**Physical factors contributing to regulation of bacterial surface motility**

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**Microbes routinely face the challenge of acquiring territory and resources on wet surfaces. Cells move**

**in large groups inside thin, surface-bound water layers, often achieving speeds of 30 μm/s within this**

**environment, where viscous forces dominate over inertial forces (low Reynolds number). The canonical**

**Gram-positive bacterium *Bacillus subtilis* is a model organism for the study of collective migration over**

**surfaces with groups exhibiting motility on length scales three orders of magnitude larger than**

**themselves within a few doubling times. Genetic and chemical studies clearly show that the secretion**

**of endogenous surfactants and availability of free surface water are required for this fast group motility.**

**Here we show that: (i) water availability is a sensitive control parameter modulating an abiotic**

**jamming-like transition that determines whether the group remains fluidized and therefore collectively**

**motile, (ii) groups self-organize into discrete layers as they travel, (iii) group motility does not require**

**proliferation, rather groups are pulled from the front, and (iv) flow within expanding groups is capable**

**of moving material from the parent colony into the expanding tip of a cellular dendrite with implications**

**for expansion into regions of varying nutrient content. Together, these findings illuminate the physical**

**structure of surface-motile groups and demonstrate that physical properties, like cellular packing**

**fraction and flow, regulate motion from the scale of individual cells up to length scales of centimeters.**

**Introduction**

   In their search for resources microbes contend with physically distinct environments ranging from soft surfaces to bulk Newtonian fluids and complex fluids like mucus. In bulk fluid environments, canonical microbes like *Escherichia coli* (1) and *Bacillus subtilis* (2, 3) ascend favorable chemical gradients via run-and-tumble chemotaxis (4). This mechanism of gradient ascent requires both flagellar-mediated motility and a complex system of phosphorylation-memory and chemical sensors on the bacterial surface, that together regulate the run-tumble transition frequency (5). In contrast, bacterial surface motility has different requirements and inputs, and different species have evolved distinct modalities of surface motion. For instance, in the predatory species *Myxococcus xanthus,* individual cells move in back-and-forth motions that employ two distinct sets of protein machinery for 'twitching' and 'gliding' motility, and cells assemble into larger motile groups that traverse surfaces as monolayers (6, 7). Many other species, including the opportunistic pathogens *Serratia marcescens* (8, 9) and *Proteus mirabilis* (10, 11), also form large groups of motile cells that are capable of rapidly expanding over surfaces, in some cases even against bulk fluid flow (12). Similarly, when present in sufficient numbers *Paenibacillus dendritiformis* form intricate fractal-like patterns on soft agar surfaces in response to lateral chemical gradients (13, 14). Even baker's yeast have been observed to exhibit group movement over fluid surfaces (15). Other species exhibit surface motility in response to non-chemical fields; for instance the cyanobacterium *Synechocystis* is phototactic, responding to incident light by asymmetrically extending and retracting pili from its surface to create a biased random walk toward a light source (16, 17). Crucial to its motion, *Synechocystis* modifies the local surface environment by secreting exopolysaccharides, and only when enough cells have participated in such surface modification can the group move toward the light source. These examples demonstrate that in response to various gradients, microbes have evolved distinct sensing capabilities and modalities of motion to acquire resources and respond to selective pressures on surfaces. Despite their differences, surface motility in all of these species (and even in abiotic systems (18)) appears to be a collective phenomenon, requiring the motion of and/or surface modification by large numbers of cells (19, 20). Uncovering biochemical and genetic factors that regulate motion is crucial for understanding how each species executes surface motility, but those factors are only part of the full picture. Physical forces that produce, regulate, and guide microbial group motion on surfaces may be relevant across species and contexts, and are thus integral to our understanding of microbial ecology in natural environments and will expand the suite of design tools for engineering microbial systems.

   Modeling plays an important role in these systems because it has the potential to connect experimental observations with physical forces and regulators of motion (21–29). Different models posit

qualitatively distinct physical mechanisms for group motility, with basal assumptions motivated by

species-specific attributes. In *P. dendritiformis* growth and spreading of cells in dendritic patterns are

modeled as a diffusion-limited conversion of nutrients into biomass at the growing tips of the dendrite

(13, 30). Cellular motion is thought to rely on chemotaxis that follows nutrient gradients that are created

by metabolic consumption of the colonies themselves (22, 30, 31). Together these effects recapitulate

many of the classes of bacterial surface patterning observed in experiments (26, 32–36).  However, within

a dense and active (e.g. swarming) group whose velocity correlations rapidly decay on the length scale of

a few cells (21), it is unclear how individuals could effectively modulate their tumble frequency and

accumulate sufficiently persistent runs to deliberately bias their random walks and hence execute run-

and-tumble chemotaxis. Indeed, previous work in *B. subtilis* (37–40) (and *Pseudomonas aeruginosa* (41),

*Salmonella enterica* (42)*,* and *Escherichia coli* (43)) shows that neither chemotaxis nor motility of

individuals are necessary for rapid and outward-directed surface motility. Other models of surface motility

(e.g. swarming or sliding) in *B. subtilis* or *P. aeruginosa* focus on the role of surfactants (23, 44) and/or

osmotic potential (24), both of which appear to be crucial for rapid surface motility in those species (37,

45–47).

While cell density is clearly an important factor in motion, in so much as many cells are required

for collective movement, and genetics implicates endogenous surfactant production as a physical

requirement for collective motion, it remains unclear how density, material flow, and dendrite structure

contribute to regulation of this ubiquitous behavior. In this work we combined high resolution video and

time-lapse microscopy with computational image processing to clarify the relative contributions of

density, flow and structure by examining the surface motility of the extensively-studied Gram-positive

bacterium *B. subitilis* (5, 24, 48, 49). Within a few hours of deposition, a small central inoculum of wild-

type cells can rapidly colonize the entire surface of a wet 10 cm agar plate via apparent swarming motility

(50). Such group motility over soft surfaces has been shown to depend on the secretion of 'surfactin', a

bacterially produced surfactant and wetting agent (22, 23, 25, 37, 51–53). Functional knockouts for

surfactin production (Δ*srf*) result in a phenotype where individual cells are still motile and chemotactic in

bulk fluid, but bacterial groups cannot move across surfaces (37, 40, 54). Localized secretion of surfactin

is thought to generate a gradient in surface tension, and thus produce collective motion via the Marangoni

force (23, 55, 56).

Building on that biophysical picture, we found evidence that the movement of *B. subtilis* over soft

surfaces was regulated by cell density, whereby groups of cells are subject to a jamming-fluidization

transition (57) that correlates with packing fraction (and thus water availability). These results

complement findings that the viscosity of bacterial suspensions depends on cell concentration, shear rate,
and level of motility (58, 59). Further, we show that such groups operate in two distinct modalities of
motion:  groups can move without growth as 'islands' that translocate independent of the parent colony,
or dendrites can extend with large-scale flow of material from the parent colony to the dendrite tip. We
used a chemotropic assay to show that when connected to a nutrient-rich parent colony, extending
dendrites can venture into nutrient-barren regions, while colonies that originate in nutrient-barren
regions are not collectively motile. These data suggest that, in addition to genetically regulated production
of surfactants, multiple physical and abiotic factors regulate motion independent of the chemotactic or
motile abilities of the constituent cells, and that groups move differently when presented with anisotropic
nutrient environments. When combined with previous work, these data suggest a model in which
individual motility and subsequent swarming are mechanisms that maintain a fluidized state on which
surface-tension gradient forces (Marangoni forces) and osmotically driven colony hydration can act to
precipitate group motion over a surface, and that a jamming-like transition, like those found in other
macroscopic granular systems (57, 60–65), may be a key regulator of whether bacterial groups are able
to move.

