

Extrinsic Noise Suppression in Micro RNA mediated Incoherent Feedforward Loops

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Abstract—MicroRNA mediated incoherent feed forward loops (IFFLs) are recurrent network motifs in mammalian cells and have been a topic of study for their noise rejection and buffering properties. Previous work showed that IFFLs can adapt to varying promoter activity and are less prone to noise than similar circuits without the feed forward loop. Furthermore, it has been shown that microRNAs are better at rejecting extrinsic noise than intrinsic noise. This work studies the biological mechanisms that lead to extrinsic noise rejection for microRNA mediated feed forward network motifs. Specifically, we compare the effects of microRNA-induced mRNA degradation and translational inhibition on extrinsic noise rejection, and identify the parameter regimes where noise is most efficiently rejected. In the case of static extrinsic noise, we find that translational inhibition can expand the regime of extrinsic noise rejection. We then analyze rejection of dynamic extrinsic noise in the case of a single-gene feed forward loop (sgFFL), a special case of the IFFL motif where the microRNA and target mRNA are co-expressed. For this special case, we demonstrate that depending on the time-scale of fluctuations in the extrinsic variable compared to the mRNA and microRNA decay rates, the feed forward loop can both buffer or amplify fluctuations in gene product copy numbers.

I. INTRODUCTION

Stochasticity in gene expression is known to be modulated by biological feedback and feed-forward network motifs [1]. In particular, recent work has found that IFFLs – a network architecture where an upstream regulator directly activates a downstream target and also indirectly inhibits it – are capable of buffering against changes in promoter activity, as well as reducing the stochasticity in gene expression [2]–[4]. Furthermore, it has also been shown that the reduction of noise is more effective at the post-transcriptional level for IFFLs [5]. In this study we focus on understanding the noise rejection properties of a particular class of post-transcriptionally regulated IFFLs, where the negative regulatory link is implemented with a microRNA (miRNA).

MiRNA are a class of non-coding, regulatory RNAs that have been linked with the post-transcriptional regulation of important biological processes including differentiation,

development and disease [6]. As was the case for transcriptionally regulated IFFLs, experimental studies with miRNA mediated IFFLs have shown that they have effective noise-suppressing and buffering properties in some parameter regimes, often displaying interesting nonlinear behaviors [2], [4]. miRNA regulation of gene expression is believed to occur through two different regulatory modes: i) by degrading mRNAs that the miRNAs bind to, and ii) by preventing efficient translation of miRNA bound mRNAs into proteins (translation-inhibition) [7]. To elucidate the mechanisms underlying these properties, several mathematical models have been proposed to describe this class of IFFLs [8]–[10]. However, these models either rely on large mathematical nonlinearities [8], or on complex network architecture [10] to explain the nonlinear noise-rejection properties of the system. Such models also often make modeling assumptions that are not justified using biological mechanisms. On the other hand, data-fitted models [9] don't explore the full parameter space, and risk to overlook important components in the machinery. Hence, a qualitative, parameter-independent model that could provide an easy biological explanation is currently missing. Here, we propose a simple quasi-linear model that utilizes mechanisms inferred from previous experimental studies [4], [7] and mild assumptions on the underlying extrinsic noise sources. Through these basic assumptions, we explain the previously-reported biological behaviors, and recapitulate the mechanisms at the origin of the noise-repression mechanism with particular attention to their biological interpretation. These insights allow us to precisely pinpoint the qualitative contribution of each cellular machinery on the overall behavior, providing a framework for understanding current biological systems and for designing synthetic systems for noise rejection.

We initially focus on the the more biologically common situation of static extrinsic noise, which arises when the mRNA degradation occur on a much faster time scale than that set by the cell cycle. We analytically obtain conditions for extrinsic noise rejection for general miRNA mediated IFFLs, demonstrating that the noise correlation alone defines the parameter space for noise rejection. Then, we demonstrate that high mRNA-miRNA binding rate could lead to an overall noise increase, and give a simple biological explanation. Moreover, we show that, contrary to what was previously reported [9], translation-inhibition is a key-player in reducing protein noise. Finally, we extend the analysis to dynamic extrinsic noise, focusing on the special case where the mRNA and miRNA are derived from the same gene (single-gene IFFL). We demonstrate that even for this simple

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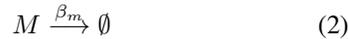
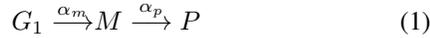
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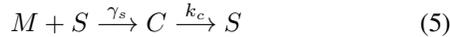
system, dynamic extrinsic noise can lead to a wide range of noise rejection regimes depending on the relative stability of the mRNA and the miRNA.

