Modeling of DNA binding to the condensin hinge domain using molecular dynamics simulations guided by atomic force microscopy

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ABSTRACT

The condensin protein complex compacts chromatin during mitosis using its DNA-loop extrusion activity. Previous studies proposed scrunching and loop-capture models as molecular mechanisms for the loop extrusion process, both of which assume the binding of double-strand (ds) DNA to the so-called hinge domain formed at the interface of the condensin subunits Smc2 and Smc4. However, how the hinge domain contacts dsDNA has remained unknown, potentially due to its conformational plasticity. Here, we conducted atomic force microscopy imaging of the budding yeast condensin holo-complex and used this data as basis for coarse-grained molecular dynamics simulations to model the hinge structure in a transient open conformation. We then simulated the dsDNA binding to open and closed hinge conformations, predicting that dsDNA binds to the outside surface when closed and to the outside and inside surfaces when open. Our simulations also suggested that the hinge can close around dsDNA bound to the inside surface. The conformational change of the hinge domain might be essential for the dsDNA binding regulation and play important roles in condensin-mediated DNA-loop extrusion.

INTRODUCTION

Proper chromosome condensation before mitosis is essential to ensure that eukaryotic genomes are correctly distributed into the two daughter cells during cell division (1–3). The condensin complex, a member of the Structural Maintenance of Chromosomes (SMC) protein family, is a key player for chromosome condensation. Indeed, purified condensin and topoisomerase II are sufficient to induce the formation of cylindrical chromatids *in vitro* (4, 5). Recent single-molecule experiments revealed that the budding yeast condensin is a molecular motor (6) that extrudes a DNA loop (7, 8), possibly through its ring structure. However, the detailed molecular mechanism of DNA loop extrusion by condensin and other members of the SMC protein family remains unclear (8–10).

The yeast condensin complex is composed of five subunits called Smc2, Smc4, Brn1, Ycg1, and Ycs4 (11). Each of Smc2 and Smc4 forms an anti-parallel coiled-coil (12) that brings together N-terminal and C-terminal regions to create a globular ATPase ‘head’ domain at one end (13). The other ends of Smc2 and Smc4 interact with each other to form the doughnut-shaped ‘hinge’ domain (14, 15). Brn1 interacts with the head domains of Smc2 and Smc4, closing the condensin ring (13, 16). Ycg1 and Ycs4 are largely composed of HEAT repeats and bind to the central region of Brn1. Ycg1 and potentially also Ycs4 bind double-strand (ds) DNA, and the crystal structure of the Ycg1/Brn1/dsDNA complex is available (13, 17). In addition, biochemical experiments and atomic force microscopy (AFM) imaging suggest that the hinge domain binds strongly to single-strand (ss) DNA (14, 18–21) and weakly to dsDNA (15, 22), respectively.

The two coiled-coils, one from Smc2 and the other from Smc4, emanate from the hinge

domain. Previous studies revealed two conformations of the hinge domain; i) a ‘closed’ conformation in which the two coiled-coils are closely juxtaposed and ii) an ‘open’ conformation in which the coiled-coils are further apart (15, 22–27). Although these conformations have been structurally characterized, there is no clear understanding about the underlying mechanism of this transition. The binding of dsDNA to the hinge domain have been proposed to regulate the inter-head engagement at the other ends of the coiled-coils, promoting ATP hydrolysis (18, 28). Previous studies have proposed ‘scrunching’ (22) and ‘loop capture’ models (29) for the molecular mechanism of the DNA loop extrusion process. The two models are not mutually exclusive and both assume dsDNA binding to the hinge domain. However, no structural information about the interaction of the hinge with dsDNA is available.

In this study, we conducted atomic force microscopy imaging of the budding yeast condensin holo-complex, and used the data to guide coarse-grained (CG) molecular dynamics (MD) simulations (30, 31) to predict the dsDNA-binding surface of the condensin hinge domain in both, open and closed conformations. First, we demonstrated that our simulations correctly predict the Ycg1/dsDNA complex structure (17), validating our CG model. Next, we conducted high-speed (HS)-AFM imaging of the whole condensin complex to obtain open hinge images (22, 32). Using this data to constrain CGMD simulation (33), we obtained a structural model for the open hinge conformation. Finally, we predicted dsDNA binding surfaces on the hinge using CGMD simulations. These simulations identifies a transient dsDNA binding surface on the inside and outside of the hinge domain, consistent with the previous biochemical (21) and HS-AFM (22) studies. It was also suggested that the hinge opening and closing regulate the dsDNA binding to the hinge domain. Our results support the model in which structural changes in the hinge domain caused by ATP hydrolysis and the accompanying dsDNA binding contribute

to the chemo-mechanical coupling in the condensin loop extrusion.

MATERIALS AND METHODS

Prediction of the Ycg1/dsDNA and Ycs4/dsDNA complex structures

Previous studies have shown that electrostatic interactions dominate protein/dsDNA interactions (34). It was also demonstrated that CGMD simulations that consider only electrostatic interactions and excluded volume effect as interactions between proteins and dsDNA could predict protein/dsDNA complex structures with a certain degree of accuracy (35–37). In this study, to validate the model, the Ycg1/dsDNA complex structure, whose crystal structure has already been solved (17), was assessed by CGMD simulations. In the crystal structure of the Ycg1/Brn1/dsDNA complex, the dsDNA fits into the Ycg1 binding cleft and is covered by Brn1. For simplicity, we neglected Brn1 and conducted the blind prediction of the Ycg1/dsDNA complex structure. As a reference for binding strength, we also predicted the Ycs4/dsDNA complex structure.

Here, we provide a brief explanation of the AICG2+ protein model used in this study (Please refer to the original work (38) for detail). In this model, each amino acid is represented by one bead located on the C_α atom position. Consecutive amino acids are connected by elastic bonds. A sequence-based statistical potential is used to model bond angles and dihedral angles. An excluded volume interaction prevents two beads from overlapping each other. The native-structure-based contact potential (Gō potential) constrains the distance between amino acid pairs that contact each other in the native Ycg1 (PDB ID: 5OQN) and Ycs4 structures. To model the native Ycs4 structure, we performed homology modeling using *Chaetomium Thermophilum* Ycs4 structure (PDB ID: 6QJ4, 47.7% sequence identity) as a template

(<https://salilab.org/modeller/>). Parameters in the AICG2+ model were decided so that the fluctuation of each amino acid in reference proteins reproduces that of all-atom simulations.