**Results**

*Bacillus subtilis* is a model motile Gram-positive bacterium, capable of sensing and responding to
its chemical environment in bulk fluid via run-and-tumble chemotaxis (2, 5, 66). On wet surfaces wild-type
*B. subtilis* rapidly move out from a central inoculum (67), apparently via collective swarming motility (50)
that requires secretion of the endogenous bio-surfactant 'surfactin'. Mutants that lack the ability to
produce surfactin (Δ*srf*) do not exhibit collective motility from their central inoculum (37) (and Movie S1).
We inoculated small, dense (OD ~ 10 - 20) droplets of wild-type *B. subtilis* on soft (~0.5% w/v agarose)
nutrient-rich surfaces. After a brief quiescent phase, groups of cells rapidly expanded over the surface in
a dendritic pattern that reached the edge of the plate (~5 cm of travel) in less than 6 hours (Movie S2).
Dendrites robustly moved outward away from the original point of inoculation into fresh territory at an
average group motility rate of ~5 μm/s and up to ~15 μm/s (SI Fig. 1). Cells within the dendrites were
highly motile, exhibiting swarming motility (37) with individual cells moving at rates up to ~30 μm/s in
highly circuitous paths. Consistent with previous work (37, 38), we confirmed that mutants lacking the
ability to run (tumble-only, Δ*CheY*), lacking the ability to tumble (run-only, Δ*CheB*), and mutants lacking
flagella (Δ*hag*) were all able to exhibit rapid surface motility, albeit with some differences in dendrite
speed and spreading pattern (SI Fig. 2, Movies S3 - S5).

In our early observations we noted transitions in the motile state of groups, wherein cells within

a *contiguous region* appeared to be either highly motile and moving as a group or immobilized on both

the individual and group levels. These distinct states of motion were frequently observed at the same

time, and cellular groups were observed transitioning between these states on the time scale of a minute,

too short to be accounted for by phenotypic changes.

**Movement is regulated by a jamming-like transition**

While surface motility of *B. subtilis* is robust to genetic manipulations of chemotaxis and motility,

it was very sensitive to gel stiffness as measured by agarose concentration (37). Below ~0.4% agarose by

weight the gel is sufficiently porous that bacteria can penetrate and swim within it, akin to canonical swim

plates (50, 68). Above ~0.7% agarose, limited water availability hinders surface motility (50), leading to

expansion across plates through the replication of sessile cells at the leading edge of growth, as seen in *P.*

*dentritiformis* (dendrite-like) or *E. coli* (growing circular colonies), both of which expand at much slower

rates (31, 69, 70). Thus, across a relatively narrow range of agarose percentages, groups of *B. subtilis*

exhibit three qualitatively distinct behaviors: swimming through agarose (below ~0.4%), rapid surface

motility (between ~0.4% and ~0.7%), and slow growth by replication (above ~0.7%).  Why does the

transition from rapid surface motility to slow proliferative growth occur over such a narrow range of

agarose stiffness, and correspondingly, water availability?

Many previous studies of *B. subtilis* characterized spreading and attendant morphological

behaviors at the length-scale of colonies (mm to cm) and on the timescale of bacterial replication, imaging

colony morphology and spreading with minutes or hours between frames (26, 32, 37, 50). These excellent

studies revealed much of what is known about spreading phenotypes, their genetic mechanisms, and

biochemical correlates. We performed high temporal and spatial resolution imaging of spreading colonies

to illuminate processes that potentiate spreading. We captured images at 30 or 60 frames per second

with spatial resolution of 5 µm/pixel or 15 µm/pixel, respectively, both of which enabled us to see intensity

variations produced by the movements of cells within the swarm. We wrote a custom image analysis script

that measured a scalar correlate of motion as a function of position through time. Briefly, the algorithm

measures the mean absolute value of local intensity fluctuations at a position across a set of $N$ (usually 5

– 7) frames and thus reports on the level of motile activity at each position through time (see Methods).

With this 'activity' filter, we were able to visualize on the scale of 10s to 100s of microns which parts of

the colony were actively motile and in which parts cells were stationary (Fig. 1A).

In wild-type cells, the rapid movement of dendrites and cellular groups was correlated with

significant increases in the quantitative measure of activity from our image analysis, indicating that

movement of individuals positively correlates with movement of the group. Conversely, regions whose

constituent cells had a low measure of activity were stationary and did not move or flow over the surface

(Fig. 1 and Movie S6). Frequently, stark boundaries in activity level formed within contiguous regions of

cells, indicating that within a genotypically and phenotypically identical population, sub-populations could

adopt two qualitatively distinct states of motion at the same time. We interpreted the high activity state

as a fluidized (low viscosity) state in which cells were swarming, and the low activity state as a 'jammed'

state where the network of contacts and forces between cells led to local immobility. We observed that

jammed regions could be fluidized if they came into contact with a fluidized region. Further, the phase

boundary between these regions actively fluctuated in time, translating fractions of a micron per second

(Fig. 1B), meaning that individual cells transitioned between motile and immotile states faster than

(potential) phenotypic changes in the local environment. Physical theories of jamming in granular

materials (61, 63, 64, 71) predict that the transition between jammed and fluidized states result from

small (technically infinitesimal) differences in packing fraction.

We took two approaches to assess the connection between packing fraction and motion. Phase

contrast video microscopy allowed us to visualize the cells as dark objects, while open spaces – even a

fraction of a cell's area – are relatively bright (SI Fig. 3A, Movie S7A).  First, we reasoned that at suitably

high magnification (60X), the average intensity in a region is directly correlated with the amount of open

space (assuming the cells are in a monolayer, see next section of Results), or in other words, brighter

regions have lower packing fractions. We hypothesized that if packing fraction was controlling the

transition, there should be a positive correlation between local intensity and the local level of activity. We

averaged both the intensity and the activity over regions ~0.8 $\mu m^2$, and then calculated (i) the histogram

of spatially correlated values between those two measurements, (ii) the mean local intensity vs. the local

activity, and (iii) the bivariate correlation coefficient (SI Fig. 3B). The mean intensity values increased

monotonically with activity and the correlation coefficient was 0.34 ($p < 10^{-6}$, relative to the null-

hypothesis of no correlation between the two signals, SI Fig. 3B). Second, we hypothesized that there

should be a negative correlation between packing fraction and flow speed – with the understanding that

the nature of the jamming transition means that the relationship between those observables is non-linear

(i.e. we would expect a significant, but weak negative correlation). After removing low frequency

background variations, we used a fixed threshold to define which pixels were cells (dark) and which were

open space (light), and averaged the resulting packing fraction over square regions 12.4 $\mu m^2$ in area. We

then used open-source particle image velocimetry software (72) to measure the local speed of cellular

flow through the same pixel groups. We calculated the histogram of spatially correlated values between

flow speed and packing fraction, with a bivariate correlation coefficient of $-0.11$ (with $p < 10^{-6}$, relative to

the null-hypothesis of no correlation between the two signals) (SI Fig. 3C). Thus both methods provide

significant support for the hypothesis that packing fraction regulates collective movement.

Next, we wanted to determine if available water (and its anticorrelate: packing fraction) could

regulate the transition between the fluidized / collectively motile state and the jammed immotile state.

We imaged expanding tips that were suddenly subjected to evaporation, and hence reduced water

availability. As water evaporated, the expanding tips exhibited a dynamic transition from a fluidized and

rapidly swarming state to a state where both individuals and the group were immotile (Fig. 2A/B and

Movie S7B). The motile-to-immotile transition did not occur at the same time across the group, and local

variations in density allowed some cells to 'rattle' in place, even after the group as a whole had become

immotile – consistent with the appearance of 'rattlers' in jammed granular systems (73, 74).  These

observations further support the hypothesis that the transition is not a phenotypic change and that local

drag and frictional forces are not preventing motion. Then, to examine the reversibility of this transition,

we took the same plates and expanding tips, resealed them to halt evaporation, and continued imaging

while water from the gel rehydrated the cells. Over a few minutes, the colony re-fluidized into domains

of high activity sub-groups, which then coalesced until the entire tip region regained fluidity and continued

to expand (Fig. 2C/D and Movie S7C). This strongly indicates that re-wetting and the corresponding

reduction in packing fraction 'reverses' the jamming-like transition back to a fluidized and collectively

motile state (61, 71, 75).