II. MODEL OF IFFL SYSTEM

We now outline the miRNA-mRNA IFFL model that we use in this study, starting from the underlying biochemical reactions. mRNA (M) and miRNA (S) are transcribed from genes \mathbf{g}_1 and \mathbf{g}_2 respectively at constant rates α_m , α_s . The process of translation then converts mRNA into protein (P) at constant rate α_p ; miRNA, by definition, is not coding for any protein. To complete the reaction network, we add the degradation processes:



The miRNA-based regulation proceeds primarily through two groups of steps namely miRNA mediated degradation and translational inhibition which affect each of these reactions separately. Many models for miRNA based degradation have been proposed [7], [11]–[13]. For the purposes of this work we rely on a previously-adopted model [4]: we assume that mRNA and miRNA form an irreversible complex C that is then degraded into the miRNA alone at rate k_c . Moreover, we assume that mRNA can be translated into protein even when bound to the miRNA, although at a lower rate. This factor accounts for the miRNA-induced translation inhibition that has been previously reported [7], [9]:



where k_1 is a real number between 0 and 1, where $k_1 = 0$ implies full translation inhibition (no protein from the complex) and $k_1 = 1$ implies that the complex has the same translation rate as the mRNA alone.

We now write a mathematical model describing 1 and 5: the species concentrations (or copy numbers) are represented with the corresponding lower case bold-faced variables which represent random processes.

$$\frac{d\mathbf{m}}{dt} = \alpha_m \mathbf{g}_1 - \beta_m \mathbf{m} - \gamma_s \mathbf{m}\mathbf{s} \quad (7)$$

$$\frac{d\mathbf{s}}{dt} = \alpha_s \mathbf{g}_2 - \beta_s \mathbf{s} + k_c \mathbf{c} \quad (8)$$

$$\frac{d\mathbf{c}}{dt} = \gamma_s \mathbf{m}\mathbf{s} - k_c \mathbf{c} \quad (9)$$

$$\frac{d\mathbf{p}}{dt} = \alpha_p (\mathbf{m} + k_1 \mathbf{c}) - \beta_p \mathbf{p} \quad (10)$$

The production rates of mRNA and miRNA are dependent respectively on the positive random variables \mathbf{g}_1 and \mathbf{g}_2 , which correspond to the average number of active genes. Genes \mathbf{g}_1 and \mathbf{g}_2 can be transcribed either dependently ($Cov(\mathbf{g}_1, \mathbf{g}_2) \neq 0$) or independently ($Cov(\mathbf{g}_1, \mathbf{g}_2) = 0$).

This model can be further simplified by assuming that the mRNA-miRNA complex formation reaches steady state at a faster pace than the other processes. Hence, we get $\mathbf{c} = \frac{\gamma_s \mathbf{m}\mathbf{s}}{k_c}$. This simplifies equations 8–10 to:

$$\frac{d\mathbf{s}}{dt} = \alpha_s \mathbf{g}_2 - \beta_s \mathbf{s} \quad (11)$$

$$\frac{d\mathbf{p}}{dt} = \alpha_p \mathbf{m}(1 + K\gamma_s \mathbf{s}) - \beta_p \mathbf{p} \quad (12)$$

Where $K = \frac{k_1}{k_c}$. In the special case where there is complete translational inhibition we observe that $K = 0$.

Steady state mRNA and protein levels of the open loop system 1 can easily be computed:

$$\mathbf{m}_{OL}^* = \frac{\alpha_m}{\beta_m} \mathbf{g}_1 \quad (13)$$

$$\mathbf{p}_{OL}^* = \frac{\alpha_p \alpha_m}{\beta_p \beta_m} \mathbf{g}_1 \quad (14)$$

Under the reasonable assumption that there is negligible variability in the cellular processes within a cell population, the extrinsic noise source is the variability of the transcription of the genes \mathbf{g}_1 and \mathbf{g}_2 . This enables us to avoid making any assumptions on the distributions of the parameter values. The coefficient of variations of the open loop system are given by:

$$\eta_{p_{OL}} = \eta_{m_{OL}} = \eta_{g_1} \quad (15)$$

The corresponding steady state levels in the general IFFL system according to 11 are:

$$\mathbf{m}^* = \frac{\mathbf{m}_{OL}^*}{1 + \frac{\gamma_s d}{\beta_m} \mathbf{g}_2} \quad (16)$$

$$\mathbf{p}^* = \frac{\mathbf{p}_{OL}^*}{\frac{1 + \frac{\gamma_s d}{\beta_m} \mathbf{g}_2}{1 + K d \gamma_s \mathbf{g}_2}} \quad (17)$$

where $d = \frac{\alpha_s}{\beta_s}$. Equations 16 and 17 highlight the separate dependency of the system equilibria on the open-loop steady states and the miRNA regulation.

III. STATIC EXTRINSIC NOISE REJECTION IN GENERAL IFFL SYSTEMS

Extrinsic Noise Measurement at Steady State for Static Noise.

Extrinsic noise arises from diversity in cell populations, such as differences in cell size, in uptake of an external signal, or in the cell cycle phase. The two random variables \mathbf{g}_1 and \mathbf{g}_2 representing gene expression level, account for the extrinsic noise in our model. We consider two cases: when the noise is static in nature (\mathbf{g}_1 and \mathbf{g}_2 are fixed distributions over time) and when it is time-varying.

