

Diverse and complex muscle spindle afferent firing properties emerge from multiscale muscle mechanics

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1 **ABSTRACT**

2 Despite decades of research, we lack a mechanistic framework capable of predicting how movement-
3 related signals are transformed into the diversity of muscle spindle afferent firing patterns observed
4 experimentally, particularly in naturalistic behaviors. Here, a biophysical model demonstrates that well-
5 known firing characteristics of muscle spindle Ia afferents – including dependence on movement history,
6 and nonlinear scaling with muscle stretch velocity – emerge from first principles of muscle contractile
7 mechanics. Further, mechanical interactions of the muscle spindle with muscle-tendon dynamics reveal
8 how motor commands to the muscle (alpha drive) versus muscle spindle (gamma drive) can cause highly
9 variable and complex activity during active muscle contraction and muscle stretch that defy simple
10 explanation. Depending on the neuromechanical conditions, the muscle spindle model output appears to
11 “encode” aspects of muscle force, yank, length, stiffness, velocity, and/or acceleration, providing an
12 extendable, multiscale, biophysical framework for understanding and predicting proprioceptive sensory
13 signals in health and disease.

14 INTRODUCTION

15 Amongst the many somatosensory receptors throughout the body that give rise to our kinesthetic
16 and proprioceptive senses, muscle spindles have a unique muscle-within-muscle design such that their
17 firing depends critically on both peripheral and central factors. Muscle spindle sensory signals are shaped
18 not only by the movements and forces on the body, but also by motor commands to both the muscle
19 (alpha drive) and to specialized muscle fibers in the muscle spindle mechanosensory region (gamma
20 drive). But, the neuromechanical interactions that lead to muscle spindle firing patterns in movement are
21 still poorly understood. Over many decades, experimental studies of muscle spindles across several
22 vertebrate species – from snakes to mammals and including humans – provide us with rich literature of
23 muscle spindle firing patterns during well-controlled conditions, revealing significant nonlinearities and
24 history-dependence with respect to muscle length and velocity (Haftel et al., 2004; Peter B. C. Matthews,
25 1981; Nichols & Cope, 2004; Arthur Prochazka & Ellaway, 2012; Uwe Proske & Gandevia, 2012).
26 Moreover, manipulations to gamma motor drive reveal that muscles spindle firing patterns do not
27 uniquely encode whole-muscle length and velocity even in tightly controlled experimental conditions
28 (Crowe & Matthews, 1964a, 1964b; Jansen & Matthews, 1962). Independent activation of alpha and
29 gamma motor neurons show that they can each profoundly alter muscle spindle sensory signals, rendering
30 a direct prediction of muscle spindle signals from recordings of movement highly unlikely. Indeed, the
31 few available recordings of muscle spindles during natural motor behaviors collectively exhibit and
32 complex relationship(s) to the biomechanics of movement that defies simple explanation (Ellaway et al.,
33 2015; A. Prochazka et al., 1976, 1977; Arthur Prochazka & Ellaway, 2012; Arthur Prochazka &
34 Gorassini, 1998; A. Taylor et al., 2000; A. Taylor & Cody, 1974; Anthony Taylor et al., 2006).

35 Developing a mechanistic framework for understanding how muscle spindle organs generate the
36 complex sensory signals during natural movements is critical for understanding muscle spindle sensory
37 signals and how they change under different behaviors, in neurological disorders, and during interactions
38 with devices or interventions to improve motor impairments. Existing computational models of muscle
39 spindle afferent firing have largely been data-driven, with well-known experimental features with

40 represented phenomenologically (A Schaafsma & Willigen, 2002; Hasan, 1983; C.-C. K. Lin & Crago,
41 2002; Mileusnic et al., 2006). Such models have high fidelity under the conditions in which the data were
42 collected, but generalize poorly to more naturalistic conditions. The anatomical arrangement of the
43 muscle and muscle spindle have also been shown in simulation to affect the forces on the muscle spindle
44 and alter muscle spindle firing properties (C.-C. K. Lin & Crago, 2002). The biomechanical properties of
45 the intrafusal muscle fibers within the muscle spindle have also been implicated in determining muscle
46 spindle receptor potentials (Blum et al., 2017; Fukami, 1978; Hunt & Ottoson, 1975b; Hunt & Wilkinson,
47 1980; Nichols & Cope, 2004; R. Poppele & Quick, 1985; U. Proske & Stuart, 1985), but intrafusal fiber
48 forces have not been recorded experimentally in conjunction with muscle spindle afferent recordings, and
49 only a few recordings of muscle spindle receptor potentials exist (Fukami, 1978; Hunt et al., 1978; Hunt
50 & Ottoson, 1975b; Hunt & Wilkinson, 1980).

51 Here, we take a biophysical modeling approach to build a muscle spindle model based on first
52 principles of intrafusal muscle contractile mechanics, along with its interaction with the muscle and
53 tendon to predict both classic and seemingly paradoxical muscle spindle Ia afferent firing during passive
54 and active conditions. We first establish a relationship between muscle spindle firing characteristics and
55 estimated intrafusal fiber force and its first time-derivative, yank (D. C. Lin et al. 2019), in passive stretch
56 conditions. We then build a generative model based on a simple representation of intrafusal muscle
57 mechanics, revealing that several of classic muscle spindle firing characteristics emerge from muscle
58 cross-bridge population kinetics. The independent effects of alpha and gamma drive on muscle spindle
59 firing were then simulated based on mechanical interactions between the muscle spindle and the muscle-
60 tendon unit. Highly variable muscle spindle activity observed during human voluntary isometric force
61 generation and muscle stretch—from silent to highly active—could be explained by as a result of
62 interactions between the muscle spindle and musculotendon unit, central activation of extrafusal (alpha
63 drive) versus intrafusal muscle fibers (gamma drive), and muscle length change. These multiscale
64 interactions may explain why muscle spindle activity may roughly approximate muscle force, length,
65 velocity, acceleration, activation level, or other variables under different movement and experimental

66 conditions. As such, our model establishes a biophysical framework to predict muscle spindle afferent
67 activity during natural movements that can be extended to examine multiscale interactions at the level of
68 cell, tissue, and limb mechanics.

69

70 RESULTS

71

72 *Muscle fiber force and yank reproduce diverse passive features of rat muscle spindle IFRs.*

73 We first demonstrated that in passive stretch of relaxed rat muscle, muscle spindle Ia afferent
74 firing rates can be described in terms of extrafusal muscle fiber force and yank. We recorded muscle
75 spindle afferent axonal potentials and computed instantaneous firing rates (IFRs) in anesthetized rats
76 while stretching the triceps surae muscles (Fig. 1). In the relaxed condition, i.e. in the absence of central
77 drive to the muscles, we assumed that extrafusal muscle fiber forces provide a reasonable proxy for
78 resistive forces of the intrafusal muscle fibers within the muscle spindle mechanosensory apparatus (Fig.
79 1). We previously demonstrated that whole musculotendon force in rat experiments is not predictive of
80 muscle spindle IFRs (Blum et al., 2019), and that only a portion of the total musculotendon force can be
81 attributed to extrafusal muscle fiber force, with the remaining nonlinear elastic component attributed to
82 extracellular tissues (Blum et al., 2019; G. A. Meyer & Lieber, 2011; G. Meyer & Lieber, 2018). We
83 modeled the extracellular tissue forces using a nonlinear elastic model and identified the parameters that
84 minimized the prediction error of muscle spindle firing rates based on estimated extrafusal muscle fiber
85 force and yank (See Methods). The initial rise in extrafusal muscle fiber force at the onset of stretch (Fig.
86 1, red trace) manifests as a large, transient yank signal (Fig. 1, blue trace) and became more apparent once
87 the extracellular tissue forces were subtracted from whole musculotendon force.

88 Like in cats (Blum et al., 2017), fine temporal details of rat Ia afferent firing in relaxed muscles
89 could be reconstructed by linear combinations of estimated extrafusal muscle fiber force and yank,
90 subjected to a threshold. We assumed average extrafusal fiber force to be proportional to intrafusal
91 muscle fiber force in the anesthetized, passive stretch condition only. The notable features of muscle

92 spindle Ia afferent firing reconstructed included initial bursts, dynamic responses during ramps, rate
93 adaptation during holds, and movement history-dependent firing (Fig. 2A-C; Fig. S1). The
94 reconstructions revealed history-dependent initial bursts in Ia afferent firing coincide with large, history-
95 dependent yank component in estimated extrafusal muscle fiber force (Fig. 2B). The dynamic Ia afferent
96 firing response during ramp stretches was primarily reconstructed by the force component (Fig. 2C) and
97 was larger during the first stretch of a repeated sequence (Fig. 2B). Rate adaptation during the hold period
98 was reconstructed by the force component in both slow and fast stretches (Fig. 2C) (P. B. C. Matthews,
99 1963a).

100 We then demonstrated that the sensitivity of Ia firing to passive muscle fiber force and yank was
101 differentially affected by two types of perturbation to the muscle spindle afferent neuron. Estimated
102 extrafusal muscle fiber force- and yank-based reconstruction of Ia firing rates was robust to experimental
103 perturbations due to either oxaliplatin chemotherapy (OX) alone or intra-axonal antidromic stimulation of
104 the afferent (STIM). While the mechanisms underlying these perturbations are undetermined, the effects
105 on firing likely involve alterations in function of ion channels in the nerve terminal function as opposed to
106 effects on properties of non-neural tissues, e.g. muscle. (Bullinger et al., 2011a; Cope et al., 2014;
107 Housley et al., 2020). Accordingly, the characteristics of the estimated intrafusal muscle fiber force and
108 yank were qualitatively similar in intact and OX rats, suggesting there was no change in muscle fiber
109 force in the OX animals. However, muscle spindles in healthy rats treated with OX maintain an initial
110 burst and dynamic response, but lack sustained firing during the hold period (Fig. 2D vs 2E; Fig. S2)
111 (Bullinger et al., 2011a). These OX Ia afferent firing phenotypes were primarily reconstructed by the yank
112 component (Fig. 2D, blue trace), with a small contribution of the force component (Fig 2D, red trace; Fig.
113 S2), suggesting a reduced sensitivity of the muscle spindle afferent to intrafusal muscle fiber force.
114 Conversely, when we perturbed the Ia afferent in healthy rats through intra-axonal electrical stimulation
115 (STIM; 500 ms duration, 100 Hz train of 30nA pulses) and applied muscle stretch immediately afterward
116 (Fig. 2F; Fig. S2), STIM Ia afferent firing phenotypes were primarily reproduced by the intrafusal fiber
117 force, with reduced sensitivity to the intrafusal fiber yank (Fig. S3). Overall, we were able to reproduce

118 perturbed Ia afferent firing data by varying only the relative weighting of force and yank. Taken together,
119 these data show that in the absence of changes in intrafusal or extrafusal muscle fiber force and yank
120 signals, the sensitivity of the muscle spindle Ia afferent to force and yank can be decoupled and therefore
121 may arise due to separate encoding or transduction mechanisms.

122 *Modifying muscle spindle Ia afferent sensitivity to force and yank generates an array of firing*
123 *phenotypes.*

124 The differential effects of oxaliplatin and axonal stimulation on spindle firing rates led us to hypothesize
125 there is a degree of independence in the transduction of force and yank in the spindle. We further tested
126 this hypothesis by simulating Ia afferent firing arising from force- and yank-based receptor driving
127 potentials (Fig 3A; see Methods). Nominal sensitivities of the model receptor current (closely related to
128 driving potential) were chosen to reproduce a typical recorded muscle spindle firing rate during passive
129 stretch (Fig 3B, green shaded box). We then generated a family of muscle spindle firing phenotypes (Fig.
130 3B, blue dots) by systematically varying the sensitivity of receptor currents (Fig. 3, thin black lines) to the
131 same muscle fiber force (Fig. 3B, vertical axis) and yank signals (Fig. 3B, horizontal axis) during the
132 same muscle stretch (Fig 3B).

133 Varying force- and yank- sensitivity generated diverse muscle spindle firing phenotypes similar
134 to those observed experimentally – including the OX and STIM phenotypes. Under passive conditions,
135 muscle spindle Ia afferent firing profiles exhibited larger initial bursts and dynamic responses as yank
136 sensitivity increased (Fig. 3B, left to right), consistent with classically-identified dynamic muscle spindle
137 firing phenotypes (Emonet-Dénand et al., 1977). High force sensitivity led to profiles with elevated
138 plateau firing (Fig. 3B, top to bottom), consistent with classically-identified static muscle spindle firing
139 phenotypes (Emonet-Dénand et al., 1977; Jansen & Matthews, 1962). The firing profiles with high yank
140 and lowest force sensitivity resembled the OX firing phenotype (Bullinger et al., 2011a) (Fig. 3B, yellow
141 shaded boxes, cf. Fig 2D) and the firing profiles with lowest yank sensitivity resembled the STIM firing
142 phenotype (Fig. 3B, red shaded boxes, cf. Fig 2D). We concluded that independently varying Ia afferent

143 sensitivity (the input-output gain) to force and yank explains a spectrum of spindle firing phenotypes.

144

145 *A muscle cross-bridge model predicts intrafusal muscle fiber force and yank underlying muscle spindle*

146 *firing rates.*

147 We next built a biophysical model of muscle spindle mechanics to more directly predict intrafusal

148 fiber force and yank during muscle stretch conditions where experimental data are not readily available.

149 History-dependent muscle forces have been simulated previously based on muscle cross-bridge

150 population cycling kinetics (K. S. Campbell & Lakie, 1998; K. S. Campbell & Moss, 2000; Kenneth S.

151 Campbell, 2014b; Kenneth S. Campbell & Moss, 2002), but no current muscle spindle model has

152 incorporated these principles, and thus none simulate history dependence in muscles spindle Ia afferent

153 firing (Hasan, 1983; C.-C. K. Lin & Crago, 2002; Mileusnic et al., 2006; Schaafsma et al., 1991). Muscle

154 spindles contain bundles of specialized “bag” and “chain” intrafusal muscle fibers (Boyd, 1962, 1976;

155 Boyd et al., 1977; M H Gladden & Boyd, 1985; A. Taylor et al., 1999). While different intrafusal muscle

156 fibers vary morphologically and mechanically, as well as in their contributions to Ia firing patterns, their

157 basic architectures are similar (Banks et al., 1997; Banks, 2005, 2015; M H Gladden & Boyd, 1985).

158 Intrafusal muscle fibers are innervated by gamma motor neurons in two polar regions containing

159 contractile muscle fibers, with an in-series non-contractile equatorial region around which the muscle

160 spindle endings are wrapped and mechanotransduction occurs. Muscle spindle Ia receptor potentials and

161 afferent firing are directly related to deformation of the equatorial regions (Boyd, 1976; Boyd et al.,

162 1977), which depends on the tensile forces in intrafusal muscle fibers.

163 Our mechanistic muscle spindle model consists of a pair of half-sarcomere muscle models

164 arranged in parallel, simulated using two-state actin-myosin population interactions (Kenneth S.

165 Campbell, 2014b). The simulated “dynamic fiber” loosely represented the putative dynamic bag1

166 intrafusal muscle fibers, thought to mediate muscle spindle initial bursts (Banks et al., 1997; U. Proske &

167 Stuart, 1985; Schäfer & Kijewski, 1974; Schäfer & Schäfer, 1969; Song et al., 2015). The simulated

168 “static fiber” loosely represented the bag2 and chain intrafusal muscle fibers of the muscle spindle
169 thought to mediate tonic **muscle spindle** firing (Fig. 4A) (R. E. Poppele et al., 1979). The dynamic fiber
170 model was designed with slower myosin attachment rates (Thornell et al., 2015b) and with a more
171 compliant passive element than the static fiber (Margaret H. Gladden, 1976) (see Methods; Fig. S5).
172 Length changes to the spindle were applied to both fibers equally. We assumed the receptor driving
173 potential of the **Ia afferent receptor ending** to be proportional to **a weighted combination of force and yank**
174 of the **intrafusal fibers based on visual inspection of intrafusal force recordings and our previous**
175 **hypotheses of intrafusal force and yank encoding (Blum et al., 2019; Fukami, 1978; Hunt & Ottoson,**
176 **1975b). This was** based on the idea that the **intrafusal mechanosensory** nuclear regions stretch linearly
177 with intrafusal force (Peter B. C. Matthews, 1981; Schaafsma et al., 1991) **and do not themselves**
178 **contribute significantly to the force profile of the intrafusal muscle fiber** (Fig. 4A, lower panel; See
179 Materials and Methods for more details).

