Stiffness and tension gradients of the hair cell's tip-link complex in the mammalian cochlea.

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SIGNIFICANCE STATEMENT: The tip-link complex of the hair cell is mechanically
tuned along the tonotopic axis of the cochlea.

ABSTRACT Frequency analysis of sound by the cochlea relies on sharp frequency tuning of
mechanosensory hair cells along a tonotopic axis. To clarify the underlying biophysical
mechanism, we have investigated the micromechanical properties of the hair cell's
mechanoreceptive hair bundle in the rat cochlea. We studied both inner and outer hair cells,
which send nervous signals to the brain and amplify cochlear vibrations, respectively. We
find that tonotopy is associated with gradients of stiffness and resting mechanical tension,
with steeper gradients for outer hair cells, emphasizing the division of labor between the two
hair-cell types. We demonstrate that tension in the tip links that convey force to the mechano-
electrical transduction channels increases at reduced Ca^{2+}. Finally, we reveal tonotopic
gradients in stiffness and tension at the level of a single tip link. We conclude that
mechanical gradients of the tip-link complex may help specify the characteristic frequency of
the hair cell.

INTRODUCTION

The cochlea –the auditory organ of the inner ear– is endowed with a few thousands of mechanosensory hair cells that are each tuned to detect a characteristic sound frequency (Fettiplace and Kim, 2014). Different frequencies are detected by different cells, which are spatially distributed in the organ according to a frequency –or tonotopic– map (Lewis et al., 1982; Greenwood, 1990; Viberg and Canlon, 2004). Despite its critical importance for frequency analysis of complex sound stimuli, determining the mechanism that specifies the characteristic frequency of a given hair cell remains a major challenge of auditory physiology.

Although certainly not the only determinant of hair-cell tuning (Fettiplace and Fuchs, 1999), we focus here on the contribution of the hair bundle, the cohesive tuft of cylindrical processes called stereocilia that protrude from the apical surface of each hair cell. The hair bundle works as the mechanical antenna of the hair cell (Hudspeth, 1989). Sound evokes hair-bundle deflections, which modulate the extension of elastic elements –the gating springs– connected to mechanosensitive ion channels. By changing the elastic energy stored in the gating springs, sound stimuli affect the channels' open probability, resulting in a mechano-electrical transduction current. Importantly, mechanical stress is conveyed to the mechano-electrical transduction channels by oblique proteinaceous tip links (Pickles et al., 1984; Kazmierczak et al., 2007) that interconnect the stereocilia near their tips. Whether or not the tip link embodies the gating spring, however, is unsure. Electron microscopy (Kachar et al., 2000) and molecular dynamics simulations (Sotomayor et al., 2010) have suggested that the tip link may be too rigid and thus that a compliant element in series with the tip link could determine the gating-spring stiffness. If so (but see (Bartsch et al., 2018)), the gating spring could reside within the assemblies of molecules that anchor both sides of the tip link to the actin core of the stereocilia (Michalski and Petit, 2015), or perhaps be associated with membrane deformations at the lower insertion point of the tip link (Powers et al., 2012). In the following, the molecular assembly comprising the tip link, the transduction channels, as well as the molecules to which they are mechanically connected will be called the 'tip-link complex'.

The operating point of the transducer lies within the steep region of the sigmoidal relation between the transduction current and the hair-bundle position (Corey and Hudspeth, 1983; Russell and Sellick, 1983; Johnson et al., 2011). This key condition for sensitive hearing is thought to be controlled by tension in the tip links at rest (Hudspeth and Gillespie, 1994; Gillespie and Muller, 2009), as well as by extracellular and intracellular calcium (Corey and

Hudspeth, 1983; Ricci et al., 1998; Fettiplace and Kim, 2014), which is thought to stabilize

the closed state of the transduction channels (Hacohen et al., 1989; Cheung and Corey, 2006).

Tip-link tension has been estimated at ~8 pN in the bullfrog's sacculus (Jaramillo and

Hudspeth, 1993) but, to our knowledge, there has been no such report in the mammalian

cochlea.

Adaptation continuously resets the mechanosensitive channels to a sensitive operating

point when static deflections of the hair bundle threaten to saturate mechanoelectrical

transduction (Eatock, 2000). Most of the available evidence indicates that movements by

molecular motors actively pulling on the tip links and calcium feedback on the open

probability of transduction channels contribute to adaptation. With mammalian cochlear hair

cells, however, the dependence of adaptation on Ca^{2+} entry has recently been the subject of

significant controversy (Peng et al., 2013; Corns et al., 2014; Peng et al., 2016; Effertz et al.,

2017). Motor forces and calcium feedback can also explain the active hair-bundle

movements, including spontaneous oscillations, that have been observed in various species

(Fettiplace and Hackney, 2006; Martin, 2008). Active hair-bundle motility may contribute to

hair-cell tuning by actively filtering and amplifying sound inputs (Hudspeth, 2008). These

findings emphasize the importance of the tip-link complex, including the transduction

channels, the tip links that convey sound-evoked forces to these channels, as well as the

molecular motors that pull on the tip links, for mechanosensitivity of the hair cell.

Electrophysiological properties of the transduction apparatus, including the activation

kinetics and the conductance of the transduction channels, as well as the kinetics of

adaptation, have been shown to vary with the characteristic frequency of the hair cell (Ricci et

al., 2003; Ricci et al., 2005; Fettiplace and Kim, 2014; Beurg et al., 2018). These

observations suggest that hair-cell tuning may depend on the transducer itself (Ricci et al.,

2005). In addition, it is a ubiquitous property of vertebrate auditory organs that the

morphology of the hair bundle varies systematically with the characteristic frequency of the

corresponding hair cell (Wright, 1984; Lim, 1986; Roth and Bruns, 1992; Tilney et al., 1992):

going from the high-frequency to the low-frequency region of the organ, the hair bundle gets

longer and comprises a progressively smaller number of stereocilia. These morphological

gradients have long been recognized as circumstantial evidence that the mechanical properties

of the hair bundle might be involved in frequency tuning (Turner et al., 1981; Flock and

Strelioff, 1984; Fettiplace and Fuchs, 1999). However, a detailed characterization of

mechanical gradients at the level of the whole hair bundle is lacking, in particular to clarify
the contribution of the tip-link complex to these gradients.

In this work, we probed passive and active hair-bundle mechanics along the tonotopic axis
of an excised preparation of the rat cochlea (Fig. 1). We worked both with inner hair cells,
which convey auditory information to the brain and are considered as the true sensors of the
organ, and outer hair cells, which are mostly dedicated to cochlear amplification of sound-
evoked vibrations (Hudspeth, 2014). We combined fluid-jet stimulation to deflect the hair
bundle, iontophoresis of a calcium chelator (EDTA) to disrupt the tip links and measure
bundle movements resulting from tension release by these links, and patch-clamp recordings
of transduction currents to infer the number of intact tip links contributing to the response.
From these measurements, we estimated the stiffness of the whole hair bundle, the
contribution of the tip links and of the stereociliary pivots to this stiffness, as well as the
resting tension in the tip links. Our results reveal mechanical gradients of the tip-link
complex according to the tonotopic map and to the division of labor between sensory inner
and amplificatory outer hair cells, providing evidence for the implication of the tip-link
complex to frequency tuning of cochlear hair cells.

RESULTS
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The hair-bundle stiffness increases along the tonotopic axis. Using a calibrated fluid jet
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(Methods; Figure 2−figure supplements 1-5), we applied force steps to single hair bundles
along the tonotopic axis of the rat cochlea, going from the very apex of the organ to mid-
cochlear locations (Fig. 1). Each hair bundle responded to a force step with a fast deflection
in the direction of the stimulus followed by a slower movement in the same direction (Fig. 2;
Figure 2−figure supplement 5A). Over the duration of the step, the deflection of the hair
bundle increased in the direction of the applied step on average by 22% for inner hair cells
and by 12% for outer hair cells, corresponding to an apparent softening of the hair bundle by
the same amount). A mechanical creep is expected from mechanical relaxation of tip-link
tension associated with myosin-based adaptation (Hudspeth, 2014). Accordingly, the creep
was strongly reduced upon tip-link disruption by EDTA treatment (Figure 2−figure
supplement 5).

