1	Dispersal increases the resilience of tropical savanna and					
2	forest distributions					
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14						

15 Abstract

Global change may induce changes in savanna and forest distributions, but the dynamics of these 16 17 changes remain unclear. Classical biome theory suggests that climate is predictive of biome distributions, such that shifts will be continuous and reversible. This view, however, cannot 18 19 explain a widely observed mismatch between climate and tree cover, which some argue results from fire-vegetation feedbacks maintaining savanna and forest as bistable states, such that, 20 21 instead, shifts will be discontinuous and irreversible. This bistable model, however, cannot reproduce the spatial aggregation of biomes. Here, we suggest that both models are limited in 22 that they ignore spatial processes, such as dispersal. We examine the contributions of dispersal to 23 24 determining savanna and forest distributions using a reaction-diffusion model, comparing results 25 qualitatively to empirical savanna and forest distributions in Africa. The diffusion model induces spatially aggregated distributions, separated by a stable savanna-forest boundary. The 26 27 equilibrium position of that boundary depends not only on precipitation but also on the curvature 28 of precipitation contours with some history dependence (although less than in the bistable 29 model). This model predicts different dynamics in response to global change: the boundary 30 continuously tracks climate, recovering following disturbances, unless remnant biome patches 31 are too small.

32 Introduction

33 Climate and land use change are expected to result in large-scale shifts in global 34 vegetation patterns (Aleman et al. 2016; Loarie et al. 2009; Malcolm et al. 2002; Salazar et al. 35 2007), leading to loss of biodiversity and ecosystem services that are vital for human livelihoods (Daily 1997). However, biosphere responses to changing climate and land use are uncertain. This 36 uncertainty stems from uncertainty in what determines global biome patterns; current biome 37 distribution models are unable to explain even simple empirical features of today's vegetation 38 patterns. Predicting changes in biome distributions with respect to global change thus requires a 39 40 better understanding of the drivers and the mechanisms by which these drivers shape global biome patterns. 41

Conceptually, the classical theory suggests that climate is the fundamental determinant of 42 vegetation pattern and that there is a one-to-one match between climate and biome (Holdridge 43 1947; von Humboldt and Bonpland 1807; Whittaker 1970), such that biomes continuously track 44 changes in climate through space. Thus, under the classical view, biome shifts are continuous 45 46 and reversible, and as such, relatively predictable. An alternative viewpoint, supported by both 47 field (Dantas et al. 2016) and remote-sensing approaches (Hirota et al. 2011; Staver et al. 2011a) suggests that a single climate can support multiple vegetation types, which are differentiated 48 49 instead by other ecological processes including chronic fires (Bond et al. 2005; Staver et al. 50 2011a; Staver et al. 2011b). Savanna and forest may be a classic example; fire experiments have 51 repeatedly shown that frequent fires can maintain savanna in regions where a closed canopy forest is climatically possible (Swaine et al. 1992; Trapnell 1959; Veenendaal et al. 2018). 52 53 Simple theoretical models that incorporate both climate and fire suggest that savanna and forest may be bistable, with substantial hysteresis in biome patterns (Beckage et al. 2009; Staver et al. 54 55 2011b; Staver and Levin 2012). In this second scenario, unlike the classical theory, vegetation 56 responses to changing climate and land use may be large and irreversible, and therefore difficult 57 to foresee.

The bistable theory, however, has its limitations as well. Although the mean-field bistable models can mechanistically explain the overlap in the climatic ranges over which savanna and forest occur (Beckage et al. 2009; Staver and Levin 2012), they cannot be used to describe spatial patterning of biomes. Most notably, they miss obvious spatial features of savanna and forest distributions: savannas are found near other savannas and forests near other forests, with a

distinct biogeographic boundary separating the two biomes (Aleman and Staver 2018). This
spatial aggregation is not an obvious outcome of mean-field models, unless they invoke an
additional assumption that the historical or paleo-distributions of savanna and forest are spatially
structured by some extrinsic process (see Aleman and Staver 2018) (*e.g.*, paleoclimate).

An alternative explanation could be that some spatial process at the savanna-forest 67 68 ecotone may spatially aggregate savanna and forest. For instance, studies show that seed dispersal from forest patches can allow recovery of nearby derived savannas (Holl et al. 2000; 69 Puyravaud et al. 1994) by clumping fragmented forest patches into bigger forest aggregates. 70 Thus, dispersal could potentially explain the observed spatial aggregation of savanna and forest. 71 72 However, only a handful of theoretical studies (Favier et al. 2004; van de Leemput et al. 2015; 73 Wuyts et al. 2018) have explicitly considered the role of dispersal in determining biome patterns at relevant spatial scales. 74

75 The problem should be tractable, however, as dispersal is among the best-studied spatial ecological processes. Traditionally, dispersal in ecology is has been studied via 76 77 reaction-diffusion equations. These equations offer a simple and analytically tractable way to 78 incorporate dispersal in modeling dynamics of populations at large spatial scales (Levin 1992; 79 Skellam 1951). Theoretical work on one-dimensional reaction-diffusion models shows that 80 coupling diffusion with a bistable model (van de Leemput et al. 2015; Wuyts et al. 2018) can yield spatially aggregated biome distributions, separated by a stable savanna-forest boundary. 81 82 Moreover, the one-dimensional reaction-diffusion model behaves, dynamically, like the classical biome theory: savanna and forest distributions continuously track changes in climate and 83 84 recovers to equilibrium following perturbations. Unfortunately, this one-dimensional diffusion model (van de Leemput et al. 2015; Wuyts et al. 2018) also reverts to the main drawback of the 85 86 classical biome theory: it, too, fails to reproduce the widely observed overlap in the climatic 87 ranges over which savanna and forest biomes occur (Hirota et al. 2011; Staver et al. 2011a). One obvious avenue for exploration is that these models that couple diffusion to models 88 for savanna-forest dynamics treat the landscape as one-dimensional (1D) (van de Leemput et al. 89 90 2015; Wuyts et al. 2018), whereas, in reality, savanna and forest dynamics play out on 91 two-dimensional (2D) landscapes. Going from 1D to 2D often gives rise to new dynamical 92 features, such as motion by mean-curvature (Allen and Cahn 1979; Chen 1992; Evans et al. 93 1992; Gandhi et al. 1999; Keener 1986; Merriman et al. 1992; Tyson and Keener 1988), that

94 could fundamentally change the dynamics of boundaries and thus their equilibrium distributions. 95 Here, we address this directly by considering a model that couples the bistable mean-field 96 vegetation structure with a diffusion process in two dimensions. In particular, we ask (i) whether 97 seed dispersal, approximated as a two-dimensional diffusion process could contribute to spatial aggregation of savanna and forest biomes at continental scales; (ii) if yes, what then determines 98 99 the equilibrium position of the savanna-forest boundary, and (*iii*) how this impacts the resilience 100 of savanna and forest biomes to perturbations and global change. Finally, (iv) we empirically test some of the key analytical predictions of our 2D reaction-diffusion model using remotely sensed 101 102 biome (Hansen et al. 2013) and climate patterns (Huffman and Bolvin 2013) in Sub-Saharan 103 Africa.

