samples the configuration with  $Q < 0.3$  till the first time it  
705 samples a configuration with  $Q > 0.7$  (in the folding direction; opposite for unfolding).  $\phi$ -  
706 Values were computed from the transition paths using the approximation:

$$707 \quad \phi(i) \approx \sum_{j:(i,j) \in \text{native}} p(q_{ij}|TP)$$

708 In which  $p(q_{ij}|TP)$  is the probability that the native contact  $q_{ij}$  between residues  $i$  and  $j$  is  
709 formed on transition paths as defined above. We also characterized the importance of  
710 individual contacts in determining the folding mechanism using  $p(TP|q_{ij})_{\text{nn}}$ , defined in Eq.  
711 1 of the main text, i.e. the probability of being on a transition path given that contact  $q_{ij}$  is  
712 formed and the protein is not yet folded. Having already calculated  $p(q_{ij}|TP)$  above,  
713 evaluating  $p(TP|q_{ij})_{\text{nn}}$  required  $p(q_{ij})_{\text{nn}}$ , the probability of a contact being formed in all  
714 non-native fragments of the trajectory, and  $p(TP)$ , the fraction of time spent on transition  
715 paths. For  $L=51$ , we obtained  $p(q_{ij})_{\text{nn}}$  directly from unbiased folding simulations, using the  
716 portion of the trajectory up to the first folding event (i.e. the first time  $Q > 0.7$ ). For  $L=31$  or

717 35, where the protein is still relatively unstable, we determined it from unfolding simulations  
718 by computing  $p(q_{ij})$  separately for the unfolded and transition-path portions of the trajectory  
719 and combining them weighted by  $p(\text{TP})_{\text{nn}}$ . We determined  $p(\text{TP})_{\text{nn}}$  via folding ( $L = 51$   
720 case) and unfolding ( $L = 31$  and  $L = 35$  cases) simulations (described above). For the  $L = 51$   
721 case,  $p(\text{TP})_{\text{nn}} = \frac{2t_{\text{TP}}}{2t_{\text{TP}} + t_{\text{F}}^{\text{mfpt}}}$ , where  $t_{\text{TP}}$  is the mean transition path time and  $t_{\text{F}}^{\text{mfpt}}$  is the mean  
722 first passage time for folding obtained from the maximum likelihood estimator  $t_{\text{F}}^{\text{mfpt}} =$   
723  $[N_{\text{fold}}t_{\text{fold}} + (N - N_{\text{fold}})t_{\text{sim}}]/N_{\text{fold}}$ , where  $N$  is the total number of trajectories ( $N = 50$ ),  $N_{\text{fold}}$   
724 is the number of trajectories folding within 4  $\mu\text{s}$ ,  $t_{\text{fold}}$  is the average folding time (of the  
725 trajectories which fold), and  $t_{\text{sim}}$  is the length of the simulations (4  $\mu\text{s}$ ). For the  $L = 31$  and  $L$   
726 = 35 cases, it is less efficient to obtain the folding time  $t_{\text{F}}^{\text{mfpt}}$  directly, therefore we estimate it  
727 based on the mean first passage time for unfolding,  $t_{\text{U}}^{\text{mfpt}}$ , from unfolding simulations.  
728  $p(\text{TP})_{\text{nn}} = \frac{2t_{\text{TP}}}{2t_{\text{TP}} + \frac{p_{\text{U}}}{p_{\text{F}}}t_{\text{U}}^{\text{mfpt}}}$ , where  $p_{\text{U}}$  and  $p_{\text{F}}$  are the equilibrium populations of the unfolded  
729 and folded respectively determined from umbrella sampling.