To quantify steady-state noise rejection in microRNA-based systems, we use the square of the coefficient of variation defined as the ratio:

$$\eta_Z^2 = \frac{\text{Var}(\mathbf{Z})}{\langle \mathbf{Z} \rangle^2} \quad (18)$$

where \mathbf{Z} is the random variable of interest, and $\text{Var}(\cdot)$ and $\langle \cdot \rangle$ are the standard notation for the variance and expected value operators. The coefficient of variation allows to compare

noise in processes that have different means. This is essential in our study, as microRNA-control reduces the level of mRNA and protein.

We use a first-order Taylor-Delta approximations to obtain the ratio and products of the two random variables (\mathbf{g}_1 and \mathbf{g}_2), which leads to the following expression for the coefficient of variation 18:

$$\eta^2\left(\frac{\mathbf{X}}{\mathbf{Y}}\right) = \frac{\text{Var}\left(\frac{\mathbf{X}}{\mathbf{Y}}\right)}{\left\langle\frac{\mathbf{X}}{\mathbf{Y}}\right\rangle^2} \quad (19)$$

$$= \frac{\frac{\text{Var}(\mathbf{X})}{\langle\mathbf{Y}\rangle^2} - 2\frac{\langle\mathbf{X}\rangle}{\langle\mathbf{Y}\rangle^3}\text{Cov}(\mathbf{X}, \mathbf{Y}) + \frac{\langle\mathbf{X}\rangle^2}{\langle\mathbf{Y}\rangle^4}\text{Var}(\mathbf{Y})}{\frac{\langle\mathbf{X}\rangle^2}{\langle\mathbf{Y}\rangle^2}} \quad (20)$$

This expression can be simplified to:

$$\eta^2\left(\frac{\mathbf{X}}{\mathbf{Y}}\right) = \eta_{\mathbf{X}}^2 + \eta_{\mathbf{Y}}^2 - 2\rho\eta_{\mathbf{X}}\eta_{\mathbf{Y}} \quad (21)$$

where ρ is the Pearson product-moment correlation coefficient between the two random variables \mathbf{X} and \mathbf{Y} .

We can now compute the coefficient of variations for the mRNA (16) and protein (17) level as:

$$\eta_m^2 = \eta_{g_1}^2 + \eta_{mm}^2 - 2\rho_m\eta_{g_1}\eta_{mm} \quad (22)$$

$$\eta_p^2 = \eta_{g_1}^2 + \eta_{pm}^2 - 2\rho_p\eta_{g_1}\eta_{pm} \quad (23)$$

where we define the miRNA-dependent noise contribution for mRNA (η_{mm}) and protein (η_{pm}) as:

$$\eta_{mm} = \frac{\gamma_s d\langle\mathbf{g}_2\rangle}{\beta_m + \gamma_s d\langle\mathbf{g}_2\rangle} \eta_{g_2} \quad (24)$$

$$\eta_{pm} = \frac{\gamma_s d\langle\mathbf{g}_2\rangle(1 - K\beta_m)}{(1 + K\gamma_s d\langle\mathbf{g}_2\rangle)(\beta_m + \gamma_s d\langle\mathbf{g}_2\rangle)} \eta_{g_2} \quad (25)$$

and ρ_m and ρ_p as:

$$\rho_m = \rho\left(g_1, \frac{\gamma_s d}{\beta_m} \mathbf{g}_2\right) \quad (26)$$

$$\rho_p = \rho\left(g_1, \frac{1 + \frac{\gamma_s d}{\beta_m} \mathbf{g}_2}{1 + K d \gamma_s \mathbf{g}_2}\right) \quad (27)$$

According to formulae 22 and 23, miRNA regulation introduces noise in the system (represented by the quantities η_{mm}^2 and η_{pm}^2), but it also cancels part of the noise components of \mathbf{g}_1 if the two processes are positively correlated ($\rho > 0$). The noise cancellation terms, $-2\rho_m\eta_{g_1}\eta_{mm}$ or $-2\rho_p\eta_{g_1}\eta_{pm}$, suggest that the noise is canceled from both sources depending on how well they correlate with each other.

Conditions for noise cancellation in miRNA-regulated IFFLs

To understand the extrinsic noise reduction properties of miRNA regulation, we studied the values of η_{mm} and η_{pm} in 22 and 23 that satisfy condition: $\eta^2 \leq \eta_{g_1}^2$, when the noise of the close-loop system is lower than for the open loop.