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105 Model Description

Here, we present a reaction-diffusion model of savanna and forest biomes that consists of two parts: the reaction term that determines how fire interacts with vegetation and climate, and the diffusion term that represents seed dispersal. Here, we first describe the reaction term, and then, following, the diffusion term.

In savanna and forest ecosystems, fire exerts strong control over tree cover (Bond et al. 110 111 2003; Bond et al. 2005) via feedbacks with vegetation. In a low tree-cover landscape, fire 112 spreads readily in the landscape (Archibald et al. 2009; Staver and Levin 2012) because sparse 113 tree cover promotes the formation of a continuous grass layer (Archibald et al. 2009; 114 Hennenberg et al. 2006; Puevo et al. 2010), in turn limiting the density of trees in the landscape 115 (Higgins et al. 2000; Prior et al. 2010; Staver et al. 2009). Meanwhile, dense tree cover shades out grasses, resulting in a discontinuous grass layer that impedes fire spread (Archibald et al. 116 117 2009; Hennenberg et al. 2006; Pueyo et al. 2010). Here, we capture these two alternative 118 feedbacks using a step fire-mortality function ϕ (see also Staver et al. 2011b; Staver and Levin 119 2012), that takes a high value (combining high fire frequency with its potential effects on forest 120 trees) at low tree-cover and a low value (representing a background mortality rate in the absence 121 of fires) at high tree-cover. Finally, we assume that in the absence of fire, the tree cover 122 accumulates logistically to some carrying capacity, with a per-capita growth rate that we normalize, without loss of generality, to precipitation P, reflecting an increase in tree growth 123 124 rates [via increased primary productivity (Lieth 1975)] with increasing precipitation.

With these simplifying assumptions, the mean-field or the reaction term can be

126 mathematically expressed as

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$$f(T,P) = T\left[P\left(1 - \frac{T}{T_k}\right) - \phi(T)\right],\tag{1}$$

where T, T_k , and P represent tree cover, local carrying capacity, and precipitation, 128 respectively. This reaction term has two important ecological features. First, in the absence of 129 fire (*i.e.*, at $\phi = 0$), the system equilibrates to a high tree-cover state. This feature of the reaction 130 term f(T, P) is consistent with long-term (50-60 years) fire experiments that show that active 131 132 fire suppression in mesic savannas can result in a closed canopy forest (Bond et al. 2005; Swaine 133 et al. 1992; Trapnell 1959). Second, in the presence of fire, the mortality rate of trees has a threshold response to the tree cover itself (fire-vegetation feedbacks), consistent with previous 134 135 empirical work (Archibald et al. 2009), because of which the equilibrium tree cover becomes 136 bimodal in some intermediate range of rainfall. Theoretically, this implies that inclusion of fire results in a potential decrease in tree cover below the system's carrying capacity by allowing for 137 multiple stable states, corresponding to savanna (T_s^*) and forest (T_F^*) , for some parts of parameter 138 space. This is also evident from the bifurcation diagram in figure 1, which shows that both 139 140 savanna and forest are stable states in the intermediate precipitation region, bounded by the two critical precipitation values (P_{SF} and P_{FS}); meanwhile, outside this rainfall region, the system 141 142 has only one stable solution corresponding to savanna and forest in low and high precipitation 143 regions, respectively. An analogous mean-field system has been thoroughly elaborated in a 144 number of papers (Staver et al. 2011b; Staver and Levin 2012; Touboul et al. 2018).

145 Next, we incorporate seed dispersal in our model following Skellam (1951). In his paper, 146 Skellam (1951) assumed that a plant disperses its propagules like a random walk process, with 147 the probability of finding a propagule highest near the parent stem and falling off with increasing 148 distance (Levin et al. 2003; Okubo and Levin 2013). This movement of plant populations, 149 although random at an organism level can be statistically approximated to a continuous Diffusion (or Laplacian) operator ∇^2 when scaled to the landscape level (Okubo and Levin 2013; Skellam 150 1951). Mathematically, ∇^2 is defined as $\left[\frac{\partial^2}{\partial x^2}\right]$ in 1D and $\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right]$ in 2D. Here, for 151 simplicity, we assume that seed dispersal is isotropic, and ignore advective effects, for example 152 153 due to wind.

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Finally, combining the reaction (mean-field) and the diffusion (spatial) components of the

155 model yields a reaction-diffusion equation:

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 $\frac{\partial T}{\partial t} = f(T, P) + D \,\nabla^2 T,$ (2)

157 where D is the diffusion coefficient that captures the rate of seed spread.

Although the reaction-diffusion approach to model plant dynamics has proven to be quite 158 159 useful because of its analytical tractability and mathematical simplicity, this approach has some 160 inherent drawbacks, such as approximating discrete variables (such as habitats) as continuous 161 (Keitt et al. 2001) and failing to consider the effects of long-range seed dispersal (Kot et al. 1996; see Appendix C), both of which have been previously shown to yield qualitatively 162 163 different results. Moreover, we also ignore fire spread as spatially explicit process [see Schertzer et al. (2015) for a more realistic way of modeling fire spread within savannas and (Cochrane 164 165 2003; Cochrane et al. 1999) for discussion of the spatial structure of fire spread at the savanna-forest boundary]; instead, we incorporate fire effects only in the reaction term. 166 167 Nevertheless, at continental scales, in the absence of appropriate continuum models, a diffusion 168 model is a reasonable place to start.