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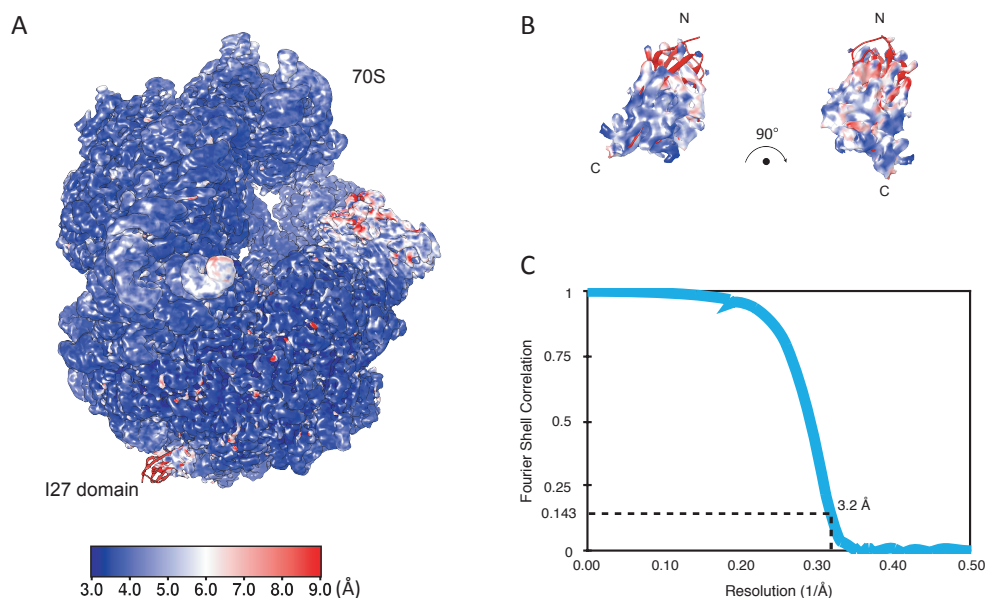
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944 **Supplementary Figures**

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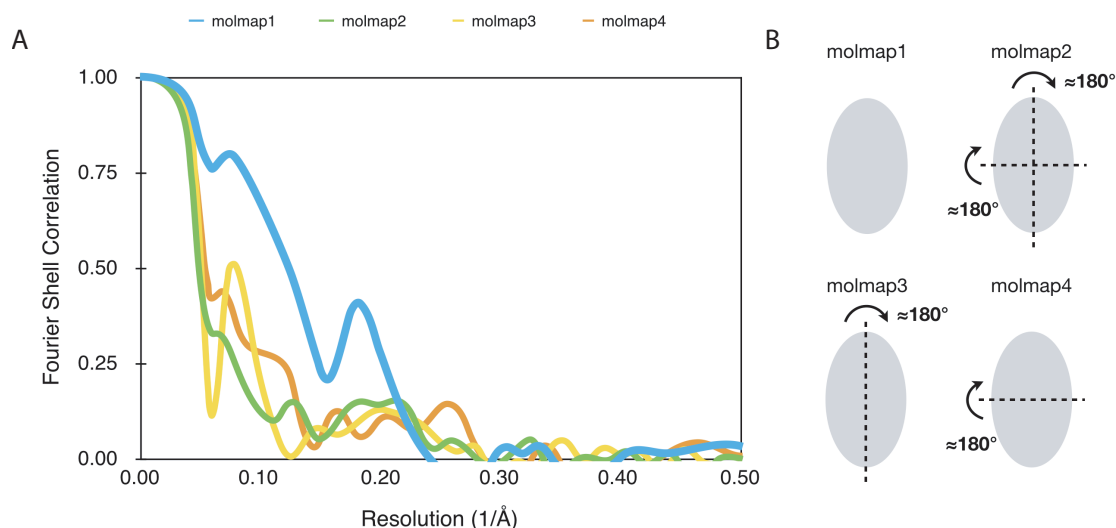
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949 **Figure S1.** Resolution of the ribosome-nascent chain complex (RNC). (A) Calculation of the  
950 local resolution using Resmap (Kucukelbir, A. et al. Nat Methods 11, 63-65, 2014). The RNC  
951 density is displayed at 1.7 RMSD. (B) local resolution of the I27 domain. The I27 domain  
952 density is displayed at 2 RMSD. N and C termini are indicated. (C) Fourier-shell correlation  
953 (FSC) curve of the refined final map of the RNC, indicating the average resolution of 3.2 Å  
954 (at 0.143).