Next, we provide a brief explanation of the 3SPN.2C DNA model (Please refer to the original work (39) for detail). In this model, each nucleotide is represented by three beads located at the positions of base, sugar, and phosphate units. Neighboring sugar-phosphate and sugar-base are connected by virtual bonds. Bond angles and dihedral angles are restrained to their values in reference B-type dsDNA. An excluded volume interaction prevents any two beads from overlapping each other. Orientation-dependent attractive potentials are applied to base pairs, cross-stacking pairs, and intra-chain-stacking pairs. The parameters are decided so that the model reproduces several types of experimental data.

When a protein forms a complex with dsDNA, the electrostatic interaction plays an important role (34). Instead of putting unit charges on the C_α atoms of charged residues, we placed partial charges on the C_α atoms of surface residues. The RESPAC algorithm (35) decided the partial charges of the surface residues to reproduce the electrostatic potential calculated for the all-atom structure model. The Debye-Hückel potential represents the interaction between charge pairs. The potential is expressed as

$$V_{DH} = \sum_{i < j}^N \frac{q_i q_j}{4\pi\epsilon_0\epsilon_k r_{ij}} e^{-r_{ij}/\lambda_D}, \quad (1)$$

where q_i is charge, ϵ_0 is the electric constant, and ϵ_k is the dielectric constant. The Debye length λ_D is expressed as

$$\lambda_D = \sqrt{\frac{\epsilon_0\epsilon_k k_B T}{2N_A e^2 I}}, \quad (1)$$

where k_B is the Boltzmann constant, T is temperature, N_A is the Avogadro's number, and I is

ionic strength. In addition to the electrostatic interaction, the excluded volume interaction is applied.

In the simulations, we put one protein molecule (Ycg1 or Ycs4) and five 18 base-pairs (bps) dsDNA molecules in the 200 Å x 160 Å x 200 Å box (Ycg1) and 210 Å x 170 Å x 210 Å box (Ycs4), respectively. The time evolution is modeled by Langevin dynamics with the friction coefficient and the time step set to 0.8435 and 0.2, respectively. The monovalent ion concentration was varied as 100, 150, 200, 250, and 400 mM. We conducted the simulations using CafeMol version 3.2 (30) (<https://www.cafemol.org>).

HS-AFM imaging of a condensin holo-complex

We conducted HS-AFM imaging to observe equilibrium dynamics of a whole condensin complex on bare mica surface in solution (40). In HS-AFM imaging, the forces that a sample imposes on a probe (tapping-cantilever) are used to form an image of a three-dimensional shape of a protein surface.

We purified condensin holo-complexes as described previously (6). Briefly, the five subunits of a condensin complex were co-overexpressed in *Saccharomyces cerevisiae* using galactose-inducible promoters. Cultures were grown at 30°C in URA/TRP dropout media containing 2% raffinose to OD₆₀₀ of 1.0. Expression was induced with 2% galactose for 8 hours. Then, cells were harvested by centrifugation, re-suspended in buffer A (50 mM Tris-HCl pH 7.5, 200 mM NaCl, 5% (v/v) glycerol, 5 mM β-mercaptoethanol, and 20 mM imidazole) containing 1x complete EDTA-free protease-inhibitor mix (Roche) and lysed in a FreezerMill (Spex). The lysate was cleared by two rounds of 20 min centrifugation at 45,000 x g at 4°C and loaded onto a

5-mL HisTrap column (GE Healthcare) pre-equilibrated with buffer A. The resin was washed with five column volumes of buffer A containing 500 mM NaCl; buffer A containing 1 mM ATP, 10 mM KCl and 1 mM MgCl₂; and then buffer A containing 40 mM imidazole to remove non-specifically bound proteins. Protein was eluted in buffer A containing 200 mM imidazole and transferred to buffer B (50 mM Tris-HCl pH7.5, 200 mM NaCl, 5% (v/v) glycerol, 1 mM DTT) using a desalting column. After the addition of EDTA to 1 mM, PMSF to 0.2 mM, and Tween20 to 0.01%, the protein was incubated overnight with 2 mL (bed volume) of pre-equilibrated Strep-Tactin high-capacity Superflow resin (IBA). The Strep-Tactin resin was packed into a column and washed with 15 resin volumes of buffer B by gravity flow. Protein was eluted with buffer B containing 5 mM desthiobiotin. The elute was concentrated by ultrafiltration and loaded onto a Superose 6 size exclusion chromatography column (GE Healthcare) pre-equilibrated in buffer B containing 1 mM MgCl₂. Peak fractions were pooled and concentrated to 4 μM by ultrafiltration. Purified proteins were analyzed by SDS-PAGE (NuPAGE 4-12% Bis-Tris protein gels, ThermoFisher Scientific).

We used a laboratory-built tapping-mode HS-AFM apparatus together with small cantilevers (BLAC10DS-A2, Olympus) designed for HS-AFM imaging (spring constant, ~0.1 N/m, resonant frequency in water, ~0.5 MHz; quality factor in water, ~1.5) (41). We deposited 2 μL of 1 nM condensin in buffer (40 mM Tris-HCl pH7.5, 125 mM KCl, 10 mM MgCl₂, 1 mM DTT, 4 mM ATP) onto the scanning mica stage of an HS-AFM apparatus. We conducted imaging at room temperature. The cantilever's free oscillation amplitude A_0 and set-point amplitude A_s were set at 1-2 nm and $\sim 0.9 \times A_0$, respectively. The scan area was 150 x 150 nm², and the imaging rate was 150 ms/frame. To define a hinge angle, first, we found out the pixel with the highest height value in a hinge globular domain. We calculated the total height values of the

parallelogram regions with a height of 10 pixels and a width of 5 pixels that extended radially from that point and selected two of them with local maximum values, respectively. The angle between the two lines passing through the center of these regions was defined to be the hinge angle.