In the next section, we explore the behavior of equation (2) using a series of simplifying 169 170 assumptions that are ecologically relevant. It may also be worth mentioning that the qualitative behavior of the equation (2) is independent of the particular details of f(T, P). However, in this 171 172 paper, we use a particular functional form of f(T, P) motivated by previous work on the subject 173 (Staver et al. 2011b; Staver and Levin 2012). We do this to compare and contrast the simulation 174 results of the previous mean-field model and its spatial counterpart, presented here.

175

176 **Methods and Results**

177 Since the mathematical literature on bistable reaction-diffusion models is scattered across various subfields of physics (Coleman 1977), mathematics (Aronson and Weinberger 1975; 178 179 Bramson 1983; Fife and McLeod 1977), and ecology (Lewis and Kareiva 1993; Murray 2001; Okubo and Levin 2013), we begin by summarizing some of the well-known results of the 1D 180 181 diffusion model in the context of savanna and forest biomes. Although some results for the 1D 182 model have been presented numerically elsewhere (Eby et al. 2017; van de Leemput et al. 2015), 183 here we provide analytical results that may yield deeper insights. These results will also provide a baseline for comparison with the 2D diffusion model that has not been discussed in the 184 185 literature.

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187 Reaction-Diffusion Model in One Dimension

188 Since one of the primary goals of the paper is to determine the spatial limits (boundaries) 189 of savanna and forest biomes, it is natural to look for solutions that naturally give rise to 190 boundaries. Based on the extensive literature on invasion biology (Hastings et al. 2005; Keitt et 191 al. 2001), we know that equation (2) has a traveling wave solution (see Fig. A1), where the 192 wavefront can be interpreted as the savanna-forest boundary. In this section, we find the velocity 193 of movement of the savanna-forest boundary as a function of system parameters, e.g., 194 precipitation. We then set the velocity to zero to find the equilibrium boundary position. Using the generalized waveform in one-dimension, T(x,t) = T(x - vt) (Aronson and 195 Weinberger 1975; Bramson 1983; Fife and McLeod 1977; Murray 2001; Okubo and Levin 196 197 2013), where v is the velocity of the savanna-forest boundary, we find that $v \propto \sqrt{D} \Delta U(P)$. 198 (3)where $\Delta U(P)$ is defined as the difference between values of the potential function at forest and 199 200 savanna state, respectively (see Appendix B.1 for calculations). Mathematically, the potential function is defined as $U(T, P) = -\int_{T_0}^T f(T, P) dT$, where f(T, P) is the mean-field growth 201 function (Nolting and Abbott 2016; Strogatz 2014). This potential function is a formal way of 202 203 defining the concept of a potential landscape that is commonly used to understand the resilience 204 of dynamical systems (Holling 1996; Strogatz 2014). In a bistable system, the potential landscape consists of two wells corresponding to the two stable states of the system (see top row 205 in Fig. 1). In the equation above, $\Delta U(P) = -\int_{T_c^*}^{T_F^*} f(T, P) dT$ is the difference between the 206 depth of potential wells corresponding to savanna and forest. 207

208 The equation above suggests that the magnitude of v is proportional to diffusion (\sqrt{D}) 209 and the difference in the depth of the potential wells (ΔU), while the direction of v is purely 210 determined by the sign of $\Delta U(P)$. Thus, in a homogeneous landscape (e.g., with constant precipitation across the whole landscape), the state with lower potential invades the one with 211 212 higher potential, except in the trivial case when the potential for both states is equal (*i.e.*, $\Delta U =$ 213 0; see top row in Fig. 1). The trivial case occurs at a unique precipitation value, which is referred 214 to as Maxwell precipitation (P_M) (Bel et al. 2012; Boettiger and Hastings 2013; Carr et al. 1984; Clerk-Maxwell 1875; Martín et al. 2015; Pomeau 1986; van de Leemput et al. 2015; Weissmann 215

216 and Shnerb 2014; Wuyts et al. 2017; Zelnik and Meron 2018). Next, to obtain the velocity of 217 movement of the savanna-forest boundary as a function of precipitation, we Taylor expand 218 $\Delta U(P)$ in equation (3) around P_M : $v \propto \sqrt{D} (P_M - P).$ 219 (4) 220 This equation implies that in a landscape with precipitation greater than P_M , forest encroaches savanna (v < 0), and conversely, that in a landscape with precipitation less than P_M , savanna 221 encroaches forest ($\nu > 0$). Only when the landscape receives precipitation exactly equal to P_M 222 is the savanna-forest boundary neutrally stable (*i.e.*, the magnitude of the small perturbations to 223 224 the boundary neither increases or decreases over time). In other words, in a homogeneous precipitation landscape with $P \neq P_M$, a stable savanna-forest boundary is not possible, under 225 226 these assumptions (van de Leemput et al. 2015). 227 228 A Precipitation Gradient and Stable Savanna-Forest Boundary 229 In the previous section, we assumed homogeneous precipitation conditions. However, at 230 continental scales, landscapes have precipitation gradients. In this section, we show how a 231 precipitation gradient can lead to a stable savanna-forest boundary (van de Leemput et al. 2015; 232 Wuyts et al. 2018). 233 To do this, we consider a 1D landscape with a linear precipitation gradient with 234 precipitation P at site x given by 235 $P(x) = g(x - x_M) + P_M,$ (5)where x_M is the spatial location receiving P_M and g is the change in precipitation per unit 236 237 distance (precipitation gradient constant). Substituting equation (5) into equation (4) we get $\Delta X = \Delta X_o \ e^{-g\beta\sqrt{D} t},$ 238 (6) where X is the position of the savanna-forest boundary, $\Delta X = X - x_M$ is the deviation of the 239 boundary from x_M , and β the natural logarithm of the proportionality constant in equation (4). 240 Equation (6) highlights two important features of the savanna-forest boundary. First, in a 241 242 1D landscape with linear precipitation gradient, the boundary equilibrates to P_M (Fig. 2). Second, if the boundary is perturbed locally (in any direction), it will recover back to x_M . 243 Moreover, the characteristic timescale of recovery is inversely proportional to \sqrt{D} and 244 precipitation gradient constant g. This suggests that the savanna-forest boundary is resilient to 245 246 local spatial perturbations (Fig. 2). Although not shown here, our numerical experiments in 1D 9

also suggest that the equilibrium distribution of savanna and forest is independent of initialconditions (see also van de Leemput et al. 2015; Wuyts et al. 2018).

