Optimal foraging of benthic diatoms evolving novel movement behaviours

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Adaptive locomotion of living organisms contributes to their competitive abilities and

help maintain their fitness in diverse environments. To date, understanding of

searching behaviours and how microscale dynamics scale up to ecosystem-level

processes remain poorly understood in ecology. Here, we investigate the motion

patterns of the biofilm-inhabiting marine diatom *Navicula arenaria* var. *rostellata* at

the two-dimensional space. We report that individual *Navicula* cells display a novel

"rotational run-and-reversal" movement behaviour at different concentrations of

dissolved silicate (dSi). Using experimental measurements of the search behaviours,

we show that translation motions of cells can be predicted exactly with a universal

model—the generalized Langevin theoretical model. Both the experimental and

theoretical results reveal quantitative agreement with an optimal foraging strategy

and show that circular and reversal behaviours both contribute to comparable spatial

diffusion properties. Our modelling results suggest that the evolving movement

behaviours of diatom cells are driven by optimization of searching their physical

surroundings, and predicted behavioural parameters coincide with the experimental

observations. These optimized movement behaviours are an evolutionarily stable

strategy to cope with environmental complexity.

Keywords: foraging behaviour, diatoms, ESS, rotational diffusion, reversal behaviour

One sentence summary: Contrary to wide belief, new experiments reveal that diatoms

can active search food in their physical surroundings by adjusting movement behaviours.

The motion behaviours of organisms in natural world is of crucial importance to their life

cycle and survival (1-4). Plants may adapt their shape and inclining position to weaken the

competition between them, grow toward a more favourable light environment and

collectively increase production per unit land area at high density (5), while most animals

and many microorganisms can actively move from one place to another to seek favourable

habitats (6), nourishment, sexual partners (7), or to escape from predators (8,9). For

example, white blood cells are attracted by inflammatory signals (10), while bacteria move

toward higher concentrations of sugars. At a large scale, foraging animals explore their

environment to search food using their senses (2). These movements appear to be random

in space and time, whereas there exist a universal principle and a unique behaviour to cope

with a stressful environment. In past decades, numerous studies addressed movement

behaviour of living organisms, in particular whether organism's movement patterns

correspond to Brownian walks or Lévy walks (11-13). Contrary to the popular consensus

that Brownian walks characterize movement of most organisms, recent studies identified

its counterpart Lévy walks/flights as a universal movement pattern at different scales,

ranging from molecular entities (14), swimming and swarming prokaryote and eukaryote

microorganisms, foraging animals and human motion (11,15,16).

These studies have focused on unravelling the universal Lévy statistics because it

accelerates the spatial diffusive potential (17,18). However, in realistic scenarios, many

species of living organisms have some characteristic movement behaviour to cope with

their complex surroundings (19,20). The most thorough studies of motility at the cellular

scale include the swimming motions of *Escherichia coli* starting from seminal work of

Howard Berg in the 1970s (21,22). It moves in almost straight runs that are interrupted by

tumbles with an average angle of about 60°, so-called 'run-and-tumble'-like behaviour.

More recently, observations of the swimming patterns of other bacteria have revealed a

variety of different movement strategies (19, 23-24). Among them, the 'run-and-reverse'

and 'run-reverse-flick' motion are two significant discoveries in the soil bacteria

Pseudomonas putida (25), *Myxococcus xanthus* (26, 27), and the marine bacteria *Vibrio*

alginoticus (20), respectively. Along with these experimental studies, a rich theoretical

literature has addressed the different motility patterns based on concepts from the theory

of random walks. These approaches are closely linked to the newly emerging research areas

of active particles and microswimmers (28,29). At the same microscopical scales, many

species of algal cells have active motility by self-propelled flagella to involve in

translational motion and vertical migration (4,6,30). However, little is known, both from

an experimental as well as a theoretical perspective, about the impact that a given

movement behaviour may have on the spatial self-organization (31) and foraging of

microalgae (4, 32).

Here, we report experimental observations of motility patterns of the marine

benthic diatom *N. arenaria* var. *rostellata* in response to varying resource availability.