**Motile groups are pulled by the front in discrete layers**

In our bright-field imaging it appeared that motile dendrites frequently traversed surfaces in

discrete layers of cells, first traversing the surface as dense monolayers, then double layers, and so forth,

up to 4 discrete layers (see Movies S6, S8 and S9). To confirm that groups migrated in discrete layers we

measured the surface height of expanding dendrites with high $Z$-resolution using an interferometric

profilometer (see Methods). In Fig. 3A/B we show the height field and a 1D profile for an expanding mono-

layer dendrite – its height across the expansion region was almost exactly 1 μm, the thickness of a single

*B. subtilis* cell. As dendrites slowed and/or encountered space constraints, the layer thickness jumped to

higher integer values, as shown in Fig. 3C/D whose dendrites transition between 0 and 3 layers of cells (4

layers visible in Movies S6 and S10). Looking back to the jamming data in Figs. 1 and 2, multiple discrete

layers are visible in the interior of the colony. Frequently, the fastest dendrites were advancing as
monolayers (the fastest moving groups in Movies S6, S8, and S9, notably tips in S10 appeared multi-
layered). In addition to being valuable *in situ* data of motile-group structure, this also meant that our
imaging of rapid motile groups (and attendant activity filtering) was characterizing the motion of all cells
in *Z*, as opposed to there being motile and immotile cells at different *Z* positions (e.g. a biofilm blocking
view of motile cells). This discrete structuring is noteworthy because relevant models of surfactant-driven
spreading take cell-layer thickness as a continuous variable, whereas these data demonstrate that
(frequently, though not strictly) motile groups move in monolayers with possible transitions to other
discrete heights, especially when near the jamming transition.

The boundary shape between bare agarose and dense cellular monolayers also contained
information. While the shape of the meniscal boundary can be influenced by surface pinning and
depinning (76), this suite of imaging data showed that such boundaries have different shape statistics near
and far from the advancing front. At the advancing front, the boundary curvature is smooth (low variation)
with an overall positive curvature whereas behind the front retraction of material causes curvature to
vary widely in a so-called 'lacunar' structure (SI Fig. 4A). Further, within the regions whose boundaries
were lacunar, voids appeared with similar distributions of boundary curvature to the external boundary
(SI Fig. 4B). Consistent with previous work on 'viscous fingering' (77) and models of surfactant-driven flow
(23, 24), these data support the hypothesis that the advancing front has positive pressure (pushes
outward, covers new territory) while the boundary behind the front experiences negative pressure (pulls
inward, retracts and reduces coverage), both consistent with cellular groups primarily being front-pulled
by gradients in surface tension, not pushed from the back by proliferation. Consequently, on one occasion
we observed a dendrite tip contacting its own surfactant field (78), resulting in an almost immediate
(within ~ 15 s) cessation of motion (SI Fig. 5 and Movie S11) despite the number of cells in the tip increasing
even after self-contact. Also consistent with a front-pull mechanism, much of the same imaging data
showed independent islands of cellular monolayers, completely detached from their parent colony,
traversing the surface independently, with characteristic positive curvature at the front and negative
curvature regions in the back (Movies S6, S8, S12 and S13).

While high numbers of cells are required for group motility, these data indicate that proliferation
alone is not generating outward pressure to drive group motion. To definitely determine this, we
measured the area of dense monolayer islands of cells as they moved over time. In Fig. 4, we show one
such island traversing an agarose surface over a distance many times its length. Rather than increasing
due to cell division and growth, the island area decreased over time as such islands tend to leave a trail of

cells immobilized on the surface by the meniscus. Therefore, during surfactin-driven surface motility

groups of cells are capable of large-scale motion without growth in the population and without connection

to their parent colony. Rather, boundary curvature analysis and the existence of independent islands

provide evidence that cells are pulled by the front, frequently in monolayers at a density near the

threshold for jamming.

**Material flows from parent colony to dendrite tip**

While our data were consistent with outward front-pressure pulling the group away from the

parent colony, multiple time-lapse experiments showed apparent rapid flow of discrete cellular layers

from parent colony to the expanding dendrite tip (e.g. see Movie S10). To confirm that these observations

were indeed flow of material over the scale of millimeters to centimeters, we performed our typical wild-

type expansion assay, but doped the initial cellular deposition with 1 µm fluorescent polystyrene beads,

at a ratio of ~ 1 bead per 100 cells (see Methods). We imaged the motion of the tracer beads at 30 fps,

which allowed us to see the trajectories of individual beads on the time scale of minutes (Movie S14, SI

Fig. 6A). The swarming motility of wild-type cells caused the motion of individual beads to be erratic, akin

to a non-thermal random walk with drift (79) especially when flow speed was slow in comparison to

cellular swim speeds (Movie S15, SI Fig. 6B). This erratic motion prohibited flow tracking in wild-type

groups, however, as we had confirmed earlier (38), mutants of *B. subtilis* unable to perform flagellar self-

propulsion (Δ*hag*) exhibit rapid surface motility with speeds and morphologies similar to wild-type cells

(see SI Fig. 2). This strain did not cause severe bead agitation and thus we were able to perform *in situ*

flow measurements using those mutants (Fig. 5 and Movie S16).

Using open-source cross-correlation software, we measured bead flow via particle image

velocimetry (72). Similar to the flow characteristics of a typical incompressible fluid, the flow speed of

dense cellular monolayers decreased in regions with wide meniscal boundaries and significantly sped up

at points of constriction (see inset Fig. 5). The similarity between incompressible 2D flow and flow of dense

cellular monolayers is consistent with the fact that cells are effectively solid steric objects densely packed

in (incompressible) water. Average flow speeds, and following the motion of individual beads, showed

that material originating in the parent colony frequently made its way to the expanding tip (Fig. 5 and

Movie S16). Thus, while our island data showed that motile groups need not be connected to the parent

colony and that cellular proliferation was not required for group motility, these data indicate that when

connected back to the parent colony, cells and material (e.g. nutrients) can be supplied to the expanding

tip of a dendrite.

Having confirmed that our bright-field data was indeed revealing large scale flow, we examined

how the speed of tip expansion into fresh territory compared with the speed of cellular flow along

established dendrites connecting the tip to the parent colony (SI Fig. 7).  Consistent with the bead-flow

data, kymograph analysis showed that when expanding dendrites are connected to their parent colony,

cells (and surrounding fluid) flow many centimeters from the colony into the expansion zone. Such flows

along an established dendrite are significantly faster (~ 10 μm/s) than the speed of tip movement (~ 3

μm/s) into fresh territory, as they must be for material to reach the tip from the colony. At this point it is

unclear how groups create these differentials in speed along a dendrite, but we offer a potential

mechanism in the Discussion.

**Motile groups use material flow to traverse nutrient poor regions**

We knew from previous work (37, 38) that neither chemotaxis nor flagellar-mediated motility is

required for group surface motility, but that surface motility was potently modulated by (e.g.) the

presence of monovalent ions (K+ (25)). This suggested that groups might respond to their chemical

environment indirectly via changes in metabolism and/or gene expression. For instance, differences in

nutrient concentration might affect the rate of synthesis and/or secretion of surfactin, and thereby affect

surfactin-mediated motility. Further, our large-scale flow data suggested that differences across the

chemical landscape that affect group surface motility could be ameliorated by flow of material from other

regions.