180 Our mechanistic model predicted a spectrum of muscle spindle firing phenotypes **identified in**
181 **passive stretch conditions** by varying the sensitivity (gain) of driving potentials to biophysically-predicted
182 intrafusal fiber force and yank (Fig. 4B). The predicted muscle spindle firing phenotypes closely
183 resembled biological firing phenotypes discussed earlier. **As a result of our chosen kinetic scheme,**
184 **history-dependent forces emerge within our simulated intrafusal muscle fibers from cross-bridge**
185 **population dynamics (Fig 4C). When the muscle is at rest, the spontaneous formation and breakage of**
186 **cross-bridges is at steady-state equilibrium (Fig. 4C, time 1). As soon as a stretch is applied, the**
187 **population of attached cross-bridges is also stretched, resulting in a short period of high-stiffness force**
188 **response, known as short-range stiffness (Fig. 4C, time 2; Lakie and Campbell 2019). Once these cross-**
189 **bridges are stretched to their limits, they break and a new attachment equilibrium is reached during the**
190 **remainder of stretch (Fig. 4C, between times 2 and 3). If the muscle is shortened back to the original**
191 **length, the cross-bridges will shorten (Fig. 4C, time 4 – note the distributions shift leftward). In this case,**
192 **the force in the fibers will decrease until the fibers are slack (Fig. 4C, time 5), at which point the cross-**
193 **bridges will shorten the muscle fibers against zero load until the slack length is reduced to zero, or**

194 another stretch is applied (Fig. 4C, time 6). If the fibers are not given enough time to return to steady-state
195 equilibrium (e.g., Fig. 4C, time 1), a stretch will result in qualitatively different, history-dependent force
196 and yank responses.

197 *A biophysical muscle spindle model predicts a variety of classic muscle spindle firing phenomena.*

198 A truly general model of a muscle spindle would predict firing properties across different
199 experimental contexts without any fundamental changes to the model structure or parameters. Using a
200 single, nominal set of parameters in our muscle spindle model, we tested whether a variety of classical
201 muscle spindle firing characteristics during passive stretch would emerge, including: fractional power
202 relationship with stretch velocity (P. B. C. Matthews, 1963a), history-dependence (Haftel et al., 2004),
203 and increased firing due to gamma motor neuron activation (Boyd et al., 1977; Crowe & Matthews,
204 1964a, 1964b). To demonstrate the ability of the mechanical signals themselves to predict the properties
205 observed in Ia afferent firing rates, we examined simulated driving potentials rather than creating
206 additional degrees of freedom in our model with a spike generating process.

207 In our model, the scaling properties of muscle spindle dynamic responses and initial bursts with
208 increasing velocity during passive muscle stretch arise from intrafusal cross-bridge kinetics. More
209 specifically, the strain dependence of the cross-bridge detachment rates produces force profiles that
210 contain linear increases in short-range stiffness and sublinear increases in force dynamic response as a
211 function of stretch velocity. In a series of constant velocity stretches relative to optimal muscle length, L_0
212 (0.1 - 1.0 L_0/s in 0.1 L_0/s increments; Fig 5A), we computed the “dynamic index”, classically defined as
213 the difference in firing rate between the end of the ramp and after 0.5 s of isometric hold, which increases
214 with stretch velocity (Fig. 5B-C; P. B. C. Matthews 1963a). We predicted a sublinear increase in dynamic
215 index with stretch velocity, emergent from intrafusal mechanics, resembling the classically-reported
216 fractional-power velocity relationship in muscle spindle firing rates (Fig. 5C Houk, Rymer, and Crago
217 1981). In the same simulations, linear scaling of initial burst amplitude with stretch acceleration was
218 predicted (Schäfer & Kijewski, 1974; Schäfer & Schäfer, 1969). Initial burst scaling was emergent from

219 intrafusal muscle fiber yank at the onset of stretch due to the elasticity of attached cross-bridges that then
220 detach rapidly after being stretched a small fraction of L_0 (Hasan & Houk, 1975; P. B. Matthews & Stein,
221 1969). To our knowledge, neither of these phenomena has been previously demonstrated to arise from the
222 same mechanistic model.

223 Our biophysical model predicted history-dependent changes in the **muscle spindle firing** initial
224 **burst and dynamic response analogous to those previously reported in acute cat, rat, and toad experiments**
225 **as well as in microneurographic recordings in awake humans (Blum et al., 2017; Day et al., 2017; Haftel**
226 **et al., 2004; U. Proske & Stuart, 1985). In our model, these features emerged from the asymmetry present**
227 **in strain-dependent cross-bridge detachment rates.** In three consecutive, identical stretches, **the**
228 **biophysical muscle spindle** model predicted an initial burst and elevated dynamic response on the first,
229 but not second stretch (Fig. 6A). In the third stretch, the amplitude of the simulated driving potentials
230 recovered asymptotically as the time interval between stretches increased to 10s (Fig. 6A), resembling the
231 recovery of spike counts during the dynamic response in rats (Fig 6B) (Haftel et al., 2004) and the
232 recovery of initial bursts as a function of hold period in toads (Fig. 6B; Proske and Stuart 1985). Our
233 model predicted an initial burst at the onset of sinusoidal muscle stretches (Fig. 6C), similar to that seen
234 in microneurographic recordings from awake humans (Fig. 6C; Day et al. 2017). **Although a cross-bridge**
235 **mechanism has been proposed previously based on experimental findings (Haftel et al., 2004; Nichols &**
236 **Cope, 2004; U. Proske et al., 1992; U. Proske & Gregory, 1977; U. Proske & Stuart, 1985), this is the first**
237 **demonstration of history dependence in a muscle spindle model.**

238

239 *A biophysical muscle spindle model predicts effects of gamma drive on muscle spindle firing.*

240 **We tested whether previously-reported effects of gamma motor neuron activity on muscle spindle afferent**
241 **firing characteristics during** passive muscle stretch conditions were also emergent from our biophysical
242 model. Dynamic and static gamma motor neurons innervate the static and dynamic intrafusal muscle
243 fibers within the muscle spindle, respectively. To simulate classic experiments in which dynamic and
244 static gamma motor neurons were electrically stimulated in anesthetized animals, we independently

245 increased the number of available actin binding sites in simulated intrafusal static and dynamic muscle
246 fibers (Fig. 7A). Consistent with many prior experimental findings, simulated dynamic gamma drive
247 increased the force and yank of the dynamic fiber during stretch, predicting increased receptor driving
248 potentials underlying initial burst and dynamic responses, with proportionately smaller increases in
249 baseline muscle spindle driving potentials (Boyd et al., 1977; Hunt & Wilkinson, 1980). In contrast,
250 simulated static gamma drive primarily increased baseline muscle spindle driving potentials, and had
251 much smaller effects on driving potentials underlying the initial burst and dynamic response during
252 stretch (Fig. 7A-B; Crowe and Matthews 1964a; Boyd et al. 1977; Crowe and Matthews 1964b). Even the
253 previously-reported increases and decreases in dynamic index as a result of dynamic versus static fiber
254 stimulation, respectively, (Fig. 7C left plot; Crowe and Matthews 1964a; 1964b) were predicted by our
255 model (Fig. 7C right plot).

256

257 *Interactions between a biophysical muscle spindle model and muscle-tendon dynamics predict*
258 *paradoxical muscle spindle firing in isometric force production*

259 To simulate how alpha and gamma drive affect muscle spindle Ia afferent firing during voluntary
260 movement, we simulated the peripheral mechanical interactions of the biophysical intrafusal muscle fiber
261 above within extrafusal muscle-tendon dynamics. We modeled the muscle spindle and extrafusal muscle-
262 tendon dynamics while the total end-to-end length was held constant to simulate isometric muscle
263 contractions (Dimitriou, 2014; Edin & Vallbo, 1990) (Fig. 8). All of the force on the tendon was
264 generated by the extrafusal muscle fibers, as the intrafusal muscle fiber was assumed to generate
265 negligible force. We constrained the end-to-end lengths of the intrafusal and extrafusal muscle fibers to be
266 the same. However, muscle fibers were allowed to go slack based on their own properties, meaning the
267 true fiber lengths were not necessarily equal when slack was present in one or more fibers. Alpha drive
268 was simulated by activation of the extrafusal muscle fiber and gamma drive was simulated by activation
269 of the intrafusal muscle fiber (Fig. 8A).

270 We first tested the individual effects of alpha and gamma drive by independently activating
271 extrafusal and intrafusal muscle fibers in gaussian activation patterns (Fig. 8B, first two columns). When
272 extrafusal muscle was activated to 50% with intrafusal activation at a nominal constant level, the stress in
273 the extrafusal fiber increased, while the stress in the intrafusal fiber decreased (Fig. 8B, first column, top
274 trace, purple and green traces, respectively). The Ia afferent firing rate simulated using an integrate-and-
275 fire neuron was abruptly silenced during extrafusal muscle contraction (Fig. 8B, first column, middle
276 trace); during this condition both the extrafusal and intrafusal muscle lengths decreased and tendon length
277 increased (Fig. 8B, First column, bottom trace, green-purple vs. orange trace). Here, the extrafusal muscle
278 shortens due to the simulated alpha drive, and the intrafusal muscle fiber goes slack, causing a decrease in
279 intrafusal fiber stress. In contrast, when intrafusal fiber activation (both static and dynamic) was increased
280 and extrafusal activation was held constant, we observed no change in extrafusal stress or force, but an
281 increase in intrafusal stress (Fig. 8B, first column, top trace, purple and green traces, respectively). The
282 increase in intrafusal stress causing an increased in simulated Ia afferent firing rate (Fig. 8B, second
283 column, middle trace) with no change in muscle fiber or tendon length (Fig. 8B, second column, bottom
284 trace).

285 When we simulated concurrent alpha and gamma drive, we observed Ia afferent firing that was
286 dependent on the relative balance of imposed shortening of intrafusal fibers by alpha drive and the
287 activation of intrafusal fibers by gamma drive (Fig. 8B, third and fourth columns). As a proof of concept,
288 we activated the intrafusal and extrafusal fibers at a “low” and “medium” level, respectively (Fig. 8B,
289 third column, top trace), which caused an increase in extrafusal muscle fiber stress, a decrease in
290 intrafusal muscle fiber stress (Fig. 8B third column, top trace, purple and green traces), and a termination
291 of Ia afferent firing (Fig. 8B, third column, middle trace), while both extrafusal and intrafusal muscle
292 fiber lengths shortened (Fig. 8B, third column, bottom trace). Here, the extrafusal shortening due to alpha
293 drive outweighed the effects of gamma drive to the intrafusal muscle fiber and silenced the Ia afferent
294 (Fig. 8B, third column). When we reversed the relative activations so intrafusal activation was “medium”
295 and extrafusal activation was “low”, both intrafusal and extrafusal muscle fiber stress was increased while

296 the muscle shortened (Fig. 8B, fourth column, top versus bottom traces), causing increased Ia afferent
297 firing (Fig. 8A, fourth column, middle trace). In this case, both extrafusal and intrafusal muscle fibers
298 were shortening, but the presence of adequate intrafusal activation maintains the force in the intrafusal
299 muscle fibers, driving the muscle spindle firing in the presence of alpha drive to the extrafusal muscle
300 fiber (Fig. 8B, fourth column).

301 The simulations just described illustrate the multiscale, complex, interactions between the
302 musculotendon and the muscle spindle in a few intuitive cases (Fig. 8C-E). Briefly, if the extrafusal
303 muscle shortening caused by alpha motor drive is not sufficiently counteracted by gamma drive, the Ia
304 afferent will reduce its firing rate or cease firing altogether (Fig. 8C). Alternatively, if the gamma drive to
305 the spindle stretches the sensory region at a greater rate than the extrafusal shortening, the Ia afferent will
306 increase its firing rate (Fig. 8D). Finally, it follows that there is a set of extra- and intrafusal states that
307 will lead to no change in Ia firing rate during isometric conditions (Fig. 8E).

308 We were inspired to use these observations to explain paradoxical Ia afferent firing rates from
309 experiments in human finger muscles (Fig. 8C; Edin and Vallbo 1990), in which muscle spindle Ia
310 afferents were shown to either increase or decrease their firing rates (Fig. 8C, bottom traces) during an
311 isometric contraction despite the same joint torque profile (Fig. 8C, top traces). We mimicked muscle
312 fiber shortening under isometric conditions by applying a ramp increase in extrafusal muscle activation
313 (Fig. 8D). As expected, increasing alpha drive decreased extrafusal muscle fiber length (Fig. 8B, black
314 trace), and increased extrafusal muscle fiber stress (Fig. 8D, purple trace). In the absence of gamma drive,
315 the intrafusal fiber stress decreased as alpha drive increased (Fig. 8D, green trace) resulting in no muscle
316 spindle firing as shown experimentally (Elek et al., 1990). When we increased gamma drive by a fixed
317 amount intrafusal fiber stress could increase (Fig 8D, pink traces), decrease (Fig 8D, green traces), or
318 experience no change (Fig 8D, purple traces) in stress depending on the amount of alpha drive. As such it
319 was possible to generate almost any level of muscle spindle firing during isometric conditions where the
320 muscle fibers shorten by varying the relative levels of alpha and gamma drive.

321 Finally, to demonstrate the complex interplay of alpha and gamma drive when muscle fibers are
322 lengthened, we applied ramp-and-hold stretches to the model musculotendon while simulating the Ia
323 afferent response under different extra- and intrafusal activation conditions. When the musculotendon was
324 passive, i.e. no alpha drive, the muscle fiber length closely followed the applied musculotendon length
325 (Fig. 9A, left panel). Conversely, when the extrafusal muscle fiber was activated during the stretch, there
326 was a marked difference in musculotendon length compared to the applied length (Fig. 9A, right panel).
327 For each of these cases, we simulated muscle spindle Ia afferent responses with and without gamma drive.
328 In the case of the passive spindle, the musculotendon stretch in the absence of alpha drive produced a
329 slightly higher muscle spindle Ia afferent response than in the presence of alpha drive (Fig. 9B, top row).
330 Specifically, the baseline firing rate, initial burst, and dynamic response were all slightly larger in the case
331 of the passive musculotendon. These differences were more exaggerated when we simulated gamma drive
332 by increasing intrafusal muscle fiber activation (Fig. 9B, bottom row). The addition of a tonic level of
333 gamma drive increased the baseline firing rate, the amplitude of the initial burst, and the dynamic
334 response to stretch in the passive extrafusal musculotendon unit (Fig. 9B bottom row, green trace).
335 Adding alpha drive decreased baseline firing only slightly, abolished the initial burst, and had little effect
336 on the dynamic response to stretch (Fig. 9B bottom row, purple trace). These simulations demonstrate the
337 importance of considering musculotendon dynamics and activation when estimating muscle spindle
338 afferent feedback in different experimental conditions.

339

340 **DISCUSSION**

341 *Summary*

342 Here, we present a multiscale biophysical model that predicts a wide range of muscle spindle Ia afferent
343 firing characteristics in both passive and active conditions, illustrating the complex central and peripheral
344 neuromechanical interactions that underlie them in naturalistic conditions. We demonstrated that a diverse
345 range of Ia afferent firing characteristics during passive stretch – well-documented in the literature but not
346 completely understood mechanistically – is emergent from first principles of intrafusal muscle contractile

347 (cross-bridge) mechanisms. Ia afferent firing in a broad range of behaviorally relevant conditions can be
348 reproduced by converting the force and yank of the intrafusal muscle fibers within the muscle spindle into
349 afferent driving potentials using a simple linear relationship. This transformation represents the yet
350 unexplained mechanotransduction process in the equatorial encoding region of muscle spindle sensory
351 organs. The multiscale mechanical interaction of the muscle spindle within the muscle tendon unit
352 revealed isometric force generating conditions in which muscle spindle firing followed intrafusal muscle
353 fiber force, which can either resemble or differ from muscle-tendon unit force or length depending on the
354 relative amounts of alpha and gamma motor neuron drive. As such, mechanical interactions within the
355 muscle-tendon unit and forces external to the muscle all contribute to firing rates of muscle spindle Ia
356 afferents; they may even create firing patterns considered to be paradoxical if viewing muscle spindles as
357 simple passive sensors of kinematic or kinetic variables. Finally, we simulated multiscale mechanical
358 interactions including movements, alpha activation of extrafusal muscle, and gamma activation of
359 intrafusal muscle to demonstrate the importance of these complex interactions when modeling
360 proprioceptive activity. Our neuromechanical framework demonstrates how muscle spindle firing can
361 reflect the complex interplay of external and self-generated forces and motion. Importantly, our model
362 provides a foundation for biophysically realistic, predictive muscle spindle models that can be extended to
363 understand how properties of tissues, muscles, neurons, and central drive affect muscle spindle sensory
364 function in health and disease.