Laboratory experiments revealed two key movement behaviours, including a unique two-

step forward-reverse motion and rotation, respectively. We found that the diatom cells

regulate the forward-backward frequencies and rotational diffusion to optimize foraging

efficiency according to the ambient dSi concentration. The reversal motion is likely driven

by spraying the extracellular polymeric substances (EPS) to create a new direction from

forward to backward (33, 34). Our theoretical model exactly captures the diffusion scaling

behaviours of mean-squared displacements (MSD) and the changes of the direction of cell

movement as observed in experimental trials. We also found that, although the reversal

events and rotational diffusion appear to be random, their rates have converged to about

0.015 and 0.0054 respectively, which corresponds to the maximized efficiency of foraging.

To our knowledge, these foraging traits have not been previously discussed or appreciated

in microalgae, and it raises interesting relative questions about the roles of these movement

behaviours on the primary productivity and biofilm formations in the benthic habitats.

Results

Movement behaviours of diatom cells. Diatom motility is involved not only in searching

mating partners, but it can also direct cells towards or away from other environmental

cues[6,7]. To characterise motility in *N. arenaria* var. *rostellata*, cell movement patterns were

measured in a series of lab experiments in various dSi concentrations using exponentially

growing cultures. Cell displacements were recorded by a Ti-E Nikon phase contrast

microscope with high temporal-spatial resolutions. Trajectory analysis revealed that they

display two main movement features during foraging: (i) *reversal behaviour* with a certain

probability (v) under a fixed dSi concentration and (ii) continuous *rotation* on their turning

angles (Fig. 1a and b, and *SI Appendix* Movie S1). The significant difference of turning

behaviours suggest that diatom cells can directly response to gradient in dSi through

adjusting the frequency of reversal events (see *SI Appendix* Fig. S1). This raises questions

about the long-term implications to the reversal and rotational behaviours. To this end, we

developed a mathematical model to show that the reversal behaviour and a minor variation

in rotational noise both lead to major changes to the trajectory patterns and ultimately

foraging efficiency.

Mathematical model. We utilize a discrete time model for the self-propelled diatom cells

moving in a two-dimensional space. Reversal events are represented by an exponential

distribution process with mean intervals derived from experimental data. Inspired by

motion behaviour predictions of self-propelled rods (28), considering an Ornstein-

Uhlenbeck process and standard Wiener processes on rotational diffusion, the stochastic

equations of a single active cell are given by

96
$$x(t + \Delta t) - x(t) = V_0 \cos(\theta) \Delta t + \sqrt{2D_r} \, dW_1, \tag{1a}$$

97
$$y(t + \Delta t) - y(t) = V_0 \sin(\theta) \Delta t + \sqrt{2D_r} \, dW_2, \tag{1b}$$

98
$$\theta(t + \Delta t) - \theta(t) = \kappa(t)\omega\Delta t + \sqrt{2D_\theta} dW_3, \tag{1c}$$

99
$$\kappa(t + \Delta t) - \kappa(t) = -2\kappa(t)B, \tag{1d}$$

where $\kappa = \pm 1$ depicts the moving in CCW and CW orbits respectively. B is a Bernoulli

random variable with success probability $\nu\Delta t$ (here ν refers the frequency of the reversal

behaviours). The components of dW_1, dW_2 and dW_3 are random variables with a standard

Wiener process. The cells tend to follow paths with angular speed ω. For the evolutionarily

stable strategy analysis, up to 1000 cells are simulated with various prescribed rotational

diffusional coefficients (noise intensity) D_θ, and fixed parameter values $V_0 = 17 \, \mu m/s$,

$D_r = 0 \, \mu m^2/s$, $\kappa = \pm 1$, $\omega = \pi/36$ rad/s (5 degree), and $\nu = 0.02 \, s^{-1}$, which are all

estimated from our experiments. Note that for noncircular motion, i.e. $\omega = 0$, model (1)

define a system of persistent random walks characterized by a diffusion coefficient $D =$

$D_r + V_0^2/(2D_\theta)$ (22, 34).

To study the ensemble behaviour, we represent the configuration of diatom cells by

the probability distribution functions $\Psi_\pm(\boldsymbol{r}, \theta, t)$ (28,36), here $\boldsymbol{r} = (x, y)$, which evolve

according to the Fokker-Planck equation associated with the Langevin equations (1). The

evolution is described by the conservation equation

$\quad \frac{\partial \Psi_\pm}{\partial t} + \nabla \cdot (\dot{r}\Psi_\pm) + \frac{\partial}{\partial \theta}(\dot{\theta}\Psi_\pm) = \nu(\Psi_\mp - \Psi_\pm),$ (2)

with $\dot{r} = \pm V_0 e(\theta) - D_r \nabla \log \Psi_\pm$, $\dot{\theta} = \pm\omega - D_\theta \frac{\partial}{\partial \theta} \log \Psi_\pm$. To demonstrate the validity

of our theoretical model (1), we obtained the expected change in orientation and mean-

square displacements (MSDs) over time from Eq. (2) and compared them directly with our

experiments and numerical simulations. Thus, we can obtain the analytical expression of

the time-dependent change in orientation and MSDs of the moving cells derived from Eq.