We wanted to test whether motile groups had differential responses to exogenously presented

nutrient gradients, and specifically whether groups would avoid negative nutrient gradients. We created

agarose plates with two distinct halves – one half containing the same rich defined medium (RDM) used

earlier, the other half containing (potassium-free) NaCl buffer osmotically matched with the RDM to

maintain the same osmotic potential (see Methods). A thin impermeable barrier separated the two

halves, initially keeping the nutrients on one side. Once the agarose set, we poured a thin (~ 1 mm) layer

of osmotically matched NaCl buffer with agarose, and thus linked the two regions into a contiguous

surface with uniform mechanical properties that allowed cells to move freely between them. The nutrient-

rich half also contained a red fluorescent tracer dye (rhodamine) that served as an approximate reporter

of nutrient diffusion over the barrier. With or without cells, these plates showed diffusion of the dye from

the nutrient rich side into the side devoid of nutrients, meaning that there was a negative concentration

gradient on the surface from the nutrient rich to the nutrient poor side. We reasoned that if wild-type

cells were responding (potentially indirectly) to the outward-facing chemical gradients on a standard plate

(i.e. with initially isotropic nutrients) or to the absence of required potassium ions on the split-plate, they

would react to the reverse nutrient gradient on these split plates and either avoid the nutrient poor region

or show reduced motility toward it.

To test this, we simultaneously inoculated such a split-plate in two positions from the same

isogenic population of wild-type cells, one colony on the nutrient rich side, one on the nutrient poor side

(Fig. 6A/B). The colony inoculated on the nutrient poor side did not exhibit collective motility and was

effectively stationary, while the colony on the nutrient rich side rapidly covered the nutrient rich zone and

simultaneously ventured into the nutrient poor zone. Initially, dendrites that expanded into the nutrient

poor zone covered the plate less densely than dendrites that never left the nutrient rich zone (Fig. 6B),

but tip speeds were similar in both zones. Ultimately, cells originating from the nutrient-rich zone

colonized the entire plate (both zones, see Movie S17). The difference in initial coverage between the

zones supports the hypothesis that groups do exhibit a (mild) chemotropic response to gradients in

nutrients, ions, or potentially their own secreted waste products. Interestingly, these data show that

connection to a parent colony in a nutrient rich zone allows cells to explore nutrient-poor regions, even if

those regions lack critical chemical species (K+ ions in wild-type (25)). This also complements findings that

bacteria can transport other bacteria as cargo during surface motility (80). We did not observe any

independently moving islands of cells in the nutrient-poor region. Our flow data offer a qualitative

mechanism for this 'scouting' ability, that large-scale flow of material from parent colony to the extending

tip brings metabolic resources, presumably to fuel, in part, surfactin production, thus allowing Marangoni

forces to continue to pull the group outward via positive tip pressure.

**Discussion**

Our data support a model in which water availability – a proxy for bacterial packing fraction – is a

sensitive control parameter for an abiotic jamming transition in the granular material that is densely

packed, expanding bacterial populations. We hypothesize that motility plays a contributory role,

encouraging a fluidized state via motility-dependent viscosity (59, 81).  When sufficient fluid is extracted

from the substrate by endogenous osmolyte secretion and cells are in a fluidized state, groups flow over

the surface driven and guided by surface-tension gradients (Marangoni forces). Conversely, when water

availability decreases, either by (e.g.) evaporation or decreasing osmotic potential, packing fraction

increases and cells transition to a jammed and immotile state. It is worth noting that while our data are

consistent with a jamming-fluidization transition, we cannot use our current data to *prove* that we are

meeting the technical requirements of such a transition. Jamming is defined by a non-crystalline network

of steric contacts and forces between particles that are 'isostatic', that is, they do not have free modes of

motion (other than rattlers (73)), but we are unaware of a modality of microscopy that permits live, *in situ*

visualization of the force and contact network in a dense bacterial group (stationary or collectively motile).

Nonetheless, the potential existence of an abiotic jamming transition that regulates group motility

through the dense packing of nematic actors offers a mechanism for why small increases in agarose gel

stiffness (~0.05%) switch groups from rapid surface migration to slow steric growth.  It also clarifies why

the secretion of both surfactants and osmolytes is required for collective motility (37, 50); the former

being crucial for force generation and the latter for maintaining fluidization. Thus we hypothesize that

colonies that are not able to achieve the fluidized state – for any reason – are constrained to expand on

the order of 100 times slower over surfaces using the forces generated by cell wall growth and cell division

(i.e. the mode of sessile growth one typically associates with single colonies growing on a petri dish).

Our examination of boundary structure (SI Fig. 4) was consistent with outward-facing pressure at

the expanding front and inward-facing pressure behind the front. However, many data sets showed

persistent flow from parent colony, or other regions behind the front, toward the advancing front, without

the appearance of lacunar structure (e.g. Movies S2 and S10).  We speculate that these differences in

boundary structure, and correspondingly flow, might reflect differences in the magnitude of force being

generated at the front and/or changes to the effective viscosity of the cellular suspension.  Differences in

force magnitude might arise from varying rates of surfactin secretion and/or processes that affect the

spreading or degradation of surfactin. Differences in effective viscosity of the cellular suspension could

arise due to changes in cellular density (including but not limited to jamming), cellular aspect ratio (9, 82),

internal shear arising from individual motility (59, 81), or processes that affect the rate of fluid extraction

(e.g. via osmolytes) or loss (e.g. via evaporation) from the substrate.  Similarly, we hypothesize that the

difference in tip speed vs. flow speed along an established dendrite (SI Fig. 7) – which allows material to

be brought from parent colony to dendrite tip – might reflect different energetic costs for motion in those

two scenarios.  A tip advancing into fresh territory uses surface tension differences to generate force, but,

given our layering data, it must also pay a line-energy cost along the meniscus on each side of the dendrite,

which to first order would appear like a constant cost in energy per unit length of movement (equivalent

to a constant reduction in overall force).  Conversely, assuming that cells flowing within an established

dendrite are using the same mechanisms for force generation, they move within the same meniscal

boundaries established by the tip, thus do not pay this energetic cost, and thus have a larger effective

force pulling them outward.

The existence of large scale flow from parent colony to dendrite tip presents a number of
biophysical questions and possibilities. How far into nutrient poor zones can expanding dendrites
venture?  How do those distances depend on the available nutrients or the position of the parent colony?
How do surface properties like local water availability and topography (83, 84) affect flow and
colonization?  How do structures on the surface guide and/or impeded collective motility?  The fact that
we observed islands of cells moving independently on nutrient-rich surfaces, whereas flow from a parent
colony was required for exploration of nutrient-poor zones, suggests that those modalities of motion
(island vs. flowing dendrite) depend on environmental conditions. For instance, an island that ventures
into a nutrient poor region might quickly deplete its resources, and thus decrease secretion of osmolytes
and surfactants required for motion (54). Likewise, the transport of materials and nutrients from a
common pool (parent colony) into discrete tendrils far from the common pool (80) spur game-theoretic
questions of how those resources are divided, whether phenotypic changes are playing a role, and – like
other collectively motile and reproductive systems (e.g. *Dictyostelium* (85, 86)) – how the specific
genomes of individuals propagate in space and time relative to their nearly isogenic kin.

Whereas individual cells in bulk fluids propel themselves and direct their motion according to the
algorithms of chemotaxis, these data and previous work strongly suggest a fundamentally different set of
rules govern group surface motility. Specifically, that collective secretion of surfactants and osmolytes
drives and directs motion, that large-scale flow plays a critical role in exploration of territories with
differing nutrient concentrations, and that bacterial packing fraction – modulated through multiple
mechanisms – is a sensitive control parameter that regulates transitions between fluidized (collectively
motile) and jammed (immotile) states of the group. Thus continued understanding and modeling of
microbial communities on surfaces must consider how these physical factors – not fully described by
genetics and biochemistry – contribute to population dynamics.