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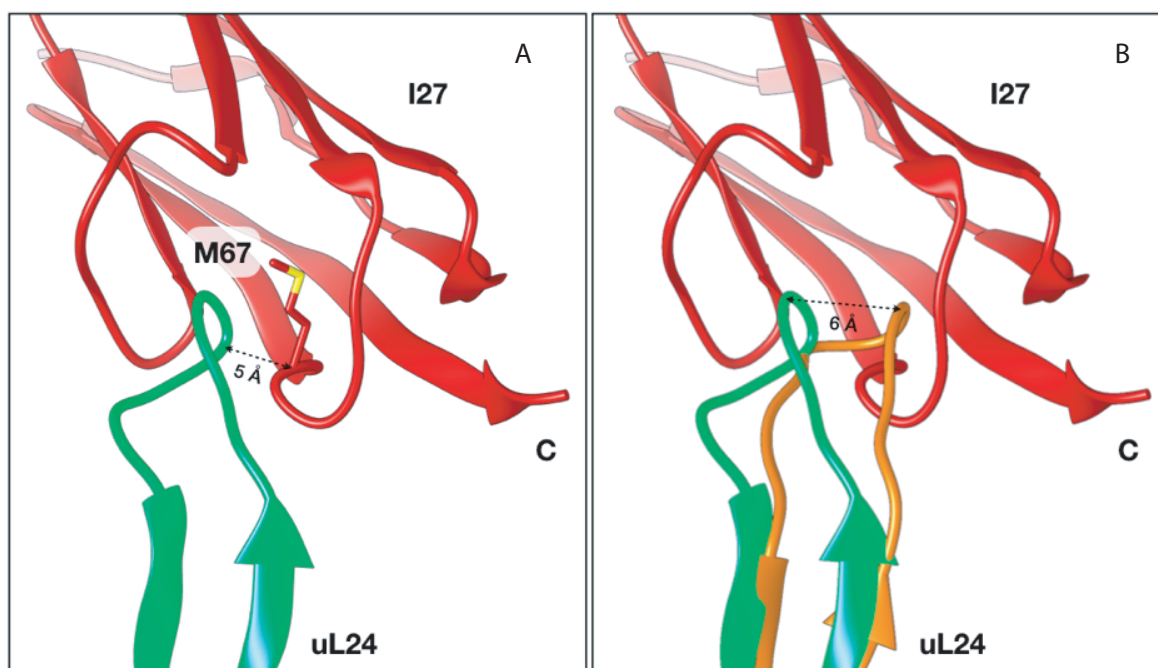
956

957 **Figure S2.** Validation of model orientation for I27 domain. To validate the orientation of the  
958 I27 domain model (PDB 1TIT) to its corresponding density, four possible orientations were  
959 tested. (A) The Fourier-shell correlations between the isolated I27 density and the map  
960 generated from the model of the final orientation (molmap1, blue) and the models fitted with  
961 the other three possibilities (molmap2, green; molmap3, yellow; molmap4, orange) were  
962 plotted. In the frequency range 0 to 0.2 (1000 to 5 Å) the correlation of molmap1 is  
963 significantly higher compared to all other orientation molmaps. (B) The illustration showing  
964 the relationship among the four model orientations. Since the density represents a flat  
965 ellipsoid, we used all four major and minor axes covering all possible orientations of the  
966 model fitting within the density.



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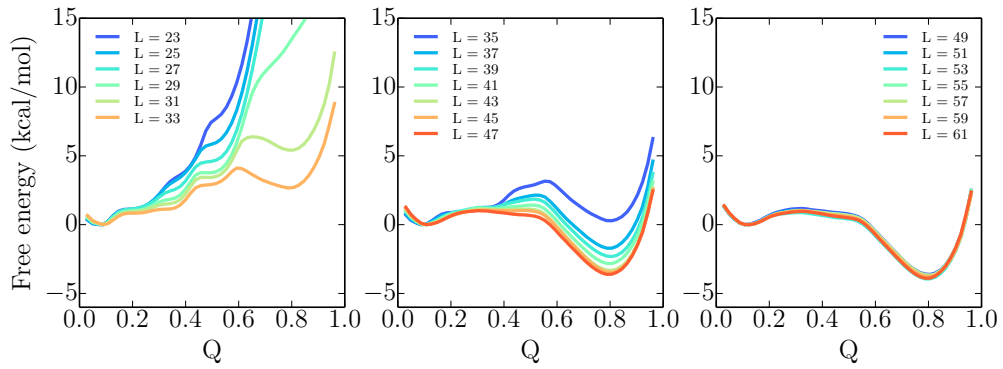


969

970 **Figure S3.** The I27 domain and a  $\beta$  hairpin in ribosomal protein uL24 close to the ribosomal  
971 exit tunnel. (A) Residue M67 in the I27 domain is located in close proximity to a  $\beta$  hairpin  
972 loop in uL24 in the cryo-EM structure of I27-TnaC[L=35] RNCs. (B) The uL24  $\beta$  hairpin in  
973 the I27-RNC (light green; re-modeled based on PDB 5NWX) is  $\sim 6$  Å shifted (distance  
974 measured via the backbone of Pro50) compared to its location in the VemP-RNC (orange;  
975 PDB 5NWX) and the TnaC-RNC (PDB 4UY8, not shown). C represents the C terminus of  
976 the I27 domain.