(2) (see *Materials and Methods* for details).

$\quad$ We can obtain the time-dependent expected change in orientation of diatom cells by

multiplying Eq. (2) with $\cos \Delta\theta$ and separately by $\kappa \sin \Delta\theta$ respectively, and then

integrating both equations over θ and r. By solving a linear system of ordinary differential

equations for $\langle \kappa \cos \Delta\theta \rangle(t)$ and $\langle \kappa \sin \Delta\theta \rangle(t)$ (*SI Appendix for the derivation details*), the

analytical prediction of direction change is given by

$\quad\quad\quad \langle \cos \Delta\theta \rangle(t) = e^{-(D_\theta + \nu)t}(\cos \sqrt{\lambda}\nu t - \frac{1}{\sqrt{\lambda}} \sin \sqrt{\lambda}\nu t),$ (3)

where $\langle \cdot \rangle = \int_\Omega dA_r \int_0^{2\pi} d\theta (\Psi_+ + \Psi_-)$ and $\lambda = \left(\frac{\omega}{\nu}\right)^2 - 1$. To compare our model with

experimental data, the qualitative results are shown in Fig. 2b where experimental data of

orientation change decay overtime are in good agreement with our prediction Eq. (3).

Theoretically, we can further obtain the analytical predictions of the time-dependent MSDs

and effective diffusion coefficient from the Fokker-Planck equation (2). Using

mathematical derivation, we can obtain the analytical expression as

$\langle \Delta x^2 \rangle (t) = 4Dt + \frac{2V_0}{v^2(\lambda+\alpha^2)^2}\left[(\lambda + 2\alpha - \alpha^2)(1 - e^{-\alpha vt}\cos\sqrt{\lambda}vt) + (\lambda - 2\alpha\lambda - \right.$

$\left. \alpha^2)\frac{1}{\sqrt{\lambda}}e^{-\alpha vt}\sin\sqrt{\lambda}vt\right],$ (4)

where $D = D_r + \frac{V_0(\alpha-1)}{2v(\lambda+\alpha^2)},$ (5)

is the effective diffusion coefficient or diffusivity with $\lambda = \frac{\omega^2}{v^2} - 1$ and $\alpha = \frac{D_\theta}{v} + 1$.

Comparison between experiment and model predictions. To validate the rationality of

model (1), we tested the model in the context of the movement scaling behaviours at lower

cell density to ignore cell encounters modifying cell movement. In the experimental setup,

low densities of individuals were used in order to minimize effects of cell-cell interactions

(i.e., the attraction pheromone ref. 7). We tracked the motion of cells in two-dimensional

space where cells can freely forage within given arena (*Materials and Methods*). The model

predicts the individuals foraging trajectories that coincide with the experimentally

observed trajectories using the experimentally determined rotational diffusion coefficient

$D_\theta \approx 5.4 \times 10^{-3}$ rad^2/s (see Fig. 1 and *SI Appendix* Movie S1 and S2). Although the

theoretical model (1) gives a constant speed rather than a fluctuating speed, it does not

affect spatial dispersal and temporal correlation. Hence, for modelling simplicity, a

constant movement behaviour is adequate to describe diatom motion.

In order to quantitatively verify our model, we conducted simulations of diatom

movements with $D_\theta = 0.1$ rad^2/s, $D_\theta = 0.0054$ rad^2/s (the experiment data), and $D_\theta =$

0.0001 rad^2/s , and comparing their MSDs and the average temporal correlation $\langle \cos \Delta\theta \rangle$

with experimental results. Fig. 2a shows the two regimes of MSDs as a function of time

intervals and extracted their slope k by linear fits of $\log MSD$ as function of $\log t$. For the

short time scales ($t < t_C = 25\ s$), the MSDs are superdiffusive dynamics with an exponent

of $k = 2$ independent of rotational diffusion, whereas for the long times scales ($t > t_C$),

the MSDs change from diffusive to subdiffusive behaviour ($k < 1$) with the simulated

strength of rotational diffusion decreasing. Figures 2a clearly indicates that decreased

rotational diffusion leads to a decline of slope k, indicating that the cell motion becomes

circular behaviours (*SI Appendix*, Fig. S2). $k = 1.0$ corresponds to random diffusive

motion without directionality, i.e. a paradigm run-and-tumble behaviour. The overall

behaviours are in good agreement with experiments, numerical simulations and theoretical

predictions. Note that the diffusion behaviour reverts to normal diffusions after a relative

long plateau for a weak rotational diffusion coefficient (see *SI Appendix*, Fig. S3).