**Methods**

*B. subtilis culture storage and preparation*

Isogenic cultures of *B. subtilis* strains were grown to OD 0.5, mixed 50% in glycerol aliquots,
individually snap frozen and stored at - 80 C. For each experiment, 200 µL aliquots were removed and
thawed, then diluted in 10 ml LB medium and grown in a shaking incubator at 37 C for 4 hours until mid-
log phase, then pelleted by centrifugation at 4000 g for 10 mins. The supernatant was removed, and the
bacterial pellets were then resuspended in 200 µL LB, yielding a culture with an approximate density OD
10, then 1.5 µL of that culture was deposited on 0.5% agarose plates with Teknova EZ-RDM media (a rich

defined medium), then dried for 10 mins in air and sealed with parafilm. For bead flow experiments, 1 µm

fluorescent polystyrene beads (Bangs Laboratories, Dragon Green) were diluted into the inoculation

culture and vortexed before deposition.

Depending on the type of time lapse experiment, plates were then either incubated for 1 hr then

imaged or imaged immediately at 37 C. For culture profilometry (height measurements), cultures were

incubated for 2 - 3 hrs at 37 C, then unsealed within 10 minutes of imaging to minimize evaporation and

jamming. For the multilayer height measurements (Fig. 3B), cultures were unsealed, transitioned to

immotility, and were then imaged between 60 to 80 minutes later.

*Plate preparation*

Plates were created by mixing 100 mL Teknova EZ-RDM and 0.5% agarose by weight and

autoclaving for 15 minutes at 121 C.  Media was cooled to 50 C in an incubator before being poured into

4 plates (25 ml each). Each plate was cooled and solidified under flame in open air for 10, 15, or 30 minutes

before immediate inoculation, creating a range of initial gel hydrations to examine the effects of water

availability on colony morphology.  For the split-plate experiment (Fig. 6) zero-nutrient buffer was made

from de-ionized water, 0.5% agarose w/v and NaCl concentration matched to the calculated osmotic

potential of the rich defined medium.

*Microscopy and Imaging*

Transmitted-light and oblique-illumination microscopy were performed using a Nikon SMZ-25

stereo zoom microscope, with a P2-SHR Plan Apo 1x objective, and a Prior ES111 OptiScan stage. High

speed images were taken using an Andor Zyla 5.5 CMOS camera. Phase contrast microscopy was

performed using a Nikon Eclipse Ti-E inverted microscope with an CFI S Plan Fluor ELWD 20x Ph1 ADM

objective (Fig. 2) or CFI S Plan Fluor ELWD 40X Ph2 ADM objective (SI Fig. 3). High speed images were

taken using an Andor iXon Ultra 888 EMCCD camera. Culture height images were captured using a Zygo

NewView 7300 optical 3D profilometer in the University of Oregon's CAMCOR facility. Low-frequency

variations in plate height were computationally removed by subtraction of a fitted a polynomial surface

to the background.

Activity overlays were generated using custom Matlab© scripts which calculated the per-pixel

sum over the absolute differences between 5 or 7 consecutive frames. Kymographs were generated using

custom Matlab scripts to identify and interpolate contours, then sum over cross-sectional slices of

intensities along the contours over time. All scripts and Matlab code are available on request.

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**References**

1.  R. Mesibov, J. Adler, Chemotaxis Toward Amino Acids in Escherichia coli. *J. Bacteriol.* **112**, 315–326 (1972).

2.  L. F. Garrity, G. W. Ordal, Chemotaxis in Bacillus subtilis: how bacteria monitor environmental signals. *Pharmacol. Ther.* **68**, 87–104 (1995).

3.  D. S. Bischoff, G. W. Ordal, Bacillus subtilis chemotaxis: a deviation from the Escherichia coli paradigm. *Mol. Microbiol.* **6**, 23–28 (1992).

4.  H. C. Berg, D. A. Brown, Chemotaxis in Escherichia coli analysed by Three-dimensional Tracking. *Nature* **239**, 500 (1972).

5.  C. V. Rao, G. D. Glekas, G. W. Ordal, The three adaptation systems of Bacillus subtilis chemotaxis. *Trends Microbiol.* **16**, 480–487 (2008).

6.  R. Balagam, O. A. Igoshin, Mechanism for Collective Cell Alignment in Myxococcus xanthus Bacteria. *PLOS Comput. Biol.* **11**, e1004474 (2015).

7.  D. B. Kearns, L. J. Shimkets, Chemotaxis in a gliding bacterium. *Proc. Natl. Acad. Sci.* **95**, 11957–11962 (1998).

8.  A. E. Patteson, P. E. Arratia, A. Gopinath, Quenching a swarm: Effect of light exposure on suppression of collective motility in swarming Serratia marcescens. *bioRxiv*, 331801 (2018).

9.  A. Rabani, G. Ariel, A. Be'er, Collective Motion of Spherical Bacteria. *PLOS ONE* **8**, e83760 (2013).

10. T. Matsuyama, *et al.*, Dynamic Aspects of the Structured Cell Population in a Swarming Colony of Proteus mirabilis. *J. Bacteriol.* **182**, 385–393 (2000).

11. H. H. Tuson, M. F. Copeland, S. Carey, R. Sacotte, D. B. Weibel, Flagellum Density Regulates Proteus mirabilis Swarmer Cell Motility in Viscous Environments. *J. Bacteriol.* **195**, 368–377 (2013).

12. B. V. Jones, Role of swarming in the formation of crystalline Proteus mirabilis biofilms on urinary catheters. *J. Med. Microbiol.* **54**, 807–813 (2005).

13. E. Ben-Jacob, I. Cohen, H. Levine, Cooperative self-organization of microorganisms. *Adv. Phys.* **49**, 395–554 (2000).

14. E. Ben-Jacob, From snowflake formation to growth of bacterial colonies II: Cooperative formation of
complex colonial patterns. *Contemp. Phys.* **38**, 205–241 (1997).

15. S. Atis, B. T. Weinstein, A. W. Murray, D. R. Nelson, Microbial Range Expansions on Liquid
Substrates. *Phys. Rev. X* **9**, 021058 (2019).

16. T. Ursell, R. M. W. Chau, S. Wisen, D. Bhaya, K. C. Huang, Motility Enhancement through Surface
Modification Is Sufficient for Cyanobacterial Community Organization during Phototaxis. *PLoS*
*Comput. Biol.* **9** (2013).

17. N. Schuergers, *et al.*, Cyanobacteria use micro-optics to sense light direction. *eLife* **5**, e12620 (2016).

18. M. Driscoll, *et al.*, Unstable fronts and motile structures formed by microrollers. *Nat. Phys.* **13**, 375–
379 (2017).

19. M. F. Copeland, D. B. Weibel, Bacterial Swarming: A Model System for Studying Dynamic Self-
assembly. *Soft Matter* **5**, 1174–1187 (2009).

20. R. M. Harshey, Bacterial Motility on a Surface: Many Ways to a Common Goal. *Annu. Rev. Microbiol.*
**57**, 249–273 (2003).

21. A. Be'er, G. Ariel, A statistical physics view of swarming bacteria. *Mov. Ecol.* **7**, 9 (2019).

22. C. Giverso, M. Verani, P. Ciarletta, Branching instability in expanding bacterial colonies. *J. R. Soc.*
*Interface* **12**, 20141290 (2015).

23. S. Trinschek, K. John, U. Thiele, Modelling of surfactant-driven front instabilities in spreading
bacterial colonies. *Soft Matter* **14**, 4464–4476 (2018).

24. S. Srinivasan, C. N. Kaplan, L. Mahadevan, A multiphase theory for spreading microbial swarms and
films. *eLife* **8**, e42697 (2019).

25. R. F. Kinsinger, M. C. Shirk, R. Fall, Rapid Surface Motility in Bacillus subtilis Is Dependent on
Extracellular Surfactin and Potassium Ion. *J. Bacteriol.* **185**, 5627–5631 (2003).