AFM fitting simulations of the hinge

Similar to the earlier reports, in the current HS-AFM imaging, we consistently observed open hinge conformations (22). To obtain the open hinge structure, we performed CG flexible-fitting MD simulations in which the restraint potential energy derived from the HS-AFM image (33) is imposed, stabilizing the open hinge structure.

The crystal structure of the hinge domain (PDB ID: 4RSI) was used as the initial structure. Using PyMol (<https://pymol.org>), we manually performed rigid body rotation and translation of Smc2 and Smc4 subunits, respectively, to roughly fit the AFM image.

Starting from this structure, we performed CG flexible-fitting MD simulations (33). The settings were similar as for the Ycg1/dsDNA complex described above, except that the AFM potential was imposed. The AFM potential was defined by

$$V_{AFM}(\mathbf{R}) = \kappa N_B T (1 - c.s.(\mathbf{R})), \quad (1)$$

where $\mathbf{R}$ collectively represents protein coordinates, k_B is the Boltzmann constant, T is the temperature (set to 300K), and $c.s.(\mathbf{R})$ is the similarity score (the cosine similarity, originally designated as a modified correlation coefficient) between the experimental AFM image and the pseudo-AFM image calculated from the simulations defined by

$$c.s.(\mathbf{R}) = \frac{\sum_p^{pixels} H_p^{(exp)} H_p^{(sim)}(\mathbf{R})}{\sqrt{\sum_p^{pixels} \left(H_p^{(exp)}\right)^2} \sqrt{\sum_p^{pixels} \left(H_p^{(sim)}(\mathbf{R})\right)^2}} \quad (2)$$

where $H_p^{(exp)}$ and $H_p^{(sim)}(\mathbf{R})$ are an experimental AFM image (a matrix listing all height values) and a pseudo-AFM image calculated from the simulated structure $\mathbf{R}$, respectively.

To decide the value of the parameter κ , we conducted 240 runs of 10^7 -steps simulations with κ set to 10, 3000, and 10000, respectively. Then, to assess how the native structures were perturbed during the fitting, we calculated the standard nativeness score, Q-score, for intra-subunit and inter-subunit contacts. Q-score is defined as the ratio of the number of bead-bead contacts formed at a certain moment in simulations to the number of contacts formed in the native protein structure. The particle pairs with a distance of less than native distance times 1.2 are considered to contact each other in CGMD simulations.

Prediction of hinge/dsDNA complex structures

To predict dsDNA binding surfaces on the hinge domain, we placed the hinge fragment and five 18-bps dsDNA molecules in a simulation box (330 Å x 330 Å x 330 Å box for open hinge and 255 Å x 255 Å x 255 Å box for closed hinge) and performed CGMD simulations. The last step of the AFM fitting simulation and the crystal structure (PDB ID: 4RSI) were used as the initial conformations for the open and closed hinge simulations, respectively. Initially, the DNA molecules were randomly placed in the box. The setup was the same as the CGMD simulations of the Ycg1 (or Ycs4)/dsDNA complex described above. In the open hinge simulations, the AFM potential with κ set to 3000 stabilized the open conformation.

RESULTS

Blind prediction of the Ycg1/dsDNA complex

As a proof-of-concept that CGMD simulations are able to correctly predict DNA-binding surfaces in the condensin complex, we first applied our approach to Ycg1. Since crystal structures of the Ycg1/Brn1/dsDNA complex have previously been reported (PDB ID: 5OQN), we can compare a blind prediction by simulations with the actual structure to assess the accuracy of simulations. For simplicity, we neglected Brn1 and conducted the blind prediction of the Ycg1/dsDNA complex structure.

We started the simulations with a single Ycg1 molecule at the center of the simulation box and five 18-base pair (bp) dsDNA molecules randomly placed around Ycg1 (Fig. 1A). In our CG model, the RESPAC algorithm (35) arranges the particle charges to reproduce the electrostatic potential calculated based on the all-atom charge arrangement. Therefore, we can predict DNA binding region more accurately than the conventional CG models (36, 37). After 1×10^7 steps of simulations, the dsDNA molecules spontaneously bound to the protein. To identify the dsDNA binding surface, we calculated the probability of protein-DNA contact for every amino-acid (Fig. 1B). At the simulated salt concentration of 100 mM, residues with high DNA-contact probability values clustered in two regions on the protein surface referred to as sites I and II in Fig. 1B. At higher salt concentration (400 mM; a condition of weaker electrostatic interactions), we only detected high contact probabilities for residues at the site I. Notably, the site I makes significant number of contacts with dsDNA in the available crystal structure (PDB ID: 5OQN).

To quantify the binding strength, we calculated the survival probability of dsDNA molecules

bound to the sites I and II during consecutive simulation steps (Fig. 1C). The dissociation rates of the sites I and II, estimated from the slopes of the fitting of the survival probabilities, were 1.2×10^{-7} (steps⁻¹) and 1.0×10^{-6} (steps⁻¹), respectively. Thus, in agreement with the experimental data, the CGMD simulations predicted the site I as the strongest dsDNA binding surface of Ycg1.

We performed conformational clustering to compare the predicted Ycg1/dsDNA complex structure with the crystal structure (PDB ID: 5OQN). First, we calculated the minimum of the distances (d_{min}^i) between the i -th protein residue of the site I and every dsDNA particle. Second, we carried out the principal component analysis of structures using d_{min}^i as coordinates. We then mapped each structure onto the first three principal component space. Finally, we counted the number of structures in a sphere with a radius of 10.0 centered on each structure in the unitless principal component space and defined the structure with the maximum value as the representative structure (Fig. 1D left). The similarity between the representative structure and the crystal structure suggests that DNA binding surfaces can be successfully predicted using the CGMD simulations. Subtle differences between the two structures are most likely due to the simplicity of our CG model and the absence of Brn1 which, in addition to Ycg1, makes contacts with dsDNA in the crystal structure.