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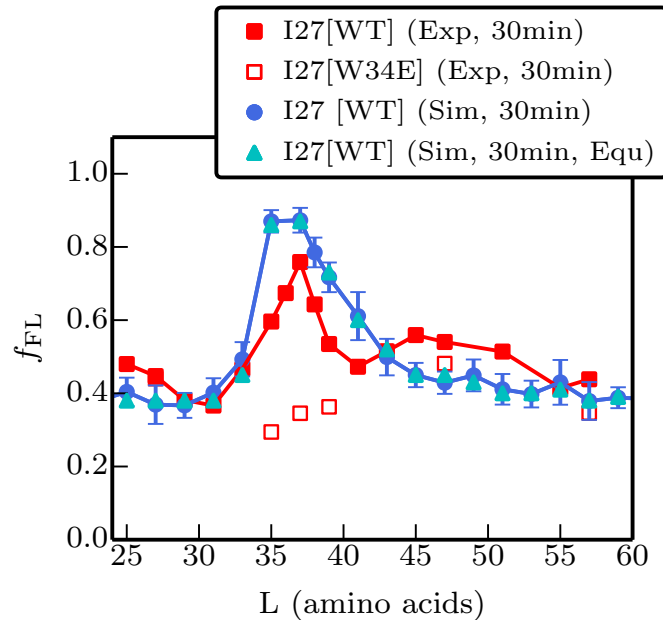
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980 **Figure S4.** Simulation free energy  $F(Q)$  projected on the fraction of native contacts,  $Q$ , for  
981 I27 folding with different linker lengths (as indicated in legend) at 291K.

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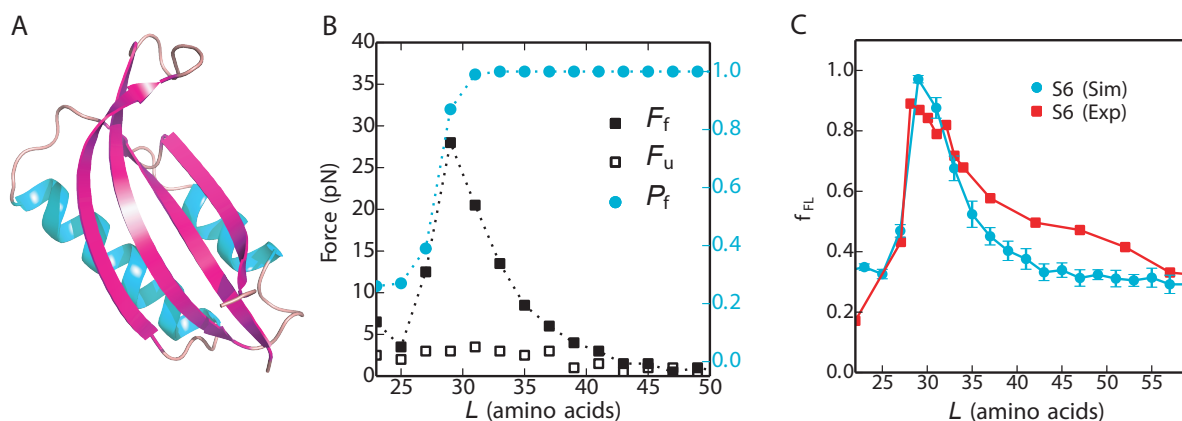


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984 **Figure S5.** Experimental (red) and simulated profiles of fraction full length protein,  $f_{FL}$ ,  
985 obtained with a 30 min incubation. Note the higher background values compared to main text  
986 Figures 1 and 3D. Force profiles calculated from simulations using full kinetic scheme and  
987 pre-equilibrium model are shown in blue circles and cyan triangles respectively. The RMSD  
988 of the  $f_{FL}$  between experiment and simulation is 0.12.