Fig. 2b shows the temporal correlation of cell orientation, where the positive correlation

coefficient suggests the positive feedback in directional persistence, while the negative one

implies a reversed motion by stochastic changes. The model accurately describes the cell

reversal and rotational behaviour in both numerical simulations and theoretical predictions

Eq. (3) as shown in Fig. 2, albeit the latter shows a slightly deviation from numerical

simulation of angular speed ω. Specially, experimental results reveal that diatom cells

display a subdiffusion to search the nutrient dSi. It is in good agreement with theoretical

predictions on both MSDs and the temporal correlation of motion orientations of cells.

Optimal searching strategies in motion behaviours. Why diatoms evolve to move like

this novel "circular run-and-reversal"? We hypothesize that it allows optimization of

searching efficiency as a fitness proxy. We start by analysing an active cell with a sense

radius r_c, blindly searching for food in an environment with a homogeneous topography.

As diatoms cruise the searching space, it continuously captures nutrients that come within

a capture radius r_c from the cells centre, as schematically shown in Fig. 3a. The amount of

leftover nutrients in each run shows a monotonous decline as a function of the area swept

by the active cell. Here, we assume that all cells use the same strategy of reversal and

rotational diffusion for the simulations. Fig. 3b plots the average amount of leftover

nutrients n, obtained from 1,000 simulated trajectories as a function of various rotational

diffusions, so that the decay rate τ of the exponential fitted is a rational indicator to evaluate

the foraging efficiency.

We use the novel rotational diffusion and reversal as the evolvable parameter. As it can

be seen in Fig. 3c, there is a remarkably peak of the foraging efficiency with evolving D_θ,

but declines when rotational diffusion passes beyond a threshold. This optimal search

strategy can be understood in quantitative terms by looking at the extreme cases. An

individual cell in the absence of such behaviour, i.e., when $D_\theta = 0, \nu = 0$, peforms a

circular walk. As D_θ decreases from the optimal peak, individuals travel probabilistically

more accurately along the circle. Thus, they visit the same area frequently unless a reversal

event occurs. As D_θ increases, more costs of energy are involved in the spatial dispersal

processes rather than nutrients uptake, i.e., more energy is consumed to spraying EPS (see

SI Appendix, Fig. S2). In contrast, there exists an evolvable ν-space with respect to the

optimal foraging efficiency, as shown in Fig.3d. The searching efficiency decreases

monotonically with increasing reversal frequency ν, but displays a wide range plateau at

low reversal frequency. Thus, they optimize dispersal benefits, defined as the effective

diffusivity, that asymptotically reaches a maximum value (Fig. 3d). Theoretically, our

foraging efficiency τ is here comparable with the theoretical effective diffusion coefficient

in Eq. (5). They predict an extreme similarity to the profile of the various rotational

diffusivity and reversal events.

Qualitatively, these optimal foraging strategies as generated by the above movement

traits are in accordance with observation of experimental data. Note that experimental

results slightly deviate from optimal values of rotational diffusion D_θ and reversal events

ν, but they only 9% and 5% underestimate comparison with the optimal values of

theoretical predictions respectively. As shown in Fig. 3c and d, by comparing the numerical

simulated data with analytical predictions, similar as the parameter τ, the effective

diffusivity can also capture the searching efficiency of rotational and reversal behaviours.

Finally, with the simulation of foraging efficiency versus various D_θ and ν, we found that

there exists an optimal region where cells have maximum foraging efficiency (Fig. 4a).

Interestingly, if we plot the experimental parameters overlapped with directly numerical

simulations and theoretical predictions in a heatmap together, which shows the simulated

foraging efficiency with respect to (D_θ, ν)-parameter space, the experimental points are in

good with the optimal strategy region of Eq. **1** on the (D_θ, ν)-parameter space (cf. Fig. 4a

and b).