To quantitatively assess similarity between the representative structure and the crystal structure, we plotted the probability of each protein residue contacting to dsDNA (Fig. 1E, orange line, 100 mM salt) and compared the probability with the DNA contact identified in the crystal structure (Fig. 1E, blue bars). Most of the residues that contact dsDNA in the crystal structure were also predicted to contact dsDNA in the representative structure obtained by our

simulation, with the exception of peaks between Ycg1 residues 450 to 650, which correspond to the site II. This correlation was strongly reduced when we randomly scrambled the charge arrangement (Fig. 1E, green line), which suggests that the surface charge distribution dictates the dsDNA binding surface of Ycg1. This result is consistent with the fact that Ycg1 residues exclusively contact the dsDNA phosphate backbone (17).

Prediction of the Ycs4/dsDNA complex

Next, we predicted the dsDNA binding surface on Ycs4, a potential dsDNA binder, to compare the binding strength of various proteins to dsDNA. As with the Ycg1 simulation, one Ycs4 was placed at the center of the simulation box and five 18 bp dsDNA molecules were randomly placed around it (Fig. S1A). In 1×10^7 steps of the simulations, dsDNA bound around the middle of the Ycs4 molecule with high probability (Fig. S1B). Visual inspection revealed that dsDNA binds to two distinct surfaces on Ycs4 (the surfaces I and II in Fig. S1C). Future studies should address if dsDNA (weakly) binds to these surfaces.

HS-AFM imaging of the condensin holo-complex

After having shown that the CGMD simulations can predict dsDNA binding surfaces in the condensin complex with a high degree of accuracy, we turned our attention to the Smc2/Smc4 hinge domain. To reveal the equilibrium hinge dynamics, we conducted HS-AFM imaging of the condensin holo-complex on bare mica surface in solution. In representative AFM images (Fig. 2A), the Smc2/Smc4 hinge domain, the coiled-coils and the ATPase head domains, as well as the two HEAT-repeat subunits Ycg1 and Ycs4 were observed, though, in most images, Smc2 and Smc4 head domains and HEAT-repeat subunits could not be clearly distinguished, similar to a

previous HS-AFM study (22). The HEAT-repeat subunits occasionally did not contact the head domains of Smc2 or Smc4 but stayed nearby (Fig. S2), suggesting that they were merely bound to the complex by the disordered Brn1 subunit, which was not visible in AFM images. The transient detachment of one of the HEAT-repeat subunits is consistent with the recently proposed flip-flop model (25).

By HS-AFM imaging, we observed transitions between open and closed hinge conformations (Fig. 2). Condensin with a closed hinge had a rod-like appearance, consistent with recent crystal and cryo-electron microscopy structures (23, 25). In the open conformation, where the coiled-coils are apart, we frequently observed a kink in the Smc2/Smc4 coiled-coils, which presumably corresponds to the ‘elbow’ region (the purple arrow in Fig. 2A), though we did not observe the structure in which the hinge domain folds back and contacts the coiled-coils, as observed in the recent cryo-EM study (25). To quantify the hinge dynamics, we plotted the time trajectory of hinge angle (Fig. 2B). The plot shows that, in most of the time, the fluctuations in hinge angle remained below $\sim 40^\circ$ and only occasionally exceeded to $50\text{--}80^\circ$ before spontaneously returning to values below $\sim 40^\circ$. Fig. 2C shows the frequency distribution of hinge angle. The fitting of this distribution required at least two Gaussian distributions, which indicated that there are at least two conformational states; a stable closed conformation (87%) and a transient open conformation (13%). These results reconcile the two conformations previously observed by different methods (22, 23, 25). Note that the relative stability of the open and closed conformations can vary depending on buffer conditions.

AFM-fitting simulation to model the open hinge conformation

To date, detailed structural information is available only for the closed hinge conformation (PDB: 4RSI) (15). To obtain a structure for the open hinge conformation, we performed CG flexible-fitting MD simulations for an AFM image of an open hinge conformation (Fig. 3A). In each step of these simulations, the pseudo-AFM image was calculated from the simulation snapshot. Then, the forces derived from the difference between the pseudo- and the experimental AFM images was used to move the CG particles so that the pseudo-AFM image in the next step was more similar to the experimental image. Eventually, we obtained the structure which reproduces the experimental AFM image.

We obtained the initial structure by rigid-body rotation and translation of the closed hinge crystal structure (Fig. 3B). In the subsequent flexible-fitting, the strength of forces derived from the AFM images can be scaled by a single scaling factor κ . The larger the value of κ is, the easier the fitting is, but the more distorted the subunit structures are. We varied κ as 10, 3000, and 10000. Each structure in the simulations are evaluated by three measures; the inter-subunit Q-score, the intra-subunit Q-score, and the cosine similarity between the pseudo- and experimental AFM images. The Q-score is a standard nativeness score defined as the ratio of the number of bead-bead contacts formed at a certain moment in simulations to the number of contacts formed in the native protein structure. The inter-subunit Q-score measures similarity to the closed hinge structure, whereas the intra-subunit Q-score measures the structural integrity of each subunit.

When κ was set to 10, the inter-subunit Q-score was larger than 0.8 in some simulations (6/240; Fig. 3C top left). This result suggests that the hinge occasionally closes in the fitting

simulation despite the presence of the AFM potential which stabilizes the open conformation. On the other hand, when κ was set to 10,000, the inter-subunit Q-score ranged from 0.6 to 0.8 in some simulations (5/240; Fig. 3C top right). The visual inspection confirmed that the structures with the inter-subunit Q-score in this range correspond to the open hinge. In this case, however, the intra-subunit Q-score is less than 0.95 in some simulations (6/240, Fig. 3C bottom right), suggesting that the subunit structures are distorted. When κ was set to 3,000, the inter-subunit Q-score was in the range from 0.6 to 0.8 in a few simulations (2/240; Fig. 3C top center). Also, the intra-subunit Q-score was almost always larger than 0.95 (239/240; Fig. 3C bottom center). These results suggest that $\kappa = 3,000$ is an appropriate value to stabilize the open hinge structure. We also tested $\kappa = 1,000$ and 5,000, and obtained similar results to those with $\kappa = 3,000$. Thus, we set κ as 3,000 for the rest of this study.