It is easy to see that a necessary condition for noise rejection is $\rho_m > 0$ and $\rho_p > 0$. At the mRNA level (eq 24), this condition is guaranteed iff $\rho(\mathbf{g}_1, \mathbf{g}_2) > 0$, which occurs when

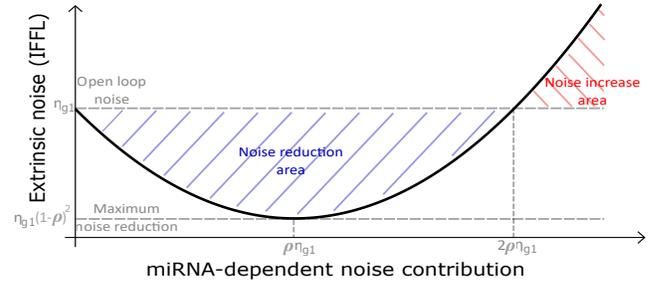


Fig. 1. Extrinsic noise rejection in microRNA based IFFLs.

the same upstream gene positively regulates both \mathbf{g}_1 and \mathbf{g}_2 , as is the case in an IFFL. Conversely, uncorrelated miRNA or coherent feed forward loops (CFFLs) do not have extrinsic noise cancellation properties at the mRNA level according to our model.

At the protein level (eq 25), two conditions need to be satisfied to have $\rho_p > 0$:

$$\rho(\mathbf{g}_1, \mathbf{g}_2) > 0 \quad (28)$$

$$1 - K\beta_m > 0 \quad (29)$$

Condition 28 is the same as for mRNA, which implies that only IFFLs can achieve noise rejection at the protein level. Condition 29 guarantees that the miRNA-mRNA complex does not translate into protein faster than mRNA alone would be degraded. If this were to happen, more protein would be produced from the complex than from the open loop mRNA. Therefore the miRNA regulation would not repress protein expression, but instead facilitate it. Hence, translation-inhibition ($K > 0$) could cancel the noise-rejection property of the IFFL.

Even if $\rho > 0$, the parabolic form of 22 and 23 implies that there is a limited range of values of η_{mm} and η_{pm} that leads to extrinsic noise rejection, before the extra noise introduced by the miRNA machinery leads to worse performance than in the open loop (Figure1). The limiting case is when $\eta_{mm} = 2\rho_m\eta_{g_1}$ and $\eta_{pm} = 2\rho_p\eta_{g_1}$, where the amount of noise that is introduced by the miRNA is equivalent to the amount of canceled noise ($\eta = \eta_{g_1}$). The minimum of the parabola corresponds to the maximum noise rejection that could be achieved. This is reached for $\eta_{mm} = \rho_m\eta_{g_1}$ and $\eta_{pm} = \rho_p\eta_{g_1}$, for which the noise rejection ratio is:

$$\eta_m^2 = \eta_{g_1}^2(1 - \rho_m^2) \quad (30)$$

$$\eta_p^2 = \eta_{g_1}^2(1 - \rho_p^2) \quad (31)$$

Thus, optimal noise rejection is achieved for perfect correlation ($\rho = 1$).

In summary, we showed that miRNA-based IFFLs are the only network architecture that can lead to noise reduction in simple miRNA mediated systems. We also demonstrated that correlation between the two types of noise (miRNA-based and mRNA) is the main factor in extrinsic noise cancellation. However, very high miRNA-related noise ($\eta_{mm} > 2\rho_m\eta_{g_1}$ or $\eta_{pm} > 2\rho_p\eta_{g_1}$) increases the system noise even at

maximum correlation ($\rho_m = 1$ and $\rho_p = 1$), as one would expect (Figure 1).

High miRNA-mRNA binding rate could lead to increased extrinsic noise

We investigated the relationship between the miRNA-mRNA binding rate (γ_s) and the extrinsic noise rejection properties of the IFFLs. We showed in the previous subsection that the amount of noise introduced by the miRNA regulation (η_{mm}^2 or η_{pm}^2) could be balanced by noise cancellation ($-2\rho_m\eta_{g_1}\eta_{mm}$ and $-2\rho_p\eta_{g_1}\eta_{pm}$). Hence, we computed the value of η_{mm} at the mRNA and the protein level and studied their dependencies to γ_s . It could be shown that η_{mm} in 24 has a monotonic dependency on γ_s :

$$\frac{d\eta_{mm}}{d\gamma_s} > 0, \forall \gamma_s > 0 \quad \lim_{\gamma_s \rightarrow \infty} \eta_{mm}(\gamma_s) = \eta_{g_2} \quad (32)$$

This equations shows that there is a maximum amount of noise that could be introduced by the miRNA regulation to the system at the mRNA level, and this is bounded by the amount of noise in the input signal \mathbf{g}_2 (specifically, by $\eta_{g_1}^2 + \eta_{g_2}^2 - 2\rho_m\eta_{g_1}\eta_{g_2}$). The mRNA extrinsic noise tends to this value in the limit, although for high values of γ_s , $\mathbf{m}^* \rightarrow 0$. In this range, the mRNA dynamics are entirely defined by fluctuations in the miRNA, as all mRNA are in a complex (Figure2(a)). Moreover, if $\eta_{g_2} < 2\rho_p\eta_{g_1}$, according to 22, noise rejection is achieved for all values of γ_s .