Invasibility analysis. Individuals optimize trade-offs between the benefits of dispersal and

the costs of energy involved in dSi detection because diatoms are the self-propulsive active

cells. For solitary individual, individual foraging behaviours are independent of the

behaviours of other individuals. Hence, the optimal foraging parameter D_θ and ν represent

an evolutionary stable strategy (ESS) of the dSi detection ability (Fig. 3c and d).

We now consider populations where individual cells may encounter each other. To

determine whether the optimal value is a long-term outcome of competition selection, or

just be exploited by free-riding strategies (37), we created a pairwise invasibility plot (PIP,

Fig. 4c) by performing an evolutionary invasibility analysis. Under a very broad range of

parameter conditions, we find that populations can convergently evolve to the ESS (Fig.

4c). The PIP reveals that diatom movement strategy of diatoms observed in our

experiments is not only an evolutionarily stable strategy but also convergence stable. Here,

we did not observe a branching to occur as the parameters of the model are changed in the

evolutionary dynamics.

Experimental evidence of movement strategies in different dSi concentrations. To

investigate the effects of different dSi concentrations on diatom foraging strategy, we

performed controlled experiments to calculate the change in reversal probability and

diffusion coefficient. The diatom cells move with low reversal probability and high

diffusion coefficient D at intermediate dSi concentrations (from 10 to 50 mg/L), whereas

low and high dSi both will lead to a decreased diffusion efficiency to cells. It is surprising

that the optimal diffusion coincides with typical dSi concentrations of the most coastal

water (Fig. 5a). The diffusion coefficient shows a monotonous decline with increased

reversal probability (Fig. 5b). This suggests that reversal behaviour is not the single trait

underlying the optimization of foraging strategy, which agrees well with our theoretical

model (1) that both reversal and rotational diffusion together control the optimal searching

strategy (Fig. 3 and 4).

Our results revealed that diatom cells adopt different movement behaviours with

changing ambient dSi concentrations. This adaptive response suggests that diatom cells are

able to sense the local dSi concentration and adjust their reversal rate to adapt to their

physical surroundings. When silicon becomes the limiting factor, diatom cells increase

searching activity to meet their dSi demand for survival and mitotic division. Consequently,

a high diffusion coefficient induces diatom cells to explore larger areas to take up dSi.

Discussion

We have observed a novel "circular run-and-reversal" behaviour in the marine biofilm-inhabiting diatom *Navicula arenaria* var. *rostellata*. The stochastic model describing this behaviour is derived from experimental observations. Our experimental and theoretical results demonstrated that the reversal and rotational behaviours play a crucial role in optimizing searching strategies. Diatoms can maximize their foraging efficiency by changing their reversal probability and rotational noise in a silicon-limited environment. Additionally, diatoms can alter their movement strategy when resource availability changes. Rather than foraging, they suppress their effective diffusivity in a nutrient-rich condition by reduced motility, contributing to reduce the costs associated with mitotic reproduction.

The searching efficiency within a low nutrient environment is strongly dependent on cell movement behaviours. Extending our results beyond dSi scavenging, there may other attractors server as the same role to impact motion behaviours of diatoms. For instance, *in silico* comparison of the experimental data led to the speculation that diatoms have a more efficient behavioural adaptation to pheromone gradients as compared to dSi (38). Together, integrating the effect of multiple attractors into evolutionary strategies remain a fascinating topic for the future research. Our model provides a universal paradigm to understand the ecologically relevant functions of movement behaviour from the perspective of foraging theory.

Especially, insight into the movement behaviour of microorganisms in aquatic environments has been generated from disciplines such as biophysics (39-41), but the focus of these studies has largely been on the statistical physical causes of behaviour and not the

ultimate cause. Cases of reversal behaviour were reported independently in different

species of marine bacteria (20, 42, 43). Recently, studies suggested that this reversal

behaviours can contribute to increase efficiency bacterial foraging (20, 41) and group social

effects (43), but similar evidence is still lacking for motile microalgae. This study

underscores the need to study the significance of these questions in other microorganisms.

Methods

Algal cell culture and images acquisition. The *Navicula arenaria* strain 0488 is

maintained in the BCCM/DCG diatom culture collection at Ghent University,

http://bccm.belspo.be/about-us/bccm-dcg. It was isolated in January 2013 from high-

nitrate intertidal flats of Paulina Schor, The Netherlands (51°21'N, 3°43'E). The isolate has

since been maintained in unialgal culture in artificial seawater medium Aquil (f/2+Si).