For outer hair cells, we found that a given series of force steps evoked hair-bundle
deflections that decreased in magnitude towards more basal locations along the tonotopic axis
(Fig. 2A). Correspondingly, the slope of a bundle's force-displacement relation (Fig. 2B),

and thus stiffness, increased with the characteristic frequency of the hair cell. The same behavior was observed for inner hair cells (Fig. 2C-D). Remarkably, the stiffness gradient was steeper (*p<0.05; Figure 2−supplement table 1) for outer hair cells than for inner hair cells (Fig. 2E). As the characteristic frequency increased from 1 to 4 kHz, the hair-bundle stiffness K_{HB} increased by 240% from 2.5±0.2 mN/m (n = 19) to 8.6±0.5 mN/m (n = 21) for outer hair cells, but only by 120% from 1.7±0.2 mN/m (n = 19) to 3.8±0.4 mN/m (n = 19) for inner hair cells. At the 15-kHz position, where stiffness was only be recorded for inner hair cells (see Methods), $K_{HB} = 5.5±0.4$ mN/m (n = 14), thus still significantly lower (***p<0.001; Figure 2−supplement table 1) than in outer hair cells at the 4-kHz position. At each cochlear position, outer hair-cell bundles were stiffer than inner hair-cell bundles, with a stiffness ratio that increased from the apex to the base of the organ.

As also observed by others using fluid-jet stimulation of cochlear hair cells (Géléoc et al., 1997; Corns et al., 2014), we measured force-displacement relations that were remarkably linear (Fig. 2B and D), showing no sign of gating compliance (Howard and Hudspeth, 1988). There are at least two possible explanations for this observation. First, the rise time (~500 µs; Figure 2−figure supplement 3) of our fluid-jet stimuli may have been too long to outspeed fast adaptation, masking gating compliance (Kennedy et al., 2003; Tinevez et al., 2007). Second, gating forces –the change in tip-link tension evoked by gating of the transduction channels (Markin and Hudspeth, 1995)− may have been too weak to affect hair-bundle mechanics under our experimental conditions (Fettiplace, 2006; Beurg et al., 2008).

Parsing out the relative contributions of gating springs and stereociliary pivots to hair-bundle stiffness. There are two contributions to the stiffness of a hair bundle: $K_{HB} = K_{GS} + K_{SP}$. First, hair-bundle deflections modulate the extension of the gating springs that control the open probability of the mechano-electrical transduction channels; we denote by K_{GS} their contribution to hair-bundle stiffness. Second, bending of the actin core of the stereocilia at the stereociliary pivots, as well as stretching horizontal lateral links that interconnect the stereocilia, provides the remaining contribution K_{SP}. Because the gating springs are in series with the tip links, disrupting the tip links affords a means to estimate both K_{GS} and K_{SP}. We used local iontophoretic application of a Ca^{2+} chelator (EDTA; Methods and Fig. 1) to disengage the Ca^{2+}-dependent adhesion of the cadherin-related molecules forming each tip link (Kazmierczak et al., 2007). From the increased magnitude of the hair-bundle response to a given mechanical stimulus (see an example in Fig. 4A), we found that the gating springs contributed up to 50% of the total hair-bundle stiffness K_{HB}. Averaging over all inner and

outer hair cells that we tested, the relative contribution of the gating springs was $r = K_{GS}/K_{HB} = 22\pm2$ % (n = 71; Figure 3–figure supplement 1), where $1 - r$ is the amplitude ratio of hair-bundle movements before and after tip-link disruption. Both inner and outer hair cells displayed a gradient of gating-spring stiffness $K_{GS} = r\, K_{HB}$ (Fig. 3A). Between the 1-kHz and the 4-kHz positions, the gating-spring stiffness increased by 520% for outer hair cells but only by 300% for inner hair cells. Similarly, the contribution $K_{SP} = (1 - r)\, K_{HB}$ of the stereociliary pivots to hair-bundle stiffness displayed tonotopic gradients for both inner and outer hair cells (Fig. 3B).

Individual gating springs are stiffer in hair cells with higher characteristic frequencies.
Both the pivot stiffness K_{SP} and the gating-spring stiffness K_{GS} are expected to vary according to hair-bundle morphology. Hair bundles get shorter and are composed of more numerous stereocilia as one progresses from the apex to the base of the cochlea (Figure 3–figure supplement 2), which ought to promote higher stiffness values. Are morphological gradients sufficient to explain the observed stiffness gradients of the hair bundle? Accounting for morphology, we write $K_{SP} = \kappa\, N_{SP}/h^2$ and $K_{GS} = k_{GS}\, N_{TL}\gamma^2$, in which h, N_{SP}, N_{TL} correspond, respectively, to the height, the number of stereocilia and the number of (intact) tip links of the hair bundle, whereas $\gamma \propto 1/h$ is a geometrical projection factor (Figure 3–figure supplements 2-3). Remarkably, the intrinsic rotational stiffness κ of a single stereocilium in outer hair cells remained the same across the positions that we explored (Fig. 3C; Figure 3–supplement table 2). Similarly, with inner hair cells, there was no significant variation of the rotational stiffness between the 1- and 2-kHz locations as well as between the 4- and 15-kHz locations, although we observed an increase by 100% between the 2- and 4-kHz locations (Figure 3–supplement table 2). Averaging over the ensembles of outer and inner hair cells that we probed, the rotational stiffness $\kappa = 1.2\pm0.2$ fN·m/rad (n = 79) in outer hair cells was about 140% higher than the value $\kappa = 0.5\pm0.1$ fN·m/rad (n = 137) measured in inner hair cells. In contrast, the intrinsic stiffness k_{GS} of a single gating spring increased by 180% from 1.3±0.4 (n = 29) to 3.7±0.7 mN/m (n = 17) in outer hair cells and by 240% from 0.5±0.2 (n = 14) to 1.7±0.3 mN/m (n = 30) in inner hair cells, for characteristic frequencies that increased from 1 to 4 kHz and from 1 to 15 kHz, respectively (Fig. 3D). Thus, morphological gradients can account for the observed gradient in pivot stiffness K_{SP}, but not for the observed gradient in gating-spring stiffness K_{GS} and in turn for the whole hair-bundle stiffness K_{HB}. The hair-bundle morphology is not the sole determinant of hair-bundle mechanics.

Tip-link tension increases along the tonotopic axis. We then estimated the mechanical tension in the tip links at rest, *i.e.* in the absence of an external stimulus. The transduction channels close when the tip links are disrupted, indicating that the channels are inherently more stable in a closed state (Assad et al., 1991; Beurg et al., 2008; Indzhykulian et al., 2013). In functional hair bundles, tip-link tension is thought to bring the operating point of the transducer within the steep region of the sigmoidal relation between the channels' open probability and the position of the hair bundle, ensuring sensitive detection of hair-bundle deflections. If there is tension in the tip links, then disrupting these links must result in a positive offset in the resting position of the hair bundle (Assad et al., 1991; Jaramillo and Hudspeth, 1993).

In response to iontophoresis of a Ca^{2+} chelator (EDTA), we observed a net positive movement ΔX_R of the hair bundle at steady state, as expected if the tip links broke and released tension (Fig. 4A). Consistent with tip-link disruption, this movement was associated with a decrease in hair-bundle stiffness, as well as with closure of the transduction channels and loss of transduction (Fig. 4B). The positive offset in resting position upon tip-link disruption was observed at all positions that we explored along the tonotopic axis of the cochlea, both for inner and outer hair cells, demonstrating that the hair bundles were indeed under tension (Fig. 5A). In addition, we observed that the magnitude of the evoked movement increased significantly (**$p<0.01$; Figure 5−supplement table 1) from 9±3 nm (n = 13) to 45±10 nm (n = 12) for outer hair cells with characteristic frequencies that increased from 1 to 4 kHz. In contrast, we observed no significant difference among inner hair cells with characteristic frequencies that varied within the 1−15-kHz range ($p>0.05$; Figure 5−supplement table1): the positive offset was 21±2 nm (n = 71) over the whole ensemble of inner hair cells.

As a result, within the range of cochlear locations that we explored, we measured a steep gradient of hair-bundle tension for outer hair cells but a comparatively weaker gradient (*$p<0.05$; Figure 5−supplement table 1) for inner hair cells (Fig. 5B). Tension $T_R = K_{SP}\,\Delta X_R$ in the hair bundle was estimated as the product of the pivot stiffness K_{SP} and the positive offset ΔX_R in resting position evoked by tip-link disruption (Methods). The hair-bundle tension increased by nearly 13-fold from 18±7 pN (n = 15) to 248±59 pN (n = 17) for outer hair cells (characteristic frequencies: 1−4 kHz) but by only 3.3-fold from 24±5 pN (n = 33) to 100±22 pN (n = 12) for inner hair cells (characteristic frequencies: 1−15 kHz). Tension in the hair bundle resulted from the summed contributions of tension in individual tip links.

Dividing the tension T_R by the average number N_{TL} of intact tip links in our recordings and
projecting the result along the oblique axis of the tip links (projection factor γ) provided
estimates of the tension $t_R = T_R/(\gamma\, N_{TL})$ in a single tip link. Remarkably, the observed
gradients in hair-bundle tension (Fig. 5B) were not only due to an increase in the number of
tip links that contributed to this tension (Figure 3−figure supplement 3), for tension in a single
tip link also showed gradients (Fig. 5C). The single tip-link tension was comparable in the
two types of cells at the 1-kHz location: 4.7±2.0 pN (n = 23) for outer hair cells and
5.9±1.3 pN (n = 43) for inner hair cells. However, at the 4-kHz location, the single tip-link
tension had increased by 620% to 34±8 pN (n = 19) in outer hair cells but only by 170% to
16±4 pN (n = 17) in inner hair cells; at the 15-kHz location, tip-link tension in inner hair cells
was 19±5 pN (n = 18). A linear regression of the relation between the single tip-link tension
and the characteristic frequency confirmed that the gradient was significantly (*p<0.05;
Figure 5−supplement table 1) steeper for outer hair cells.