We plotted the time trajectories of the cosine similarity between the pseudo- and experimental AFM images (*c. s.* in eq. 2) and inter-subunit Q-score (Fig. 3D and E). Fig. 3D shows that *c. s.* increases right after the start of the simulation. Conversely, the inter-subunit Q-score takes longer until it suddenly reaches 0.7 (Fig. 3E). These plots suggest that the two subunits are initially roughly positioned according to the AFM potential and subsequently, the inter-subunit contacts are recovered according to the native structure-based potential.

In the AFM image, Smc2 and Smc4 subunits could not be distinguished (Fig. 3A). In the initial structure, Smc2 and Smc4 were placed in the region I and II, respectively, by a random choice. When the position was exchanged (Smc2 in region II and SMC4 in region I), we obtained almost the same, but the upside-down structure. This result suggests that the AFM potential dictates the hinge angle but not the orientation of the molecule.

From our simulations, we obtained a structure for the open hinge conformation that reproduces the experimental AFM images (Fig. 3F). Superposition of a part of globular domains (residue 505 to 676 of Smc2 and 671 to 835 of Smc4) of the open (simulated structure) and closed (crystal structure PDB ID: 4RSI) hinge structures illustrates that the Smc2 coiled-coil has the same orientation in both conformations (Fig. 3G). Therefore, this result suggests that the Smc4 coiled-coil changes its orientation drastically upon the hinge conformational change.

Prediction of dsDNA binding surfaces on the hinge domain

To predict dsDNA binding surfaces on the hinge domain, we performed CGMD simulations of dsDNA with the closed hinge structure (PDB:4RSI) and the open hinge structure obtained by the AFM fitting simulation (Fig. 3F). During the simulation, we set κ to 3,000 to stabilize the open hinge structure. We then placed five 18-bps dsDNA molecules randomly in the simulation box.

After starting the open and closed hinge simulations, we observed spontaneous dsDNA association with the hinge (Fig. 4A and B). In the open hinge simulation, dsDNA bound both to the ‘inside’ region between the coiled-coils and to the opposite ‘outside’ region (Fig. 4A and C). This result is consistent with previous suggestions for dsDNA and ssDNA binding surfaces based on the conservation of positively charged residues in the hinge domain of SMC proteins (15, 21, 22). As expected, dsDNA could not access the inside region in the closed hinge simulation (Fig. 4B and D). Instead, our simulation predicts that dsDNA weakly binds to two distinct sites on the outside region. The surface is accessible even in the folded-back state of the condensin holo-complex observed in the previous cryo-EM study (25). The comparison of Fig. 4C with D suggests that these two binding sites in the closed hinge merge into one in the open hinge.

We calculated the electrostatic potential on the hinge surface using the APBS plugin of PyMol (<https://pymol.org>) to see the difference of electrostatic potential between the open and closed hinge structures (Fig. 4E and F). Fig. 4E (left) shows that positively charged surfaces inside the hinge are exposed upon the hinge opening. Comparison of Fig. 4E (right) with F (right) shows that two positively charged surfaces in the closed hinge merge into one region in the open hinge. Together, the hinge opening is required to make the inner dsDNA binding site accessible and to merge two small positively charged outside surfaces into a single large surface.

In this study, we performed CGMD simulations of Ycg1/dsDNA (Fig. 1), Ycs4/dsDNA (Fig. S1), and hinge/dsDNA complexes (Fig. 4). To compare the DNA-binding strength predicted for the hinge domain to that predicted for the Ycg1 and Ycs4, we plotted the survival probabilities of DNA molecules bound to the different DNA-binding surfaces at 400 mM salt concentration (a weak binding condition) (Fig. 4G). The simulations showed that DNA bound most strongly to Ycg1 (dissociation rate constant 1.2×10^{-7} steps⁻¹), followed by Ycs4 (7.7×10^{-7} steps⁻¹) and the hinge domain (open outside 4.1×10^{-6} steps⁻¹, closed outside 4.9×10^{-6} steps⁻¹, open inside 5.4×10^{-6} steps⁻¹). The fastest dissociation from the inside of the hinge domain suggests that, although dsDNA can potentially access the open hinge inside, any interaction might only be very transient. The ability to detect transient DNA binding, which might be missed in biochemical and structural studies, is the strength of MD simulations.

At the 400 mM salt concentration, dsDNA binding at the hinge domain was too weak to compare the binding strengths between the different surfaces. Therefore, we repeated the simulation at the 200 mM salt for the open and the closed hinge with dsDNA to calculate the survival probabilities (Fig. 4H). The result shows that dsDNA binding to the outside region of

the open hinge (5.4×10^{-7} steps⁻¹) is slightly stronger than the inside region (1.2×10^{-6} steps⁻¹) or the outside region of the closed hinge (1.4×10^{-6} steps⁻¹), further confirming that dsDNA binding to the inside region of the hinge domain might only be transient. Note that the dsDNA binding to any region of the open and closed hinges is weaker than the binding to Ycg1.

We also tested whether the hinge can close while dsDNA remains bound to the inside surface. To accomplish this, we performed 10 mM salt simulations without the AFM fitting potential using the open hinge with dsDNA bound to the inner site as an initial structure. In these simulations, the hinge closed within 2.1×10^6 steps on average without dsDNA dissociation (79/80 simulations) (Fig. 5A and B). Thus, we obtained the closed hinge structure with dsDNA trapped inside.

Next, we performed 400 mM salt simulations using the closed hinge structure with dsDNA trapped inside as an initial structure and observed the closed hinge slid off from one end of dsDNA. The survival probability plots (Fig. 5C) show that dsDNA binding to the closed hinge inside (9.2×10^{-8} step⁻¹) is significantly stronger than the open hinge inside (5.4×10^{-6} step⁻¹) and is comparable to the binding strength calculated for Ycg1 (1.2×10^{-7} step⁻¹). This result indicates that dsDNA can access the inside of the open hinge and might then induce the closing of the hinge, which considerably stabilizes dsDNA binding.

