

# Printing, Characterising, and Assessing Transparent 3D Printed Lenses for Optical Imaging

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## 1 Abstract

1 High-quality lens production has involved subtractive manufacturing methods for centuries. These methods  
2 demand specialist equipment and expertise that often render custom high-grade glass optics inaccessible. We  
3 aimed to develop a low-cost, accessible, and reproducible method to manufacture high-quality three-  
4 dimensional (3D) printed lenses using consumer-grade technology. Various planoconvex lenses were produced  
5 using a consumer-grade 3D printer and low-cost spin coating setup, and printed lenses were compared to  
6 commercial glass counterparts. A range of mechanical and optical methods are introduced to determine the  
7 surface quality and curvature of 3D printed lenses. Amongst others, high-resolution interference reflection  
8 microscopy methods were used to reconstruct the convex surface of printed lenses and quantify their radius of  
9 curvature. The optical throughput and performance of 3D printed lenses were assessed using optical  
10 transmissivity measurements and classical beam characterisation methods. We determined that 3D printed  
11 lenses had comparable curvature and performance to commercial glass lenses. Finally, we demonstrated the  
12 application of 3D printed lenses for brightfield transmission microscopy, resolving sub-cellular structures over  
13 a 2.3 mm field-of-view. The high reproducibility and comparable performance of 3D printed lenses present great  
14 opportunities for additive manufacturing of bespoke optics for low-cost rapid prototyping and improved  
15 accessibility to high-quality optics in low-resource settings.

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## 16 Introduction

17 Additive manufacturing, in particular 3D printing, has resulted in a range of innovative open hardware  
18 solutions for optical imaging applications. The variety of printing modalities, commitment to open sharing by  
19 users, low barrier-to-entry, and rapid evolution of 3D printing technologies has provided new opportunities to  
20 make imaging more accessible, particularly in low-resource settings [1–4]. Open microscopy initiatives have  
21 aimed to provide improved access to 3D printed microscopy hardware for use in the field or in rapid clinical  
22 diagnostics, such as the OpenFlexure project [5–8]. However, such initiatives routinely focus on 3D printing the  
23 mechanical parts of the microscope, such as the chassis, focusing assemblies, or specimen mounts, mainly using  
24 fused deposition modelling (FDM) printing to manufacture parts from heated plastic filaments. The optical  
25 elements in these applications still use glass objectives or plastic camera lenses. There is a need for accessible  
26 and robust methods to produce high-quality optical elements.

27 Optical imaging has relied on the use of ground glass lenses to manipulate light for centuries[9–12].  
28 Manufacturing such glass lenses is a subtractive process that can be prohibitively expensive for custom optics,  
29 requires specialist equipment and expertise, and produces a product that is delicate and easily damaged.  
30 Additive manufacturing, specifically 3D printing, has the potential to mitigate each of these barriers. Recently,  
31 additive manufacturing methods using injection moulding, magnetorheological, and molten glass printing have  
32 successfully produced bespoke optical elements [13–16]. In the latter, glass filaments are heated and extruded  
33 as in FDM printing, annealed, and cured in a kiln to produce the final product. However, this process is  
34 prohibitively costly owing to the specialist equipment and expertise required to manipulate glass in such ways.  
35 Resin-based printing techniques provide a viable solution to create bespoke optical elements for use in open  
36 hardware and imaging applications.

37 A variety of resin-based printing methods are available, mainly based on stereolithography (SLA) techniques.  
38 Two-photon polymerisation is routine for microfabrication of lenses, particularly for microlens arrays and x-ray  
39 imaging optics [17–19]. However, the requirement for specialist ultrashort pulsed laser sources and expensive  
40 printing instrumentation serves as a barrier to entry for these methods. Masked SLA (MSLA) printing is one of  
41 the most accessible SLA techniques [20]. This method uses a proprietary mix of methacrylate-based resin that is  
42 available in transparent and opaque forms and photoinitiated cross-linkers to polymerise under irradiation  
43 using 405 nm light. The illumination pattern is provided by an array of UV light emitting diodes, which can be  
44 collimated to provide homogeneous illumination over the entire build area. A liquid-crystal display (LCD)  
45 screen then masks the structure of discrete layers of the print design, changing in unison with the axial position  
46 of a buildplate mounted to a motorised stage to create a printed 3D structure [21,22]. Recent studies have  
47 produced MSLA 3D printed lenses for spectrophotometry applications [23] and described methods to quantify  
48 their material properties [24]; however, the optical performance of 3D printed lenses remains uncharacterised.  
49 Current advances in consumer-grade MSLA printer technology enable printing with a lateral resolution of up  
50 to 18  $\mu\text{m}$  and an axial resolution of up to 10  $\mu\text{m}$ .