Note that we have estimated here the stiffness and tension of a single tip-link complex
along the hair bundle's axis of bilateral symmetry, which corresponds to the natural axis of
mechanical stimulation in the cochlea. If the horizontal projection of tip links were actually
oriented at an angle with respect to this axis, stiffness and tension ought to be larger along the
tip-link axis than the values reported here. A systematic change in tip-link orientation along
the tonotopic axis would also affect the gradients. These effects are expected to be small in
inner hair cells, which are endowed with nearly linear hair bundles, but could be relevant for
the V-shaped hair bundles of outer hair cells (Figure 3−figure supplement 2 and (Pickles et
al., 1987)).

Tip-link tension first increases upon Ca^{2+} chelation. The dynamic response to an
iontophoretic step of EDTA, and thus to a decrease of the extracellular Ca^{2+} concentration,
was biphasic. The hair bundle first moved in the negative direction (arrowhead in Fig. 4A),
before the directionality of the movement reverted and the bundle showed the positive
movement associated with tip-link disruption. The negative movement was associated with
an increased inward current of similar time course (Fig. 4B). Within the framework of the
gating-spring model of mechanoelectrical transduction (Corey and Hudspeth, 1983; Markin
and Hudspeth, 1995), this observation is readily explained if the evoked decrease in the
extracellular Ca^{2+} concentration resulted in an increase in gating-spring tension, which both
pulled the hair bundle in the negative direction and led to the opening of the transduction
channels.

The magnitude of the negative movement at the peak showed no significant gradient and
was similar between inner and outer hair cells, with an average magnitude of $\Delta X_{Ca} =$
-26 ± 2 nm over the whole ensemble of hair cells (n = 83; Fig. 6A).   However, because
morphological gradients (Figure 3–figure supplement 2) resulted in gradients of pivot
stiffness K_{SP} (Fig. 3B), the maximal increase $\Delta T = -K_{SP}\,\Delta X_{Ca}$ in hair-bundle tension was
larger for hair cells with higher characteristic frequencies (Fig. 6B), as was the maximal
tension t_{max} that a single tip link sustained before tip-link disruption (Fig. 6C).   Going from
1-kHz to 4-kHz locations, this maximal tip-link tension displayed a gradient from 14 ± 4 pN
(n = 30) to 54 ± 12 pN (n = 23) in outer hair cells and from 15 ± 3 pN (n = 49) to 44 ± 8 pN
(n = 39) in inner hair cells; at the 15-kHz location, the maximal tension was not significantly
different than at the 4-kHz location in inner hair cells.

When immersing the hair cells in low-Ca^{2+} saline, the negative movement was always
followed by tip-link disruption and could thus not be observed twice with the same hair
bundle. However, in six different preparations for which the hair bundle was immersed in
saline with a higher Ca^{2+} concentration (500 µM) than usual (20 µM), we were able to
preserve the integrity of the tip links and demonstrate that the negative movements could be
reversible (Fig. 6D). Under such conditions, we observed that the absolute magnitude and the
speed of the negative movement increased with the magnitude of the iontophoretic current.
Notably, the hair bundle reached a new steady-state position when the iontophoretic step was
long enough (Fig. 6E), suggesting that resting tension in the tip links could be modulated by
the extracellular Ca^{2+} concentration, with higher tensions at lower Ca^{2+} concentrations.

DISCUSSION
Tonotopy of the mammalian cochlea is known to be associated with gradients of hair-bundle
morphology (Wright, 1984; Lim, 1986; Roth and Bruns, 1992; Tilney et al., 1992), as well as
of electrophysiological properties of the transduction apparatus (Ricci et al., 2003; Ricci et al.,
2005; Fettiplace and Kim, 2014; Beurg et al., 2018). The work presented here reveals that
tonotopy is also associated with gradients of intrinsic mechanical properties of the hair cell's
tip-link complex. Specifically, by dissecting the relative contributions of the tip links and of
the stereociliary pivots to the micromechanical properties of the hair bundle, we found that the
gating springs that control the open probability of the mechanoelectrical transduction channels
are stiffer (Fig. 3D) and subjected to higher mechanical tension (Fig. 5C) in hair cells that
respond to higher characteristic frequencies. In return, our data raises the possibility of a

mechanical role of the tip-link complex in the process that sets the characteristic frequency of the hair cell.

The stiffness K_{HB} of the whole hair bundle displayed steeper gradients than those expected using the rough estimate $K_{HB} \propto N_{SP}/h^2$ from morphological changes (Figure 3–figure supplement 2) in length h and number of stereocilia N_{SP}. Computing the stiffness ratio between the most basal and the most apical cochlear location that we were able to probe, the measured stiffness ratios (Fig. 2E) were ~50% and ~70% larger than those expected from morphology for outer and inner hair cells, respectively. We interpret this result as the consequence of intrinsic gradients of the single gating-spring stiffness (Fig. 3D). Further emphasizing mechanical regulation at the level of the tip-link complex, we also observed that the rotational stiffness of a single stereocilium was nearly uniform across the cochlear locations that we tested, especially in outer hair cells (Fig. 3C). Stiffness gradients of hair bundles with disrupted tip links are thus entirely determined by morphology, in contradistinction to those observed with intact hair bundles.

Role of hair-bundle maturation in mechanical gradients. Our experiments were performed with hair cells from juvenile animals (P7-P10), before the onset of hearing. Hair-cell maturation progresses from base to apex in the cochlea (Wu and Kelley, 2012), which may thus have affected our estimates of mechanical gradients of the tip-link complex. However, because 92% percent of our recordings were performed at P8 or later (Methods), the tip-link complex ought to be nearly mature in our experiments, at least in outer hair cells (Roth and Bruns, 1992; Waguespack et al., 2007; Beurg et al., 2018). In inner hair cells, we cannot exclude that maturation of the hair-bundle morphology was still proceeding at the most apical cochlear positions explored in our study (Peng et al., 2009). Maturation sharpens the apex-to-base gradient of bundle length (Roth and Bruns, 1992); based on bundle morphology only, we would expect to underestimate stiffness gradients with immature inner hair cells.

Within the tip-link complex, transmembrane channel-like protein isoforms 1 and 2 (TMC1 and TMC2) are thought to be essential components of the transduction channels (Fettiplace and Kim, 2014; Pan et al., 2018). TMC2 is only transiently expressed after birth; TMC1 is expressed later than TMC2 but is fundamental to mechanoelectrical transduction of mature cochlear hair cells (Kawashima et al., 2011; Kim and Fettiplace, 2013). Stereociliary expression levels of TMC1, as well as their tonotopic gradients, were recently shown in mice to be nearly mature by P6, both in inner and outer hair cells, but TMC2 may still be present in apical inner hair cells until about P13 (Beurg et al., 2018). Because TMC2 confers larger

327 Ca^{2+} permeability to the transduction channels (Kim and Fettiplace, 2013), the Ca^{2+} influx at

328 rest in the inner hair cells at P8-P10 may have been larger than at more mature developmental

329 ages, possibly lowering tip-link tension (Fig. 6) and steepening its gradient.

How stiffness gradients may contribute to the tonotopic map. We observed that the hair-

bundle stiffness increased by 240% over two octaves (1–4 kHz) of characteristic frequencies

for outer hair cells and a similar increase but over 4 octaves (1–15 kHz) for inner hair cells

(Fig. 2E). Whether or not stiffness would continue increasing along the same gradient

towards more basal locations of the cochlea is unknown. If it were the case, we would expect

a base-to-apex stiffness ratio of ~40 for outer hair cells, which is comparable to the base-to-

apex ratio of characteristic frequencies in the rat cochlea (range: 0.5-50 kHz; (Viberg and

Canlon, 2004)), but only of ~6 for inner hair cells. The interplay between the stiffness and

mass of a hair bundle could in principle help specify the preferred frequency of vibration of

the hair cell through passive mechanical resonance with sound stimuli (Frishkopf and

DeRosier, 1983; Holton and Hudspeth, 1983; Manley et al., 1988; Freeman and Weiss, 1990;

Gummer et al., 1996). The resonance frequency $\omega_C = \sqrt{k/m}$ of a spring-mass system is

given by the square root of the system's stiffness k divided by the mass m; it thus increases

with stiffness, but relatively slowly. Assuming for simplicity that the bundle's mass remains

nearly the same along the tonotopic axis (Tilney and Tilney, 1988), two orders of magnitude

in frequency must be produced by a 10,000-fold increase in stiffness, corresponding to much

steeper gradients than those reported here.