249 Thus, the 1D diffusion model reproduces what the bistable biome theory (Beckage et al. 250 2009; Staver et al. 2011b; Staver and Levin 2012) could not: that spatial interactions, overlaid on 251 a large-scale precipitation gradient, can result in the spatial aggregation of savanna with savanna and forest with forest, separated by a stable savanna-forest boundary (Fig. 2). Moreover, the 252 253 model also predicts that biome shifts are reversible provided the climatic conditions are restored. 254 Unfortunately, the 1D diffusion model also predicts that the spatial limits of savanna and forest 255 biomes are solely determined by Maxwell precipitation (P_M) . In other words, this model fails to 256 produce overlap in the rainfall ranges of savanna and forest biomes (see also Eby et al. 2017; van 257 de Leemput et al. 2015; Wuyts et al. 2018), observed in the empirical data (Dantas et al. 2016; 258 Hirota et al. 2011; Staver et al. 2011a; Staver et al. 2011b).

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260 *Reaction-Diffusion in Two Dimensions*

Above, we assumed a one-dimensional landscape. This assumption, however, may not be realistic for understanding distribution of savanna and forest biomes, since it is somewhat obvious to observe that their dynamics are better described on a two-dimensional landscape. In this section, we show that adding a second dimension can qualitatively change the equilibrium position of the savanna-forest boundary, which can explain the overlap in the rainfall ranges over which biomes occur.

To incorporate the second dimension in the model, we use a 2D polar representation of the Laplacian operator in equation (2). Following the same analytical approach as in the 1D case, we show that

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$|\Delta P| \propto \sqrt{D} |\kappa_{Mc}|,\tag{7}$

where ΔP represents the difference between the precipitation at the savanna-forest boundary and the Maxwell precipitation contour (P_{Mc} , notationally distinct from the Maxwell point in 1D P_M ; see Appendix B.2 for calculations), $|\kappa_{Mc}|$ the absolute curvature of P_{Mc} (a measure of roundness; see Fig. A2), and *D* the diffusion constant (as above). This equation describes the local deviation of the savanna-forest boundary from the location of the P_{Mc} in terms of the difference in the precipitation; to obtain the deviation in terms of absolute distance, we multiply ΔP by *g* (precipitation gradient; see Eq. 5).

In plain terms, this means that, in a 2D landscape, the location of the boundary between savanna and forest is not determined only by precipitation but also crucially depends on the geometrical shape of the precipitation contours (specifically, of the Maxwell precipitation contour P_{Mc}). When the Maxwell precipitation contour is a straight line (*i.e.*, where $|\kappa_{Mc}| = 0$), the system behaves like a 1D model, and the savanna-forest boundary coincides with P_{Mc} (Fig. 3A). However, for an arbitrarily shaped P_{Mc} ($|\kappa_{Mc}| \neq 0$), the boundary deviates from the P_{Mc} depending upon the local curvature of P_{Mc} (Fig. 3B).

285 Ecologically, curvature effects described in equation (7) arise because of source-sink 286 dynamics (Pulliam 1988; Pulliam 2000) at the savanna-forest boundary. When P_{Mc} is a straight line, the inflow and outflow of seeds are balanced, thus resulting in a stable savanna-forest 287 288 boundary that coincides exactly with P_{Mc} . However, if P_{Mc} is curved, the balance between inflow and outflow of seeds is disrupted. For example, when the shape of P_{Mc} is such that there 289 290 are more forest neighbors than savanna neighbors surrounding a point in the landscape with P =291 P_{MC} (upper part of Fig. 3B), the inflow of seeds will be higher than their outflow. This creates a net positive inflow of seeds, resulting in a higher growth rate of trees. Forests expand, pushing 292 293 the boundary into savanna region till the added growth rate of trees due to a higher influx of 294 seeds is compensated by reduced growth rate due to a decrease in precipitation. Conversely, 295 when a point on P_{Mc} is surrounded by more savanna patches than forest patches (lower part of Fig. 3B), there will be a net positive outflow of seeds, which will favor savanna expansion. 296 297 Similar to the previous case, the boundary will move into forest regions until the reduced growth 298 rate of trees due to a lower influx of seeds is balanced by increased growth rate due to an 299 increase in precipitation.

300 The reaction-diffusion model, presented above, however, has some assumptions that are 301 likely to be violated in real-world: (1) that dispersal is local (because of diffusion 302 approximation), (2) that vegetation dynamics have no demographic or external noise, and (3) that 303 the reaction part has a well defined potential function. As a robustness check, we relax these 304 assumptions one by one, and numerically test their consequences for the theoretical results 305 presented above. First, we find that incorporating long-range dispersal (via fat tail dispersal 306 kernels) does not change equilibrium biome distributions (see Appendix C), presumably because 307 fire vegetation feedbacks prevent tree establishment far away from the source even when a seed arrives there (Barton and Turelli 2011; Bates et al. 1997; Kot et al. 1996). Second, we find that 308

309 adding noise makes the boundary increasingly rough with increasing noise; however, the 310 location of the boundary at a coarser scale does not move appreciably from its equilibrium 311 position predicted from the deterministic 2D diffusion model (see Fig. D1). Thirdly, and finally, 312 we consider a two-dynamical-variable reaction-diffusion system where a potential function 313 cannot be defined. In such a system, the position of the boundary, in addition to the control 314 parameter (e.g., precipitation), is dependent on the ratio of the two diffusion constants (Fig. D2); 315 the system still exhibits curvature effects in equation (7). Therefore, in no case did we find that 316 violating the above assumptions qualitatively changed dynamics.

317 To summarize, curvature effects in the 2D diffusion model can phenomenologically 318 reproduce the overlap in the precipitation ranges over which savanna and forest biomes occur, 319 missing from the 1D diffusion model, while simultaneously retaining the spatial aggregation 320 property of biomes (Fig. 3D). Moreover, the 2D diffusion model suggests that this precipitation 321 overlap is not maintained by hysteresis, a defining feature of a bistable biome theory. Instead, our simulations and analytical calculations suggest that in a landscape with a monotonic 322 323 precipitation gradient, hysteresis is unlikely. Below we discuss an ecological scenario under 324 which hysteresis may reappear.