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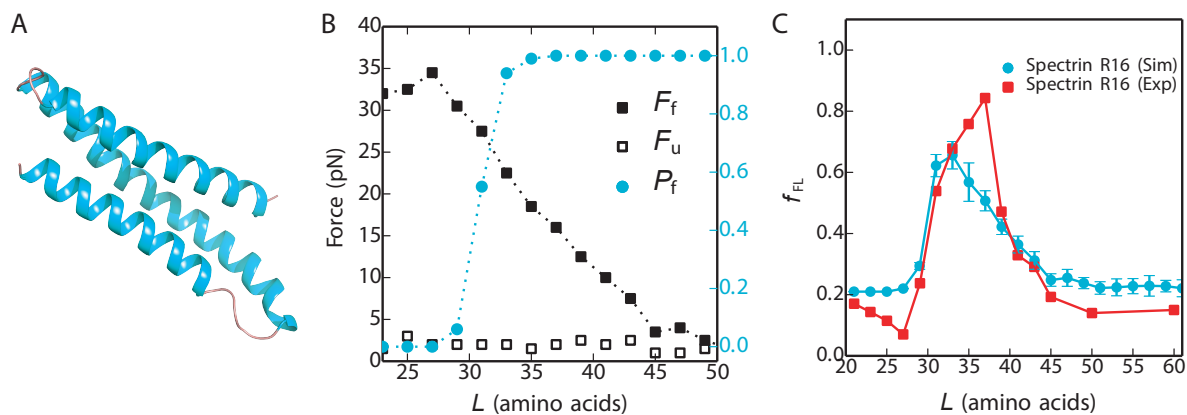


991

992 **Figure S6.** (A) Native structure of protein S6 (pdb code: 2KJV (78)) (B) Average forces  
 993 exerted on the AP by the unfolded state ( $F_f$ , filled black symbols) and folded state ( $F_u$ ,  
 994 empty black symbols) of S6 at different linker lengths  $L$ . The average fraction folded S6 for  
 995 different  $L$ ,  $P_f$ , is shown in cyan on the right axis. (C) Experimental (red) and simulated  
 996 (cyan) force profiles for cotranslational folding of S6 based on pre-equilibrium kinetic  
 997 solution.

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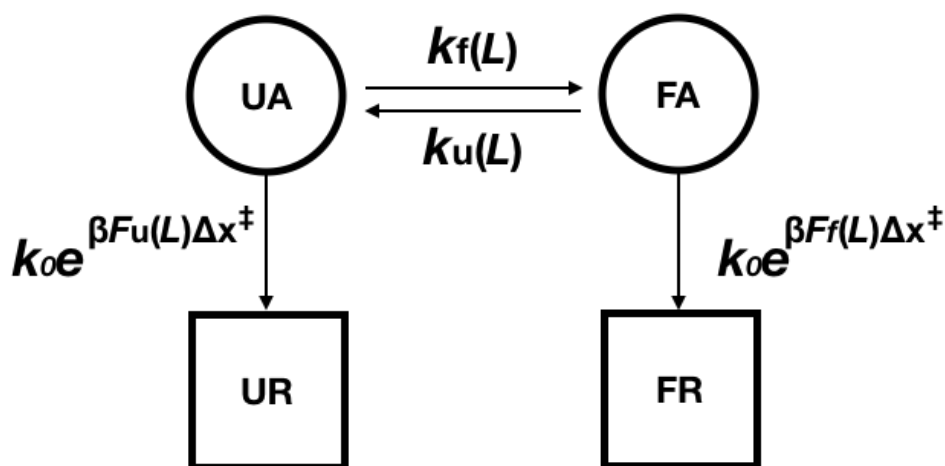
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1001 **Figure S7.** (A) Native structure of Spectrin R16 domain (PDB 1AJ3 (79)). (B) Average  
 1002 forces exerted on the AP by the unfolded state ( $F_f$ , empty black symbols) and folded state ( $F_u$ ,  
 1003 filled black symbols) of R16 at different linker lengths  $L$ . The average fraction folded R16  
 1004 for different  $L$ ,  $P_f$ , is shown in cyan on the right axis. (C) Experimental (red) and simulated  
 1005 (cyan) force profiles for cotranslational folding of R16 based on the pre-equilibrium kinetic  
 1006 solution.