Alternatively, it has been proposed that the hair bundle could actively resonate with sound

as the result of spontaneous oscillations (Martin et al., 2001; Hudspeth, 2008). Within this

framework, the characteristic frequency is set by the frequency of the oscillator, which is

expected to increase with the stiffness of the hair bundle (Vilfan and Duke, 2003; Tinevez et

al., 2007; Martin, 2008; Barral et al., 2018). Notably, the relation may be steeper than that

resulting from a passive spring-mass system, possibly approximating a linear dependence

(Hudspeth et al., 2010). In this case, the stiffness gradient observed here (Fig. 2E) for outer

hair cells, but not for inner hair cells, could be steep enough to be a major determinant of the

tonotopic map.

Functional role of tension gradients. Tip-link tension is thought to control the open

357 probability of the transduction channels, with higher tension promoting opening of the

358 channels (Hudspeth and Gillespie, 1994). On this basis, a gradient of tip-link tension

(Fig. 5C) ought to result in a gradient of open probability. Yet, it has been shown in outer hair cells that the channels' open probability –the operating point of the transducer– remains remarkably uniform along the tonotopic axis, near a value of ½ (Johnson et al., 2011). To explain this observation, we note that the tension gradient for outer hair cells is associated with a gradient of single-channel conductance (Beurg et al., 2006; Beurg et al., 2015; Beurg et al., 2018). As a consequence, the magnitude of the Ca^{2+} influx into transducing stereocilia is expected to increase with the characteristic frequency of the hair cell. Manipulations that affect the extracellular or the intracellular Ca^{2+} concentration indicate that the transduction channels close at increased Ca^{2+} concentrations (reviewed in (Fettiplace and Kim, 2014)), possibly because the channels are harder to open when the Ca^{2+} concentration is higher near the channel's pore (Cheung and Corey, 2006). Thus, the gradient of tip-link tension reported here (Fig. 5C) may compensate for the effects of the conductance gradient on the open probability: channels with higher conductance impart higher Ca^{2+} influxes (closing the channels) but are also subjected to higher tension (opening the channels), perhaps maintaining an optimal operating point for the transducer at all cochlear locations.

Tension in the tip links is thought to be produced actively by pulling forces from molecular motors interacting with the actin core of the stereocilia at the upper insertion point of the tip link (Gillespie and Muller, 2009). The observed tension gradient in turn implies that, towards basal cochlear locations, there are more motors or that each motor exerts higher forces than near the apex. Notably, the tip links of inner-hair-cell bundles were found to bear less tension than those of outer hair cell (Fig. 5B-C). This property qualitatively makes sense, for the open probability of the transduction channels is thought to be smaller in inner hair cells than in outer hair cells (Russell and Sellick, 1983). There is also no, or only a weak, gradient of the single-channel conductance in inner hair cells (Beurg et al., 2006; Beurg et al., 2018), which parallels the relatively weak gradient of tip-link tension observed here.

Tip-link tension may be high enough to alter tip-link conformation and affect gating-spring stiffness. The tip link is composed of the association of two cadherin-related proteins, cadherin-23 and protocadherin-15 (PCDH15) (Kazmierczak et al., 2007). Molecular dynamics simulations have suggested that a bend between extracellular cadherin (EC) repeats 9 and 10 of PCDH15 may confer some compliance to otherwise rigid tip links (Araya-Secchi et al., 2016). Tensions higher than ~10 pN are predicted to evoke complete unbending of EC9-10, resulting in significant stiffening of the tip link. Assuming that PCDH15 in the tip link forms a dimer (Kazmierczak et al., 2007; Ge et al., 2018) and that tip-link tension is

equally shared by the two filaments, our estimates of tip-link tension (Fig. 5C) are compatible
with a contribution of the bending elasticity of EC9-10 to gating-spring stiffness at the apex of
the rat cochlea, especially in inner hair cells. In outer hair cells, as one progresses from the
very apex towards more basal cochlear locations, tension may quickly become too high to
allow a bent conformation in EC9-10. At the 4-kHz location, we estimated a resting tip-link
tension of ~35 pN. Taking the measured unfolding forces of Ig domains in titin as a reference
(Rief et al., 1997), tip-link tension might actually be high enough to evoke unfolding of EC
domains, at least under resting conditions or at physiological loading rates. Recent evidence
suggests that unfolding a various number of EC domains may contribute to a gradation of
gating-spring stiffness (Bartsch et al., 2018; Bartsch and Hudspeth, 2018).

Notably, the estimated gradients of gating-spring tension (Fig. 5C) were associated with
gradients of gating-spring stiffness (Fig. 3D): stiffer gating springs are subjected to more
resting tension. Strain stiffening is a common phenomenon associated with the entropic
elasticity of macromolecules, including the tip-link component PCDH15 (Bartsch et al.,
2018), as well as with filamentous protein networks (Bustamante et al., 1994; Rief et al.,
1997; Kang et al., 2009). A tension gradient may thus in part explain the existence of the
observed gradient of gating-spring stiffness. Alternatively, the gating-spring stiffness could
vary if the gating spring were composed of a variable number of compliant molecules
operating in parallel and connected to a single tip link.

Tip-link tension depends on calcium. Upon iontophoretic application of a Ca^{2+} chelator
(EDTA), before tip-link disruption, we observed that the hair bundle first moved in the
negative direction and that this movement was associated with a concomitant opening of the
transduction channels (Fig. 4). Calcium acts as a permeant channel blocker of the
transduction channels (Fettiplace and Kim, 2014). Lowering the extracellular Ca^{2+}
concentration is thus expected to increase the magnitude of the current flowing through open
transduction channels but not to produce hair-bundle movements, at least as the result of
block release only. A decrease of the extracellular Ca^{2+} concentration also promotes opening
of the transduction channels (Hacohen et al., 1989; Johnson et al., 2011). Within the
framework of the gating-spring model of mechanoelectrical transduction, channel opening
must reduce gating-spring extension and in turn tension, fostering *positive* movements of the
hair bundle. Thus, the observed *negative* movements cannot result from internal forces
associated with channel gating. Instead, our observations are readily explained if the evoked
reduction of extracellular Ca^{2+} concentration resulted in an increase of tip-link (and thus

gating-spring) tension. If tip-link tension at rest is set by myosin molecular motors that pull

on the tip links (Hudspeth and Gillespie, 1994), then the motor force must increase at

decreased Ca^{2+} concentrations.

Alternatively, a negative deflection would also be produced if lowering the extracellular

Ca^{2+} concentration evoked stiffening (Martin et al., 2003; Beurg et al., 2008) or shortening of

the gating springs, increasing tension in the tip links. However, this mechanism would result

in a negative steady-state offset of the bundle's resting position (Fig. 6E) only if there were no

or little myosin-based adaptation –usually called 'slow adaptation'(Hudspeth and Gillespie,

1994; Peng et al., 2013)– to relax the tension change and return the hair bundle near its initial

resting position. Because we observed the mechanical correlate of slow adaptation in

response to force steps (Figure 2−figure supplement 5), a Ca^{2+}-dependent regulation of tip-

link tension via a change of gating-spring stiffness or length appears unlikely. This inference

may be relevant to recent data implicating the stereociliary membrane as a key regulator of

transduction-channel gating. Calcium ions have indeed been proposed to interact

extracellularly with the local lipid environment of the transduction channels, promoting lower

values of the channels' open probability (Peng et al., 2016). Our recordings (Fig. 4B and

Fig. 6D-E) are consistent with this hypothesis but a change in membrane mechanics wouldn't

explain the observed steady-state increase in tip-link tension (Fig. 6E) if slow adaptation

happens in these cells (Figure2− figure supplement 5).

Interestingly, depolarization of rat outer hair cells was previously shown to evoke positive

movements of the hair bundle (Kennedy et al., 2006). Both depolarization and chelation of

extracellular Ca^{2+} are expected to reduce the intracellular Ca^{2+} concentration in the vicinity of

the transduction channel's pore. Yet, the directionality of active hair-bundle movements is

opposite in the two studies, suggesting that the hair bundle can operate in two regimes

(Tinevez et al., 2007). In the first regime (Kennedy et al., 2006), the response to Ca^{2+} changes

is dominated by gating forces (Howard and Hudspeth, 1988) so that the resting tension in the

tip links is nearly the same before and after application of the stimulus. In the other regime

(our study), Ca^{2+}-evoked changes of the resting tension in the tip links (Fig. 6) dominate

gating forces. In the chicken cochlea, depolarization of the hair cell was reported to evoke

negative movements of the hair bundle (Beurg et al., 2013), a directionality in agreement with

that found here (Fig. 4A). In addition, it has been shown in the bullfrog's sacculus (Tinevez

et al., 2007) and the turtle's cochlea (Ricci et al., 2002) that the response of different hair cells

to a given Ca^{2+} change can be of either directionality and that the directionality of the

response for a given hair cell can even be reversed by applying a position offset to the hair

bundle. The two regimes of active hair-bundle motility can thus potentially coexist within the same hair cell, but only if gating forces are strong enough (Tinevez et al., 2007). We measured force-displacement relations that were remarkably linear (Fig. 2B and D), showing no sign of gating compliance (Howard and Hudspeth, 1988). This observation confirms that gating forces were relatively weak under our experimental conditions, although others have shown that gating compliance can be measured with mammalian cochlear hair cells (Russell et al., 1992; Kennedy et al., 2005).