The optimal noise rejection is reached for:

$$\frac{\gamma_s d\langle \mathbf{g}_2 \rangle}{\beta_m + \gamma_s d\langle \mathbf{g}_2 \rangle} \eta_{g_2} = \rho_m \eta_{g_1} \quad (33)$$

This equation shows that the mRNA noise is reduced when the miRNA-induced degradation can compete with the natural mRNA degradation β_m . Increasing or reducing the value of β_m requires a higher or lower value of γ_s to reach the minimum of the parabola. At the protein level, translation inhibition introduces a non-trivial nonlinear dynamic. In fact, η_{pm} is not monotonic on γ_s , and it reaches a maximum for:

$$\gamma_s^{max} = \sqrt{\frac{\beta_m}{K}} \frac{1}{d\langle \mathbf{g}_2 \rangle} \quad (34)$$

$$\eta_{pm}(\gamma_s^{max}) = \frac{1 - \sqrt{\beta_m K}}{1 + \sqrt{\beta_m K}} \eta_{g_2} \quad (35)$$

We notice that the maximum value of $\eta_{pm}(\gamma_s^{max})$ is positive iff $1 - \beta_m K > 0$, which is assured by the noise-rejection condition 29. Equation 25 also shows that the value of η_{pm} as $\gamma_s \rightarrow \infty$ is 0. Hence the maximum amount of noise that could be introduced by the miRNA regulation to the system at the protein level is reached for $\eta_{pm}(\gamma_s^{max})$ (Figure 2(b)). This quantity is less or equal than η_{g_2} , satisfying the equality for $K = 0$ (complete translation-inhibition), when the dynamic of the protein noise is equivalent to the mRNA one.

Depending on the values of 35, there are three possible

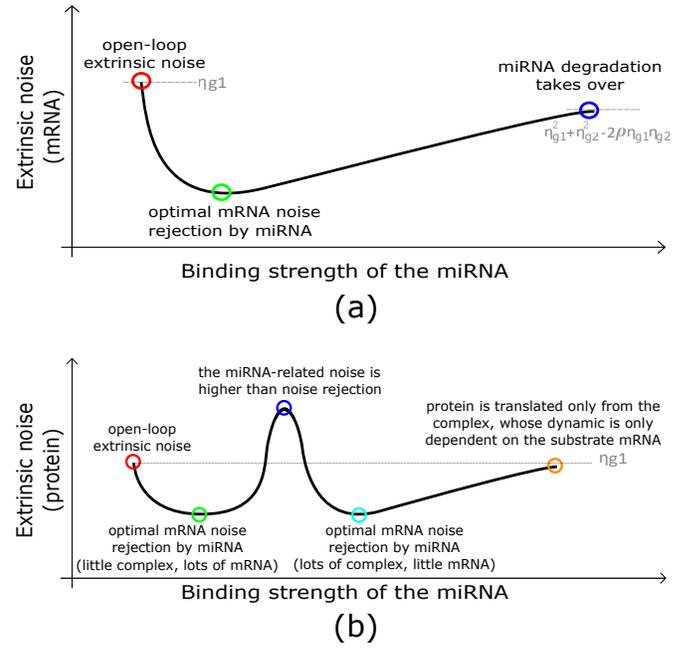


Fig. 2. High miRNA-mRNA binding rate has a nonlinear effect on the overall extrinsic noise. (a) Noise rejection at the mRNA level for a miRNA-based IFFL. miRNA-mRNA binding strength drives noise reduction from 0 (red circle) to a minimum (green circle, maximum noise cancellation). Higher binding rates introduce more miRNA noise in the mRNA spectrum till it reaches its maximum (blue circle). (b) Noise rejection at the protein level for a miRNA-based IFFL. Protein is generated from both complex and mRNA (assuming $K \neq 0$): for intermediate values of γ_s , this leads to a biphasic distribution, which increases the overall noise (dark blue circle).

scenarios:

$$0 < \eta_{pm}(\gamma_s^{max}) < \rho_p \eta_{g_1} \quad (36)$$

$$\eta_{pm}(\gamma_s^{max}) \leq \rho_p \eta_{g_1} \quad (37)$$

$$\eta_{pm}(\gamma_s^{max}) > \rho_p \eta_{g_1} \quad (38)$$

If 36 is satisfied, then there is only one optimal binding rate γ_s , which is sub-optimal for the given distribution \mathbf{g}_1 and \mathbf{g}_2 . If 37 is satisfied, then optimal noise rejection is reached for a unique value of γ_s . If 38 is satisfied, then optimal noise rejection is reached for two different values of γ_s , and noise increases between the two (local maximum). Moreover, if is bigger than $2\rho_p\eta_{g_1}$, then this local maximum is higher than the open loop noise. In all three scenarios, the IFFL noise tends to the open-loop noise as $\gamma_s \rightarrow \infty$ (Figure2(b)).

In summary, the protein noise follows the mRNA noise closely for low binding rates. As the binding rate increases, the miRNA-mRNA complex becomes dominant and competes with mRNA for translation to protein: these two separate translation events lead to a bimodal protein population ('threshold effect'), which increases the total extrinsic noise (maximum highlighted in blue in Figure 2(b)). At higher binding rate, the miRNA-mRNA complex mostly dominates translation: protein noise is then dependent on the miRNA noise, reaching a second minimum (highlighted in light blue in Figure 2(b)). For even higher values of γ_s , protein translation occurs mainly through the miRNA-mRNA complex,

because little free mRNA is left. In this regime, the protein noise is independent on miRNA, since at steady state the complex is dependent only on the amount of the substrate mRNA. Hence, η_p reaches the open-loop mRNA noise η_{g_1} as $\gamma_s \rightarrow \infty$.