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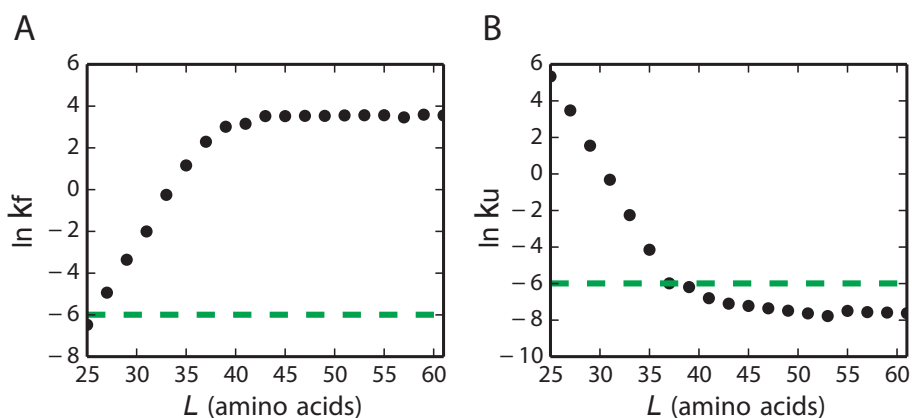


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1009 **Figure S8.** Schematic for the full kinetic model which describes the escape pathway of I27  
 1010 from the ribosome.  $k_f$  and  $k_u$  are the linker length-dependent folding and unfolding rates  
 1011 respectively. UA: I27 is unfolded and arrested by ribosome. FA: I27 is folded and arrested by  
 1012 ribosome. UR: I27 is unfolded and released from ribosome. FR: I27 is folded and released  
 1013 from ribosome.

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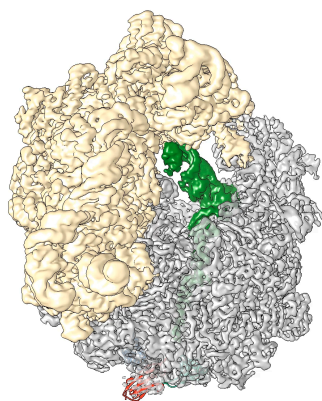


1016

1017 **Figure S9.** Dependence of folding rate  $k_f$  (left) and unfolding rate  $k_u$  (right) on the length of  
 1018 the linker between the AP and I27. Rates determined directly from simulations have been  
 1019 scaled so that  $k_f$  and  $k_u$  at large linker lengths are equal to the experimental values determined  
 1020 for the isolated protein. The green dashed line indicate the force-dependent escape rate of  $\sim$   
 1021  $2.4 \times 10^{-3} \text{ s}^{-1}$  obtained at the force of 20 pN.

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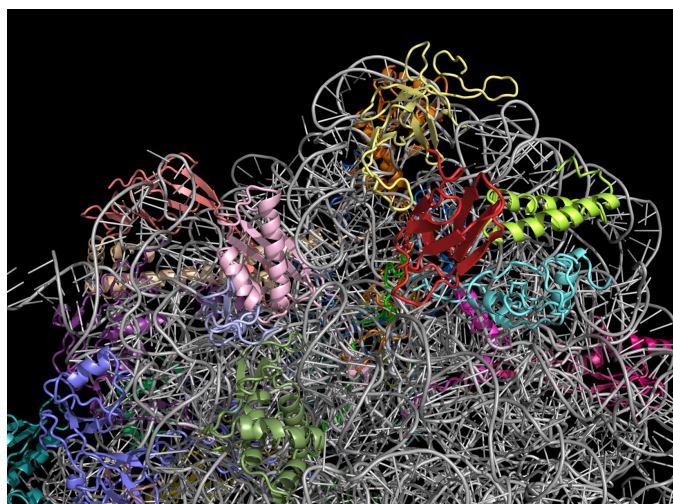
1024

1025 **Supporting Video S1.** Cryo-EM density of ribosome and I27 (one static frame of the video).  
1026 Video showing cryo-EM map for I27[L=35] RNCs. 30S in yellow, 50S and I27 domain in  
1027 grey, tRNA and nascent chain in green, the model (PDB 1TIT) of I27 domain in red.

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1032 **Supporting Video S2.** MD folding simulation (one static frame of the video). Video showing  
1033 an unbiased 1.8  $\mu$ sec fragment of an MD trajectory of I27 folding and unfolding at linker  
1034 length  $L=35$ . Ribosomal 23s rRNA is shown in white cartoon mode, ribosomal proteins  
1035 uL24, uL29, uL23 are shown in yellow, lime and cyan respectively. I27 and linker are in red  
1036 and green cartoon mode respectively.

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