Mechanical gradients reflect the division of labor between inner and outer hair cells. Stiffness (Fig. 2E) and tension (Fig. 5B) gradients were steeper for outer hair cells, which serve primarily as mechanical amplifiers of sound-evoked vibrations, than for inner hair cells, the true sensors of the inner ear (Hudspeth, 2014). Other properties, such as the height of the hair bundle (Wright, 1984; Lim, 1986; Roth and Bruns, 1992) or the conductance of the transduction channels (Beurg et al., 2006; Beurg et al., 2018), show a similar behavior. Thus, the division of labor between inner and outer hair cells may impart more stringent regulatory constrains to outer hair cells to tune their mechanoreceptive antenna according to the local characteristic frequency of the cochlear partition. However, the exact contribution of the hair bundle to frequency tuning remains unsure and, more generally, the mechanism that specifies the characteristic frequency remains a fundamental problem in auditory physiology. This may be in part because frequency selectivity cannot be ascribed to one element only, for instance the passive resonant property of the basilar membrane that was characterized in the pioneering work of von Bekesy (Von Békésy and Wever, 1960). Various models of cochlear mechanics instead indicate that the characteristic frequency emerges from an active dynamic interplay between somatic electromotility of outer hair cells (Ashmore, 2008) and the micromechanical environment, including the basilar and tectorial membranes ,as well as the hair bundle (Nobili and Mammano, 1996; Hudspeth et al., 2010; Ó Maoiléidigh and Jülicher, 2010; Meaud and Grosh, 2011; Hudspeth, 2014; Reichenbach and Hudspeth, 2014). Mechanical tuning of the inner constituents of the cochlear partition appears to happen at many scales: from the mesoscopic scale of the basilar and tectorial membranes, to the cellular scale of the hair bundle and hair-cell soma, down to the molecular scale of the hair cell's transduction apparatus. Our work demonstrates that tonotopy is associated, in addition to other factors, with stiffness and tension gradients of the tip-link complex.

METHODS

Experimental preparation.   All experimental procedures were approved by the Ethics committee on animal experimentation of the Institut Curie; they complied with the European and French-National Regulation for the Protection of Vertebrate Animals used for Experimental and other Scientific Purposes (Directive 2010/63; French Decree 2013-118). Experiments were performed on excised cochlear coils of Sprague Dawley rats (Janvier Labs) between postnatal day 7 and 10 (P7–P10), with 8% of the cells at P7, 75% at P8–P9 and 17% at P10. The dissection of the cochlea followed a published procedure (Kennedy et al., 2003). In short, we cracked open the bony shell covering the cochlear tissue, unwound the cochlear tube from the modiolus, removed the stria vascularis, and gently peeled the tectorial membrane. Apical or middle turns of the organ of Corti were positioned under strands of nylon fibres in the experimental chamber. We recorded from inner hair cells at 4 positions along the longitudinal axis of the cochlea (Fig. 1A), corresponding to fractional distances of 5%, 10%, 20%, and 50% from the cochlear apex. According to the tonotopic map in this species (Viberg and Canlon, 2004), these cells were tuned at characteristic frequencies of 1, 2, 4, and 15 kHz, respectively. We also recorded from outer hair cells but only at the first three positions along the tonotopic axis. We have attempted to record from outer hair cells farther towards the cochlear base, in particular at the 7 kHz location. However, our success rate was too low to get reliable estimates of tip-link tension and gating-spring stiffness, the primary goal of our work. Among possible reasons, basal outer hair cells might not withstand immersion in low-Ca^{2+} (~20 µM) saline, a condition that we used to disrupt the tip links. In addition, the hair bundles of outer hair cells are smaller towards the base; monitoring their movements from direct imaging on photodiodes (see 'Microscopic apparatus' below) thus gets harder because the contrast of the image is lower.

The tissue was bathed in a standard saline containing 150 mM NaCl, 6 mM KCl, 1.5 mM $CaCl_2$, 2 mM Na-pyruvate, 8 mM glucose and 10 mM Na-HEPES. In some experiments, we used a low-Ca^{2+} saline containing 150 mM NaCl, 6 mM KCl, 3.3 mM $CaCl_2$, 4 mM HEDTA, 2 mM Na-pyruvate, 8 mM glucose, and 10 mM Na-HEPES. As measured with a Ca^{2+}-sensitive electrode, this solution had a free Ca^{2+} concentration of 22 µM, similar to that found in rat endolymph (Bosher and Warren, 1978). All solutions had a pH of 7.4 and an osmotic strength of 315 mOsm·kg^{-1}. Experiments were performed at a room temperature of 20–25°C.

Microscopic Apparatus.   The preparation was viewed through a ×60 water-immersion objective of an upright microscope (BX51WI, Olympus). The tips of individual hair bundles

were imaged at a magnification of ×1,000 onto a displacement monitor that included a dual photodiode. Calibration was performed before each recording by measuring the output voltages of the monitor in response to a series of offset displacements of the photodiode. For hair-bundle movements that did not exceed ±150 nm in the sample plane, the displacement monitor was linear.

Iontophoresis of a Ca^{2+} Chelator. We used iontophoresis to apply the calcium chelator EDTA in the vicinity of a hair bundle (Fig. 1B) and disrupt its tip links (Assad et al., 1991; Jaramillo and Hudspeth, 1993; Marquis and Hudspeth, 1997). Coarse microelectrodes were fabricated from borosilicate capillaries with a pipette puller (P97, Sutter Instrument); their resistance was 1 MΩ when filled with 3 M KCl and immersed in the same solution. In experiments, the electrodes were filled with a solution containing 100 mM EDTA and 25 mM KCl. The electrode's tip was positioned at ∼3 μm from the hair bundle. A holding current of +10 nA was continuously applied to counteract the diffusive release of EDTA from the electrode. The stimulus consisted of a −100-nA current step on top of the holding current, resulting in a net iontophoretic current of −90 nA. To facilitate tip-link disruption upon EDTA iontophoresis, the cochlear tissues were immersed in low-Ca^{2+} saline (~20-μM Ca^{2+}).

Mechanical stimulation and stiffness measurements. The hair bundles of inner and outer hair cells were mechanically stimulated using a fluid-jet device (Kros et al., 1992; Géléoc et al., 1997; Johnson et al., 2011). Pipettes were pulled from borosilicate glass (TW150-F, World Precision Instruments); their tip diameter was adjusted within a range of 5−10 μm. Fluid flow through a pipette was driven by a voltage command to a piezoelectric disk (Murata 7BB-27-4). Any steady-state flow coming in or out of the pipette was nulled by changing the hydrodynamic pressure inside the fluid-jet pipette; the hydrodynamic pressure was adjusted with a syringe connected to the body of the fluid-jet device. The fluid-jet pipette was positioned on the abneural side of the bundle along the hair bundle's axis of mirror symmetry (Fig. 1C). Fluid coming out the pipette thus deflected the hair bundles towards the shortest stereociliary row, closing the ion channels that mediate mechanoelectrical transduction. This direction of bundle movement is defined as the negative direction in this paper; conversely, positive movements were directed towards the tallest row of stereocilia, fostering opening of the transduction channels. Mechanical stimuli were applied as 100-ms paired-pulse steps (Fig. 2 and Figure 2−figure supplement 3), or 60-Hz sinusoids (Fig. 4) with the magnitude of driving voltages varying between 0 and 60 V.

For stiffness measurements, we measured hair-bundle movements evoked by 100-ms force

steps (Fig. 2; see the force-calibration procedure below).   The bundle displacement was

measured 5–10 ms after the onset of the step stimulus; the stiffness was given by the slope of

the relation between the force (noted F in the following) and the displacement of the bundle's

tip. These measurements were performed in standard saline.