Translation-inhibition modulates the functional range of miRNA extrinsic noise rejection

To address the importance of translation-inhibition for the miRNA-induced noise rejection, we analyzed the dependency of our results in the previous section to perturbation of the translation-inhibition parameter K . Equations 34 and III show that for a fixed value of γ_s , the performance of the miRNA-IFFL could be tuned by increasing or decreasing the translation-inhibition parameter K or the mRNA degradation β_m . However, tuning these parameters have opposite effects on optimality: γ_s^{max} is reduced if K is increased, and it increases if β_m is increased. This is not surprising since the two mechanisms are competing for mRNA degradation.

As shown in Figure 3, increasing K changes the location and

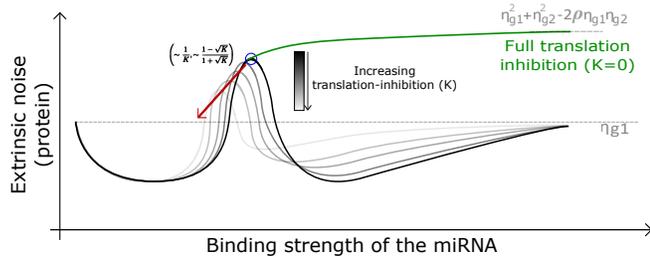


Fig. 3. Translation-inhibition modulates the functional range of miRNA extrinsic noise rejection. Upon binding to mRNA, miRNA interferes with the protein translation machinery (translation-inhibition, $K \neq 0$). This effect modulates the extrinsic noise peak (blue circle) both in intensity (y-axis) and in its corresponding binding strength (x-axis). For high values of K (low translation-inhibition), overall noise rejection occurs at a wider range of γ_s values rather than for full translation inhibition ($K = 0$).

height of the noise peak both in intensity (y-axis) and in its corresponding binding strength (x-axis). On the other hand, if $K = 0$, the dynamic follows the mRNA noise closely and reaches its maximum at $\eta_{g_1}^2 + \eta_{g_2}^2 - 2\rho_p \eta_{g_1} \eta_{g_2}$ as $\gamma_s \rightarrow \infty$ (green line in Figure 3).

Hence, translation-inhibition modulates the noise-rejection range of the miRNA-mRNA binding rate γ_s . For high-noise miRNA, full translation-inhibition ($K = 0$) reduces protein noise over a larger range of γ_s (since $\eta_{g_1}^2 + \eta_{g_2}^2 - 2\rho_p \eta_{g_1} \eta_{g_2} < \eta_{g_1}$), while little translation-inhibition is preferred for low-noise miRNA.

IV. STATIC AND DYNAMIC EXTRINSIC NOISE REJECTION IN sgFFL SYSTEMS

A. Static noise rejection by sgFFLs

A very common and functionally important sub-class of IFFLs is the single gene feed forward loop or sgFFL [14], [15]. sgFFLs are a special class of IFFLs where the mRNA and the miRNA are co-expressed by the same gene [4], [16]. Typically, the miRNA is located within an intron of the gene

encoding the mRNA. The mRNA and the miRNA are first transcribed together into one transcript and then separated from each other by splicing [17], [18]. Mathematically, this can be expressed as $\mathbf{g}_1 = \mathbf{g}_2 = \mathbf{g}$, which simplifies expression 22 to

$$\eta_m = \eta_g - \eta_{mm} = \eta_g - \frac{\eta_g}{\frac{\beta_m}{\gamma_s d(\mathbf{g})} + 1} \quad (39)$$

As a result of this simplification, the mRNA noise η_m can be reduced arbitrarily close to 0 by increasing γ_s , since $\eta_{mm} \rightarrow \eta_g$ as $\gamma_s \rightarrow \infty$. Extrinsic variations of mRNA levels would be canceled out by equivalent variations of miRNA: each cell would then be reduced to the same mRNA steady state, defined by the fixed biochemical rates of degradation. This behavior is conserved at the protein level if there is full translation inhibition. If one assumes that some of the complex gets translated, then there is an optimal value of γ_s for which extrinsic noise is minimized, and the protein concentration is not reduced to 0. This value is equivalent to:

$$\gamma_s^{opt} = \frac{1}{d(\mathbf{g})} \sqrt{\frac{\beta_m}{K}} \quad (40)$$

which is the sgFFL case of 34. This value is a ratio between mRNA that escapes miRNA degradation (naturally degraded at rate β_m or translated as miRNA-mRNA complex) and total amount of miRNA ($d(\mathbf{g})$).