We then assessed the importance of amino acid residues on the inner binding site of the hinge domain. We performed a PSI-BLAST search using the amino acid sequence of *Saccharomyces cerevisiae* condensin as a query and obtained a dataset containing 1,000 related sequences from various species. We calculated a conservation score for each residue and plotted the distribution of these scores (Fig. 6A). The frequency distribution of the conservation scores (Fig. 6A) has two

peaks, one around 0.0 and the other around 0.9. The amino acids with a high conservation score may play an essential role in the condensin function. The hinge structure in which each amino acid is colored according to its conservation score (Fig. 6B) reveals that the positively charged amino acids (arginine and lysine) on the inside have relatively high conservation scores (> 0.8), fully consistent with the previous studies (19, 21). Thus, these residues likely have an essential role in the dsDNA binding.

DISCUSSION

The purpose of this study is to predict the dsDNA binding surfaces on the condensin hinge domain which has been suggested in literature (14, 15, 18–22). In this study, we conducted HS-AFM imaging of the budding yeast condensin holo-complex and used this data as basis for coarse-grained molecular dynamics simulations to model the hinge structure in a transient open conformation. We then simulated the dsDNA binding to open and closed hinge conformations, predicting that dsDNA binds to the outside surface when closed and to the outside and inside surfaces when open. The simulations also show that the coiled-coils might be able to close around dsDNA bound to the inside region of the hinge and thereby stabilize the dsDNA binding.

Soh *et al.* showed that the yeast condensin hinge domain binds to dsDNA. However, they did not address the role of this binding *in vivo*. Recently, Nunez *et al.* created a mutant (3A-hng) in which three amino acids involved in the dsDNA binding of the bacterial SMC hinge were replaced with alanine (42). The mutations do not significantly affect translocation of the bacterial SMC complex along dsDNA, indicating that the hinge is irrelevant for DNA translocation. However, they observed that the 3A-hng head domain binds to dsDNA more stably than the

wild-type protein, suggesting that this binding acts as a replacement for the hinge-dsDNA binding. It has not been explored if the dsDNA binding to the yeast condensin hinge is essential *in vivo*. Thus, *in vivo* experiments should be performed using mutant yeast condensin in the near future.

Uchiyama *et al.* revealed that ssDNA binds to the outside surface of the hinge (21). In this binding, the initially buried binding surface should be exposed by the hinge conformational change. Consistently, the current study suggested that dsDNA binds to the hinge outside. The *in vivo* function of the ssDNA and dsDNA binding should be investigated in the future.

In this study, we used CGMD simulations to predict the protein/dsDNA complex structures. The CGMD simulations increase the calculation speed with some compromise on accuracy compared with all-atom (AA) MD simulations. The CG model neglects the protein-dsDNA hydrophobic interaction, based on the assumption that the electrostatic interaction is dominant in the protein-dsDNA interaction (34). However, the prediction could be more accurate if we introduce hydrophobic interactions into the CG model (43). Also, the predicted CG structure could be refined using AAMD simulations (44). It is expected that improvement of the CGMD simulation method will lead to a more accurate prediction of a protein/DNA complex structure.

Previous studies proposed the scrunching (22) and the loop capture models (29) (not mutually exclusive) as molecular mechanisms for the loop extrusion, both of which assume the dsDNA binding to the hinge domain. The simulations in this study support a model in which dynamic structural changes in the hinge regulate the dsDNA binding. The previous study suggested that the structural change around the head domain induced by ATP hydrolysis is propagated along coiled-coils to the hinge domain and leads to hinge opening. Based on this model and results

from the current study, we speculate that the ATP hydrolysis might control the binding of dsDNA to the hinge domain. Thus, the hinge closing and opening, and the accompanying dsDNA binding and dissociation might underlie the molecular mechanism of chromatin loop extrusion by condensin (29, 45).

AUTHOR CONTRIBUTIONS

T.T. designed the study. H.K. designed and conducted molecular dynamics simulations, AFM experiments, and data analysis. N.K. designed and conducted AFM experiments. S.B. and T.T. purified condensin complexes. All authors discussed the findings and co-wrote the manuscript.

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REFERENCES

1. Naumova,N., Imakaev,M., Fudenberg,G., Zhan,Y., Lajoie,B.R., Mirny,L.A. and Dekker,J. (2013) Organization of the Mitotic Chromosome. *Science*, **342**, 948–953.
2. Gibcus,J.H., Samejima,K., Goloborodko,A., Samejima,I., Naumova,N., Nuebler,J., Kanemaki,M.T., Xie,L., Paulson,J.R., Earnshaw,W.C., *et al.* (2018) A pathway for mitotic chromosome formation. *Science*, **359**, eaao6135.
3. Batty,P. and Gerlich,D.W. (2019) Mitotic Chromosome Mechanics: How Cells Segregate Their Genome. *Trends in Cell Biology*, **29**, 717–726.
4. Shintomi,K., Takahashi,T.S. and Hirano,T. (2015) Reconstitution of mitotic chromatids with a minimum set of purified factors. *Nat Cell Biol*, **17**, 1014–1023.
5. Shintomi,K., Inoue,F., Watanabe,H., Ohsumi,K., Ohsugi,M. and Hirano,T. (2017) Mitotic chromosome assembly despite nucleosome depletion in *Xenopus* egg extracts. **356**, 1284–1287.
6. Terakawa,T., Bisht,S., Eeftens,J.M., Dekker,C., Haering,C.H. and Greene,E.C. (2017) The condensin complex is a mechanochemical motor that translocates along DNA. *Science*, **358**, 672–676.
7. Ganji,M., Shaltiel,I.A., Bisht,S., Kim,E., Kalichava,A., Haering,C.H. and Dekker,C. (2018) Real-time imaging of DNA loop extrusion by condensin. *Science*, **360**, 102–105.
8. Kim,E., Kerssemakers,J., Shaltiel,I.A., Haering,C.H. and Dekker,C. (2020) DNA-loop extruding condensin complexes can traverse one another. *Nature*, **579**, 438–442.
9. Davidson,I.F., Bauer,B., Goetz,D., Tang,W., Wutz,G. and Peters,J.-M. (2019) DNA loop extrusion by human cohesin. *Science*, **366**, 1338–1345.
10. Golfier,S., Quail,T., Kimura,H. and Brugués,J. (2020) Cohesin and condensin extrude DNA loops in a cell-cycle dependent manner. *eLife*, **9**, e53885.
11. Hassler,M., Shaltiel,I.A. and Haering,C.H. (2018) Towards a Unified Model of SMC Complex Function. *Current Biology*, **28**, R1266–R1281.
12. Matityahu,A. (2018) A new twist in the coil: functions of the coiled-coil domain of structural