Applying and measuring forces with the fluid jet.   We describe here how we calibrated the

hydrodynamic drag force F_D applied to the hair bundle by a fluid jet by using a flexible glass

fiber of known stiffness as a reference.   The method is based on a published procedure

(Géléoc et al., 1997) that we refined to account for the non-uniform velocity field of the fluid

(Figure 2−figure supplement 1). Using a generalized Stokes equation (Leith, 1987), the drag

force can be written as $F_D = 6\pi\eta R_{HB}\, U$, in which η is the viscosity of the surrounding fluid

and R_{HB} is the effective hydrodynamic radius of the bundle. The effective radius R_{HB} was

approximated by that of a prolate ellipsoid of short axis h and long axis W, which correspond

to the bundle's height and width, respectively. For a fluid flow perpendicular to the axis of

rotational symmetry of the ellipsoid, this yields:

$$R_{HB} \cong 4h \Big/ \left\{ 3 \left[\phi/(\phi^2 - 1) + \left((2\phi^2 - 3) \ln\left(\phi + \sqrt{\phi^2 - 1}\right) \right)/(\phi^2 - 1)^{3/2} \right] \right\}, \quad (1)$$

in which $\phi = W/h$ represents the aspect ratio of the ellipsoid (Happel and Brenner, 2012).

We note that the hydrodynamic radius given by Equation 1 is exact only for an ellipsoid

immersed in an infinite volume of fluid, whereas the hair bundle instead stands erect at the

apical surface of the hair cell; this expression is thus an approximation. Figure 3−supplement

table 1 and Figure 3−figure supplement 2 recapitulate the values of parameters h and W that

we used to model inner and outer hair-cell bundles along the tonotopic axis of the rat cochlea,

as well as the resulting values of R_{HB}. The effective velocity $U \cong \int_{-W/2}^{W/2} v_X(x,y)\, dy / W$ was

estimated by computing the mean of the velocity field $v_X(x,y)$ of the fluid over the width W

of the hair bundle. Here, $v_X(x,y) = \vec{v} \cdot \vec{e}_X$ is the projection of the fluid velocity $\vec{v}$ on the axis

of mechanosensitivity (axis X) of the hair bundle; its value is estimated along the axis (Y)

perpendicular to axis X for a bundle positioned at a distance x from the mouth of the fluid-jet

pipette (Figure 2−figure supplement 1A). Using bead tracers, we found that the velocity

profile $v_X(x,y)$ obeyed (Schlichting, 1933)

$$v_X(x,y) = V_{\max}(x)/(1 + (y/A(x))^2)^2, \quad (2)$$

where $V_{\max}(x)$ and $A(x)$ characterize, respectively, the maximal speed and the lateral extension of the velocity field at position x (Figure 2–Figure supplement 1B-D). By integrating the velocity profile, we obtain an expression for the force

$$F_D = 6\pi\eta R_{\mathrm{HB}}\,\beta_{\mathrm{HB}}\,V_{\max}\,, \qquad (3)$$

where $\beta_{\mathrm{HB}} = \beta(w) = \frac{1}{2w}(w/(1+w^2) + \tan^{-1}w)$ is a constant that depends on the normalized width of the hair bundle $w = W/(2A)$. Thus, calibrating the force F_D is equivalent to calibrating the maximal fluid velocity $V_{\max}$.

To estimate $V_{\max}$, we measured the force $\bar{F}_D \cong 6\pi\eta R_{\mathrm{F}}\,\bar{U}$ applied by the same jet on a calibrated glass fiber, whose longitudinal axis was oriented perpendicularly to that of the fluid-jet pipette. Given the diameter D_{F} of the fiber, the effective hydrodynamic radius of a cylindrical fiber was calculated as $R_{\mathrm{F}} = 2L/[3(\ln(L/D_{\mathrm{F}}) + 0.84)]$ (Tirado and Torre, 1979). Because the conical fluid jet intersected the fiber over a length $L > W$, the effective fluid velocity $\bar{U} \cong \int_{-L/2}^{+L/2} v_X(x,y)\,dy/L = \beta_{\mathrm{F}}\,V_{\max}$ for the fiber was smaller than the effective velocity U for the hair bundle, where $\beta_{\mathrm{F}} = \beta(L/(2A)) < \beta_{\mathrm{HB}}$. In practice, we used $L(x) = 2\,x\,\tan\alpha + D_{\mathrm{FJ}}$, where α is the half-aperture of the conical fluid jet that was visualized using a dye (Coomassie Brilliant Blue; Figure 2–figure supplement 2) and D_{FJ} is the diameter of the mouth of the fluid-jet pipette. We noticed that $L(x) \cong 2A(x)$ (Figure 2–figure supplement 1; figure supplement 2). We used this property to estimate $\beta_{\mathrm{HB}} \cong \beta(W/L)$ and $\beta_{\mathrm{F}} \cong \beta(1)$ without having to measure A directly in every experiment.

In experiments, the projected horizontal distance between the tip of the fluid-jet pipette and the hair bundle or the fiber was fixed at $x \cong 8\,\mu\mathrm{m}$ (mean ± SD: 7.8±0.6 µm; range: 5.9–8.8 µm). Flexible fibers of diameters $D_{\mathrm{F}} = 0.7$–1.5 µm and stiffness $k_{\mathrm{F}} = 0.2$–2 mN/m were fabricated and calibrated as described before (Bormuth et al., 2014); their effective hydrodynamic radii varied within a range of $R_{\mathrm{F}} = 2.5$–3.2 µm. A fluid jet of given magnitude elicited a force $\bar{F}_D = k_{\mathrm{F}}\,\Delta X$, where ΔX is the measured deflection of the fiber. The relation between the force $\bar{F}$ applied to a fiber and the voltage command to the fluid-jet device was linear; its slope provided the calibration constant C (Figure 2–figure supplement 3). When stimulating a hair bundle, a voltage command V_C to the fluid-jet device thus elicited a force $F_D \cong GCV_C$, where $G = F_D/\bar{F}_D = (\beta_{\mathrm{HB}}R_{\mathrm{HB}})/(\beta_{\mathrm{F}}R_{\mathrm{F}})$. We used $G = 1.4\pm0.1$ (mean ± SD; range: 1.27–1.65) for inner hair cells and $G = 1.3\pm0.1$ (mean ± SD; range: 1.12–1.47) for outer hair cells. Thus, we estimate that the force applied on the hair bundle was 30–40% higher than that measured on the calibration fiber using the same jet of fluid. In practice,

we calculated G in each experiment from the geometrical parameters of the fluid-jet pipette, the calibration fiber, and the hair bundle. We noticed that $L(x) \cong 2A(x)$; at a distance $y = L/2$ from the center of the fluid jet ($y = 0$), the fluid velocity is expected to be 25% of the maximal value. Thus, some of the moving fluid was not taken into account in our estimates of the force acting on the fiber, resulting in an underestimation. However, taking $L(x) = 4A(x)$ resulted in an increase of G by only 5% while the fluid velocity dropped to 4% of the maximal value at the edge of the fluid cone (in $y = L/2 = 2A$).

All the forces reported in this work correspond to the effective force F that one would have to apply at the bundle's tip to evoke the same bundle deflection X as the hydrodynamic drag force F_D actually exerted by the fluid jet and distributed over the height of the hair bundle. Correspondingly, the stiffness of the hair bundle was calculated as the slope of the force-displacement relation $F(X)$. In the following, we develop an approximate description of hair-bundle mechanics to relate the effective force F to the drag force F_D. Modelling the hair bundle as a pivoting solid (Kozlov et al., 2007; Kozlov et al., 2011), we write the torque produced by the fluid jet as $\Gamma = \int_0^h f(z)\, z\, dz = \kappa_{HB}\, \theta$. Here, h is the height of the hair bundle, f is the drag force per unit (vertical) length experienced by the hair bundle at a vertical distance z from the apical surface of the hair cell, κ_{HB} is the rotational stiffness of the hair bundle, and $\theta \cong X/h$ is the pivoting angle corresponding to the deflection X measured at the bundle's top. With these notations, the total drag force on the hair bundle reads $F_D = \int_0^h f(z)\, dz$. Assuming that the drag force per unit length $f \propto v(z)$ is proportional to the local fluid velocity $v(z)$ and that the condition of no-slip at the apical surface of the hair cell ($v(z = 0) = 0$) imposes a linear velocity profile $v(z) = v(h)\, z/h$, we take $f \propto z$. Injecting this expression into the integrals defining the torque Γ and the total drag force F_D, we find $\Gamma = \frac{2}{3} F_D h$. As a result, we identify the effective force $F = \frac{2}{3} F_D$ that must be applied at the bundle's top to evoke the same torque as the fluid jet; the effective force F is smaller than the total drag force F_D exerted by the fluid jet, corresponding to a ratio $F/F_D = 2/3$.