B. Dynamic noise rejection by sgFFLs

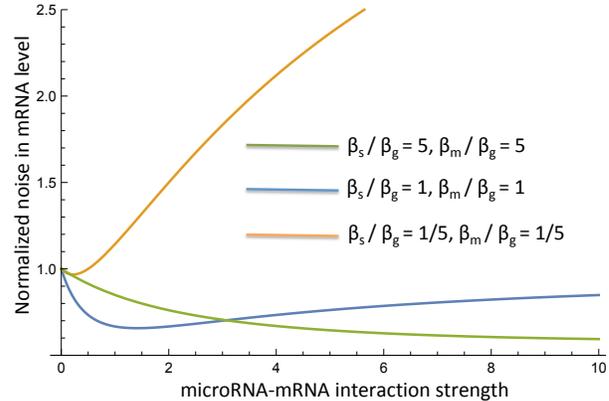


Fig. 4. Dynamical noise rejection by miRNA mediated IFFLs. Plot of noise in mRNA levels η_m as given by (47) as a function of the mRNA-miRNA interaction strength for different values β_g (time-scale of extrinsic fluctuations), β_m (mRNA decay rate) and β_s (miRNA decay rate). While low values of β_g cause η_m to monotonically decrease with γ_s (green line), values of β_g comparable or higher than β_m and β_s result in a non-monotonic profile for η_m (blue and orange lines), where mRNA noise is minimized at an optimal interaction strength. Note that in this plot mRNA noise levels are normalized by their value in the open loop system with no mRNA-miRNA interaction.

Having so far studied the effect of static extrinsic noise on sgFFL circuits, we now consider the scenario of dynamic noise by allowing

$$\mathbf{g}_1(t) = \mathbf{g}_2(t) = \mathbf{g}(t) \quad (41)$$

to be a stochastic process modeled as an Ornstein-Uhlenbeck (OU) process

$$d\mathbf{g}(t) = \beta_g (\langle \mathbf{g} \rangle - \mathbf{g}(t)) dt + \eta_g \langle \mathbf{g} \rangle d\mathbf{w}(t), \quad (42)$$

where $\langle \mathbf{g} \rangle$ is the steady-state mean level of $\mathbf{g}(t)$, $\mathbf{w}(t)$ is a Wiener process, η_g and β_g are the magnitude (coefficient of variation) and time-scale of fluctuations in $\mathbf{g}(t)$, respectively. This stochastic process drives the sgFFL circuit described earlier by the following system of different equations

$$\frac{d\mathbf{m}}{dt} = \alpha_m \mathbf{g} - \beta_m \mathbf{m} - \gamma_s \mathbf{m} \mathbf{s} \quad (43)$$

$$\frac{ds}{dt} = \alpha_s \mathbf{g} - \beta_s s. \quad (44)$$

For simplicity, in this subsection we will only focus on noise at the mRNA level, and hence ignore the protein dynamics. Note that the limit $\beta_g \rightarrow 0$, where fluctuations in $\mathbf{g}(t)$ are significantly slower than the mRNA-miRNA dynamics, recovers the scenario of static noise that we have up till now considered. A key question of interest is as we now vary β_g what miRNA-mRNA interaction strengths provide the most efficient noise buffering?

To address these questions, we investigate the statistical moments of the joint stochastic processes $\{\mathbf{g}(t), \mathbf{m}(t), \mathbf{s}(t)\}$. Readers are referred to [19], [20] for details on deriving moment dynamics for hybrid systems of the form (42)-(44) that couple stochastic and ordinary differential equations. It turns out that the nonlinear product term $\mathbf{m} \mathbf{s}$ in (43) makes computations non-trivial due to unclosed moment dynamics - the time evolution of the lower-order moments always depends on high-order moments. Moving forward, one general approach is to exploit closure schemes that approximate higher-order moments as nonlinear functions of lower-order moments [21]–[35]. Here we use an alternative approach based on linearizing the nonlinearity

$$\mathbf{m} \mathbf{s} \approx \langle \mathbf{s} \rangle \mathbf{m} + \langle \mathbf{m} \rangle \mathbf{s} - \langle \mathbf{m} \rangle \langle \mathbf{s} \rangle \quad (45)$$

where $\langle \mathbf{m} \rangle$ and $\langle \mathbf{s} \rangle$ denote the mean steady-state levels for the mRNA and miRNA, respectively. Using this approximation in place of $\mathbf{m} \mathbf{s}$ in (43) result in a linear dynamical system driven by an OU process, and standard theory can now be applied to compute moments. More specifically, let $\boldsymbol{\mu}$ denote a vector consisting of all the first and second order moments of $\{\mathbf{g}(t), \mathbf{m}(t), \mathbf{s}(t)\}$, then its time evolution is given by the linear system

$$\dot{\boldsymbol{\mu}} = \mathbf{a} + \mathbf{A}\boldsymbol{\mu}, \quad (46)$$

where vector \mathbf{a} and matrix \mathbf{A} depend on model parameters [19], [20]. Steady-state analysis of (46) yields a complicate formula for the steady-state noise (coefficient of variation) in mRNA levels η_m , as given by (47) on top of the next page. As expected, in the limit of static noise ($\beta_g \rightarrow 0$), (47) reduces to (39), where η_m^2 monotonically decreases with increasing mRNA-miRNA interaction strength γ_s .