- maintenance of chromosome (SMC) proteins. *Current Genetics*, **64**, 109–116.
13. Hassler,M., Shaltiel,I.A., Kschonsak,M., Simon,B., Merkel,F., Thärichen,L., Bailey,H.J., Macošek,J., Bravo,S., Metz,J., *et al.* (2019) Structural Basis of an Asymmetric Condensin ATPase Cycle. *Molecular Cell*, **74**, 1175-1188.e9.
 14. Hirano,M. (2002) Hinge-mediated dimerization of SMC protein is essential for its dynamic interaction with DNA. *The EMBO Journal*, **21**, 5733–5744.
 15. Soh,Y.-M., Bürmann,F., Shin,H.-C., Oda,T., Jin,K.S., Toseland,C.P., Kim,C., Lee,H., Kim,S.J., Kong,M.-S., *et al.* (2015) Molecular Basis for SMC Rod Formation and Its Dissolution upon DNA Binding. *Molecular Cell*, **57**, 290–303.
 16. Schleiffer,A., Kaitna,S., Maurer-Stroh,S., Glotzer,M., Nasmyth,K. and Eisenhaber,F. (2003) Kleisins: A Superfamily of Bacterial and Eukaryotic SMC Protein Partners. *Molecular Cell*, **11**, 571–575.
 17. Kschonsak,M., Merkel,F., Bisht,S., Metz,J., Rybin,V., Hassler,M. and Haering,C.H. (2017) Structural Basis for a Safety-Belt Mechanism That Anchors Condensin to Chromosomes. *Cell*, **171**, 588-600.e24.
 18. Hirano,M. and Hirano,T. (2006) Opening Closed Arms: Long-Distance Activation of SMC ATPase by Hinge-DNA Interactions. *Molecular Cell*, **21**, 175–186.
 19. Griesse,J.J., Witte,G. and Hopfner,K.-P. (2010) Structure and DNA binding activity of the mouse condensin hinge domain highlight common and diverse features of SMC proteins. *Nucleic Acids Research*, **38**, 3454–3465.
 20. Griesse,J.J. and Hopfner,K.-P. (2011) Structure and DNA-binding activity of the *Pyrococcus furiosus* SMC protein hinge domain. *Proteins*, **79**, 558–568.
 21. Uchiyama,S., Kawahara,K., Hosokawa,Y., Fukakusa,S., Oki,H., Nakamura,S., Kojima,Y., Noda,M., Takino,R., Miyahara,Y., *et al.* (2015) Structural Basis for Dimer Formation of Human Condensin Structural Maintenance of Chromosome Proteins and Its Implications for Single-stranded DNA Recognition. *J. Biol. Chem.*, **290**, 29461–29477.
 22. Ryu,J.-K., Katan,A.J., van der Sluis,E.O., Wisse,T., de Groot,R., Haering,C.H. and Dekker,C. (2020) The condensin holocomplex cycles dynamically between open and collapsed

- states. *Nat Struct Mol Biol*, **27**, 1134–1141.
23. Anderson,D.E., Losada,A., Erickson,H.P. and Hirano,T. (2002) Condensin and cohesin display different arm conformations with characteristic hinge angles. *The Journal of Cell Biology*, **156**, 419–424.
 24. Kamada,K., Su’etsugu,M., Takada,H., Miyata,M. and Hirano,T. (2017) Overall Shapes of the SMC-ScpAB Complex Are Determined by Balance between Constraint and Relaxation of Its Structural Parts. *Structure*, **25**, 603-616.e4.
 25. Lee,B.-G., Merkel,F., Allegretti,M., Hassler,M., Cawood,C., Lecomte,L., O’Reilly,F.J., Sinn,L.R., Gutierrez-Escribano,P., Kschonsak,M., *et al.* (2020) Cryo-EM structures of holo condensin reveal a subunit flip-flop mechanism. *Nat Struct Mol Biol*, **27**, 743–751.
 26. Shi,Z., Gao,H., Bai,X. and Yu,H. (2020) Cryo-EM structure of the human cohesin-NIPBL-DNA complex. *Science*, **368**, 1454–1459.
 27. Krepel,D., Cheng,R.R., Di Pierro,M. and Onuchic,J.N. (2018) Deciphering the structure of the condensin protein complex. *Proc Natl Acad Sci USA*, **115**, 11911–11916.
 28. Hirano,M. (2001) Bimodal activation of SMC ATPase by intra- and inter-molecular interactions. *The EMBO Journal*, **20**, 3238–3250.
 29. Marko,J.F., De Los Rios,P., Barducci,A. and Gruber,S. (2019) DNA-segment-capture model for loop extrusion by structural maintenance of chromosome (SMC) protein complexes. *Nucleic Acids Research*, **47**, 6956–6972.
 30. Kenzaki,H., Koga,N., Hori,N., Kanada,R., Li,W., Okazaki,K., Yao,X.-Q. and Takada,S. (2011) CafeMol: A Coarse-Grained Biomolecular Simulator for Simulating Proteins at Work. *J. Chem. Theory Comput.*, **7**, 1979–1989.
 31. Takada,S., Kanada,R., Tan,C., Terakawa,T., Li,W. and Kenzaki,H. (2015) Modeling Structural Dynamics of Biomolecular Complexes by Coarse-Grained Molecular Simulations. *Acc. Chem. Res.*, **48**, 3026–3035.
 32. Eeftens,J.M., Katan,A.J., Kschonsak,M., Hassler,M., de Wilde,L., Dief,E.M., Haering,C.H. and Dekker,C. (2016) Condensin Smc2-Smc4 Dimers Are Flexible and Dynamic. *Cell Reports*, **14**, 1813–1818.