The diatom cells were grown using a standard protocol. One months before the

experiment, the diatom was acclimated to 2000 Lux light intensity with a light dark cycle

of 12:12 hours (INFORS HT Multitron pro, Switzerland). A 100 ml flask suspension was

grown on a shaker at 20°C rotating with 100 rpm. Mid–exponential–phase cells were

diluted with filtered autoclaved seawater and introduced into the test chamber for

observations.

In order to obtain long time data and prevent water evaporation, we put the suspension

into a transparent, sterilized polymer coverslip microchannel (μ-Slide I 0.2; Ibidi,

Martinsried, Germany). The channel length is 50 mm, height is 0.2 mm, and channel area

is about 2.5 cm^2. The transparent channel can provide a static state hydrodynamic

environment. We used Ti-E Nikon phase contrast microscope to record the pictures. We

recorded a series of continuous data under the 4-fold objective for 4 hours at a rate of 4

frames per second in the centre of the channel.

Numerical simulations. The parameters of the simulation correspond to Fig. 3c and d, Fig.

4a. For each D and v, 1000 trajectories of 600 sec have been simulated using a time-step

$\delta t = 0.1$ sec with the parameter values $V_0 = 17\ \mu m/s$, $D_r = 0\ \mu m^2/s$, $\kappa = \pm 1$, $\omega =$

$\pi/36$ rad/s. The 'cells' move in a homogeneous space with randomly distributed nutrients

($n = 1000$). At each step, the 'nutrients' that come within a capture radius r_c from the cell

centre will be removed.

Calculation of time-dependent orientation correlation and MSDs of moving cells. We

computed the average temporal correlation as follows: $\langle \cos \Delta\theta \rangle (\Delta t) = [\langle v(t + \Delta t) \cdot$

$v(t) \rangle - \langle v \rangle^2]/(\langle v^2 \rangle - \langle v \rangle^2)$, and caculated the mean square displacement via $MSD(\Delta t) =$

$\langle |r(t + \Delta t) - r(t)|^2 \rangle$.

Data avialability

The supplementary movies S1 and S2 are available from author homepage

(http://www.quan-xinglab.org/publications.html). The experimental data and Matlab code

used to generate Figure 1-5 are archived in a dryad public repostory

(https://datadryad.org/).

Acknowledgments

This paper is a product of the project "Coping with deltas in transition" within the

Programme of Strategic Scientific Alliances between China and The Netherlands (PSA)

financed by the Chinese Ministry of Science and Technology (2016YFE0133700), and the

National Natural Science Foundation of China (41676084).

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Lists of figure captions:

Fig. 1: Experimental observations and theoretical model predictions of the main

features of reversal behaviours of diatom *Navicula arenaria* var. *rostellata*. (**a**) A

typical cell trajectory containing circling run and reversal behaviours captured with a

microscopy at 4 frames per second (see *SI Appendix* Movie S1) for 5 min. (**b**) Cropping of

the partial trajectory of the panel (a) depicts a reversal behaviour, where the running from

CCW switches to CW through the reversal behaviours, and vice versa. The arrows indicate

the moving direction of the cell. (**c**) Experimental data showing the movement velocity

before and after a reversal occurrence; for clarity, not all speeds are shown here. (**d**) and

(**e**) The spatial trajectory and the reversal event predictions are given by model (1) with

parameters value $V_0 = 17\ \mu m/s$, $D_\theta = 0.0054$ rad^2/s, $\nu = 0.02$, and $\omega = \pi/36$ rad/s.

Colorbars in panel (a, d) depict the time (see *SI Appendix* Movie S2 for details).

Fig. 2: Laboratory measured and theoretical predicted diffusion behaviours of

diatom cells. (**a**) Mean squared displacement (MSD) for three different values of the

rotational diffusion coefficient D_θ obtained by performing a numerical simulation of model

(1) and comparison with the experiments (circles symbols), respectively. By decreasing the

strength of rotational diffusion in the model, the dependence of the MSD on time becomes

consistent with confined diffusivity from ballistic behaviours similarly to cage-effect

emergence after the characteristic times ($\sim 25\ s$). Parameters are $\omega = \pi/36$ rad/s, $\nu =$

0.02, and $V_0 = 17\ \mu$m/s. The dashed lines are a guide to the eye to mark the change of the

scaling law, and solid line predicted by Eq. (4). (**b**) Correlation of measured and predicted

changes in the direction of cells moving. Experimental data ($\square$ symbols) have error bars

representing lower and upper SD. Corresponding predictions (solid lines and dashed lines)

are given by Eq. (**3**) and numerical simulations obtained from model (1).