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326 Critical Patch Size Effects in a Landscape with Non-Monotonic Precipitation Gradient

In the previous section, we assumed a monotonic gradient in precipitation. However, precipitation gradients in real-world landscapes are not always monotonic. As such, a landscape can have a complex distribution of precipitation with high precipitation regions intermittently distributed in low precipitation regions, and vice versa. In the following section, we describe how this feature of precipitation gradients can potentially lead to hysteresis.

But before we do that, we first consider a simpler case of a homogeneous precipitation model for analytical insight. Based on the theoretical works of Bradford and Philip (1970a,b), it can be shown that in a homogeneous precipitation landscape with bistable dynamics, the fate of an invasion process by a particular vegetation state into another is dependent on two factors:

- precipitation P and the initial patch area of the invading state A. An invading patch of
- 337 vegetation smaller than a critical patch area $(A < A_c)$ will not be able to expand even though that
- 338 vegetation state is climatically favourable, *i.e.*, the state which has lower potential (Bradford and
- Philip 1970a; Bradford and Philip 1970b; Holmes et al. 1994; Oxtoby 1998; Skellam 1951). Our

calculations suggest that near P_M , the critical patch size in a homogeneous landscape can be approximated as

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$$A_c \propto \frac{D}{(\Delta P_H)^{2'}} \tag{8}$$

where ΔP_H is the difference between the precipitation of the homogeneous landscape and P_M . 343 To fully understand the implications of equation (8), consider an initial savanna landscape with 344 precipitation just above P_M . Although in this landscape forest is more favourable than savanna 345 346 because of lower potential (Eq. 3), the forest state will only be able to invade if there is an initial patch of forest that has an area greater than A_c (Eq. 8). This is because a small patch of forest 347 348 has a high perimeter-to-area ratio (Skellam 1951), such that the accumulation of trees is slow because seed inflow per unit area from forest patches is low, preventing forest expansion. By the 349 350 same token, a landscape in a forest state with precipitation just below P_M would require a large patch of savanna to overcome high levels of seed rain from neighboring forest patches. 351

The same phenomenon also applies to a landscape with precipitation gradients. Consider an initial savanna landscape with spatially varying precipitation patterns such that the whole landscape has rainfall less than P_M , except in the center where the rainfall is just above P_M . Since the whole landscape was initialized with savanna, the center of the landscape will remain in a savanna state, unless the central region is initialized with forest patch of area greater than A_c . Conversely, a similar argument holds for an all forest landscape with a low-rainfall island in the center.

This suggests that the vegetation state of small and isolated patches in intermediate rainfall regions depends on the availability of a nucleation center, suggesting that the characteristic biome state in those areas might be contingent on historical biome patterns, thus exhibiting hysteresis. And more importantly, this analysis suggests that critical patch size (Eq. 8), in addition to curvature effects (Eq. 7), can also explain overlap in the rainfall ranges of savanna and forest biomes.

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366 Curvature and Critical Patch Size Effects in Empirical Systems

In this section, we test some of the key analytical predictions of our 2D reaction-diffusion
model—particularly those concerning curvature (Eq. 7) and critical patch size effects (Eq.
8)—using real savanna and forest distributions in Sub-Saharan Africa. As described above, a 2D

370 reaction-diffusion approximation of savanna-forest dynamics predicts that the location of the

371 savanna-forest boundary with respect to precipitation should vary depending upon the local 372 curvature of the boundary (Eq. 7; Fig. 3B) in such a way that the difference between the 373 precipitation at the boundary and P_{Mc} is linearly proportional to the local curvature of the 374 boundary. As such, plotting the absolute curvature of the boundary (ignoring its convexity) as a 375 function of boundary precipitation should yield a V-shaped curve, the vertex of which should 376 correspond to P_M . Indeed, using African tree cover and mean annual precipitation (MAP) data 377 (see Data Analysis in Online Appendix D), we show that the current distribution of savanna and 378 forest in Africa is consistent with this prediction (black line in Fig. 4). The location of the vertex of that curve provides an estimate of $P_M = 1508 \pm 84$ mm MAP, consistent with previous 379 380 empirical work (Staal et al. 2016 found $P_M = 1580$ mm MAP). The confidence interval for this estimate was determined by calculating P_M for various combination of parameter values of (i) 381 382 boundary tree cover (73-80%) used for identifying boundary location, and (ii) arc length of the 383 boundary (100-1000 km) used to estimate curvature. For more details on estimating P_M , see figure D3 and sensitivity analysis in Online Appendix D. 384

385 Next, we compare the above estimate of P_M , by estimating P_M with an alternative method. This involves simulating the potential distribution of savanna and forest using a 2D 386 387 reaction-diffusion model with current biome distributions as the initial condition. To do this, we 388 simulated spatial distributions of savanna and forest using present-day precipitation patterns for 389 various combinations of P_M and D (see Data Analysis in Online Appendix D), and, using a genetic algorithm (Scrucca 2013), selected those parameter values (P_M and D) that yielded the 390 'best fit' to the current distribution of biomes. Here, we refer to 'best fit' as maximizing pixel by 391 392 pixel match between simulated and empirical savanna and forest distribution (see Data Analysis in Online Appendix D). This procedure yielded an estimate of P_M (= 1538 mm MAP) that lay 393 394 within the expected precipitation range obtained from the curvature analysis in figure 4 (see Fig. 395 D4). These large-scale simulations also reproduced empirically observed biome distributions in 396 Sub-Saharan Africa surprisingly well for such a simple model (Fig. 5A), except for small regions 397 in the Bateke Plateau in Congo and Western Africa. In the Bateke Plateau, empirically observed 398 savannas may be maintained because shallow sandy soils that reduce effective soil moisture 399 (White 1986) or may alternatively be anthropogenic. Meanwhile, it is well established that 400 savannas in Western Africa are a result of historical deforestation (Adams and Faure 1997; 401 Aleman et al. 2017).