To test our force-calibration procedure for the fluid jet, we performed a control experiment using hair-cell bundles from the sacculus of the frog (strain "Rivan92" of *Rana ridibunda* (Neveu, 2009)). Details of the experimental preparation have been published elsewhere (Tinevez et al., 2007). In a similar preparation, the hair bundles have been shown to be very cohesive (Kozlov et al., 2007). As a result, the stereocilia are constrained to move as a unit in response to point-force application at the kinociliary bulb with a flexible fiber, allowing for

well-controlled stiffness measurements (Martin et al., 2000; Bormuth et al., 2014). We compared stiffness estimates from stimulation with fluid jets and flexible fibers (Figure 2–figure supplement 4). From a sample of 13 hair bundles, we found a stiffness of 0.40±0.17 mN/m (mean±SD) using fluid-jet stimulation and of 0.42±0.16 mN/m (mean±SD) using fiber stimulation; the mean values of the stiffness estimates were not statistically different (paired-sample t-test; ***p<0.001). The good agreement between the two methods gave us confidence that the procedure that we used to calibrate the fluid jet is satisfactory.

Electrophysiological recordings. We used the patch-clamp technique to measure mechano-electrical transduction currents. Borosilicate patch pipettes were filled with an intracellular solution containing 142 mM CsCl, 3.5 mM $MgCl_2$, 1 mM EGTA, 5 mM Na_2-ATP, 0.5 mM Na_2-GTP and 10 mM HEPES (pH=7.3, 295 mOsmol/kg). When immersed in standard saline, these pipettes had a resistance of 1.5–3 MΩ. A patch pipette was inserted in the organ of Corti through a pre-formed hole in the reticular lamina and approached parallel to the hair cell rows towards the soma of a target hair cell. During the approach, standard saline was abundantly perfused to protect the Ca^{2+}-sensitive tip-links from EGTA. Hair cells were whole-cell voltage clamped at a holding potential of −80 mV; transduction currents were low-pass filtered at 1-10 kHz (Axopatch 200B; Axon Instruments). No correction was made for the liquid-junction potential. The series resistance was always below 10 MΩ and was compensated up to 70%. To disrupt the tip links with EDTA iontophoresis, the solution bathing the cells was changed to low-Ca^{2+} saline after the cell was patched; the solution change was performed either with a perfusion or with a Picospritzer (Picospritzer III, Parker).

Scanning electron microscopy. Cochleae from P8 rats were processed with osmium tetroxide/thiocarbohydrazide, as previously described (Furness et al., 2008). Samples were analysed by field emission scanning electron microscopy operated at 5 kV (Jeol JSM6700F). The number of stereocilia in inner and outer hair-cell bundles was estimated from electron micrographs at each of the cochlear locations where we performed mechanical and electrophysiological measurements (Figure 3–figure supplement 2; Figure 3−supplement table 1).

Estimating the number of intact tip links in a hair bundle. We performed patch-clamp recordings of the transduction current I_{MAX} elicited at saturation by large hair-bundle deflections (Figure 3–supplement figure 3). In inner hair cells, the number of intact tip links $N_{TL} = I_{MAX}/I_1$ was calculated by dividing the saturated current I_{MAX} for the whole hair bundle by the published estimate $I_1 = 35.4$ pA for the transduction current flowing through

the tip of a single transducing stereocilium (Beurg et al., 2009); electron microscopy has

indeed shown that there is precisely one tip link per stereocilium in an intact hair bundle

(Pickles et al., 1984; Kachar et al., 2000). We used the same value of I_1 at all cochlear

locations (Beurg et al., 2006; Beurg et al., 2018). Given the magnitude $i = 15$ pA of the

current flowing through a single transduction channel, there was on average $I_1/i = 2.36$

transduction channels per transducing stereocilium (Beurg et al., 2006). In outer hair cells,

there is no direct estimate of I_1. However, the unitary current i was shown to increase (Beurg

et al., 2006; Beurg et al., 2015; Beurg et al., 2018) from 8.3 pA to 12.1 pA when the hair

cell's characteristic frequency increases from 4 kHz to 14 kHz (Beurg et al., 2006). All these

currents were measured under the same experimental conditions as ours, in particular using a

−80 mV holding potential and with the hair cells immersed in a standard saline containing

1.5 mM Ca^{2+}. Assuming a linear relation between the unitary current and the position of the

hair cell along the tonotopic axis of the cochlea (Beurg et al., 2015; Beurg et al., 2018), we

inferred the unitary currents at other cochlear locations. We then assumed that the average

number of transduction channels per tip link was 2.36, as estimated in inner hair cells (Beurg

et al., 2009). The number of intact tip links was then calculated as $I_{MAX}/(2.36\,i)$. We

performed this measurement for 10 hair cells at each cochlear location, both for inner and

outer hair cells, to calculate the average number of intact tip links in any given hair cell. In

these experiments, the hair cells were immersed in standard saline.

Recent measurements in the mouse cochlea have revealed that unitary currents may

represent an ensemble average over multiple conductance states, raising the possibility that

these currents are produced by a few (up to 5) identical transduction channels that gate

cooperatively (Beurg et al., 2018). This finding does not affect our estimates, because the

current that flows through a single stereocilium stays the same, whether or not it results from

cooperative gating of multiple channels or from gating of an effective channel endowed with

the same conductance as the total conductance of the group.

Signal Generation and Acquisition. All signals were generated and acquired under the

control of a computer running a user interface programmed with LabVIEW software (version

2011, National Instruments). Command signals were produced by a 16-bit interface card at a

sampling rate of 25 kHz (PCI-6733, National Instruments). A second interface card (PCI-

6250, National Instruments) conducted signal acquisition with a precision of 16 bits.

Sampling rates for signal generation and acquisition varied within the range 2.5−25 kHz. All

signals were conditioned with an eight-pole Bessel antialiasing filter adjusted to a low-pass

half-power frequency at half the sampling rate of signal acquisition.

Statistical significance. Unless otherwise noted, all results are quoted as mean ± standard

error of the mean (*n*) with a number *n* of cells of at least 10 per group. G-Power analysis

ensured that this number was sufficient to achieve a signal-to-noise ratio of 1–1.5, with 80%

power at a 5% significance level. We performed a one-way ANOVA to assay statistical

significance of the measured mean-value variation of a given property, *e.g.* the hair-bundle

stiffness, between the different cochlear locations for inner (IHC) or outer (OHC) hair cells.

We also used two-tailed unpaired Student's *t*-tests with Welch's correction when comparing

mean values between two groups of a given hair-cell type (IHC or OHC) with different

characteristic frequencies or between the two cell types (IHC/OHC) with a given

characteristic frequency. Stars correspond to p-values with *p < 0.05, **p < 0.01, and

***p < 0.001, whereas 'n.s.' (p > 0.05) indicates non-significant differences. To determine

whether variables estimated from the product of M independent variables X_i ($i = 1 .. M$) had

means that were statistically different, we first calculated the standard error of the mean

$\sigma_{\prod X_i}$ of the product and the effective number of degrees of freedom ν_{eff} of the product.

Defining $\overline{X_i}$ as the mean value of the variable X_i over n_i measurements, s_i the standard

deviation and $\sigma_i = s_i/\sqrt{n_i}$ the standard error of the mean, the standard error of the mean for

the product was calculated as $\sigma_{\prod X_i} = \prod \overline{X_i} \sqrt{\sum \left(\frac{\sigma_i}{\overline{X_i}}\right)^2}$ and the effective number of degrees of

freedom associated with the product was calculated using the Welch-Satterthwaite

approximation as $\nu_{\text{eff}} = \left[\sigma_{\prod X_i}/\prod \overline{X_i}\right]^4 / \sum \frac{1}{n_i-1}\left(\frac{\sigma_i}{\overline{X_i}}\right)^4$. Finally, we characterized tonotopic

gradients by performing weighted linear regressions, in which the weight applied to each data

point was given by the inverse of the squared standard error of the mean. We then applied a *t*-

test on the resulting coefficients to determine whether the observed difference between the

gradients measured with inner and outer hair cells was statistically significant. The results of

all statistical analyses are listed in supplement tables associated to the main figures.

Acknowledgements:

We thank Benoît Lemaire and Rémy Fert from the machine shop of the Curie Institute, Prof. Walter Marcotti for help in the design of our fluid-jet device, Jérémie Barral for a critical reading of the manuscript and Christine Petit for fruitful discussions. This research was supported by the French National Agency for Research (ANR-11-BSV5 0011 and ANR-16-CE13-0015) and by the Labex Celtisphybio ANR-10-LABX-0038 part of the Idex PSL. M.T. is alumnus of the Frontiers in Life Science PhD program of Université Paris Diderot and thanks the Fondation Agir pour l'Audition for a doctoral fellowship.

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FIGURES

Figure 1: Hair-bundle stimulation along the tonotopic axis of the rat cochlea.