26. D. Kaiser, Bacterial Swarming: A Re-examination of Cell-Movement Patterns. *Curr. Biol.* **17**, R561–
R570 (2007).

27. J. D. Partridge, R. M. Harshey, Swarming: Flexible Roaming Plans. *J. Bacteriol.* **195**, 909–918 (2013).

28. T. A. Witten, L. M. Sander, Diffusion-Limited Aggregation, a Kinetic Critical Phenomenon. *Phys. Rev.*
*Lett.* **47**, 1400–1403 (1981).

29. J. Y. Wakano, S. Maenosono, A. Komoto, N. Eiha, Y. Yamaguchi, Self-Organized Pattern Formation of
a Bacteria Colony Modeled by a Reaction Diffusion System and Nucleation Theory. *Phys. Rev. Lett.*
**90**, 258102 (2003).

30. I. Golding, Y. Kozlovsky, I. Cohen, E. Ben-Jacob, Studies of bacterial branching growth using reaction–
diffusion models for colonial development. *Phys. Stat. Mech. Its Appl.* **260**, 510–554 (1998).

31. E. Tamar, M. Koler, A. Vaknin, The role of motility and chemotaxis in the bacterial colonization of
protected surfaces. *Sci. Rep.* **6**, 19616 (2016).

32. N. C. Caiazza, R. M. Q. Shanks, G. A. O'Toole, Rhamnolipids Modulate Swarming Motility Patterns of
Pseudomonas aeruginosa. *J. Bacteriol.* **187**, 7351–7361 (2005).

33. C. J. Ingham, E. B. Jacob, Swarming and complex pattern formation in Paenibacillus vortex studied
by imaging and tracking cells. *BMC Microbiol.* **8**, 36 (2008).

34. J. K. Parrish, L. Edelstein-Keshet, Complexity, Pattern, and Evolutionary Trade-Offs in Animal
Aggregation. *Science* **284**, 99–101 (1999).

35. E. B. Steager, C.-B. Kim, M. J. Kim, Dynamics of pattern formation in bacterial swarms. *Phys. Fluids*
**20**, 073601 (2008).

36. A. Marrocco, *et al.*, Models of Self-Organizing Bacterial Communities and Comparisons with
Experimental Observations. *Math. Model. Nat. Phenom.* **5**, 148–162 (2010).

37. D. B. Kearns, R. Losick, Swarming motility in undomesticated Bacillus subtilis. *Mol. Microbiol.* **49**,
581–590 (2003).

38. R. Fall, D. B. Kearns, T. Nguyen, A defined medium to investigate sliding motility in a Bacillus subtilis
flagella-less mutant. *BMC Microbiol.* **6**, 31 (2006).

39. S. Mukherjee, D. B. Kearns, The structure and regulation of flagella in Bacillus subtilis. *Annu. Rev.*
*Genet.* **48**, 319–340 (2014).

40. R. A. Calvo, D. B. Kearns, FlgM Is Secreted by the Flagellar Export Apparatus in Bacillus subtilis. *J.*
*Bacteriol.* **197**, 81–91 (2015).

41. T. S. Murray, B. I. Kazmierczak, Pseudomonas aeruginosa Exhibits Sliding Motility in the Absence of
Type IV Pili and Flagella. *J. Bacteriol.* **190**, 2700–2708 (2008).

42. A. Toguchi, M. Siano, M. Burkart, R. M. Harshey, Genetics of Swarming Motility in Salmonella
enterica Serovar Typhimurium: Critical Role for Lipopolysaccharide. *J. Bacteriol.* **182**, 6308–6321
(2000).

43. M. Burkart, A. Toguchi, R. M. Harshey, The chemotaxis system, but not chemotaxis, is essential for
swarming motility in Escherichia coli. *Proc. Natl. Acad. Sci. U. S. A.* **95**, 2568–2573 (1998).

44. H. Du, *et al.*, High Density Waves of the Bacterium Pseudomonas aeruginosa in Propagating Swarms
Result in Efficient Colonization of Surfaces. *Biophys. J.* **103**, 601–609 (2012).

45. A. Yang, W. S. Tang, T. Si, J. X. Tang, Influence of Physical Effects on the Swarming Motility of
Pseudomonas aeruginosa. *Biophys. J.* **112**, 1462–1471 (2017).

46. R. Chen, S. B. Guttenplan, K. M. Blair, D. B. Kearns, Role of the σD-Dependent Autolysins in Bacillus
subtilis Population Heterogeneity. *J. Bacteriol.* **191**, 5775–5784 (2009).

47. A. Be'er, *et al.*, Paenibacillus dendritiformis Bacterial Colony Growth Depends on Surfactant but Not
on Bacterial Motion. *J. Bacteriol.* **191**, 5758–5764 (2009).

48. E. Ward, *et al.*, Organization of the flagellar switch complex of Bacillus subtilis. *J. Bacteriol.*,
JB.00626-18 (2018).

49. J. R. Kirby, T. B. Niewold, S. Maloy, G. W. Ordal, CheB is required for behavioural responses to
negative stimuli during chemotaxis in Bacillus subtilis. *Mol. Microbiol.* **35**, 44–57 (2000).

50. D. B. Kearns, A field guide to bacterial swarming motility. *Nat. Rev. Microbiol.* **8**, 634–644 (2010).

51. L. W. Schwartz, R. V. Roy, Some results concerning the potential energy of interfaces with
nonuniformly distributed surfactant. *Phys. Fluids* **13**, 3089–3092 (2001).

52. Á. T. Kovács, Bacillus subtilis. *Trends Microbiol.* **27**, 724–725 (2019).

53. T. E. Angelini, M. Roper, R. Kolter, D. A. Weitz, M. P. Brenner, <em>Bacillus subtilis</em> spreads by
surfing on waves of surfactant. *Proc. Natl. Acad. Sci.* **106**, 18109 (2009).

54. W.-J. Ke, Y.-H. Hsueh, Y.-C. Cheng, C.-C. Wu, S.-T. Liu, Water surface tension modulates the
swarming mechanics of Bacillus subtilis. *Front. Microbiol.* **6** (2015).

55. S. Srinivasan, N. C. Kaplan, L. Mahadevan, Dynamics of spreading microbial swarms and films.
*bioRxiv* (2018) https:/doi.org/10.1101/344267 (February 5, 2019).

56. M. Fauvart, *et al.*, Surface tension gradient control of bacterial swarming in colonies of
Pseudomonas aeruginosa. *Soft Matter* **8**, 70–76 (2012).

57. H. H. Wensink, *et al.*, Meso-scale turbulence in living fluids. *Proc. Natl. Acad. Sci.* **109**, 14308–14313
(2012).

58. H. M. López, J. Gachelin, C. Douarche, H. Auradou, E. Clément, Turning Bacteria Suspensions into
Superfluids. *Phys. Rev. Lett.* **115**, 028301 (2015).

59. A. Sokolov, I. S. Aranson, Reduction of Viscosity in Suspension of Swimming Bacteria. *Phys. Rev. Lett.*
**103**, 148101 (2009).

60. M. Burel, S. Martin, O. Bonnefoy, Jamming/flowing transition of non-Brownian particles suspended
in a iso-density fluid flowing in a 2D rectangular duct. *EPJ Web Conf.* **140**, 03086 (2017).

61. E. Woldhuis, V. Chikkadi, M. S. van Deen, P. Schall, M. van Hecke, Fluctuations in flows near
jamming. *Soft Matter* **11**, 7024–7031 (2015).

62. K. A. Dahmen, Y. Ben-Zion, J. T. Uhl, A simple analytic theory for the statistics of avalanches in
sheared granular materials. *Nat. Phys.* **7**, 554–557 (2011).

63. M. E. Cates, J. P. Wittmer, J.-P. Bouchaud, P. Claudin, Jamming, Force Chains, and Fragile Matter.
*Phys. Rev. Lett.* **81**, 1841–1844 (1998).