402 Next, to check whether the results of large-scale simulation in figure 5A were dependent 403 on the initial conditions – at least theoretically possible because of critical patch size effects, as 404 described above - we simulated the vegetation distribution for two more initial conditions: 'all 405 savanna' and 'all forest' in Sub-Saharan Africa (Fig. 5B-C), using the best fit parameter values estimated above (P_M and D). Simulations with 'all savanna' initial conditions (Fig. 5B) matched 406 those using current distributions as initial conditions (see Fig. 5A). However, 'all forest' initial 407 conditions produce substantially different biome patterns in the Southern Congo and Ethiopian 408 Highlands (Fig. 5C). 409

410 We propose that the critical patch area (A_c) requirement can potentially explain why the 411 simulations over-predict the forest extent in the Southern Congo (Fig. 5C) and under-predict the forest extent in Ethiopian Highlands (Fig. 5B). Since both of these regions are disconnected from 412 413 the main forest cluster by savanna vegetation, biome distributions in these regions are dependent 414 on the availability of historical nucleation centers (or initial conditions; Eq. 8). Based on our 415 simulations we suspect that Southern Congo and Ethiopian Highlands were historically occupied 416 by savannas and forests, respectively, which resulted in their present distribution. Although this 417 claim is currently hard to test due to lack of reliable long-term paleo-records from these regions 418 (however, see Elenga et al. 1994; Jolly et al. 1998), historical vegetation reconstructions for the early 20th century (Aleman et al. 2017; White 1986) are consistent with the theoretical 419 predictions of the model. 420

421

422 *Comparisons with Alternative Models*

423 Our calculations show that novel dynamical features of the 2D diffusion model — such 424 as spatial aggregation (Eq. 6), curvature effects (Eq. 7), and critical patch size (Eq. 8) — can 425 qualitatively explain many empirical features of savanna and forest distributions that previous 426 biome distribution models could not. In this section, we investigate whether these dynamical 427 features improve upon the predictions from previously proposed models of biome distribution. 428 To do this, we simulate the distribution of savanna and forest in Sub-Saharan Africa using three 429 alternatives (see Data Analysis in Online Appendix D). First, (a) we consider a 'one-climate 430 one-biome' model in which the savanna-forest boundary is determined by a unique precipitation contour. This model is analogous to the classical biome theory (Fig. 6A). Next, (b) we consider a 431 432 model in which the local vegetation dynamics in each patch are governed by mean-field bistable

model and the neighbouring patches do not interact. In this model, we randomly initialize the
landscape with savanna and forest patches (Fig. 6B); note, however, that this test does not
consider the possibility that initial conditions could be spatially structured, leading to spatial
structure in biome distributions today. Finally, (*c*) we consider a 2D reaction-diffusion model,
already described at length above (Fig. 6C). The diffusion model incorporates both bistability
vegetation dynamcis and dispersal.

We measure whether these models can – with parameter optimization – reproduce three components of biome distribution: overlap in the rainfall ranges of biomes, the spatial aggregation of savanna with savanna and forest with forest, and the match between the simulated and actual distribution of biomes (see Fig. 6, Fig. D5, and Table 1). Note, again, that tuned parameters do not necessarily correspond to demographic rates, etc., that might be measured empirically; note also that the three model alternatives we propose here are not exhaustive.

In the 'one-climate one-biome' model, the precipitation cutoff between savanna and 445 446 forest was found to be 1583 mm MAP (see Table 1). Whereas, in the other two models the rainfall ranges of savanna and forest showed considerable overlap between 1000 mm and 2000 447 448 mm. Meanwhile, the one-climate one-biome and 2D diffusion models show a high probability of 449 spatial aggregation (above 90%) that is missing in the mean-field bistable model (below 68%). 450 Thus, of the three models, only the 2D diffusion model can reproduce both spatial aggregation 451 and overlap in the rainfall ranges of biomes. Therefore, it is not surprising that the 2D diffusion 452 model also outperforms (97% accuracy) other models in terms of predicting the spatial distribution of biomes in Sub-Saharan Africa. 453

In summary, all models, except in a one-climate one-biome model, reproduce at least some overlap in the rainfall ranges of biomes (Table 1 and Fig. D5). Of the two remaining models, the mean-field bistable model fails to reproduce the spatial aggregation of biomes (Fig. 6 and Table 1). This leaves us with the 2D reaction-diffusion model, which reproduces not only the climatic overlap in the limits of biomes, and the spatial aggregation in biome distributions, but also the overall biome distributions in Sub-Saharan Africa with remarkable accuracy (see Table 1).

461

462 Discussion

463 In this paper, we develop and analyze a reaction-diffusion model to examine the contributions

of dispersal to the distribution and resilience of tropical savanna and forest biomes. The model
assumes that the local mean-field dynamics of biomes are governed by non-linear fire-vegetation
feedbacks and that adjacent savanna and forest patches interact spatially through seed dispersal.

467 We find that the model reproduces empirical features missing from existing biome 468 distribution models. Specifically, the 2D reaction-diffusion model simultaneously reproduces 469 both overlap in the climatic ranges of biomes, as well as spatial aggregation of savanna with 470 savanna and forest with forest. As before, we find that fire-vegetation feedbacks may substantially expand savanna distributions at the expense of forests, but that in a spatial context 471 this does not necessarily translate into bistable vegetation distributions. Instead, the equilibrium 472 473 position of the savanna-forest boundary is determined by a combination of three factors: (a) 474 climate (via impacts on the relative depth of potential wells for each biome), (b) source-sink 475 dynamics (via local curvature of the Maxwell precipitation contour), and, occasionally, by (c) 476 availability of historical nucleation centers (which contributes an element of hysteresis to 477 distribution dynamics, albeit more limited than that described before). These theoretical 478 predictions are empirically consistent with observations of the curvature of the savanna-forest 479 boundary, and large-scale simulations which show that the 2D diffusion model can — with 480 parameter optimization-reproduce empirically observed patterns of savanna and forest 481 distributions in Sub-Saharan Africa.