51 We aimed to create a robust method to manufacture and characterise the optical performance of transparent 3D  
52 printed high-quality bulk optics using a consumer-grade printer and commercially available resin. We used a

53 low-cost MSLA 3D printer and employed a spin-coating method render the lenses transparent by minimising  
54 surface imperfections and reduce scattering and refraction from step structures originating from the printing  
55 process (the so-called 'staircase effect'). We printed a range of planoconvex lenses of various optical  
56 prescriptions, we characterised their surface profile in comparison to their commercial glass counterparts, and  
57 we assessed their optical performance.

## 58 Methods

### 59 3D printing

60 All 3D printing was conducted using a consumer-grade Mars 3 printer (Elegoo, China) fitted with a magnetic  
61 build plate attachment for easy print removal (Sovol, China). Clear resin (Clear Resin V4; Formlabs, USA) was  
62 used as the substrate for all prints. The print settings were optimised using the Cones of Calibration test print  
63 (TableFlip Foundry) and the print quality was verified using the Ameralabs Town test print (MB Labsamera,  
64 Lithuania) (Supplementary Figure 1). All lens prints were conducted using the parameters detailed in  
65 Supplementary Information, with the planar side of the lenses printed directly on the buildplate surface. Figure  
66 1 provides a graphical overview of the manufacturing process for 3D printed lenses.

67 Planoconvex lens designs (each with a lens diameter of 12.5 mm) were acquired from the manufacturer  
68 (Thorlabs, USA or Edmund Optics, USA) as .STEP files. Files were imported into Fusion 360 (v2.0.16985;  
69 Autodesk, USA) and the polygon count was increased to the maximum available before being exported as .STL  
70 files. Files were then imported into LycheeSlicer (v5.2.201; Mango3D, France), where the print parameters were  
71 applied, and the print files were exported as .CTB files. The .CTB lens files were printed with a raft surrounding  
72 the lens perimeter which provided a handling surface to avoid touching the lens surface. This increased the total  
73 diameter of the print to 25 mm, and ensured the print would fit in a standard 25 mm diameter optical mount  
74 (LMR1S; Thorlabs, USA). Four planoconvex lens prescriptions of different focal lengths ( $f$ ) were selected, based  
75 on commercial glass counterparts;  $f = +12.5$  mm (37-385; Edmund Optics, USA),  $f = +19.9$  mm (LA1074; ThorLabs,  
76 USA),  $f = +35.0$  mm (37-791; Edmund Optics, USA), and  $f = +49.8$  mm (LA1213; ThorLabs, USA).

77 Following printing, the lenses were removed from the magnetic build plate and washed with neat isopropanol  
78 (10592921; FisherScientific, USA) for 9 minutes in a Mercury X washing station (Elegoo, China). The prints were  
79 removed and carefully air-dried using a compressed Ultra Pure Duster (Thorlabs, USA) before a post-print UV  
80 curing step for 20 minutes in a Mercury X curing station (Elegoo, China).

### 81 Spin Coating and Lens Preparation

82 The printed lenses were rendered optically transparent by spin coating a thin layer of resin over both the curved  
83 and planar surfaces, minimising layer artefacts and surface structures from the printing process. For the convex  
84 surface, a spin coater (L2001A3-E463; Ossila, UK) was fitted with a custom 3D printed chuck with a 25 mm  
85 diameter well to accommodate a printed lens. The printed lens was cleaned again prior to coating with 100%  
86 isopropanol, airdried using compressed air, and placed in the chuck. 100  $\mu$ l of Clear UV Resin (4<sup>th</sup> Generation;  
87 VidaRosa, China) was deposited on the apex of the lenses (50  $\mu$ l for  $f = +49.8$  mm to accommodate the shallower  
88 curvature) and spun for 10 seconds at 2000 rpm. The coated lenses were stored in darkened conditions for 30

89 minutes to allow the liquid resin to settle and were subsequently cured for 20 minutes using the Mercury X  
90 curer as described above.