365

366 *A biophysical muscle model connects many independent observations with a simple hypothesis*

367 Our approach of starting with biophysical principles of muscle force generation enabled us to
368 develop a muscle spindle model that predicts nonlinear muscle spindle firing properties beyond the
369 conditions that it was designed to replicate, without any change in model parameters. We simulated
370 muscle cross-bridge kinetics to produce the muscle fiber force and yank that drive history-dependent
371 responses to a repeated muscle stretch stimulus, including the initial burst. Other muscle spindle models
372 do not exhibit history-dependence and – by design – either lack (Mileusnic et al., 2006) or always

373 generate (Hasan, 1983) an initial burst at the onset of stretch. In our model, the initial burst emerges from
374 the number of cross-bridges that are formed after a period of rest, temporarily increasing the stiffness of
375 the fibers before being stretched to their mechanical limit. This number depends on the prior and current
376 fiber length, velocity, acceleration, and activation of the intrafusal fibers, and thus leads to the modulation
377 of the initial burst of the spindle afferent with these variables. Similarly, the fractional power relationship
378 between muscle spindle dynamic responses and stretch velocity (Hasan, 1983; C.-C. K. Lin & Crago,
379 2002) has previously been imposed by including it explicitly within the model definition. While it can be
380 simulated using a phenomenological muscle model with a built-in fractional power function, prior models
381 fail to account for the history dependence of the velocity sensitivity and thus fail to replicate muscle
382 spindle firing patterns outside of tightly controlled laboratory conditions. **Similarly, the adaptation of**
383 **muscle spindle firing rates when held at a constant length after stretch follows the decrease in force of the**
384 **intrafusal muscle fiber when it is held at a constant length (Blum et al., 2017; Husmark & Ottoson, 1971).**
385 Here, history dependence, the nonlinear relationship to stretch velocity, the scaling of the initial burst to
386 acceleration, and firing rate adaptation are all implicit, emergent properties of intrafusal cross-bridge
387 force and yank under different conditions – no explicit formulation of these relationships is necessary.

388 Even the apparent differential effects on Ia afferent encoding by static and dynamic fusimotor
389 activity are accounted for by the mechanics of muscle force generation (Crowe & Matthews, 1964a,
390 1964b). We designed our model with slow cycling myosin in the “dynamic” fiber and fast cycling myosin
391 in the “static” fiber, modeled after nuclear bag and chain fibers, respectively (Thornell et al., 2015b). In
392 doing so, the dynamic and static fusimotor effects on Ia afferent stretch sensitivity were closely predicted,
393 emerging from muscle biophysics and a simple representation of mechanotransduction (Banks et al.,
394 1997). By taking a structural approach to the interactions between intrafusal and extrafusal muscle fibers
395 the effect of gamma motor drive were more robustly simulated (C.-C. K. Lin & Crago, 2002).

396 Variations in force and yank sensitivity account for both overall firing characteristics as well as
397 “static-dynamic” variability across Ia afferents during passive stretch. The few available recordings of
398 muscle spindle receptor potentials during muscle stretch exhibit fast initial transients similar to muscle

399 fiber yank, as well as a sustained potential that resembles muscle force and **its decay following a stretch**
400 **(Hunt & Ottoson, 1975a)**. We show that sweeping a range of sensitivities to intrafusal force and yank
401 predicts a continuum of firing phenotypes resembling those seen previously in both healthy and perturbed
402 conditions (Housley et al., 2020; Vincent et al., 2016). Further, each set of sensitivities can predict muscle
403 spindle firing across different stretch velocities and displacements. Whether a similar principle describes
404 the firing of the second type of muscle spindle afferent, group II afferents, remains to be tested.

405

406 *Multiscale interactions are necessary to consider for natural behavior*

407 Explicit simulation of the mechanical arrangement of the intrafusal muscle fibers within the
408 muscle-tendon unit enabled us to demonstrate the profound effects that central drive to alpha and gamma
409 motor neurons can have on muscle spindle firing during active force and muscle stretch conditions. Our
410 work emphasizes the importance of modeling the anatomical arrangement of the muscle shown in some
411 prior models (C.-C. K. Lin & Crago, 2002; Mileusnic et al., 2006). With the addition of biophysical force
412 generation mechanics we now provide a computational testbed to better explore the long-discussed idea
413 that muscle spindle firing rates are driven by properties of intrafusal muscle fiber forces (Peter B. C.
414 Matthews, 1981), and determined by both the central drive to intrafusal and extrafusal muscle fibers, as
415 well as interactions with external forces on the muscle. When simulating the mechanical interaction of the
416 muscle spindles within the muscle tendon unit, our model reveals that muscle spindle firing can depend
417 critically on the activation, force, and length of the intrafusal fiber within the muscle spindle, which
418 shapes the degree to which muscle spindle firing resembles extrafusal muscle length, velocity, force or
419 other biomechanical variables (Bewick & Banks, 2015; Kruse & Poppele, 1991).

420 The multiscale model presented here provides a mechanistic framework that can explain the rich
421 diversity of movement-related biomechanical signals in naturalistic behaviors. There are many such
422 scenarios we have not simulated that our model may provide insights for. For example, the mechanical
423 interactions between intrafusal and extrafusal muscle fibers (Fig. 8, 9) may explain the complex
424 relationships between Ia afferent signals in locomotion to muscle length and force. During the swing

425 phase of the cat locomotor step cycle, relaxed ankle extensor muscles are stretched by activity in ankle
426 flexor muscles, and extensor Ia firing rates appear to follow muscle velocity and/or length (A Prochazka
427 & Gorassini, 1998; Arthur Prochazka & Gorassini, 1998), which closely resemble extrafusal and
428 intrafusal muscle force and yank (Blum et al., 2017) in passive conditions. However, the same muscle
429 spindle Ia afferent also fires in early stance when there is no change in muscle length, but the muscle is
430 both active due to alpha drive, and loaded due to the carrying the weight of the animal: this firing is likely
431 due to intrafusal force resulting from the interactions between gamma drive, alpha drive, and the external
432 load (Fig. 8, green traces). The same muscle spindle afferent ceases firing in midstance, when the
433 intrafusal muscle is possibly activated by gamma drive, but not enough to prevent the muscle spindle
434 from falling slack as the extrafusal muscle fiber shortens due to alpha drive (Ellaway et al., 2015; Arthur
435 Prochazka, 2015; Arthur Prochazka & Gorassini, 1998) (Fig. 8, purple traces). We did not simulate such
436 experiments, but doing so would provide an excellent opportunity to further improve our framework for
437 natural movements.

438 Complex biomechanical interactions may further explain muscle spindle afferent activity that
439 have heretofore proven challenging model and to understand mechanistically. Ia afferent firing has been
440 shown in humans to be highly dependent on resistive loads applied to an attempted movement (Dimitriou,
441 2014; Edin & Vallbo, 1990; Vallbo, 1974). An interpretation of this observation could be that, in coupled
442 agonist-antagonist systems, Ia afferents under alpha-gamma coactivation or beta motor control signal the
443 effects of loads on a joint originating outside their parent muscles. For example, in an unimpeded
444 concentric contraction, the antagonist Ia afferent stretches and signals the effect on the joint torque and
445 angle caused by the agonist contraction. If, in a similar scenario, a load prevents the agonist from
446 shortening, there would be no change in the antagonist Ia afferent activity, correctly signaling the lack of
447 joint movement. Simultaneously, the agonist Ia afferents would increase their firing due to concomitant
448 fusimotor drive, signaling the impeding effects of the external load on the joint. In both cases, the Ia
449 afferent, whether in an agonist or antagonist muscle, would fire when interacting with a force that
450 originates outside the parent muscle. Whether the antagonist load or reciprocal inhibition of the fusimotor

451 drive between antagonist muscle pairs explains these results, Ia afferents in opposing muscles may work
452 together to signal joint-level information about external loads (Dimitriou, 2014). Simulating such
453 intermuscular interactions is likely necessary to predict proprioceptor activity during movement, and
454 would require multiple muscle and muscle spindle models to be coupled.

455

456 *Limitations of our current model and future directions*

457 The core framework presented provides proof of concept for more elaborate and biophysically accurate
458 model of muscle spindle firing necessary to predict muscle spindle signals during naturalistic movement
459 conditions in health and disease. Importantly, our model can be used as a platform to enable the
460 development of computational models capable of testing the effects of a wide variety of multiscale
461 mechanisms on muscle spindle firing, including architectural arrangement of the muscle spindle within
462 the muscle (Maas et al., 2009), intra- and extrafusal muscle myosin expression, more complex muscle
463 force-generating mechanisms, extracellular matrix stiffness, mechanosensory encoding mechanisms, and
464 biophysical neural dynamics that could all be affected by aging and disease.

465 Our primary focus in developing our mechanistic model was to use the simplest two-state cross-
466 bridge model of actin-myosin interactions to qualitatively, but not quantitatively, reproduce history
467 dependence in its stretch responses (Figure 6A) as seen in biological Ia afferents (Figure 2B). It should be
468 noted that, despite the discrepancies in the more abrupt stretch-shorten “triangle” stretches (Figure 6A),
469 our model was able to closely reproduce the human Ia afferent response (Figure 6B, C) during the slow
470 sinusoidal stretch, which may be more physiologically relevant. However, we acknowledge the apparent
471 shortcomings of our chosen model in reproducing the triangle responses from anesthetized rat. With our
472 approach, it is impossible to separate the elastic contributions from endomysial tissue from those in
473 myofilaments, from structures such as titin. It remains to be seen which aspects of complex muscle
474 spindle properties will improve the model’s predictive capabilities, but it may be necessary to account for
475 muscle-specific (Kenneth S. Campbell, 2014b) kinetic schemes or to replace the simple linear model of
476 titin with more sophisticated properties (Nishikawa et al., 2012) to capture calcium-dependent hysteresis

477 in regulation of titin stiffness and a range of other complex processes of muscle force generation (Kenneth
478 S. Campbell, 2016, 2017; Mann et al., 2020). We relied on the similarly timed short-range elastic
479 component of extrafusal force response to predict the putative intrafusal mechanics-driven Ia afferent
480 stretch responses. Additional sources of history-dependence could include history dependence in
481 extrafusal muscle fibers and non-contractile tissues, or the inertial of the muscle mass when rapidly
482 accelerated. Another possibility is that we failed to account for the complex process(es) responsible for
483 the occlusion between multiple transduction zones in the primary afferent (Banks et al., 1997; Mileusnic
484 et al., 2006), in which the firing rate may be more dependent on the bag1 fiber during the initial stretch
485 and become more dependent on to bag2 and chain fibers in later stretches.

486 A comprehensive understanding of the mechanisms underlying muscle spindle
487 mechanotransduction, i.e. translation of mechanical stimuli to receptor potentials, remains elusive.
488 Lacking sufficiently detailed information from the literature, we represented the entire neuromechanical
489 transduction process by a set of constant gains. However, we do recognize that the dynamics of yet
490 uncharacterized viscoelastic properties of the equatorial regions of intrafusal muscle fibers, ion channels
491 in muscle spindle afferent endings (Bewick & Banks, 2015; Carrasco et al., 2017), and intrinsic neural
492 dynamics (e.g. those underlying neural history dependence, spike-frequency adaptation, and occlusion
493 Banks et al., 1997) of the afferent plays a significant role in the resulting muscle spindle firing signals.
494 Finally, we focused only on the muscle spindle Ia afferent, but the anatomical arrangement of our model
495 could also provide a biophysical framework for understanding muscle spindle group II afferents that are
496 located at the junction between the polar and equatorial regions of intrafusal muscle fibers, as well as and
497 Golgi tendon organ Ib afferent located at the extrafusal musculotendinous junction. In future studies, it
498 would also be beneficial to adopt the approach of Mileusnic and colleagues (Mileusnic et al., 2006) and
499 optimize the parameters of our model to match the careful experimental conditions of specific muscles
500 and animal models.

501

502 *Conclusion*

503 Ultimately, it is clear that the complex interplay between central drive, internal, and
504 environmental forces cannot be ignored when decoding muscle spindle firing patterns, especially in
505 voluntary movement and during interactions with the environment. Such neuromechanical interactions are
506 key in interpreting the afferent signals and percepts in a variety of conditions where mismatches in
507 extrafusal and intrafusal mechanical state due to fatigue or pre-conditioning alter the perception of both
508 force and position (Uwe Proske et al., 2014; Uwe Proske & Gandevia, 2012), or the role of muscle
509 spindles in perception of force and weight (Jami, 1992; Luu et al., 2011; Savage et al., 2015). Our
510 multiscale biophysical model provides a new platform by which we may move beyond the idea that
511 muscle spindles can be characterized as simple passive encoders of biomechanical signals during
512 movement enabling a more sophisticated understanding of muscle spindles as proprioceptive sensors,
513 their role in neural control of movement, and how neurological disorders disrupt sensorimotor systems.

FIGURES

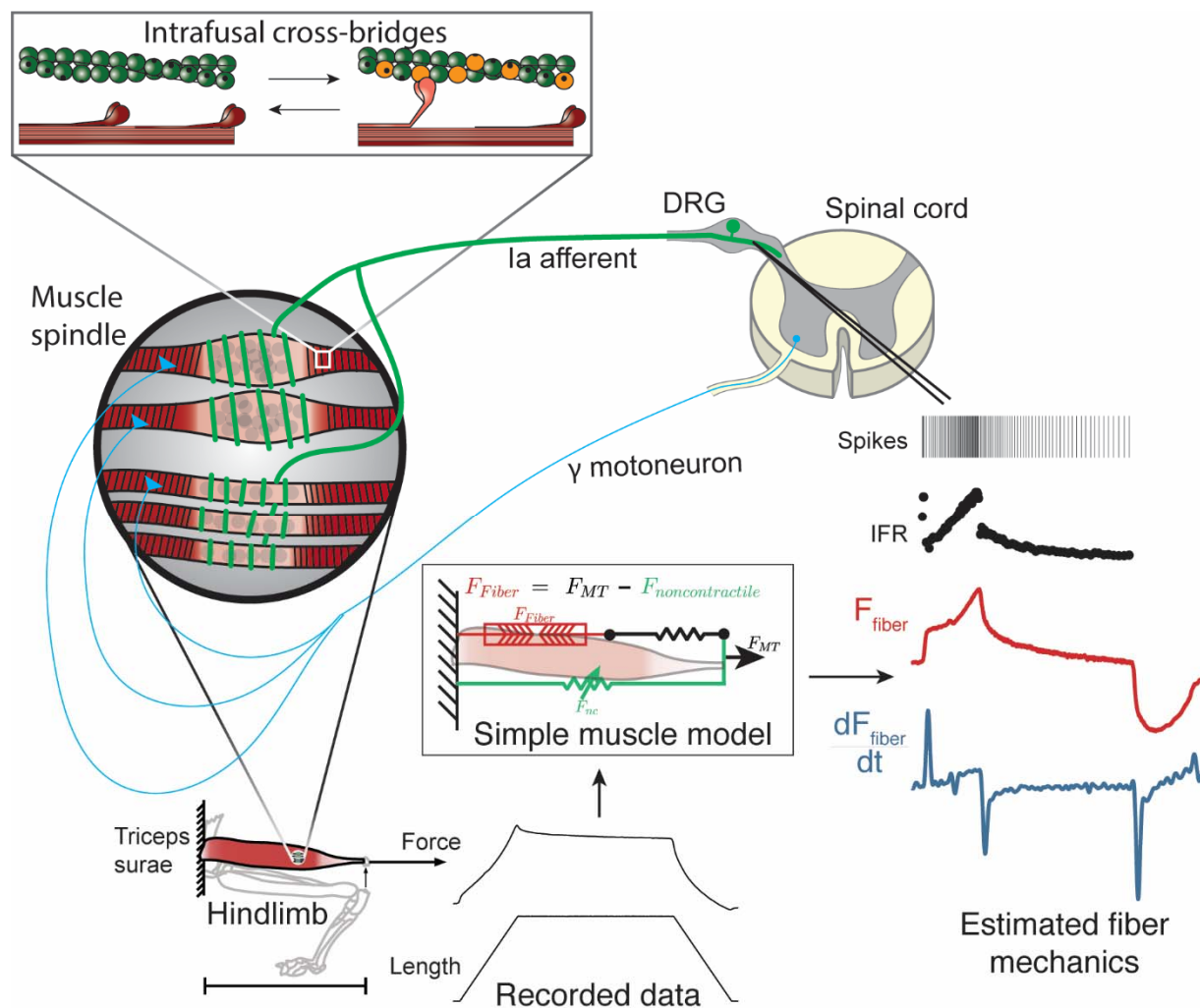


Figure 1 Overview of methodologies to test hypothesis that muscle spindle Ia afferent firing rates follow intrafusal muscle fiber force due to cross-bridge interactions. Ia afferent firing rates were recorded from dorsal rootlets during stretches of the triceps surae muscle in anesthetized rats. Muscle fiber forces were estimated by subtracting noncontractile forces from measured whole musculotendon force. The exponential rise in force with stretch was assumed to arise from non-contractile tissue in parallel with the muscle-tendon unit with exponential stiffness (Blum et al. 2019). The remaining estimated muscle fiber force and yank exhibited similar temporal characteristics to the muscle spindle IFR. Intrafusal muscle fiber force and yank were then simulated using a cross-bridge based model to predict muscle spindle IFRs.