482 These findings have direct implications for how we think of the stability and resilience of 483 tropical biomes. Classical biome theories suggest that perturbations to biome distributions should 484 be easily reversible since vegetation tracks climate directly (Holdridge 1947; Schimper 1902; 485 von Humboldt and Bonpland 1807; Whittaker 1970). By contrast, more recent work has suggested that fire-vegetation feedbacks can stabilize savanna as an alternative to forest in some 486 487 areas, such that perturbations to biome distributions may not be easily reversible (Beckage et al. 488 2009; Staver et al. 2011b; Staver and Levin 2012). Here, we show that combining a spatial 489 dispersal process with an underlying bistable model radically alters stability predictions: biome 490 recovery after perturbation becomes much more likely, even if fire-vegetation feedbacks do 491 modify vegetation (which they probably do; see Bond et al. 2005). In this scenario, biome 492 transitions may be regionally predictable and reversible, even if they are locally abrupt. 493 However, there is a notable caveat to this prediction. In isolated rainfall islands,

494 vegetation distributions may exhibit hysteresis; analogously, if remnant vegetation patches are

495 reduced below a critical area, recovery of the boundary may be impossible, resulting in a 496 permanent loss of vegetation. As a result, extensive historical forest loss, *e.g.*, in West Africa, 497 coastal Kenya and Tanzania, and the Ethiopian highlands, may be irrecoverable without direct 498 intervention, since remnant forest patches may be too small for forest to recolonize successfully 499 (Aleman and Staver 2018). This also raises contrasting concerns about proposed afforestation 500 plans in mesic savannas of the Southern Congo (Veldman et al. 2015). These isolated mesic 501 savannas might be historically maintained as a stable alternative biome state (Aleman et al. 502 2017); proposed afforestation practices (Veldman et al. 2015) in these regions could trigger a 503 permanent shift in ecosystem state from savanna to forest (Fig. 5C), which may lead to loss of 504 endemic biodiversity in mesic savannas (Bond 2016) and wastage of scarce management 505 resources.

The results of the reaction-diffusion model presented herein should be interpreted with 506 507 caution, however. For starters, we here incorporate only a subset of important spatial processes, 508 notably ignoring the long-range spread of fire within savannas (Schertzer et al. 2015) and local 509 fire spread at the savanna-forest boundary (Cochrane 2003; Cochrane et al. 1999) both of which 510 may be significant (note, however, that we have included fire effects in the reaction term of the model). However, ongoing analytical work on a more thorough set of models that examine fire 511 512 effects at the boundary between savanna and forest (Durrett and Ma 2018) suggest that the 513 phenomenological results presented herein may be applicable more broadly: scaling limits to 514 those models appear also to be characterized by traveling waves, with the occurrence of 515 stationary savanna-forest boundaries only in landscapes that include a gradient in rainfall 516 (Durrett and Ma 2018). Notably, however, long-range fire effects seem to change predictions 517 somewhat (Li et al., in review), resulting in the emergence of stable savanna-forest mosaics even 518 under homogenous climatic conditions (Schertzer et al. 2015).

Another major question surrounds the problem of time-scales of ecological processes. Here, we have considered only the equilibrium distribution of biomes, ignoring the speed of equilibration. Modern climate change is sufficiently rapid (Karl and Trenberth 2003), and so associated with extreme climatic events (Jentsch et al. 2007; Katz and Brown 1992), that biome responses to ongoing anthropogenic global change are unlikely to be dominated by these local spatial processes. This may result in transient mismatches between climate and equilibrium vegetation (Webb 1986), which may be persistent from timescales ranging from decades to

526 millennia depending on the speed of ecological dynamics (Hastings 2004; Hastings et al. 2018;

- 527 Hastings and Higgins 1994). Therefore, understanding how fast biomes respond to changing
- 528 climate (empirically from the paleo-records) and using this information to incorporate dispersal
- 529 into existing non-spatial biosphere models (Bond and Keeley 2005; Bond et al. 2003; Moncrieff
- et al. 2014; Scheiter and Higgins 2009; Scheiter et al. 2013) will be critical to generating
- 531 informative predictions for the effects of anthropogenic global change on biome distributions.

Projections show that rapidly changing climate (Lewis et al. 2011; Nepstad et al. 2004) and land-use change (Aleman et al. 2016; Cochrane and Laurance 2002) are expected to result in large-scale biome shifts, which may yield huge economic and ecological losses. Here, we argue that, except in a few cases, dispersal can, in general, increase the resilience of tropical savanna and forest biomes to natural and anthropogenic disturbances (see also van de Leemput et al. 2015). However, recovery from disturbance could be slow, due to slow dynamics of biomes and

- anthropogenic or natural dispersal barriers.
- 539

540 Author Contributions

541 NG and ACS designed this model based on a concept from ACS. NG and VG implemented the

542 diffusion and integro-differential equation models. NG performed simulations and data analyses.

- 543 ACS and NG co-wrote the manuscript with feedback from VG and SAL. All the authors
- 544 contributed ideas and discussions.
- 545

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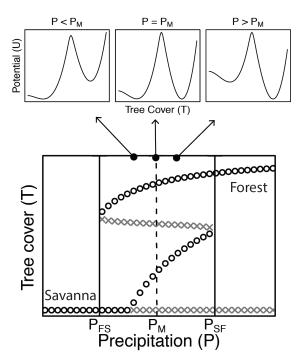
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804 FIGURES

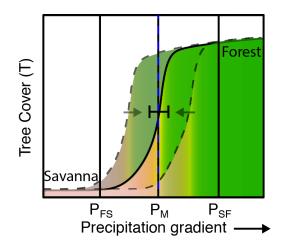


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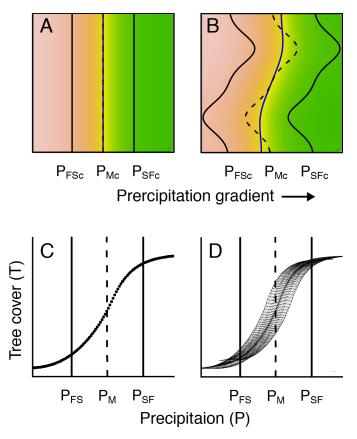
Figure 1: Potential functions of different vegetation configurations (top) and bifurcation diagram of the mean-field reaction term (bottom). In the intermediate precipitation region (bounded by the critical precipitation values P_{FS} and P_{SF}), the bifurcation diagram shows that the system can exist in both savanna and forest states, depending upon the initial conditions. In the bistable region, the depth of the potential function (U in the top row) corresponding to savanna and forest states depends on the precipitation value. Both savanna and forest states have equal potential at a unique precipitation value, referred to as Maxwell precipitation (P_M). Below (above) P_M ,

savanna (forest) state has a deeper potential than forest. In the bottom panel, stable (unstable)

equilibrium points are marked as black circles (dark-grey crosses).