Fig. 3: Predicting optimal foraging strategy in spatial randomized nutrient targets

and theoretical results. (**a**) Schematic representation (not to scale) of diatom cells

blindly searching for randomly distributed nutrient resources (dots). The cells placed in

two-dimensional space move with constant speed $\boldsymbol{V_0}$ and variable orientation described

model (1). The capture radius $\boldsymbol{r_c}$ is about 20 $\boldsymbol{\mu}$m size (dashed circle area). (**b**) The

distinct exponential function, $\boldsymbol{n(t) = Ae^{-\tau t}}$, with the decay rate τ describes the foraging

efficiency of diatom movement strategy with respect to various $\boldsymbol{D_\theta}$ and $\boldsymbol{v}$. (**c, d**) The

efficiency of captured nutrients as a function of $\boldsymbol{D_\theta}$ and $\boldsymbol{v}$, respectively. The efficiency

values are averaged over 1000 trajectories at stable capture rates with various $\boldsymbol{\omega}$, where

the plot is scaled to the maximum value at $\boldsymbol{D_\theta = 0.3}$ and $\boldsymbol{v = 0.0001}$ respectively. The

solid lines represent an analytical prediction of effective diffusivity from theory Eq. **4**,

coinciding with directly numerical simulations of model (1). The gray shaded area

represents mean $\pm$ SD around the experimentally measured values of $\boldsymbol{D_\theta}$ and $\boldsymbol{v}$ on

species *Navicula arenaria* var. *rostellata*, respectively.

Fig. 4: Theoretical and experimental results implicate the emergence of the foraging

efficiency for various behavioural strategies. (**a**) Heatmap of foraging efficiency with

respect to (D_θ, v)-parameter space obtained from spatial randomly distributed nutrient

targets and constant movement speed for $\omega = \pi/36$ rad/s and $V_0 = 17$ μm/s. The optimal

foraging occurs over a window of behavioural parameters of v and D_θ, for instance, the

yellow areas are the optimal foraging regions. The boundaries of the optimal regions

change sharply with increasing reversal rate (white dashed lines with intervals $\Delta\tau = 0.1$).

In the low reversal rate limit, there are nonlinear effects of the rotational diffusion on

diatom foraging. The colored-solid dots correspond to the experimentally measured

rotational diffusion coefficients versus reversal rate on diatom *Navicula arenaria* var.

rostellata and the colorscale indicates the scaled foraging efficiency, τ from 0 to 1.0. (**b**)

Theoretical prediction of Eq. (**5**) on the effective diffusivity as functions of the rotational

diffusivity and reversal events. It shows a similar spatial profile comparison with directly

numerical simulations. (**c**) Pairwise invasibility plot (PIP) indicating that the movement

behavioural strategy of rotational diffusivity evolves toward a stable point 0.2 (vertical

dashed line). For a range of resident (*x*-axis) and mutant (*y*-axis) movement strategies, the

PIP indicates whether a mutant has a higher (red) or a lower (green) fitness than the resident.

Plus and minus symbols indicate strategy combinations resulting in positive and negative

invasion fitness, respectively. Here, the PIP shows that the rotational diffusivity with 0.2

is the sole evolutionarily stable strategy (ESS). Parameters: $\nu = 0.02$, and $\omega =$

$\pi/36$ rad/s.

Fig. 5: Reversal behaviour and diffusivity depend on the ambient dSi concentration.

(**a**) The diffusivity of diatom cells maximizes at an ambient dSi concentration of about 30

mg/L and declines at low and high dSi concentrations. The probability of reversal events

shows a sharply increase when dSi goes beyond 60 mg/L, but it maintains a plateau at low

dSi. (**b**) Diffusion coefficient, showing a nonmonotonic relationship with increased

reversal events, which have a maximized dispersal coefficient about $\nu = 0.02$ coincident

with theoretical predictions. The grayscale rectangle indicates the dSi concentrations in

most coastal ecosystems.

Figure 1.

Figure 2.

Figure 3.

Figure 4.

Figure 5.