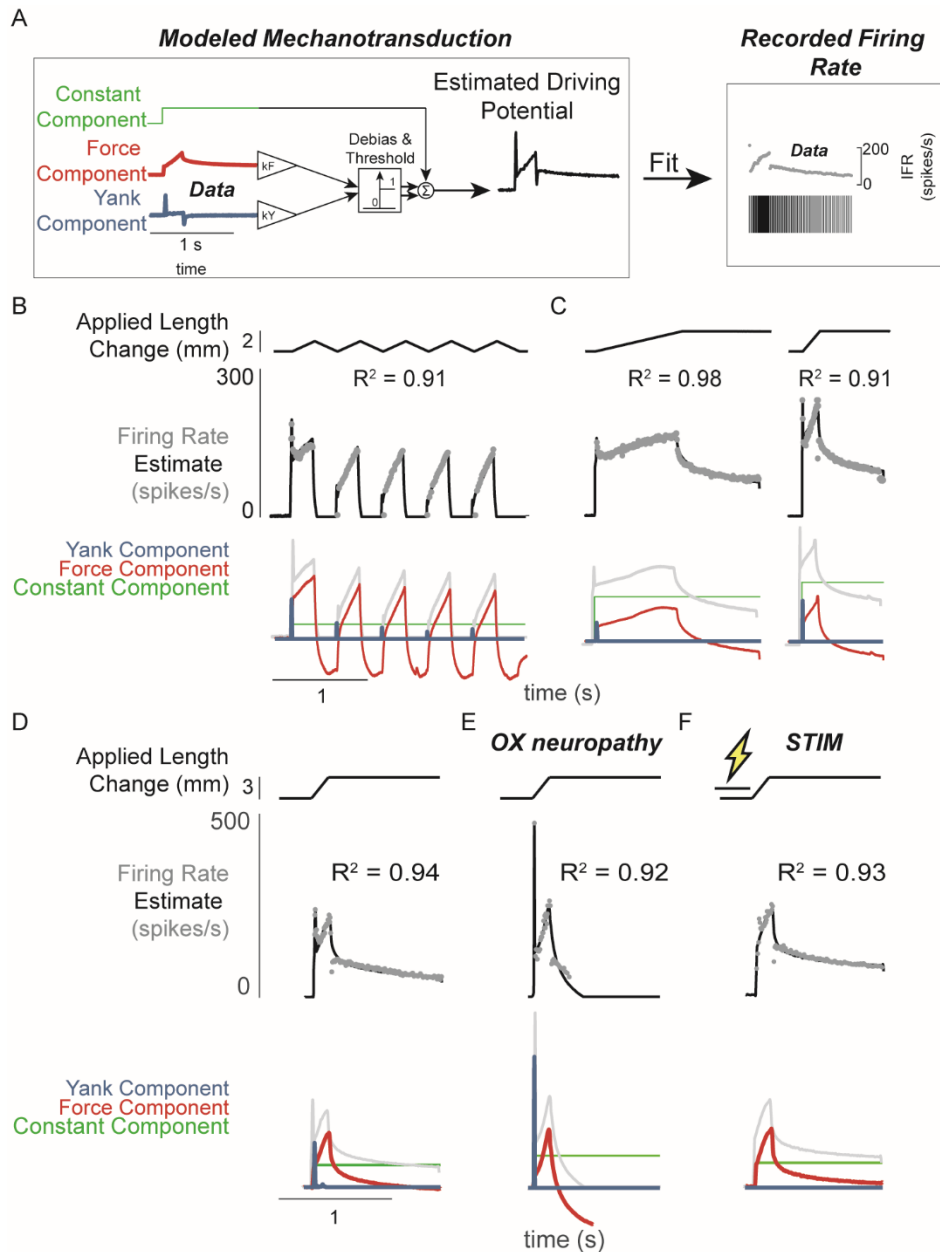


Figure 2 Muscle spindle driving potentials estimated from independent contributions of experimentally-derived muscle fiber force and yank. A) Estimated driving potentials were derived from linear combinations of muscle fiber force and yank, half-wave rectified, and compared against recorded muscle spindle Ia afferent firing rates. Weights of each component were optimized to match recorded spiking dynamics. B) Recorded muscle spindle Ia afferent firing rates (gray dots) in history-dependent conditions, having non-unique relationships to muscle length and velocity, were reproduced using muscle fiber force and yank (black lines). Notably, the initial burst and increased firing during ramp in the first stretch were attributed to increased muscle fiber yank and to greater force during the first stretch, respectively. C) Likewise, muscle fiber force and yank could also account for the temporal dynamics of Ia afferent firing in response to both slow and fast stretches. D) This model permits independent manipulation of the force and yank contributions to muscle spindle firing rates. As such, we can

explain the altered muscle spindle Ia afferent firing patterns in E) oxaliplatin chemotherapy-induced sensory neuropathy as a loss of force sensitivity, and F) after antidromic electrical stimulation of the axon as loss of yank sensitivity.

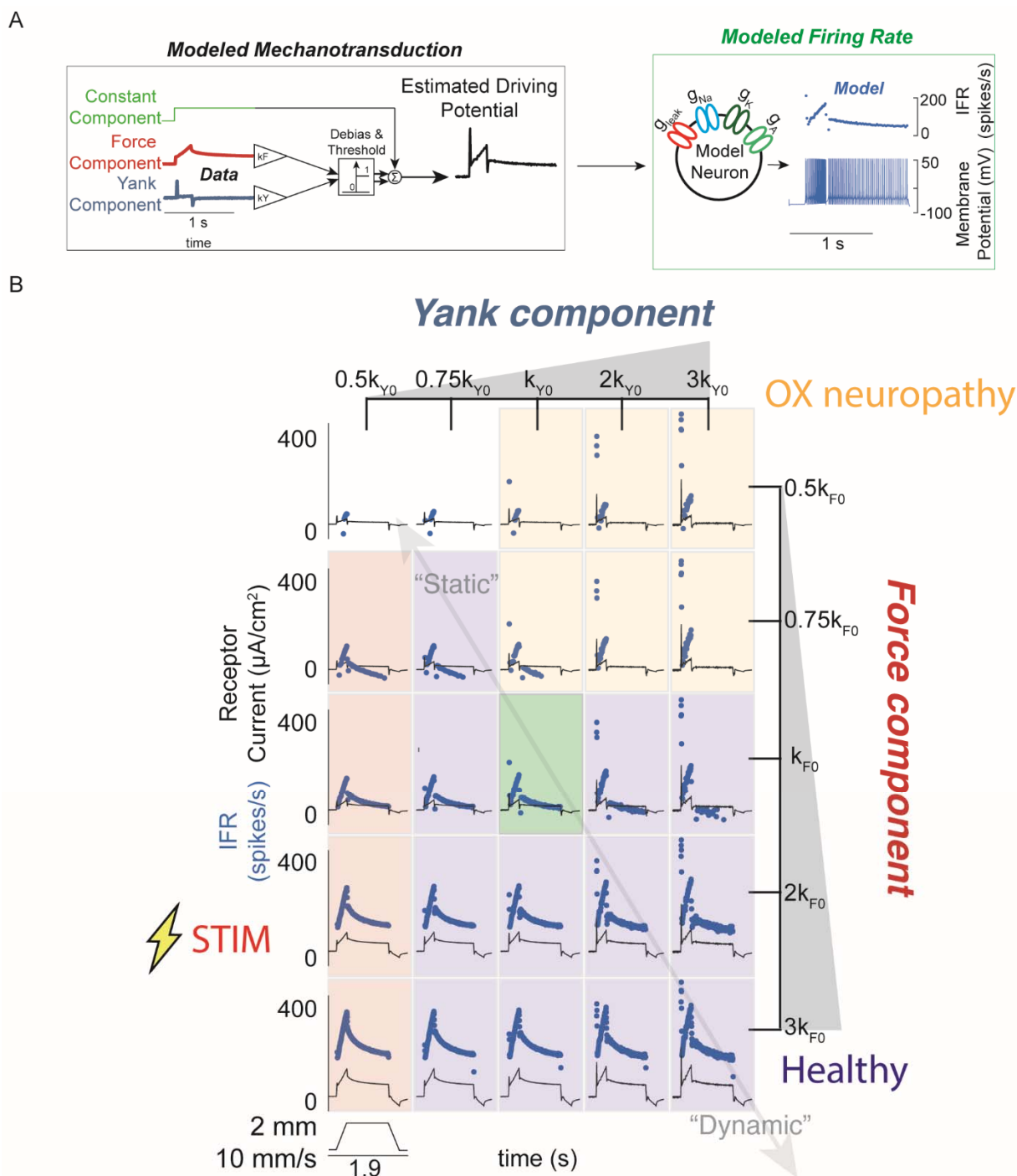


Figure 3 Spectrum of passive muscle spindle firing phenotypes accounted for by varying sensitivity to experimentally-derived muscle fiber force and yank. A) Sensitivity to experimentally-derived force and yank was systematically varied for a single stretch and resultant driving potentials were input to a Connor-Stevens model neuron to generate firing patterns. B) Nominal force and yank weights were identified to recreate experimentally-recorded muscle spindle response to a representative stretch (green box). Increasing sensitivity to yank (left to right) generated larger initial bursts and dynamic responses during the ramp, and resembled responses from oxaliplatin-treated specimens at the highest yank and lowest force

sensitivities (orange boxes, compare to Figure 2E). Increasing sensitivity to force (top to bottom) generated higher firing rates during the hold period and resembled Ia afferent firing responses after axonal stimulation at the lowest yank and highest force sensitivities (red boxes, compare to Figure 2F). Varying the weights of the force and yank sensitivities could recreate the spectrum of healthy muscle spindle firing profiles reported classically (purple boxes).

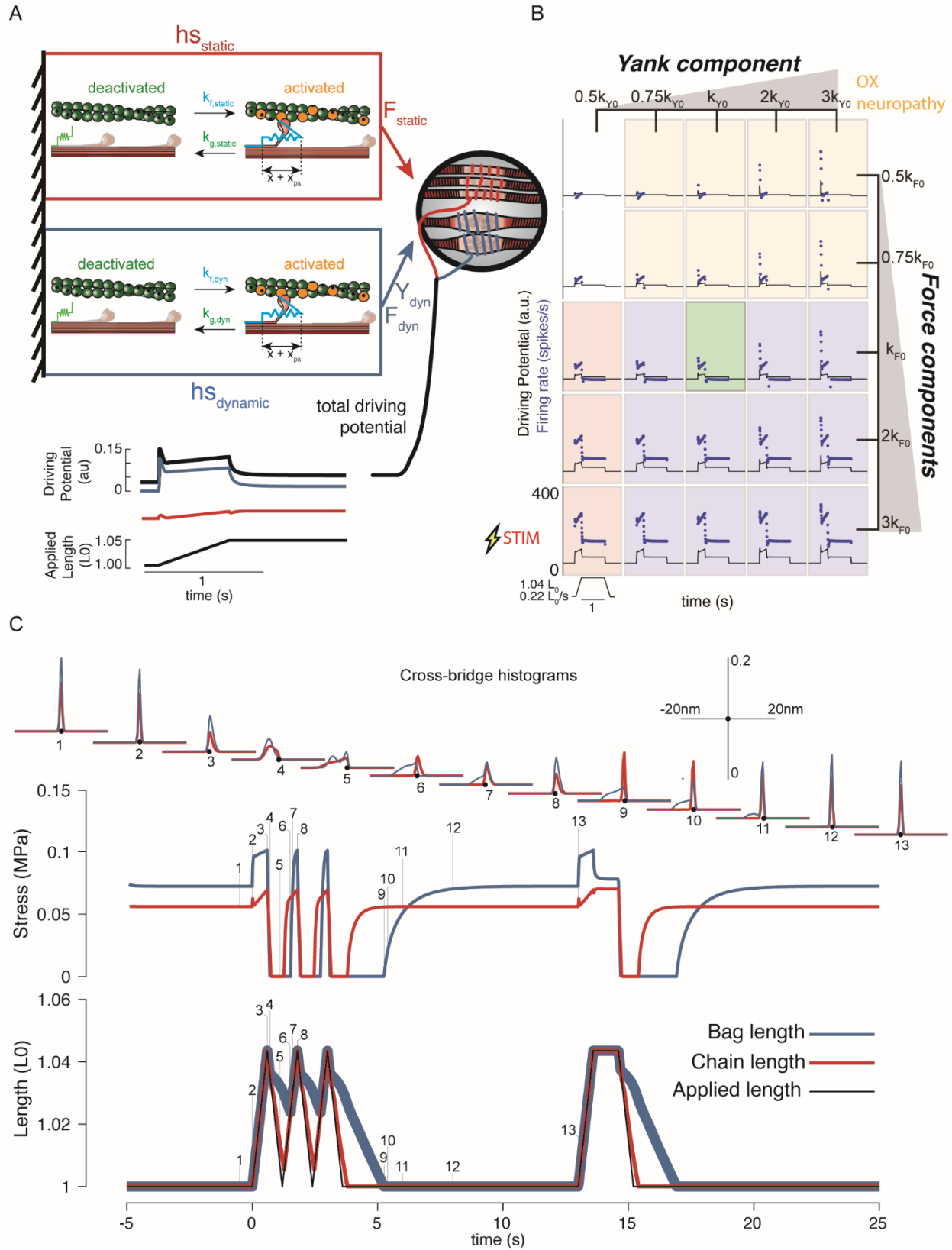


Figure 4 Generative model of muscle spindle driving potentials based on simulated muscle cross-bridge kinetics. A) The muscle spindle model consists of two muscle fibers in a parallel mechanical arrangement, loosely representing intrafusal bag and chain fibers. B) During stretch, force and yank of the “dynamic” fiber is linearly combined with force of the “static” fiber, with different proportions generating driving potentials consistent with “dynamic” and “static” muscle spindle firing response phenotypes. C) A population of myosin cross-bridges and their relative displacement and velocity with respect to active actin binding sites was simulated during three consecutive ramp-stretch-release stretches. The distribution of cross-bridge lengths relative to actin binding sites is shown at different timepoints of imposed kinematics (numbered graphs and timepoints). The length of the dynamic and static fibers (lower trace) and the stress in the dynamic and static fibers (middle trace) is shown. Deviations between applied length and muscle fiber length occur due to muscle fiber slack.