In the limit of no mRNA-miRNA interaction, the mRNA noise level is given by

$$\lim_{\gamma_s \rightarrow 0} \eta_m^2 = \eta_g^2 \frac{\beta_m}{\beta_m + \beta_g} \quad (48)$$

and is an increasing function of the mRNA decay rate β_m . Intuitively, as we decrease the mRNA half-life, it becomes less efficient in averaging out upstream fluctuations in $\mathbf{g}(t)$, and more noise from $\mathbf{g}(t)$ propagates to the mRNA level. In the limit of strong mRNA-miRNA interaction,

$$\lim_{\gamma_s \rightarrow \infty} \eta_m^2 = \eta_g^2 \frac{\beta_g}{\beta_s + \beta_g}, \quad (49)$$

and comparing (50) and (48) reveals an intriguing result: for sufficiently fast fluctuation in $\mathbf{g}(t)$, such that, $\beta_g^2 > \beta_m \beta_s$

$$\lim_{\gamma_s \rightarrow \infty} \eta_m^2 > \lim_{\gamma_s \rightarrow 0} \eta_m^2, \quad (50)$$

i.e., strong mRNA-miRNA feed forward interaction will actually amplify noise compared to the open loop system. This counter intuitive result can be understood from the fact that while the incoherent feed forward loop cancels external noise, it also decreases the mRNA half-life which causes more noise propagation from \mathbf{g} to \mathbf{m} . The balance of these two effects results in an interesting set of behaviors illustrated in Fig 4 which plots (47) as a function of γ_s . In particular, we see

- For slow fluctuations in $\mathbf{g}(t)$, mRNA noise level decreases with increasing γ_s as expected from our analysis of static noise (green line).
- When time-scale of fluctuations in $\mathbf{g}(t)$ are comparable to mRNA and miRNA half-lives, η_m^2 first decrease with increasing γ_s to reach a minimum, and then increase with γ_s to create a U-shape profile (blue line).
- For sufficient fast fluctuations in $\mathbf{g}(t)$, the U-shape become shallower and shifts to the left such that η_m^2 is mostly an increasing function of γ_s (orange line).

In summary, our results show that if the time-scale of fluctuations in the external variable is somewhat comparable to the system dynamics, then introducing noise statically in the system can lead to erroneous results. Given the space constraints, we have only focused on the noise at the mRNA level in sgFFLs, and future work investigating dynamic noise propagation at the protein level in general mRNA-miRNA feed forward circuits is clearly warranted.

V. CONCLUSION AND FUTURE WORK

Here we have identified the regimes for noise rejection at the mRNA and the protein level for miRNA-mediated IFFLs. In the case of static extrinsic noise, we investigated how the different modes of post-transcriptional regulation by miRNAs affect noise levels. Interestingly, we found that IFFLs are the only microRNA mediated process with extrinsic noise reduction capability. Furthermore, we found that due to translational inhibition, the regimes for noise rejection are quite different at the level of mRNA and protein. While ordinarily this added non-linearity leads to an increase in noise, under certain biologically realistic circumstances, this effect can be reversed leading to a greater range of noise reduction at the protein level. We also mathematically verify the experimental observation [4] that high mRNA-miRNA binding rate can lead to an increase in extrinsic noise.

$$\frac{\eta_m^2}{\eta_g^2} = \frac{\beta_m^2(\beta_m + \beta_s)(\beta_s + \beta_g) + \beta_m^2(\beta_s + 3\beta_g)\gamma_s d(\mathbf{g}) + 3\beta_m\beta_g(\gamma_s d(\mathbf{g}))^2 + \beta_g(\gamma_s d(\mathbf{g}))^3}{(\beta_s + \beta_g)(\beta_m + \gamma_s d(\mathbf{g}))(\beta_m + \beta_s + \gamma_s d(\mathbf{g}))(\beta_m + \beta_g + \gamma_s d(\mathbf{g}))} \quad (47)$$

For the case of dynamic extrinsic noise, we limit our investigation to the special case of co-expressed mRNA and miRNA (sgFFL system). Our results show counter-intuitive effects if the time-scale of fluctuations in $\mathbf{g}(t)$ are comparable or faster than the time scale of mRNA and miRNA turnover. In particular, while for static extrinsic noise the mRNA noise is predicted to monotonically decrease with increasing mRNA-miRNA interaction strength, in the case of dynamic noise (for all other model parameters fixed) there exists an optimal feed forward strength that minimizes mRNA noise.

Future work will be aimed at extending this analysis to intrinsic noise. Such a study will enable a more comprehensive understanding of the noise rejection properties of such systems. Another promising direction of future work relates to studying the noise modulation by more complicated miRNA mediated network motifs, particularly those involving multiple correlated or uncorrelated miRNAs.

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