816 Figure 2: Equilibrium (solid sigmoidal curve) and transient (dash sigmoidal curve) tree cover 817 along a linear spatial precipitation gradient in a 1D landscape. The two solid vertical lines correspond to the two critical points (P_{FS} and P_{SF}), and the vertical black dashed line to the 818 Maxwell precipitation (P_M) . The plot suggests that spatial interactions coupled with a large-scale 819 820 gradient in precipitation can result in the spatial aggregation of savanna and forest, separated by 821 a stable savanna-forest boundary (indicated by the blue vertical line, in this case coincident with P_{M}). This boundary is resilient to perturbations and always recovers back to its equilibrium 822 position after a disturbance. This model, however, fails to reproduce the non-deterministic 823 824 relationship between biome and precipitation, observed in the empirical data. The simulations were initialized with random initial condition (see Numerical Methods in Online Appendix D). 825



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Figure 3: Simulated equilibrium distribution of tree cover (A and B) and its relationship with 827 828 precipitation (C and D) in a 2D landscape with a monotonic precipitation gradient. The columns show the results for two geometries of precipitation contours: linear (A and C) and curved (B and 829 830 D). Simulations suggest that in a 2D landscape, the equilibrium position of the boundary is not only determined by P_M , but also depends on the curvature of the Maxwell precipitation contour 831 κ_{Mc} . When $\kappa_{Mc} = 0$ (linear P_{Mc}), the boundary aligns with P_{Mc} (A). This situation is 832 analogous to the one-dimensional model in figure 2. However, when $\kappa_{Mc} \neq 0$ (arbitrary shaped 833 P_{Mc}), the boundary deviates from P_{Mc} according to equation (7) (B). These curvature effects 834 can reproduce the non-deterministic relationship between biome and precipitation (D), missing in 835 the one-dimensional model (Fig. 2). The simulations were initialized with random initial 836 837 condition (see Numerical Methods in Online Appendix D).

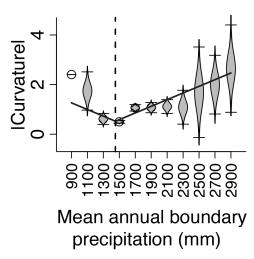


Figure 4: Empirical response of local savanna-forest boundary curvature to mean annual precipitation at the boundary in sub-Saharan Africa. Results show that absolute curvature vs. boundary precipitation exhibit V-shaped relationship (black line), consistent with our theoretical prediction (eq. 7). In theory, the vertex of V corresponds to P_M , loosely corresponding to results from the extensive sensitivity analysis that estimates $P_M = 1508 \pm 84$ MAP (Fig. D3; see Data Analysis in Online Appendix D).

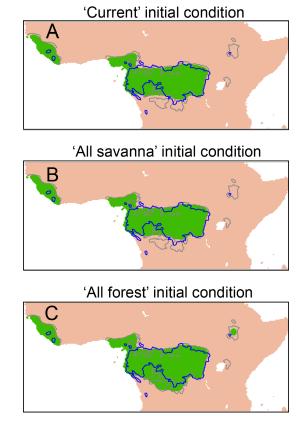


Figure 5: Simulated distributions of savanna and forest in Africa, initialized with the current 846 distribution of biomes (A), all savanna (B), and all forest (C). Blue lines correspond to the 847 observed present-day savanna-forest boundary, and the grey line represents P_{Mc} (with $P_M =$ 848 1538 mm MAP). The large-scale simulations in (A) and (B) matched, and reproduced the 849 current distribution of savanna and forest, except the edaphic savannas on the Bateke Plateau in 850 851 Congo and deforested areas in western Africa. However, the simulations in (C) significantly 852 overpredicted the extent of forest in the Southern Congo. This region can climatically support 853 both savanna and forest depending upon the historical vegetation state of the region (see Data 854 Analysis in Online Appendix D).

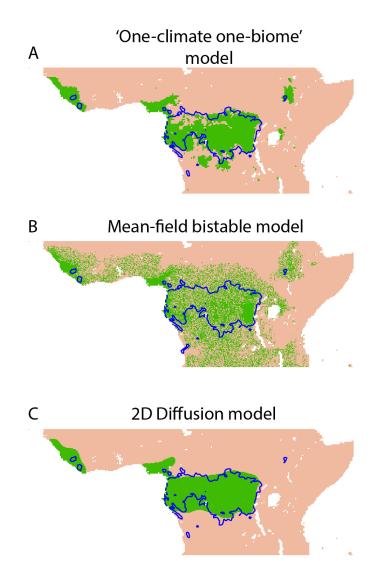


Figure 6: Simulated distribution of savanna and forest biomes from three biome distribution models: 'one-climate one-biome' model (A), mean-field bistable model (B), 2D reaction-diffusion model (C). Blue lines correspond to the present-day savanna-forest boundary. These results indicate that the 2D reaction-diffusion model can, with tuning, describe the quantitative distribution of biome patterns. This model reproduces both spatial aggregation and overlap in rainfall ranges of biomes and is also the best predictor of biome patterns in Central Africa (see Table 1 and Data Analysis in Online Appendix D).

Table 1: Summary statistics of the simulated distribution of biomes in Central Africa using three alternative models. The performance of models was evaluated on three aspects: overlap in the rainfall ranges of biomes (columns 2 and 3), spatial aggregation of savanna with savanna and forest with forest (columns 4 and 5), and match between the simulated and actual distribution of biomes (columns 6, 7, and 8). Note that we excluded the contributions of the deforested regions in Western Africa and edaphic savannas of Bateke Plateau while calculating the goodness of fit (see Data Analysis in Online Appendix D).

	Rainfall range (mm MAP)		Spatial aggregation		Goodness of fit (%)		
Model	Savanna	Forest	Probability of finding savanna next to a savanna	Probability of finding forest next to a forest	Correctly classified	Incorrectly classified as savanna	Incorrectly classified as forest
'One-climate one-biome' model	Below 1583	Above 1583	0.96	0.90	93.43	3.09	3.48
Mean-field bistable model	Below 2000	Above 1000	0.68	0.36	78.90	5.49	15.61
2D diffusion model	Below 2042	Above 1058	0.97	0.96	97.00	1.24	1.76