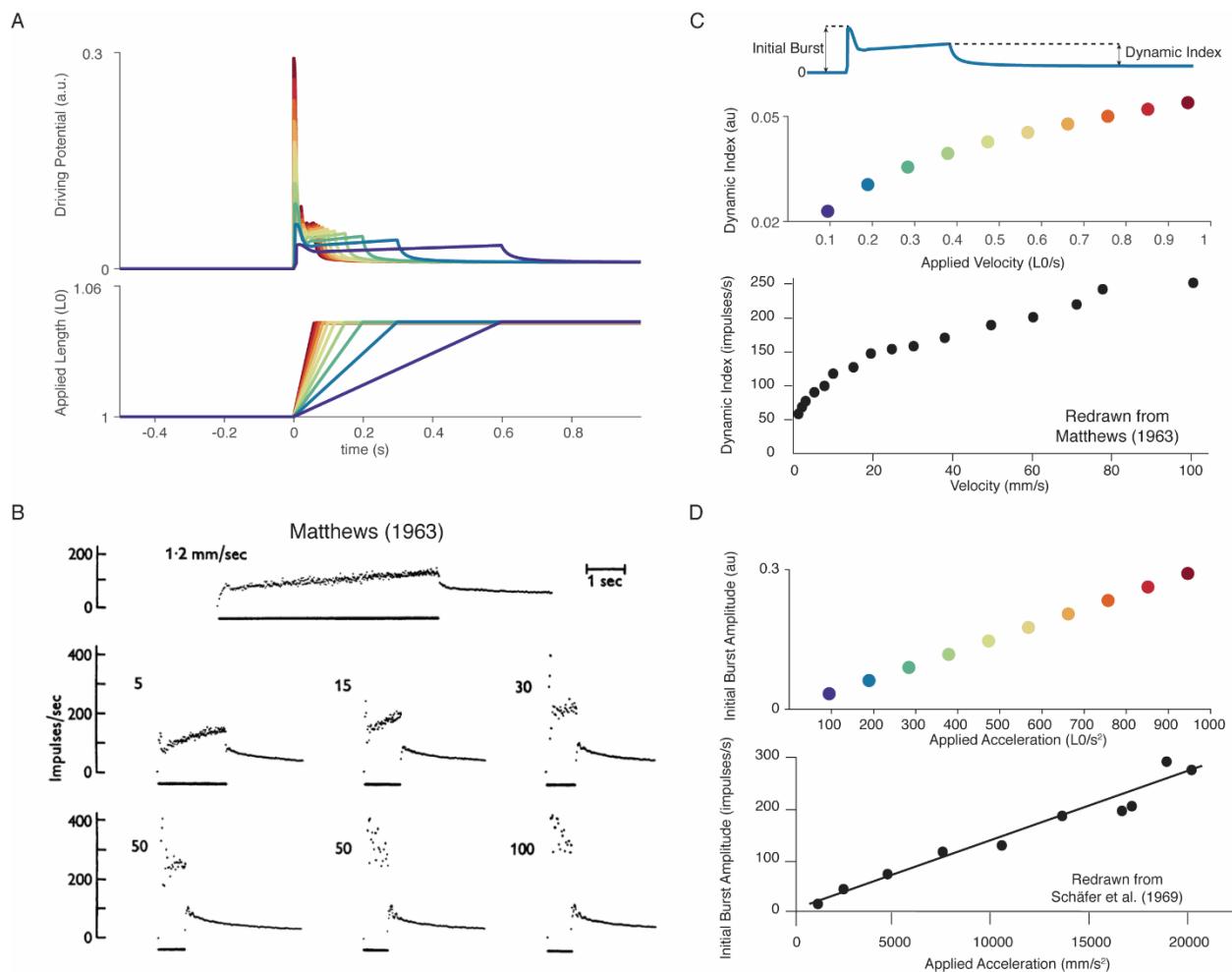


Figure 5 Muscle spindle firing dynamics emerge from cross-bridge mechanisms during simulated muscle stretch. Simulations assume a low level of muscle cross-bridge cycling consistent with relaxed muscle. Length displacements were imposed on the muscle fiber. A) Predicted driving potentials (upper traces) during ramp stretches of varying velocity and acceleration (lower traces). B) Classical data showing Ia afferent firing modulation with different stretch velocity (Matthews 1963). C) Dynamic index emergent from cross-bridge mechanisms. Dynamic index is defined classically as the ratio of firing rate at the end of the ramp phase and the firing rate 0.5 seconds into the hold phase (upper diagram). The muscle spindle model exhibits a sublinear relationship between dynamic index and stretch velocity (middle plot – colors correspond to A), similarly to classical findings (bottom plot). E) Linear acceleration scaling of initial burst emergent from cross-bridge mechanisms. Initial burst amplitude is defined as the difference between peak firing rate during initial burst and baseline. Muscle spindle model exhibits linear scaling with stretch acceleration at stretch onset (top plot), which is consistent with classical findings (bottom plot – Schäfer 1969). Figures reproduced from Matthews (1963) and Schäfer et al. (1969).

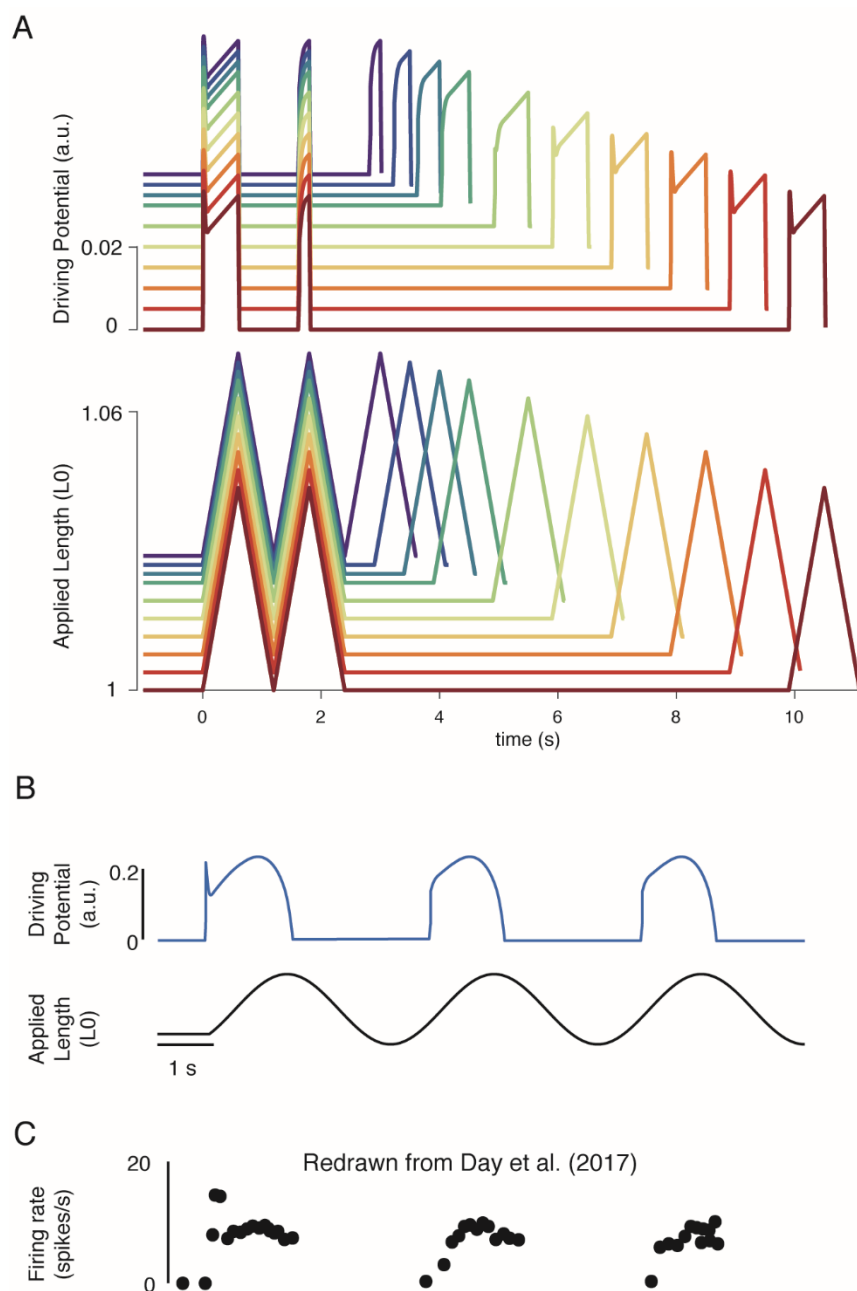


Figure 6 Emergent movement history-dependence of muscle spindle. A) When repeated stretch-shorten length cycles were applied, a larger response was predicted if the length was held constant prior to stretch (bottom plot – all traces). An abolished initial burst and reduced dynamic response were predicted in the second stretch, immediately applied after the length was returned to the initial value (top plot – all traces). In the third stretch, recovery of the initial burst was dependent on the time interval between the second and third stretch, with the effect saturating between 5-10 s (top plot, recovery from violet to red traces). This finding predicts the results found by Haftel et al. (2004) in rat muscle spindles. Similarly, dynamic response recovered gradually with time interval between second and third stretch (top plot, recovery from

violet to red traces). This finding predicts the results found by Proske and Stuart (1985) in toad muscle spindles. B) Sinusoidal displacement imposed from rest elicited a history-dependent initial burst in the predicted muscle spindle driving potential at the onset of stretch, resembling data from C) human muscle spindles recorded during the application of sinusoidal motion to the ankle in relaxed conditions (Day et al. 2017).

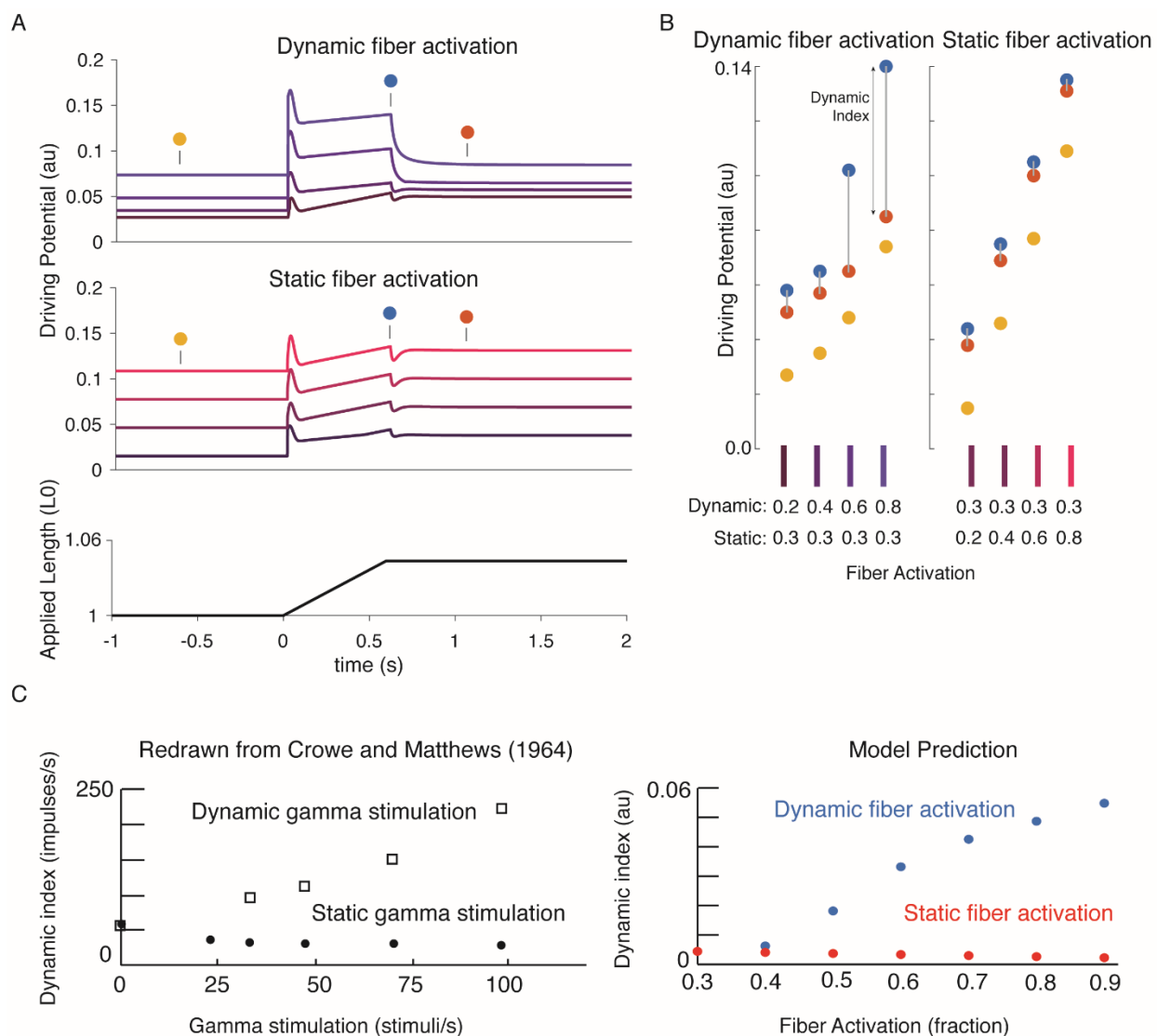


Figure 7 Changes in muscle spindle sensitivity caused by central drive emergent from interactions between dynamic and static fibers. Fusimotor activity was imposed on either the dynamic or the static fiber by increasing the number of active actin binding sites in the appropriate fiber. A) Simulated dynamic fiber activation increased the driving potential predominantly during the ramp, with smaller increases during the background and hold period (top traces). Simulated static fiber activation predominantly increases the driving potential rate during the background and hold period, with only modest increases in during the ramp. B) Emergent scaling of the dynamic index with dynamic (increase in dynamic index) and static fiber activation (decrease in dynamic index) resembled trends reported previously in the literature with C) dynamic index increasing with bag fiber activation, and dynamic index decreasing with chain fiber activation, respectively. Figures adapted and reproduced from Crowe and Matthews (1964) with permission.

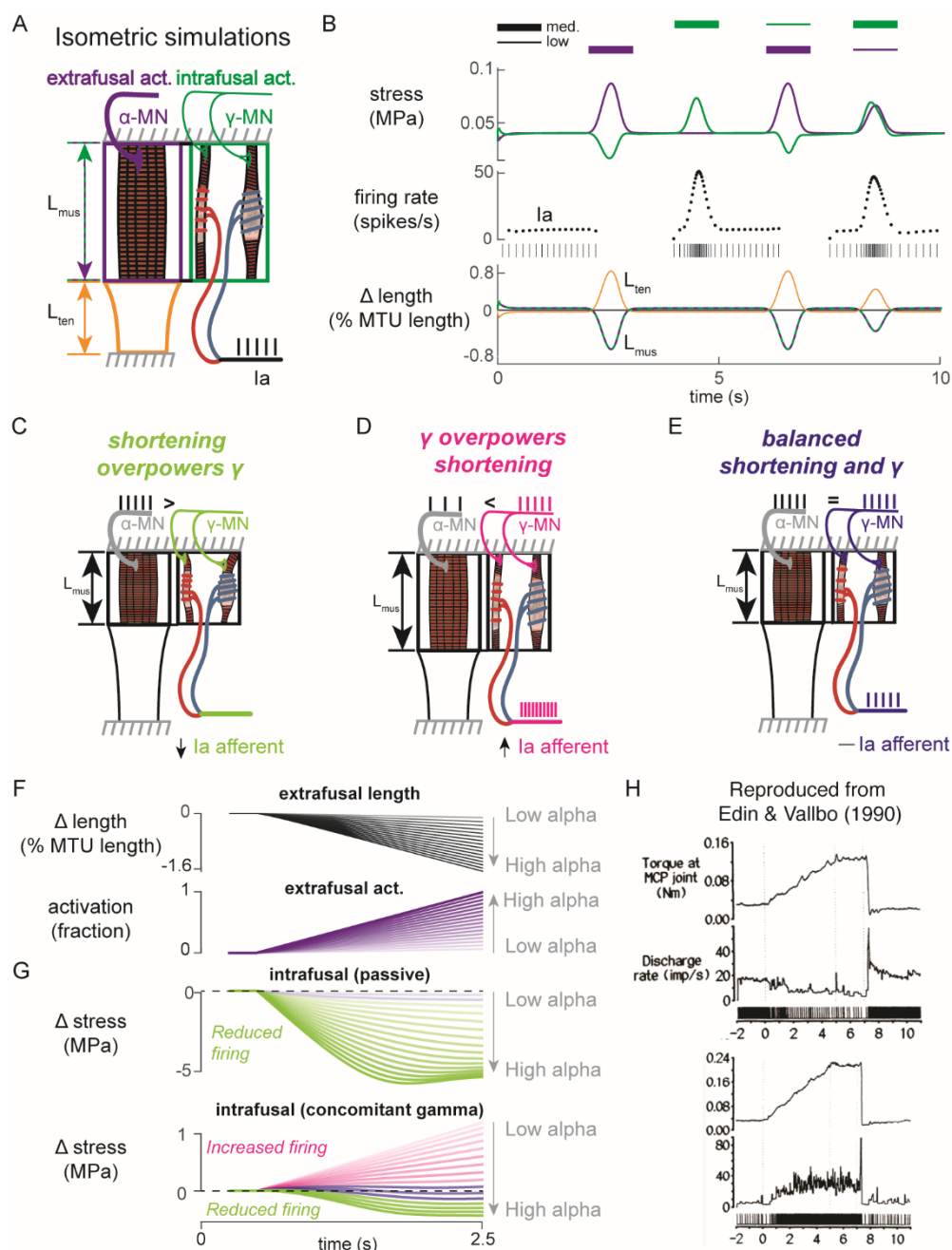


Figure 8

Paradoxical muscle spindle firing activity during isometric musculotendon conditions emergent from multiscale muscle mechanics. A) Framework of isometric musculotendon simulations. We simulated an in-series muscle-tendon with ends fixed to its initial length. The extrafusal muscle length was shared with the intrafusal muscle fibers, which were activated independently from the extrafusal muscle. B) Decoupling of intra- and extrafusal force during short isometric contractions. Different combinations of extrafusal (purple bars) and intrafusal (green bars) activation were simulated. Top row contains intra- (green) and extrafusal (purple) stress during these conditions. Middle row is the simulated Ia afferent firing rate, generated by using our mechanistic model simulated driving potential combined with a leaky integrate-and-fire neural

model. Extrafusal activation shortens the extrafusal muscle, and in the absence of sufficient intrafusal activation, silences the Ia afferent firing. Intrafusal activation by itself causes the Ia afferent to fire more. When both intra- and extrafusal muscle are activated concomitantly, the relative weighting of activity determines whether the Ia afferent will increase or decrease its firing rate. C-E) Schematic representation of the multi-scale mechanics responsible for Ia afferent behavior in these simulations. When the effects of extrafusal shortening overpower the effect of fusimotor activity on the spindle, the Ia afferent firing rate decreases, or stops altogether (left, green). When fusimotor activity overpowers the effect of extrafusal shortening on the spindle, the Ia afferent firing rate increases (center, pink)). In theory, this suggests there are combinations of fusimotor activity and shortening that will result in no net change to the Ia afferent firing, despite the dynamics of the muscle (right, blue). F) A spectrum of musculotendon isometric conditions arising from the interplay of extrafusal shortening and intrafusal activation. We activated the musculotendon model with ramp activations and simulated the intrafusal muscle force with and without concomitant gamma activation. When increasing alpha motor neuron drive (purple traces; directionally indicated in all plots by an arrow), extrafusal muscle shortened more (top plot). G) The passive spindle always decreased its force because of the shortening imposed by the surrounding extrafusal muscle (top plot). When the intrafusal muscle was activated with concomitant drive, similar to the largest amplitude ramp of extrafusal muscle, we observed the spectrum of responses predicted in C-E: the intrafusal force either decreased (green), stayed relatively constant (purple), or increased (pink). H) Experimental examples of one muscle spindle decreasing its firing rate (top traces) and another increasing its firing rate (bottom traces) during an isometric task in human. Figure reproduced from Edin and Vallbo (1990) with permission.