64. A. J. Liu, S. R. Nagel, Jamming is not just cool any more. *Nature* **396**, 21–22 (1998).

65. I. R. Peters, *et al.*, Dynamic jamming of iceberg-choked fjords. *Geophys. Res. Lett.* **42**, 1122–1129
(2015).

66. J. Henrichsen, Bacterial surface translocation: a survey and a classification. *Bacteriol. Rev.* **36**, 478–
503 (1972).

67. N. Morales-Soto, *et al.*, Preparation, Imaging, and Quantification of Bacterial Surface Motility
Assays. *J. Vis. Exp. JoVE* (2015) https:/doi.org/10.3791/52338 (October 7, 2019).

68. A. J. Wolfe, H. C. Berg, Migration of bacteria in semisolid agar. *Proc. Natl. Acad. Sci.* **86**, 6973–6977
(1989).

69. L. Jauffred, R. Munk Vejborg, K. S. Korolev, S. Brown, L. B. Oddershede, Chirality in microbial biofilms
is mediated by close interactions between the cell surface and the substratum. *ISME J.* **11**, 1688–
1701 (2017).

70. C. Beloin, A. Roux, J.-M. Ghigo, Escherichia coli biofilms. *Curr. Top. Microbiol. Immunol.* **322**, 249–
289 (2008).

71. C. Reichhardt, C. J. O. Reichhardt, Aspects of jamming in two-dimensional athermal frictionless
systems. *Soft Matter* **10**, 2932–2944 (2014).

72. W. Thielicke, E. Stamhuis, PIVlab – Towards User-friendly, Affordable and Accurate Digital Particle
Image Velocimetry in MATLAB. *J. Open Res. Softw.* **2**, e30 (2014).

73. M. Skoge, A. Donev, F. H. Stillinger, S. Torquato, Packing hyperspheres in high-dimensional
Euclidean spaces. *Phys. Rev. E* **74**, 041127 (2006).

74. S. Atkinson, F. H. Stillinger, S. Torquato, Detailed characterization of rattlers in exactly isostatic,
strictly jammed sphere packings. *Phys. Rev. E Stat. Nonlin. Soft Matter Phys.* **88**, 062208 (2013).

75. K. Roeller, J. Blaschke, S. Herminghaus, J. Vollmer, Arrest of the flow of wet granular matter. *J. Fluid
Mech.* **738**, 407–422 (2014).

76. M. Hennes, J. Tailleur, G. Charron, A. Daerr, Active depinning of bacterial droplets: The collective
surfing of Bacillus subtilis. *Proc. Natl. Acad. Sci.*, 201703997 (2017).

77. O. Praud, H. L. Swinney, Fractal dimension and unscreened angles measured for radial viscous
fingering. *Phys. Rev. E* **72**, 011406 (2005).

78. B. L. James, J. Kret, J. E. Patrick, D. B. Kearns, R. Fall, Growing Bacillus subtilis tendrils sense and
avoid each other. *FEMS Microbiol. Lett.* **298**, 12–19 (2009).

79. Y. Wu, H. C. Berg, Water reservoir maintained by cell growth fuels the spreading of a bacterial
swarm. *Proc. Natl. Acad. Sci.* **109**, 4128–4133 (2012).

80. A. Finkelshtein, D. Roth, E. Ben Jacob, C. J. Ingham, Bacterial Swarms Recruit Cargo Bacteria To Pave
the Way in Toxic Environments. *mBio* **6** (2015).

81. D. Lopez, E. Lauga, Dynamics of swimming bacteria at complex interfaces. *Phys. Fluids* **26**, 071902
(2014).

82. B. Ilkanaiv, D. B. Kearns, G. Ariel, A. Be'er, Effect of Cell Aspect Ratio on Swarming Bacteria. *Phys.*
*Rev. Lett.* **118**, 158002 (2017).

83. E. S. Gloag, *et al.*, Micro-Patterned Surfaces That Exploit Stigmergy to Inhibit Biofilm Expansion.
*Front. Microbiol.* **7** (2017).

84. M. Werb, *et al.*, Surface topology affects wetting behavior of Bacillus subtilis biofilms. *Npj Biofilms*
*Microbiomes* **3**, 1–10 (2017).

85. D. A. Brock, T. E. Douglas, D. C. Queller, J. E. Strassmann, Primitive agriculture in a social amoeba.
*Nature* **469**, 393–396 (2011).

86. G. D. Palo, D. Yi, R. G. Endres, A critical-like collective state leads to long-range cell communication
in Dictyostelium discoideum aggregation. *PLOS Biol.* **15**, e1002602 (2017).

**Figure 1. Isogenic cells exist in motile and immotile phases within the same colony. (A)** Time-lapse imaging of

wild-type *B. subtilis* spreading dendritically on a nutrient-rich soft-agarose surface. We applied a computational

image filter that uses local intensity fluctuations over time to report on the degree of movement. Green regions

are high motile activity and stationary regions are left gray. Transient motile regions and fluctuating boundaries

between motile and stationary regions emerged and dissipated on timescales shorter than (potential) phenotypic

changes (see Movie S1). **(B)** Kymographs along the colored dashed lines in (A) (time is downward). Transitions

from large-scale motility (green) to immobilized (jammed) states are denoted by the red arrows.  Regions could

transition from jammed back to fluidized if they came in contact with a fluidized region (see boundary movements

in yellow). At later times, once the entire population had become immotile, cells grew 'upward' into bilayers that

frequently coalesced.  See Movie S6.

**Figure 2. Water availability controls a jamming-like transition**. **(A)** Time-lapse imaging of the tip of a wild- type *B. subtilis* dendrite on a nutrient-rich low-agarose surface.  At *t* = 0 the plate was unsealed and its surface was exposed to air.  Subsequent evaporation reduced water availability at the surface, causing both collective motion (green color) and the computational measure of mean motile activity **(B)** to drop as the cellular group entered the jammed state. Across the entire dendrite tip collective motility halted. **(C - D)** The same plate was then immediately resealed and imaged over 20 minutes. The punctuated drop in evaporation rate allowed cells to retain more water from the gel. Fluidized domains of motile cells subsequently emerged and coalesced until the entire dendrite was, once again, collectively motile. In both sets of images, the rectangular structure toward the top is a cellular bilayer (see Fig. 3). Noise in the activity measurement (B) was (in part) due to manual refocusing required by the retraction of the surface as water evaporated. Scale bar is 200 μm. See Movies S7A and S7B.

**Figure 3. Dendrites expand in discrete layers.** The left-hand column shows height images of wild-type *B. subtilis*
dendrites on soft agarose taken using an interferometric profilometer (see Methods); the image backgrounds
have been computationally flattened. **(A)** Advancing dendrites flowed downward (white arrow) and the fronts of
those dendrites exhibited overall positive curvature (red dashed lines, inset) indicating outward pressure. Behind
the advancing front, boundaries between agarose and cells exhibited a 'lacunar' structure (e.g. inset, cyan dashed
ellipse), where the mode of the curvature distribution was negative (SI Fig. 3). Voids appeared behind the front,
consistent with negative pressure in those regions overcoming meniscal forces. Both voids and lacunar structure
are consistent with viscous fingering and cells de-wetting and being pulled to the front. **(B)** A one-dimensional
profile of the surface height along the black line in (A). Here the advancing front is a densely packed cellular
monolayer, whose height is one cell thickness as shown by the horizontal grey dashed lines (1 μm apart). **(C)** Same
imaging modality as in (A), now after dendrites have jammed on the surface and then grown for ~ 1 hr. The original
flow direction is indicated by the white arrow. **(D)** A one dimensional profile of the surface height along the black
line in (C). Within contiguous cellular regions, bacterial film height is discretized into layers one cell thick; here the
profile shows transitions from monolayer, to bilayer, to trilayer. Discretization ceased above four layers, when
presumably cells move in any orientation, rotating about their long axis within the entire vertical structure (see
also Movies S6, S8 and S9).