(A) Schematic representation of the tonotopic axis of the rat cochlea. Recordings were made at locations marked by black disks, corresponding to characteristic frequencies (in kHz) increasing from the apex to the base of the cochlea as indicated on the figure and to fractional distances from the apex of 5%, 10%, 20%, and 50%. The rat cochlea was typically 10 mm long. Adapted from (Viberg and Canlon, 2004). **(B)** Schematic layout of the experimental pipettes around a given outer hair cell. We combined fluid-jet stimulation of single hair bundles, iontophoresis of a Ca^{2+} chelator (EDTA), patch-clamp recordings of transduction currents, and perfusion of low-Ca^{2+} saline. **(C)** Schematic representation of the fluid-jet pipette and of a hair bundle (left) and micrograph of a fluid-jet pipette ready to stimulate an outer hair cell of the rat cochlea (right). A positive (negative) deflection of the hair bundle, as defined on the drawing, was elicited by fluid suction (ejection) into (from) the pipette, promoting opening (closure) of the transduction channels. The horizontal projected distance

between the mouth of the pipette (blue vertical line) and the hair bundle (green vertical line)

was set at ~8 µm.

FIGURE 2: Stiffness gradients of the hair bundle.

(A) Hair-bundle movements (top) in response to series of force steps (bottom) for outer hair

cells (OHC) with characteristic frequencies of 1, 2 and 4 kHz (from left to right). **(B)** Force-

displacement relations for the data shown in **A**, with black disks, white disks and black

triangles corresponding to characteristic frequencies of 1, 2 and 4 kHz, respectively. **(C)**

Hair-bundle movements (top) in response to series of force steps (bottom) for inner hair cells

(IHC) with characteristic frequencies of 1, 2, 4, and 15 kHz. **(D)** Force-displacement

relations for the data shown in **C**, with black disks, white disks, black triangles, and white

squares corresponding to characteristic frequencies of 1, 2, 4, and 15 kHz, respectively. **(E)**

Stiffness (K_{HB}) of a hair bundle under control conditions as a function of the characteristic

frequency (CF) for inner (white disks) and outer (black disks) hair cells. Each data point in **E**

is the mean ± standard error of the mean (SEM) with the number of cells indicated between

brackets.

The following source data is available for figure 2:

Source data 1. Statistical significance.

Source data 2. Hair-bundle stiffness of inner and outer hair cells as a function of the

characteristic frequency.

The following figure supplements are available for figure 2:

Figure supplement 1. Velocity field of a fluid jet.

Figure supplement 2. Geometrical characteristics of a fluid jet.

Figure supplement 3. Rise time and linearity of the fluid-jet stimulus.

Figure supplement 4. Test of fluid-jet calibration in the frog's sacculus.

Figure Supplement 5. Mechanical creep during a force step.

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FIGURE 3: Stiffness gradients of the gating springs and of the stereociliary pivots.

Stiffness **(A)** of the contribution of the gating springs ($K_{GS} = r\,K_{HB}$), **(B)** of a hair bundle

after tip-link disruption, corresponding to the contribution of the stereociliary pivots ($K_{SP} =$

$(1 - r)\,K_{HB}$), **(C)** of a single stereociliary pivot ($\kappa = K_{SP}\,h^2/N_{SP}$), and **(D)** of a single gating

spring ($k_{GS} = K_{GS}\,\gamma^2/N_{TL}$) as a function of the characteristic frequency (CF) for inner (white

disks) and outer (black disks) hair cells. These stiffnesses were calculated from measured

values of the hair-bundle stiffness K_{HB} (Fig. 2), the amplitude ratio $1 - r$ of hair-bundle

movements before and after tip-link disruption (Figure 3–figure supplement 1), the hair-

bundle height h and the number of stereocilia N_{SP} (Figure 3-figure supplement 2), and the

average number N_{TL} of intact tip links (Figure 3–figure supplement 3). Each data point is the

mean ± SEM; SEMs were calculated as described in the Methods.

The following source data is available for figure 3:

Source data 1. Morphological parameters of inner and outer hair-cell bundles.

Source data 2. Statistical significance.

Source data 3. Gating-spring contribution to the hair-bundle stiffness.

Source data 4. Hair-bundle morphology along the tonotopic axis.

Source data 5. Transduction currents and number of intact tip links along the tonotopic axis.

The following figure supplements are available for figure 3:

Figure supplement 1. Gating-spring contribution to the hair-bundle stiffness.

Figure supplement 2. Hair-bundle morphology along the tonotopic axis.

Figure supplement 3. Transduction currents and number of intact tip links along the

tonotopic axis.

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FIGURE 4: **Mechanical and electrical response of a hair bundle to fluid-jet stimulation and fast calcium chelation.**

(A) An iontophoretic step of a calcium chelator (EDTA; top) elicited a biphasic movement of the hair bundle from an inner hair cell (bottom): the hair bundle first moved in the negative direction (arrow head) and then in the positive direction. After iontophoresis, the position baseline was offset by $\Delta X_R = +78$ nm with respect to the resting position at the start of the experiment. A sinusoidal command to a fluid jet (middle) evoked hair-bundle movements (bottom) that increased in magnitude, here by 50%, after application of the iontophoretic step. Repeating the iontophoretic step elicited no further movement and the response to fluid-jet stimulation remained of the same magnitude. A similar behaviour was observed with 101 inner and 44 outer hair cells. (B) An iontophoretic step of EDTA (top) also elicited biphasic variations of the transduction current: the inward current first increased (arrow head) and then decreased. Before application of the calcium chelator, fluid-jet stimulation evoked a transduction current of 1.5-nA peak-to-peak magnitude; the open probability of the transduction channels was near ½. The transduction current was abolished by the iontophoretic step. Outer hair cell at the 4-kHz location; the same behaviour was observed

1045   with 17 outer hair cells.   In **A-B**, the command signal to the fluid-jet device was a 60-Hz

1046   sinusoid and we applied a −100-nA iontophoretic step on top of a +10-nA holding current.

1047   The hair bundles were exposed to 20-µM Ca^{2+}. In **B**, the dashed line indicates the current for

1048   which the transduction channels are all closed.

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FIGURE 5: Gradients in tip-link tension at rest.

Offset ΔX_R in the resting position of a hair bundle resulting from tension release in the tip

links **(A)**, tip-link tension $T_R = K_{SP} \Delta X_R$ in the hair bundle **(B)** and tension $t_R = T_R/(\gamma\, N_{TL})$

along the oblique axis of a single tip link **(C)** as a function of the characteristic frequency

(CF) for inner (white disks) and outer (black disks) hair cells. The hair-bundle tension T_R **(B)**

was calculated as the product of the stereociliary-pivot stiffness K_{SP} shown in Fig. 3B and the

data shown in **A**; this tension is estimated along the bundle's horizontal axis of mirror

symmetry. The single tip-link tension t_R was then deduced from the projection factor γ and

the average number N_{TL} of intact tip links in a hair bundle (Figure–figure supplement 3).

Each data point in **A** is the mean ± SEM with the number of cells indicated between brackets;

in **B-C**, mean values and SEMs were calculated as described in the Methods.

The following source data is available for figure 5:

Source data 1. Statistical significance.

Source data 2. Offset in the resting position of a hair bundle upon tip-link disruption.

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FIGURE 6: Tensioning of the tip links at decreased Ca²⁺ concentrations.

The amplitude of the negative hair-bundle movement ΔX_{Ca} (**A**), of the maximal increase

$\Delta T = -K_{SP} \Delta X_{Ca}$ in hair-bundle tension (**B**), and of the maximal tension $t_{max} = t_R +$

$\Delta T/(\gamma N_{TL})$ in a single tip link (**C**) are plotted as a function of the hair cell's characteristic

frequency (CF). The tension increase in **B** was calculated from the stiffness K_{SP} of the

stereociliary pivots (Fig. 3B) and the data shown in **A**. The single tip-link tension t_{max} was

then deduced from the tension at rest t_R in a single tip link (Fig. 5C), the projection factor γ

(Figure 3–figure supplement 2) and the average number N_{TL} of intact tip links (Figure 3–

figure supplement 3). (**D**) Current-step commands (top) applied to an iontophoretic pipette

containing the Ca²⁺ chelator EDTA evoked reversible negative movements of the hair bundle

(bottom). (**E**) When the stimulus (top) was long enough, the hair bundle position could

reach a steady state (bottom), corresponding to higher resting tension in the tip links. In **A-C**,

the hair bundles were immersed in low-Ca²⁺ saline, for which EDTA iontophoresis led to tip-

link disruption. Positions and tensions were estimated at the point of polarity reversal of the

hair-bundle movement (see Fig. 4A), thus at the initiation of tip-link disruption, where the

hair bundle reached its largest deflection in the negative direction and tension was thus

maximal. Black and white disks correspond to outer and inner hair cells, respectively. The

error bars in **A** represent ±SEM with numbers of cells indicated between brackets; in **B-C**,

mean values and SEMs were calculated as described in the Methods. In **D-E**, the hair bundles

were immersed in a saline containing 500-µM Ca^{2+}; this higher Ca^{2+} concentration preserved

the integrity of the tip links upon EDTA iontophoresis.

The following source data is available for figure 6:

Source data 1. Statistical significance.

Source data 2. Negative movement ΔX_{Ca} of the hair-bundle before tip-link disruption.

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