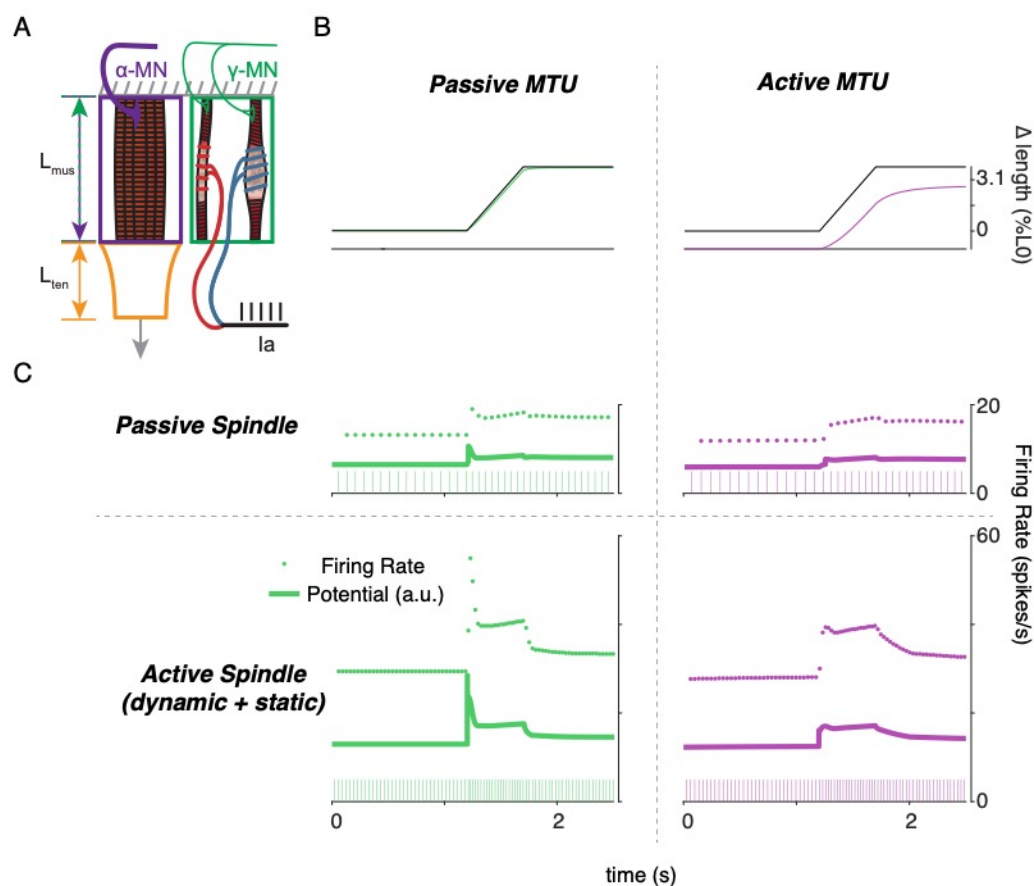


Figure 9

Complex interplay of movement, alpha, and fusimotor drive underlie muscle spindle Ia firing rates. A) Framework of free-end musculotendon simulations. We simulated an in-series muscle-tendon with one end fixed and one end free to move. The extrafusal muscle length was shared with the intrafusal muscle fibers, which were activated independently from the extrafusal muscle. B) Passive musculotendon has a relatively compliant muscle, leading to a MTU length change being primarily applied to the muscle, not the tendon. Extrafusal activation creates a relatively stiff muscle, leading to a smaller length change for the same MTU length applied. C) Musculotendon activation changes the coding properties of both passive and active muscle spindles. When the MTU is passive, a passive muscle spindle has an initial burst.

514

515 **METHODS**

516 *Animal care*

517 All procedures and experiments were approved by the Georgia Institute of Technology's
518 Institutional Animal Care and Use Committee. Adult female Wistar rats (250–300 g) were studied in
519 terminal experiments only and were not subject to any other experimental procedures. All animals were
520 housed in clean cages and provided food and water ad libitum in a temperature- and light-controlled
521 environment in Georgia Institute of Technology's Animal facility.

522

523 *Terminal physiological experiments*

524 Experiments were designed to measure the firing of individual muscle afferents in response to
525 controlled muscle stretch using conventional electromechanical methods applied in vivo as reported by
526 Vincent et al. 2017 (see Fig. 1). Briefly described, rats were deeply anesthetized (complete absence of
527 withdrawal reflex) by inhalation of isoflurane, initially in an induction chamber (5% in 100% O₂) and
528 maintained for the remainder of the experiment via a tracheal cannula (1.5–2.5% in 100% O₂).
529 Respiratory and pulse rates, P_{co2}, P_{O2}, blood pressure, and core temperature were maintained within
530 physiological ranges variously by injecting lactated ringers subcutaneously and by adjusting anesthesia
531 and radiant heat. Rats were secured in a rigid frame at the snout, lumbosacral vertebrae, and distally on
532 the left femur and tibia (knee angle 120°). Dorsal roots S1-L4 were exposed by laminectomy and
533 supported in continuity on a bipolar recording electrode. Triceps-surae muscles in the left hindlimb were
534 dissected free of surrounding tissue and their tendon of insertion was cut and attached to the lever arm of
535 a force and length-sensing servomotor (Model 305C-LR, Aurora Scientific Inc.). Triceps surae nerves
536 were placed in continuity on monopolar stimulating electrodes, and all other nerves in the left hindlimb
537 were crushed to eliminate their contribution to muscle-stretch evoked action potentials recorded in dorsal
538 roots. Exposed tissues were covered in pools of warm mineral oil Surgical and recording preparation

539 followed by data collection (see below) lasted for up to 10 hrs. At the conclusion of data collection, rats
540 were killed by exsanguination preceded either by overdose with isoflurane inhalation (5%).

541 Data collection centered on trains of action potentials recorded from individual primary sensory
542 neurons responding to length perturbations of the parent triceps surae muscles. Action potentials were
543 recorded intracellularly from axons penetrated by glass micropipettes in dorsal root. Sensory axons were
544 identified as triceps surae when responding with orthodromic action potentials evoked by electrical
545 stimulation of the corresponding muscle nerves. Axons selected for further study were classified as group
546 Ia when (a) isometric twitch contractions interrupted muscle stretch-evoked firing, (b) each cycle of high
547 frequency (>200Hz), low amplitude (80 μ m) muscle vibration evoked action potentials, and (c) muscle
548 stretch produced an initial burst firing of high frequency (>100Hz) firing at the onset of rapid muscle
549 stretch.

550 Ia afferents were studied for their firing responses to two forms of muscle stretch beginning from
551 and returning to $L_{r(\text{rest})0}$, which was the muscle length corresponding to 10 grams passive muscle force:
552 (1) ramp-hold-release stretch (20mm/s constant velocity ramp and release with intervening hold to 3 mm
553 length) repeated every 4 secs and (2) sequential triangular ramp-release stretches to 3 mm at constant
554 velocity (4 mm/s). Records of action potentials, muscle force and length were digitized and stored on
555 computer for later analysis.

556 In the 11 Ia afferents for which the initial pseudolinear model analyses were performed, a range
557 of 6-99 stretch trials with varying maximum length and velocity were achieved depending on the
558 recording stability. To ensure sufficient information for statistical measures, we required that stretch trials
559 have at least 50 recorded action potentials in order to be included in statistical analyses. Stretch trials
560 where spikes were not discriminable were excluded. These criteria yielded suitable datasets for 11
561 individual afferents from 5 animals for pseudolinear model analyses and 6 individual afferents from 5
562 animals for the axonal stimulation analyses. We also included 3 afferents from 3 rats treated with
563 oxaliplatin (Bullinger et al., 2011b).

564

565

566

567 *Muscle fiber force estimation*

568 To isolate the component of recorded musculotendon force arising from the muscle fibers (used
569 as a proxy for intrafusal muscle force), we assumed an idealized musculotendon mechanical arrangement
570 (Fig. 1). In summary, we assumed there was noncontractile passive connective tissue arranged
571 mechanically in parallel with the muscle fibers and removed it analytically, as previously described
572 (Blum et al., 2019).

573 Briefly, we assumed the noncontractile tissue acted as a nonlinear spring of the form:

574

$$F_{nc} = k_{lin}(L - L_0) + Ae^{k_{exp}(L-L_0)}$$

575 where k_{lin} , k_{exp} , and A are greater than or equal to zero. Once parameters were selected by the
576 optimization procedure, the estimated noncontractile tissue forces were subtracted from the recorded
577 force to estimate the muscle fiber force, which was fit to the IFRs.

578

579 *Pseudolinear models for predicting firing responses*

580 We predicted spiking responses using pseudolinear combinations of either recorded
581 musculotendon length-related (length, velocity, and acceleration) or force-related (estimated muscle fiber
582 force and yank) variables (Fig. S1). The relative weights and offsets for each variable in a model were
583 optimized to minimize the squared error between the model prediction and Ia spike rates on a per-trial
584 basis.

585 For both the force- and length-related models, we fit the estimated IFR for each model to the IFR
586 of the afferent for each stretch trial included in our analyses (for all 20 afferents presented in this study).
587 The model parameters, consisting of a weight (k_i) and offset (b_i) for each force- or length-related variable

588 included in the sum, were found via least-squares regression using MATLAB's optimization toolbox
589 (*fmincon.m*) and custom scripts.

590 We observed a peak to peak delay from the whole musculotendon yank and the initial burst,
591 likely caused by delayed force transmission to the intrafusal fibers from the tendon (Blum et al., 2019). A
592 time delay (λ_j) was determined by shifting the timestamp of the variables forward relative to the IFR data
593 to be fit (note: this time delay was 0 for all variables except yank, to account for the apparent delay
594 between the onset of muscle force response and the onset of the spiking response). The general form of
595 the models was:

$$IFR_{j,n}(t) = \left(\sum_{i=1}^n k_i \cdot ([x_i(t - \lambda_j)] + b_i) \right) \quad (1)$$

596 where the IFR estimate of the j th model for the n th perturbation was estimated by a sum of n force- or
597 length-related variables, offset by a single value, b_i , and scaled by a gain, k_i . $[]$ denote positive values of
598 the argument. Model estimates for IFR were related to the recorded IFR of the m th afferent by the
599 equation:

$$IFR_{j,n}(t) + e(t) = IFR_{m,n}(t) \quad (2)$$

600 Error, $e(t)$, was minimized by finding the set of parameters for each model that minimizes a measure
601 related to $e(t)^2$.

602

603 *Antidromic Axonal stimulation dataset*

604 To test whether the force and yank components could arise from separate mechanosensitive
605 mechanisms, another set of experiments was performed on 6 additional afferents in 4 animals. Each
606 dataset consisted of three ramp-hold-release stretches repeated at 4 sec intervals. The second of three
607 stretches was preceded by antidromic action potentials in a 50 Hz train lasting 500 ms. These action
608 potentials were isolated exclusively to the Ia afferent under study by injecting suprathreshold current

609 pulses (0.5ms) directly into the axon through the micropipette located in the dorsal root. The first and
610 third stretches were unconditioned and served as bookend controls

611 For each trial in these 6 afferents, we found the best-fit prediction for the force-related model
612 using the parameter optimization described earlier. For the pre- and post-stimulation control trials, we
613 first fit the model without a yank component, and then refit the model with a yank component. For the
614 trials in which the electrical stimulus was applied, the yank component was set to be zero and the force
615 and constant components were optimized as described before.

616 We performed one-way ANOVA on model performance (R^2), yank sensitivity (k_Y), force
617 sensitivity (k_F), and the constant component (C) across 5 groups of model fits: pre-stimulus control trials
618 without (1) and with (2) yank sensitivity, stimulus trials (3), and post-stimulus trials without (4) and with
619 (5) yank sensitivity. We used the Tukey-Kramer method to examine all pairwise comparisons between
620 groups.

621

622 *Oxaliplatin dataset*

623 We used data collected previously to test whether force and yank components were altered by
624 oxaliplatin chemotherapy alone (Bullinger et al., 2011b). The effect of oxaliplatin on sensory coding of
625 Ia afferents has been well-documented, so we analyzed three afferents from different animals. We fit the
626 muscle fiber force-related model (described above) to three stretch trials for each afferent (3mm, 20
627 mm/s). We performed one-way ANOVA on model performance (R^2) between model fits with and without
628 yank for each afferent to test the significance of the yank component on model performance.

629

630 *Applying estimated fiber force-related driving potentials to model neuron*

631 To test the feasibility of the force, yank, and constant components of the muscle fiber force-
632 related model as mechanical signals encoded by the muscle spindle receptor, we applied a range of
633 combinations of components to a conductance-based model neuron (based on the Connor-Stevens model;
634 see next section) and examined the resulting firing rates. We first estimated the muscle fiber force and

635 yank, as described previously, and varied the relative gains of these signals before adding them with a
636 constant component. Once the components were added together, they were half-wave rectified, and
637 applied to the model neuron as a stimulus current.

638 Model neuron sensitivities to these components were tuned until the model instantaneous firing
639 rate was within 10 spikes/s for initial burst, dynamic response, and final plateau. We treated the parameter
640 values which produced this response as the nominal values for the model. The relative sensitivities of the
641 model neuron to force and yank component were then swept from 10-600% of their respective nominal
642 values. We then compared the resulting changes in predicted firing rates with different phenotypical
643 muscle spindle responses observed from these and other experiments.

644

645 *Responses of muscle spindle Ia afferents to stretch*

646 Consistent with prior studies, all Ia afferents exhibited initial bursts at onset of applied stretch,
647 followed by a dynamic response during constant velocity stretch, and a period of rate adaptation during
648 the subsequent isometric hold period. When repeated ramp-release length-changes were applied to the
649 muscle, an initial burst and dynamic response was present during the first ramp, but the initial burst was
650 absent and dynamic response was reduced during subsequent stretches—a phenomenon in Ia afferents
651 known as history-dependence (cf. Haftel et al. 2004).

652 The population of 11 Ia afferents considered for the first analysis varied in sensitivity to stretch
653 length, velocity, and acceleration. More dynamic afferents, as quantified by dynamic index (P. B. C.
654 Matthews, 1963a) typically had relatively large spike responses during positive velocity stretch. More
655 static afferents exhibited more firing during the plateau phase of stretch, with relatively smaller dynamic
656 indices. The population of afferents also exhibited a range of initial burst amplitudes in response to
657 stretch. There was no clear relationship between the dynamic index and initial burst amplitudes for a
658 given afferent. Despite the differences in sensitivity amongst the afferent population, the waveforms of

659 afferent responses to the same stretch stimuli contained the same features (i.e. all afferents exhibited
660 initial bursts, dynamic responses, and rate adaptation to varying degrees).

661

662 *Conductance-based model neuron for reproducing spiking activity*

663 To demonstrate the plausibility of force- and yank-related ionic currents caused by stretch, we
664 used a modified Connor-Stevens conductance-based model neuron to model the transformation of graded
665 receptor potentials into action potentials by the afferent (Connor & Stevens, 1971). The model neuron
666 contained a fast sodium, delayed rectifier potassium, transient potassium, and leak conductances
667 implemented in Simulink using built-in differential equation solvers (ode23s.m).

668

669 *Intrafusal muscle model*

670 All intrafusal muscle model code and scripts used to generate the data presented in this study are
671 available at <https://github.com/kyleplum/MechanisticSpindleManuscript.git>. If interested in using the
672 model as implemented in this study, a description of how to use it is also available there.