33. Niina,T., Fuchigami,S. and Takada,S. (2020) Flexible Fitting of Biomolecular Structures to Atomic Force Microscopy Images via Biased Molecular Simulations. *J. Chem. Theory Comput.*, **16**, 1349–1358.
34. Privalov,P.L., Dragan,A.I. and Crane-Robinson,C. (2011) Interpreting protein/DNA interactions: distinguishing specific from non-specific and electrostatic from non-electrostatic components. *Nucleic Acids Research*, **39**, 2483–2491.
35. Terakawa,T. and Takada,S. (2014) RESPAC: Method to Determine Partial Charges in Coarse-Grained Protein Model and Its Application to DNA-Binding Proteins. *J. Chem. Theory Comput.*, **10**, 711–721.
36. Marcovitz,A. and Levy,Y. (2011) Frustration in protein-DNA binding influences conformational switching and target search kinetics. *Proceedings of the National Academy of Sciences*, **108**, 17957–17962.
37. Marcovitz,A. and Levy,Y. (2013) Weak Frustration Regulates Sliding and Binding Kinetics on Rugged Protein–DNA Landscapes. *J. Phys. Chem. B*, **117**, 13005–13014.
38. Li,W., Terakawa,T., Wang,W. and Takada,S. (2012) Energy landscape and multiroute folding of topologically complex proteins adenylate kinase and 2ouf-knot. *Proc Natl Acad Sci USA*, **109**, 17789–17794.
39. Freeman,G.S., Hinckley,D.M., Lequieu,J.P., Whitmer,J.K. and de Pablo,J.J. (2014) Coarse-grained modeling of DNA curvature. *Journal of Chemical Physics*, **141**, 165103.
40. Ando,T. (2018) High-speed atomic force microscopy and its future prospects. *Biophys Rev*, **10**, 285–292.
41. Kodera,N., Yamamoto,D., Ishikawa,R. and Ando,T. (2010) Video imaging of walking myosin V by high-speed atomic force microscopy. *Nature*, **468**, 72–76.
42. Vazquez Nunez,R., Ruiz Avila,L.B. and Gruber,S. (2019) Transient DNA Occupancy of the SMC Interarm Space in Prokaryotic Condensin. *Molecular Cell*, **75**, 209-223.e6.
43. Onufriev,A.V. and Case,D.A. (2019) Generalized Born Implicit Solvent Models for Biomolecules. *Annu. Rev. Biophys.*, **48**, 275–296.

44. Shimizu,M. and Takada,S. (2018) Reconstruction of Atomistic Structures from Coarse-Grained Models for Protein–DNA Complexes. *J. Chem. Theory Comput.*, **14**, 1682–1694.
45. Diebold-Durand,M.-L., Lee,H., Ruiz Avila,L.B., Noh,H., Shin,H.-C., Im,H., Bock,F.P., Bürmann,F., Durand,A., Basfeld,A., *et al.* (2017) Structure of Full-Length SMC and Rearrangements Required for Chromosome Organization. *Molecular Cell*, **67**, 334-347.e5.

FIGURE LEGENDS

Figure 1

(A) The initial structure of CGMD simulations of the Ycg1/dsDNA complex. dsDNA is shown in white and magenta. Ycg1 is shown in green. **(B)** DNA contact probabilities mapped on the Ycg1 structure. **(C)** A survival probability plot for the Ycg1/dsDNA binding. **(D)** Representative simulation (left) and the crystal (right; PDB ID: 5OQN) structures of the Ycg1/dsDNA complex. **(E)** Probabilities of each Ycg1 amino acid contacting with dsDNA. The orange and green lines represent the data from the simulations with the native and shuffled charge distribution. The vertical light blue lines represent the amino acid contacting dsDNA in the crystal structure.

Figure 2

(A) The cartoon and the representative AFM images of the condensin holo-complex with various hinge angles. **(B)** The time trajectory of hinge angle. The blue and green lines represent data with and without median filter, respectively. **(C)** The frequency distribution of hinge angle. The data is fit by a linear combination (dark blue) of two Gaussian distributions (light blue).

Figure 3

(A) The AFM image of the open hinge. One pixel corresponds to 1.875 nm. **(B)** The crystal hinge structure and the initial structure of the AFM fitting simulation. The initial structure was

obtained by a manual rigid body transformation of the crystal structure. **(C)** Inter-subunit Q-scores and intra-subunit Q-scores. The scores were calculated from simulations with the fitting strength parameter κ set to 10, 3000, and 10000. **(D)** The time trajectory of the cosine similarity (c.s.) between an experimental AFM image and the pseudo-AFM image calculated from the simulations. **(E)** The time trajectory of intra-subunit Q-score. **(F)** The open hinge structure obtained with the AFM fitting simulation overlaid with a grid representation of the reference AFM image. **(G)** The crystal structure (light blue) superimposed on the simulated open hinge structure.

Figure 4

(A, B) The representative structures from the CGMD simulations of the open (A) and closed hinge (B) with dsDNA. **(C, D)** dsDNA contact probabilities mapped on the open (C) and closed (D) hinge structures. **(E, F)** Surface charges mapped on the open (E) and closed (F) hinge structures. **(G, H)** The survival probability plots for the Ycg1/dsDNA, Ycs4/dsDNA, and hinge/dsDNA binding in 400 (G) and 200 (H) mM salt, respectively.

Figure 5

(A) The structures of the 0-th, 1.6×10^6 -th, and 1.0×10^7 -th steps of the CGMD simulation of hinge/dsDNA complex without the AFM fitting potential. **(B)** The time trajectories of inter-subunit Q-score and hinge angle from a CGMD simulation. The inset illustrates the hinge angle of a hinge structure. **(C)** The survival probability plot for hinge/dsDNA and Ycg1/dsDNA binding.

Figure 6

(A) The frequency distribution of conservation probability of each amino acid in the hinge domain (PDB ID: 4RSI). **(B)** The open hinge structure in which each particle is colored according to its conservation probability.

Figure S1

(A) The initial structure of CGMD simulations of the Ycs4/dsDNA complex. **(B)** The

representative structures of the CGMD simulation. **(C)** DNA contact probabilities mapped on the Ycs4 structure.

Figure S2

The representative AFM images of the condensin holo-complex. In these images, one of the HEAT repeat subunits pointed by the red arrow repeatedly dissociates from and associates to the head domains of Smc2 and Smc4.

Figure 1

Figure 2

Figure 3

Figure 4

Figure 5

Figure 6

Figure S1

Figure S2