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**Figure 4. Groups can move independently without growth. (A)** Four snapshots of an independent monolayer
'island' of wild-type *B. subtilis* cells moving to the right (white arrow) on soft agarose. Over the course of imaging,
the group moved at an average speed of ~ 6 μm/s and the total amount of movement was ~ 700 μm, necessitating
movement of the viewing frame (grey data in B). Common to our observations, such islands leave a trail of cells
immobilized on the surface behind them (orange highlighted regions). **(B)** The area of the island decreased over
time (time labels correspond between A and B). When combined with our data showing movement in monolayers,
these data demonstrate that movement does not require proliferation. See Movie S13.

**Figure 5. Dendrite expansion corresponds with large-scale flow.** Fluorescent 1 μm tracer beads were diluted with the cellular suspension and deposited onto the agarose surface. We examined the motion of beads in Δ*hag* mutant cells (lacking flagella); this minimized bead agitation due to individual cellular motility (see SI Fig. 5 and Movies S14 and 15 for wild-type). The beads were imaged in an expanding dendrite millimeters long on the time scale of minutes; the positions of individual beads are shown colored by time. Here the dendrite is connected to its parent colony, material flow moved beads from the parent colony into the dendrite. The accompanying movie (Movie S16) shows individual beads starting in the colony and moving to the dendrite tip, demonstrating that flow generated by pulling at the front can move material long distances compared to cell size, from the colony to the tip of a growing dendrite. The inset shows the average flow velocity (orange arrows) and average flow speed through a region with varying width. Similar to a 2D incompressible fluid, constrictions accelerated flow, whereas flow slowed at wider points along the dendrite.

**Figure 6. Surface motility in the presence of an exogenous nutrient gradient. (A)** Soft agarose plates were poured
with a barrier between nutrient-rich media doped with a fluorescent tracer dye (rhodamine) on the right half of
the plate and osmotically balanced zero-nutrient buffer on the left, with a top layer of balanced zero-nutrient
buffer to create a contiguous surface. The spatial distribution of the dye confirmed nutrient diffusion with a
negative nutrient gradient from right to left. Wild-type *B. subtilis* colonies were deposited on either side of the
barrier. **(B)** The colony inoculated in the nutrient-rich zone (right, red circle) spread dendritically across the divided
plate into the zero-nutrient zone, effectively traversing along a negative nutrient gradient. An identical inoculum
in the zero-nutrient zone (left, green circle) exhibited no collective motility. Scale bar is 1 cm. **(C)** Some dendrites
remained in the nutrient-rich zone (dashed lines) while other dendrites ventured into the zero-nutrient zone (solid
lines). Both classes of dendrites exhibited similar speeds, and flow moved material along dendrites in both regions
(see also Movie S17). The small arrows on the time axis are when each dendrite entered the zero-nutrient zone.
Colors are matched between (B) and (C).

0 seconds 254

1 mm

Average Speed (um/s)
0 10 20 30 40 50 60

parent colony

**Figure 1. Isogenic cells exist in motile and immotile phases within the same colony. (A)** Time-lapse imaging of

wild-type *B. subtilis* spreading dendritically on a nutrient-rich soft-agarose surface. We applied a computational

image filter that uses local intensity fluctuations over time to report on the degree of movement. Green regions

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retain more water from the gel. Fluidized domains of motile cells subsequently emerged and coalesced until the
entire dendrite was, once again, collectively motile. In both sets of images, the rectangular structure toward the
top is a cellular bilayer (see Fig. 3). Noise in the activity measurement (B) was (in part) due to manual refocusing
required by the retraction of the surface as water evaporated. Scale bar is 200 µm. See Movies S7A and S7B.

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**Figure 3. Dendrites expand in discrete layers.** The left-hand column shows height images of wild-type *B. subtilis*
dendrites on soft agarose taken using an interferometric profilometer (see Methods); the image backgrounds
have been computationally flattened. **(A)** Advancing dendrites flowed downward (white arrow) and the fronts of
those dendrites exhibited overall positive curvature (red dashed lines, inset) indicating outward pressure. Behind
the advancing front, boundaries between agarose and cells exhibited a 'lacunar' structure (e.g. inset, cyan dashed
ellipse), where the mode of the curvature distribution was negative (SI Fig. 3). Voids appeared behind the front,
consistent with negative pressure in those regions overcoming meniscal forces. Both voids and lacunar structure
are consistent with viscous fingering and cells de-wetting and being pulled to the front. **(B)** A one-dimensional
profile of the surface height along the black line in (A). Here the advancing front is a densely packed cellular
monolayer, whose height is one cell thickness as shown by the horizontal grey dashed lines (1 μm apart). **(C)** Same
imaging modality as in (A), now after dendrites have jammed on the surface and then grown for ~ 1 hr. The original
flow direction is indicated by the white arrow. **(D)** A one dimensional profile of the surface height along the black
line in (C). Within contiguous cellular regions, bacterial film height is discretized into layers one cell thick; here the
profile shows transitions from monolayer, to bilayer, to trilayer. Discretization ceased above four layers, when
presumably cells move in any orientation, rotating about their long axis within the entire vertical structure (see
also Movies S6, S8 and S9).

**Figure 4. Groups can move independently without growth. (A)** Four snapshots of an independent monolayer

'island' of wild-type *B. subtilis* cells moving to the right (white arrow) on soft agarose. Over the course of imaging,

the group moved at an average speed of ~ 6 µm/s and the total amount of movement was ~ 700 µm, necessitating

movement of the viewing frame (grey data in B). Common to our observations, such islands leave a trail of cells

immobilized on the surface behind them (orange highlighted regions). **(B)** The area of the island decreased over

time (time labels correspond between A and B). When combined with our data showing movement in monolayers,

these data demonstrate that movement does not require proliferation. See Movie S13.

**Figure 5. Dendrite expansion corresponds with large-scale flow.** Fluorescent 1 μm tracer beads were diluted with the cellular suspension and deposited onto the agarose surface. We examined the motion of beads in Δ*hag* mutant cells (lacking flagella); this minimized bead agitation due to individual cellular motility (see SI Fig. 5 and Movies S14 and 15 for wild-type). The beads were imaged in an expanding dendrite millimeters long on the time scale of minutes; the positions of individual beads are shown colored by time. Here the dendrite is connected to its parent colony, material flow moved beads from the parent colony into the dendrite. The accompanying movie (Movie S16) shows individual beads starting in the colony and moving to the dendrite tip, demonstrating that flow generated by pulling at the front can move material long distances compared to cell size, from the colony to the tip of a growing dendrite. The inset shows the average flow velocity (orange arrows) and average flow speed through a region with varying width. Similar to a 2D incompressible fluid, constrictions accelerated flow, whereas flow slowed at wider points along the dendrite.

**Figure 6. Surface motility in the presence of an exogenous nutrient gradient. (A)** Soft agarose plates were poured

with a barrier between nutrient-rich media doped with a fluorescent tracer dye (rhodamine) on the right half of

the plate and osmotically balanced zero-nutrient buffer on the left, with a top layer of balanced zero-nutrient

buffer to create a contiguous surface. The spatial distribution of the dye confirmed nutrient diffusion with a

negative nutrient gradient from right to left. Wild-type *B. subtilis* colonies were deposited on either side of the

barrier. **(B)** The colony inoculated in the nutrient-rich zone (right, red circle) spread dendritically across the divided

plate into the zero-nutrient zone, effectively traversing along a negative nutrient gradient. An identical inoculum

in the zero-nutrient zone (left, green circle) exhibited no collective motility. Scale bar is 1 cm. **(C)** Some dendrites

remained in the nutrient-rich zone (dashed lines) while other dendrites ventured into the zero-nutrient zone (solid

lines). Both classes of dendrites exhibited similar speeds, and flow moved material along dendrites in both regions

(see also Movie S17). The small arrows on the time axis are when each dendrite entered the zero-nutrient zone.

Colors are matched between (B) and (C).