673 To simulate the hypothesized history-dependent mechanisms of intrafusal muscle fibers, we used
674 a computational model of cross-bridge cycling. We implemented a model in MATLAB based on a
675 simplified structure of the model developed by Campbell (K. S. Campbell & Lakie, 1998; K. S. Campbell
676 & Moss, 2000; Kenneth S. Campbell, 2014b; Kenneth S. Campbell & Moss, 2002) (Supplemental Fig.
677 5). *Instead of simulating the coupled dynamics between myosin heads and actin binding sites (Kenneth S.
678 Campbell, 2014b), we focused on the thick filament kinetics and greatly simplified the thin filament. In
679 brief, the force in each half sarcomere was calculated as a sum of two components: an active component
680 generated by the cycling activity of a population of myosin heads and a passive component generated by a
681 simulated linear spring modelling the contributions of titin in the half sarcomere:*

682

$$Total\ Force = \int_{-\infty}^{+\infty} (k_{cb} \rho f(x) (x + x_{ps}) dx) + k_{pas}(l_{hs} - l_0)$$

683

684 The model calculates the force of a half sarcomere at each time step by adding the forces generated from
 685 each myosin head attached to an actin binding site (active force) with the elastic force of titin (passive
 686 force). The active force is calculated as the fraction of attached myosin heads, f , multiplied by the
 687 number density, ρ , unit stiffness of a single attached actin-myosin cross-bridge, k_{cb} , multiplied by the
 688 length of the cross-bridges, $x + x_{ps}$ (where x_{ps} represents the additional displacement of an attached
 689 cross-bridge required to generate a “powerstroke” force when multiplied by k_{cb}), integrated across cross-
 690 bridge lengths. The passive force is calculated as the length of the half sarcomere, l_{hs} , relative to a
 691 reference length, l_0 , multiplied by a linear stiffness, k_{pas} .

692 To simplify the coupled dynamics of Campbell, we decided to control intrafusal activation
 693 directly instead of ionic calcium concentration:

$$\frac{df_{act}(t)}{dt} = u(t),$$

694 Where the change in the fraction of activated actin binding sites, $df_{act}(t)$, is determined by user input
 695 $u(t)$.

696 The myosin cycling dynamics were modelled using a two-state system in which each myosin
 697 head could either be attached as a cross-bridge with length x , or detached. The numbers of myosin heads
 698 in each of these states were governed by the system of partial differential equations, as in Myosim:

$$\frac{\delta A(x, t)}{\delta t} = k_f(x)D(t) - k_g(x)A(x, t)$$

$$\frac{\delta D(t)}{\delta t} = \int_{-\infty}^{+\infty} k_g(x)A(x, t)dx - \int_{-\infty}^{+\infty} k_f(x)D(t)dx$$

699 where $A(x, t)$ is the number of attached myosin heads attached to actin binding sites at time t extended by
 700 length x and $D(t)$ is the number of myosin heads in the detached state. The rate equations $k_f(x)$ and

701 $k_g(x)$ are the forward and reverse rates of myosin attachment, respectively, and are both functions of
702 cross-bridge length, x . The partial differential equations were simplified into a system of ODEs by
703 solving the time-dependent equations simultaneously at each cross-bridge length. The resulting system of
704 ODEs was solved using a built-in Matlab ODE solver (*ode23.m*). Thick filament dynamics are coupled to
705 the activation of the thin filament by the number of cross-bridges participating in cycling:

$$n(t) = \int_{-\infty}^{+\infty} A(x, t) dx + D(t)$$

706 where the fraction of bound cross-bridges is defined by the thick filament:

$$f_{bound}(t) = \int_{-\infty}^{+\infty} A(x, t) dx,$$

707 and the number of detached cross-bridges is defined by the thin filament activation and $f_{bound}(t)$:

$$D(t) = f_{act}(t) - f_{bound}(t).$$

708 There are two modes of operation for these muscle models: length control model and slack mode.
709 The model switches between the two modes automatically, based on whether each intrafusal fiber is taut
710 (force greater than zero) or slack (force equal to zero).

711 Length control mode is used by the model when the half sarcomere length is equal to the end-to-
712 end length of the sarcomere (i.e., the half sarcomere is not slack). When in length control mode, the
713 command length is applied directly to the half sarcomere, and the system of differential equations is
714 solved as described earlier. The half sarcomere is updated in the following order: 1) The length of the half
715 sarcomere and calcium concentration are each updated by the change in command length and change in
716 command calcium concentration applied to the model, respectively, 2) the amount of filament overlap is
717 updated, 3) thin filament dynamics are updated, 4) the population of myosin heads are allowed to evolve,
718 and 5) the population of myosin heads is moved according to the change in length of the half sarcomere.

719 Slack mode contains several extra steps for solving the equations governing the model. At every
720 time step, the model uses an iterative search (*fzero.m*) to find the length at which the force in the current-
721 state half sarcomere would reach zero and compare it to the command length. If the command length is
722 smaller than the “slack length” calculated, the model enters into slack mode; otherwise, the model stays in
723 length control mode.

724 Once the model enters into slack mode, the model first evolves the cross-bridge distribution in
725 time based on the system of differential equations described earlier. The model then repeats the iterative
726 search to find the slack length for the half sarcomere given the new cross-bridge distribution and shifts it
727 to the slack length for the half sarcomere. At the next time step, the model repeats the comparison
728 between the command length and the slack length for the current state of the half sarcomere. If the
729 command length is greater than the slack length, the model returns to length control mode; otherwise, the
730 model stays in slack mode.

731

732 *Intrafusal muscle model parameters*

733 Model parameters were either chosen to match the default parameters from Campbell (Kenneth S.
734 Campbell, 2014a), so the model would exhibit history-dependence at the time-course measured in this
735 work, or based on the limited qualitative information we have regarding intrafusal muscle fibers (Thornell
736 et al., 2015b). All simulations used the same set of model parameters (Supplemental Table S1).

737 Myosin attachment and detachment rates equations, $k_f(x)$ and $k_g(x)$, were selected such that the
738 force response of the model would exhibit history-dependent features consistent with observations in both
739 permeabilized muscle fibers and instantaneous firing rates of muscle spindle Ia afferents. The bag and
740 chain fibers were differentiated by relatively slower myosin attachment kinetics in the bag fiber and a
741 more compliant passive elastic element (M H Gladden & Boyd, 1985; Margaret H. Gladden, 1976;
742 Thornell et al., 2015a). Proske (U. Proske et al., 1992; Uwe Proske et al., 2014; Uwe Proske & Gandevia,
743 2012) hypothesized that history-dependent muscle spindle IFRs (and corresponding perceptual errors) are

744 caused by a population of cross-bridges within the intrafusal muscle that are unable to “keep up” with the
745 rate of shortening during an imposed movement, causing the intrafusal muscle fibers to fall slack. To
746 model this behavior, we selected rate equations that would produce relatively slow cross-bridge
747 reattachment during shortening, but would retain other desired characteristics, such as short-range
748 stiffness.

749

750 *Model of muscle spindle responses to stretch of intrafusal muscle*

751 To model the transformation of intrafusal muscle fiber stress into a firing waveform, we used a
752 pseudolinear combination of force and its first time-derivative, yank, based on previously published
753 observations (Blum et al., 2017). Our model consists of two intrafusal muscle fiber models, a “static”
754 fiber and a “dynamic” fiber based on observations that muscle spindle primary afferent responses to
755 stretch consist of two components (Banks et al., 1997; Blum et al., 2017; Boyd, 1976; Boyd et al., 1977;
756 Hasan, 1983; Jami et al., 1982; Jami & Petit, 1979; Lewis & Proske, 1972). For these simulations, each
757 muscle fiber model used identical parameters, but the contribution of each fiber to the neural firing rate
758 varied. The equation describing the contribution of each fiber to the total firing rate is:

$$r(t) = r_{dynamic}(t) + r_{static}(t), \quad (15)$$

759 where the total firing rate of the afferent, $r(t)$, is a sum of the dynamic and static fiber components, or
760 $r_{dynamic}(t)$ and $r_{static}(t)$, respectively. The static component was defined as:

$$r_{static}(t) = k_{Fs} F_s(t), \quad (16)$$

761 where k_{Fs} is a constant and $F_s(t)$ is the total force in the static fiber. The dynamic component was defined
762 as:

$$r_{dynamic}(t) = [k_{Fd} F_d(t) + k_{\dot{F}d} \dot{F}_d(t)], \quad (17)$$

763 where k_{Fd} and $k_{\dot{F}d}$ are constants, and $F_d(t)$ and $\dot{F}_d(t)$ are respectively the force and yank of the cycling
764 cross-bridges in the dynamic fiber. We used default k_{FS} , k_{Fd} , and $k_{\dot{F}d}$ values of 1, 1, and 0.03,
765 respectively, unless otherwise noted.

766 The static and dynamic fibers are arranged in perfect mechanical parallel and were allowed to be
767 activated independently. Thus, the actions of the dynamic and static fibers could be simulated
768 simultaneously or sequentially.

769 *Occlusion between dynamic and static components*

770 To account for the evidence of so-called “occlusive interaction” between dynamic and static
771 branches of the muscle spindle Ia afferent ending, we used a nonlinear summation of the static and
772 dynamic components. Previous models have used complete occlusion (Hulliger et al., 1977; C.-C. K. Lin
773 & Crago, 2002) but we used a partial occlusion based on more recent findings (Banks et al., 1997). With
774 occlusion, the total firing rate of the model Ia afferent becomes:

$$r(t) = f_{occ}r_{dynamic}(t) + r_{static}(t), \quad r_{dynamic} \geq r_{static} \quad (18)$$

$$r(t) = r_{dynamic}(t) + f_{occ}r_{static}(t), \quad r_{static} > r_{dynamic} \quad (19)$$

775 where f_{occ} is an occlusion factor limiting the contribution of either component to the overall firing rate.
776 This parameter was set to 0.3 (unitless) for all simulations unless otherwise noted (Banks et al., 1997).

777 *Dynamic response simulations*

778 To demonstrate the ability of our model to produce the classical fractional power relationship
779 between the dynamic response of muscle spindle Ia firing rates and ramp velocity (Hasan, 1983; Houk et
780 al., 1981; P. B. C. Matthews, 1963b), we applied a series of ramp-hold stretches to the model with each
781 fiber’s proportion of available binding sites set to 0.3 (Fig. 6A). The ramp stretches consisted of a pre-
782 stretch isometric hold period, followed by a constant velocity stretch that varied linearly between trials
783 from $0.079L_0/s$ to $0.79L_0/s$, and another isometric hold period at its new length ($1.059L_0$). The duration of

784 stretch was shortened proportionally to the stretch velocity to ensure the same total length was applied in
785 each trial.

786 *Time-history dependence simulations*

787 To demonstrate the unique ability of our model to vary its own sensitivity to stretch based on the
788 history of movement applied to the muscle (Haftel et al., 2004), we applied series of triangular ramp-
789 release stretches with each fiber's activation set to 0.3. Each series consisted of three stretch-shorten
790 cycles, with a $1.047L_0$ amplitude and stretch and shorten velocities of $0.079L_0/s$. The first two cycles were
791 applied sequentially with no pause between them, whereas the third sequence was applied after a varied
792 isometric hold period at L_0 ranging from 0 – 10 s.

793 To demonstrate the robustness of our mechanistic model to produce history dependent responses
794 similar to those observed in awake humans, we approximated stimuli from microneurography studies in
795 humans from the lower limb (Day et al., 2017). We applied a $1.042 L_0$ sinusoidal length change at 1.57
796 Hz to the model at the baseline activation to mimic the passive manipulation of the ankle in the study
797 (Day et al., 2017). To roughly match the predicted driving potential to the firing rate of the spindle, we
798 used k_{FS} , k_{Fd} , and $k_{\dot{F}d}$ values of 1.8, 2, and 0.15, respectively.

799

800 *Tonic gamma activation simulations*

801 To demonstrate the effects of muscle activation on the firing response of our model (Fig. 7), we
802 applied a range of activations to the static and dynamic fibers (Emonet-Dénand et al., 1977). We varied
803 the activation levels of the static and dynamic fibers independently, between 0 – 1.0, before applying a
804 $1.047L_0$ ramp-hold stretch at a constant velocity of $0.079L_0/s$. We used k_{FS} , k_{Fd} , and $k_{\dot{F}d}$ values of 1.5,
805 0.8, and 0.03, respectively, for these simulations in order to better visualize the effects of gamma
806 activation on the predicted driving potential.

807 *Musculotendon stimulations*

808 To more completely address the complex interplay between alpha and fusimotor drive and the
809 effect therein on the afferent outputs of the spindle, we created a model of a musculotendinous unit,
810 comprised of a single half-sarcomere model mechanically in series with a linear elastic element to model
811 the tendon. To balance the forces between the muscle and tendon, a force balancing operation was
812 performed by searching iteratively for the pair of muscle and tendon length changes that would result in
813 equal tendon and muscle forces.

814 To demonstrate the independent effects of alpha and fusimotor drive on the muscle spindle Ia
815 afferent, we modeled a series of isometric activation pulses. Each pulse was based on a Gaussian
816 waveform (*gausswin.m*) 1 second long, with a width factor of 0.3. Medium activation was defined as 50%
817 peak activation, and low activation was defined as 25% peak activation for both extra- and intrafusal
818 muscle. We simulated the musculotendon responses to these pulses and used the extrafusal muscle length
819 as the length input to the intrafusal muscle fibers. We calculated the resulting firing rate using the
820 previously defined r , from the force and yank of intrafusal models. As a visual aide, we generated a
821 possible Ia afferent spike response by driving a leaky integrate-and-fire neural model with r as the input.

822 Next, we wanted to demonstrate the opposing effects of extrafusal muscle shortening and
823 increasing fusimotor drive to mimic the effects of putative alpha-gamma coactivation during isometric
824 conditions. To accomplish this, we emulated the microneurography work of Edin and Vallbo (Edin &
825 Vallbo, 1990) by applying ramp increases to extrafusal activation with concomitant fusimotor ramps. We
826 ran 16 simulations with increasing 2 second activation ramps with peak activation at the end of the ramp
827 ranging from 25-100% extrafusal activation in linear increments. The intrafusal peak activation was 100%
828 for all 16 simulations.

829 Finally, to demonstrate the combined interactions of alpha and fusimotor drive with
830 musculotendon movements, we ran a set of simulations varying levels of extra- and intrafusal activation
831 during an imposed stretch of the musculotendon. For each simulation, we applied a 3.8% L_0 , 7.7% L_0/s
832 ramp-hold stretch to the musculotendon under low (10%) and medium (60%) tonic extrafusal activation.
833 For each of these conditions, we also simulated low (30%) and high (70%) tonic dynamic and static

834 intrafusal activation to highlight the effects of extrafusal activation on Ia afferent coding. As with the
835 isometric pulse simulations, we visualized Ia afferent spiking using a leaky integrate-and-fire neuron.

836

837 *Adaptation of previously published figures*

838 Previously published results that were used for comparison with our model predictions were
839 redrawn in Adobe Illustrator. Only single data points and lines were approximated by tracing over their
840 apparent geometric centroids. These data were redrawn for aesthetic purposes only and were not used for
841 any quantitative comparisons. Any comparison of data from these studies with the present study were
842 performed using the original manuscripts.

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SUPPLEMENTAL FIGURES AND LEGENDS

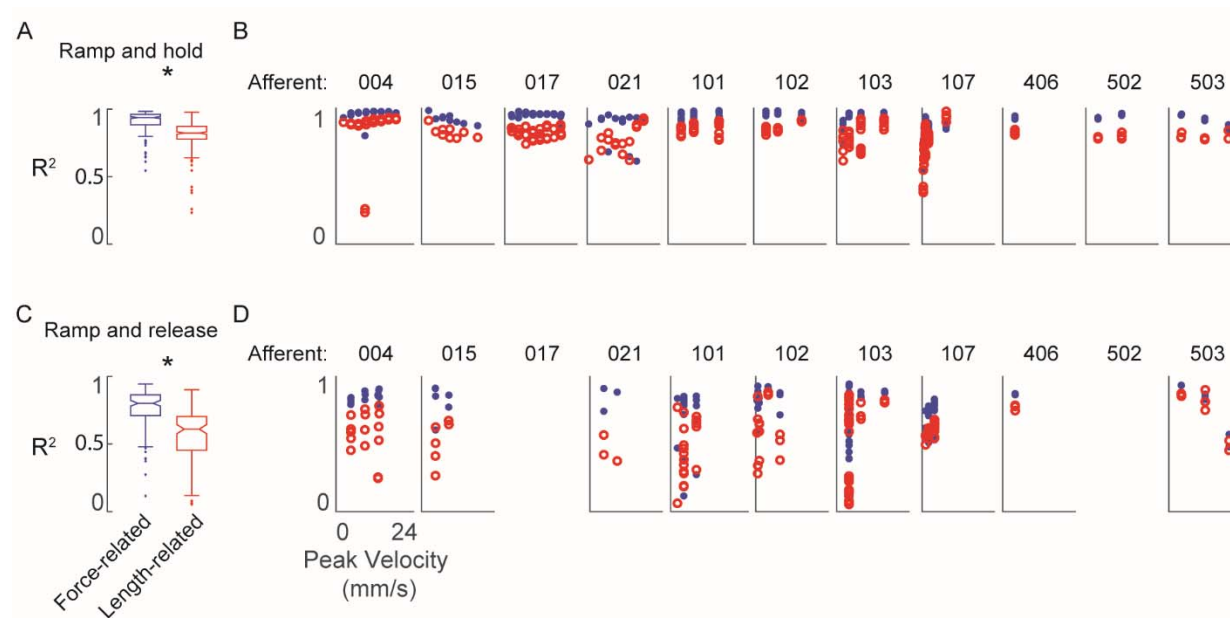


Figure S1 Goodness-of-fit to measured muscle spindle Ia afferent firing rates using estimated muscle fiber force and yank (blue) compared to kinematics model (red) as baseline comparison. A) Box and whisker plots of model goodness of fit (R^2) for all ramp-and-hold trials using estimated fiber force-based (blue) versus length-based (red) models. * denotes a significant difference between the model goodness of fit based on cumulative squared errors; $p < 0.05$, one-way ANOVA. B) R^2 values for all individual trials of ramp-and-hold and within each afferent plotted versus stretch velocity for force-based (blue dots) versus length-based (red dots) models. C-D) Same as A-B, only for repeated ramp-release trials.

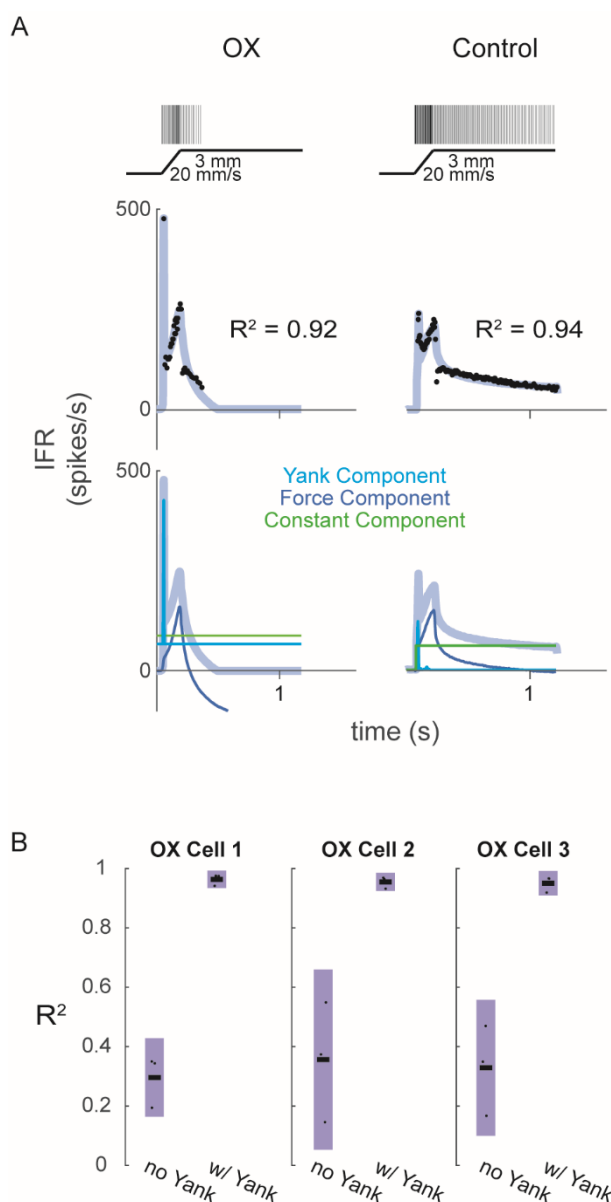


Figure S2

Fits of muscle spindle firing rates before and after oxaliplatin-induced neuropathy using estimated muscle fiber force and yank. (A) Example Ia afferent responses to stretch and force-based model fits. Left column shows a typical Ia afferent response to a ramp-hold stretch applied to the triceps surae of a rat treated with oxaliplatin chemotherapy (Bullinger et al. 2011). Right column shows typical response of Ia afferent in response to same ramp-hold stretch. Raster plots indicate times at which action potentials are recorded and are shown above imposed stretches. Below are the IFR and corresponding model fits shown above the same model fits with their respective components. (B) Variance of muscle spindle Ia afferent responses accounted for by force related model with (right bar in each plot) and without (left bar in each plot) yank for 3 Ia afferents from 3 oxaliplatin-treated rats. Black horizontal bars represent the means, blue bars represent the standard deviations, and black dots represent the data points from each trial (3 trials per afferent).

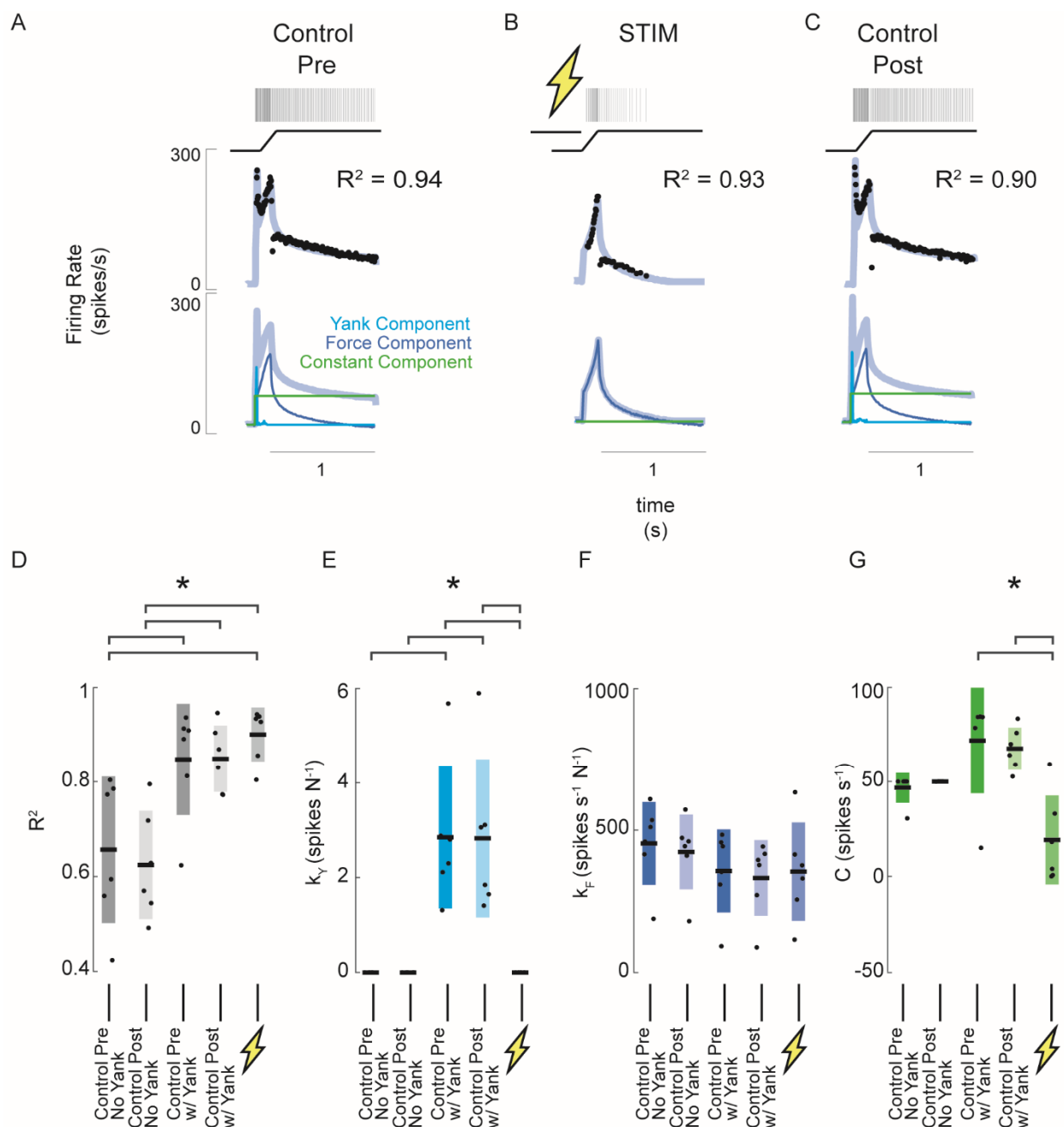


Figure S3

Estimated muscle fiber force-related model predicts changes in muscle spindle encoding caused by axonal stimulation. (A) Example of pre-stimulus control trial consisting of a 3 mm ramp-hold stretch at 20 mm/s. Raster represents the times of recorded action potentials in response to the ramp-hold stimulus shown directly below it. Black dots represent the afferent IFR corresponding to the raster. The gray-blue line represents the model prediction using force, yank, and constant components. Below this are the same model prediction as above (grey-blue) and model prediction components (blue – force component; cyan – yank component; green – constant component). (B) Example of stimulus trial consisting of the same stretch applied as A, with a train of antidromic action potentials induced by suprathreshold current injection into the Ia axon prior to stretch. The model fit in this trial represents the best fit without the yank component.

Notice the quality of fit does is roughly equal in this trial with that of the pre-stimulus trial in A, but no yank component was necessary. (C) Example of post-stimulus control trial. Same stretch and model components were used as in A. (D) Goodness of fit (R^2) of force-related model for 6 afferents subjected to the same trials as shown in A-C. The first two columns show R^2 values of the model fits without yank components for the pre- and post-stimulus control trials, respectively, for all 6 afferents. The third and fourth columns from the left show R^2 values for the model fits with yank components for the pre- and post-stimulus control trials, respectively, for the same 6 afferents. The fifth column shows R^2 values for the model fits for the stimulus trials. This model did not use a yank component. (E) Model sensitivity to yank for the same model fits and afferents as D. (F) Model sensitivity to force for the same model fits and afferents as D and E. (G) Model constant component for the same model fits and afferents as D-F. Brackets above plots indicate significant differences between the means ($p < 0.05$).

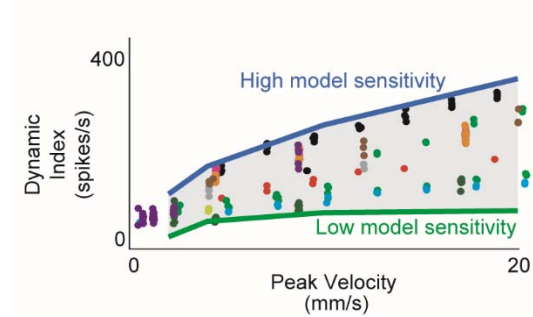


Figure S4

Estimated muscle fiber force model predicts inter-afferent variability of healthy afferent firing properties across perturbation velocity and acceleration. (A) Same dynamic index data shown in Figure 3-2 (colored dots) with predicted range of dynamic indices from model (gray shaded area). Nominal simulations were performed with the same force and yank model sensitivities for 4 stretch trials (2, 4, 10, 20 mm/s) from the same animal. Model parameter sweeps were performed for force and yank sensitivities for each of the 4 trials from 0.1 to 5 times the nominal value for each parameter. Green line represents the minimum dynamic indices from each of the 4 parameter sweeps from trials for which there was a spiking response (usually corresponding to low force and yank sensitivities). Blue line represents maximum dynamic indices from each of the 4 parameter sweeps (usually corresponding to high force and yank sensitivities). Gray shaded area represents plausible space of simulated dynamic index.

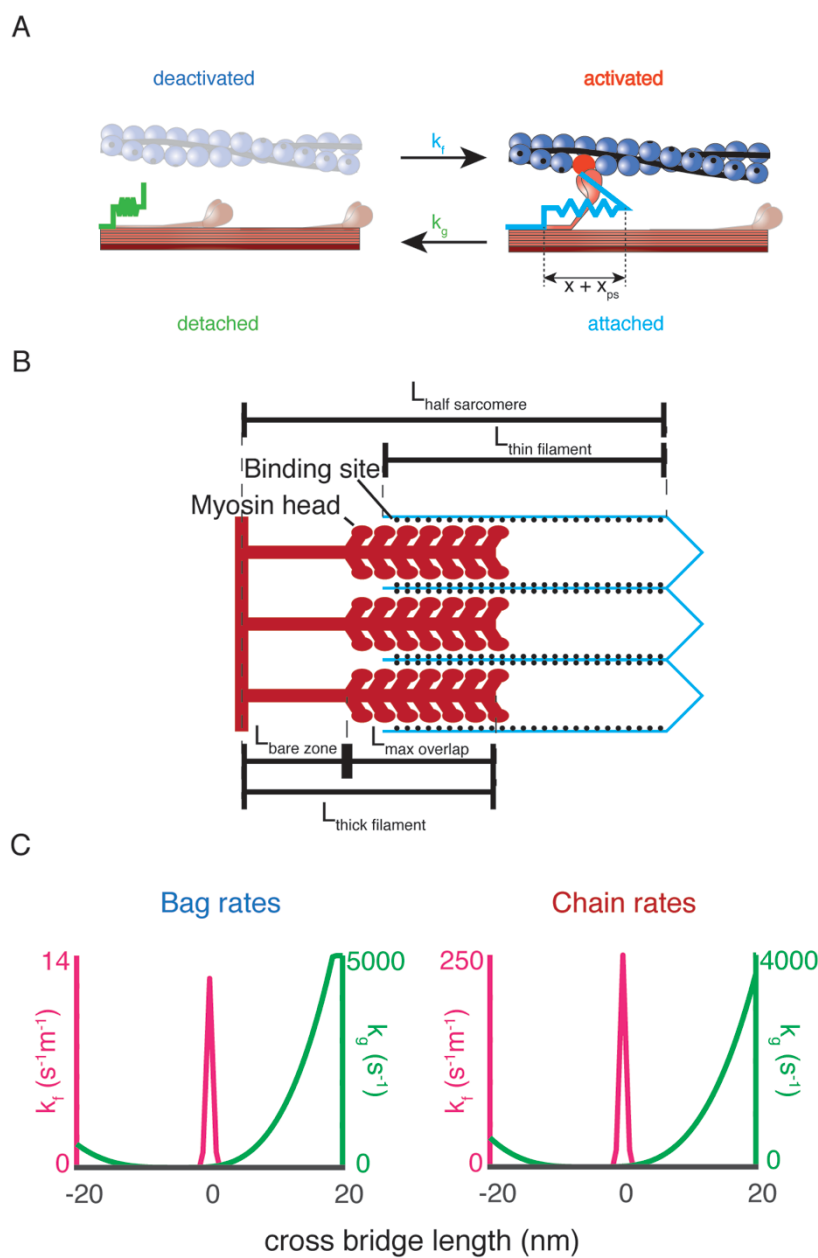


Figure S5

Biophysical intrafusal muscle models. A) Two-state dynamic system of cross-bridge cycling. A population of detached cross-bridges attaches at rate $k_f(x)$, and the population of attached cross-bridges detaches at rate $k_g(x)$. When a cross-bridge is formed at length x , an additional “powerstroke” length, x_{ps} , is applied to the cross-bridge to generate a contractile force. B) Schematic of length variables accounted for in muscle model. The amount of overlap between the myosin heads of the thick filament (red) and actin binding sites of the thin filament (blue lines, black dots) is the relevant variable for the simulations. The fraction of overlap is simply the difference between the total length of the thick and thin filaments with the length of the half sarcomere, expressed as a fraction of the maximum potential overlap. C) Rate equations for myosin dynamics. The rate at which detached myosin

heads will attach as a function of the length, $k_f(x)$ of attachment is a Gaussian function, centered around 0 nm (shown in blue). The rate at which attached cross-bridges will detach as a function of their length $k_g(x)$ is an offset polynomial function (shown in orange).

Parameter	Value (dynamic, static fiber)	Units	Description
k_{cb}	1	mN m ⁻¹	Unit cross-bridge stiffness
x_{ps}	2.5	nm	Unit power stroke distance
$L_{thick\ filament}$	815	nm	Length of thick filament
$L_{thin\ filament}$	1120	nm	Length of thin filament
$L_{bare\ zone}$	80	nm	Length of bare zone
$c_{filament}$	0.5	-	Filament compliance factor
ρ_{cb}	6.9×10^{16}	m ⁻²	Cross-bridge number density
l_0	1050, 1200	nm	Passive force reference length
k_{pas}	100, 250	N m ⁻² nm ⁻¹	Passive force linear stiffness

Table S1

Constant parameters used in both dynamic and static intrafusal muscle fiber models. These parameters did not change in any simulation presented in this study. When two values are presented, they represent the respective values for the dynamic and static fibers.