91 For the planar surface, the spin coater was fitted with a 76 mm x 26 mm microscope slide chuck and a clean  
92 microscope slide placed in the chuck. A thin resin layer was created by depositing 100  $\mu$ l of Clear Resin (v4;  
93 Formlabs, USA) on the slide and spinning for 10 seconds at 2000 rpm. The cleaned planar slide of the printed  
94 lens was carefully placed on to the resin-coated slide to avoid introducing air bubbles and was placed in a  
95 vacuum chamber (2 L,; Bacoeng, USA) fitted with a vacuum pump (Capex 8C; Charles Austen Pumps Ltd, UK).  
96 The assembly was maintained under a vacuum on 0.9 bars for 30 minutes before curing for 20 minutes as  
97 described above. The melded lens-slide combination was stored at -20°C for 3 minutes and carefully levered  
98 from the slide, relying on the differential thermal expansion of the glass and resin to remove the lens from the  
99 microscope slide.

## 100 Spin-coat Thickness Measurements

101 To determine the thickness of the spin coated layers, coumarin-30, a non-polar green-emitting organic  
102 fluorophore, was prepared as a 10 mM stock in 100% isopropanol and mixed with 1 ml of Clear UV Resin (4<sup>th</sup>  
103 Generation; VidaRosa, China) at a final concentration of 100  $\mu$ M. Lenses from each of the four test prescriptions  
104 were spin coated and cured as described above.

105 The thickness of the fluorescent spin coated layer was measured by acquiring a 3D image stack using a confocal  
106 laser scanning microscope. An Olympus IX81 inverted microscope coupled to an FV1000 confocal laser scanning  
107 unit (Olympus, Japan) was used for imaging. Excitation of fluorescence was performed using a 488 nm argon  
108 laser (GLG3135; Showa Optronics, Japan) and fluorescence emission from coumarin-30 was detected by a  
109 photomultiplier tube (PMT) with a detection spectral window of 500 nm to 550 nm. Coated lenses were placed  
110 with the curved surface in contact with a Type 1.5 coverglass and imaged using a 10 $\times$ /0.4 numerical aperture  
111 (NA) objective lens (Olympus, Japan). All images were acquired at the axial Nyquist sampling rate for the  
112 imaging objective ( $\Delta z = 1.53 \mu$ m). The thickness of the fluorescent layer was measured using a linear plot profile  
113 of the fluorescence intensity in orthogonal ( $x,z$  and  $y,z$ ) views of the 3D image stack using FIJI (v1.53t) [25].  
114 Analysis was conducted using three replicate printed lenses with each focal length.

## 115 White Light Interferometry and Stylus Profilometry

116 Methods for measuring the surface profile of planoconvex lenses include non-contact white light interferometry  
117 and contact stylus profilometry, with contact measurement approaches often only reporting the curvature of a  
118 linear trace instead of providing a 3D reconstruction of the specimen topology. Non-contact surface profiles  
119 were obtained using a white light interferometer (Wyko NT1100; Veeco Instruments Inc, USA) which used  
120 coherent light to generate interference fringes which are axially shifted through the optical surface, providing  
121 two-dimensional surface roughness and uniformity measurements. The interferometer was used in a vertical  
122 scanning interferometry configuration where an internal translator axially scanned in one direction during the  
123 measurement as the in-built camera detector recorded each frame. The non-contact approach provides an  
124 approximate 300 x 200  $\mu$ m surface area measurement using a 20x objective, with a colour gradient to indicate  
125 height ( $\Delta z$ ) changes as well as read-out line profiles to provide sub-nanometre-scale surface roughness across

126 the  $x$  and  $y$  axes of the specimen. The contact approach utilised a stylus profiler (Alpha-Step IQ; KLA Corp.,  
127 USA) with 5  $\mu\text{m}$  tip diameter for one-dimensional surface topology. The stylus profilometry technique provides  
128 millimetre-range one-dimensional measurements with sub-nanometre height resolution for curvature and  
129 roughness analysis.

## 130 Interference Reflection Microscopy (IRM)

131 The underpinning theory of IRM relating to imaging of plano-convex lenses has been explained elsewhere  
132 [26,27]. Briefly, the printed lens specimens were placed convex side down on a Type 1.5 coverglass, bridging  
133 the stage insert of an IX81 inverted microscope coupled to an FV1000 confocal laser scanning unit (Olympus,  
134 Japan). An 80/20 beamsplitter was used in place of the dichroic filter in the confocal microscope, which  
135 facilitated a configuration to detect reflected light from the specimen plane. A 458 nm argon laser (GLG3135;  
136 Showa Optronics, Japan) provided incident light, which was reflected from refractive index boundaries at the  
137 specimen plane (i.e., coverglass-air and air-lens interfaces). Reflected light from each interface coincided, leading  
138 to constructive and destructive interference depending on the optical path difference of the two reflected beams  
139 (Figure 2a). The resulting image provided a 2D projection of the 3D topography of the specimen surface, where  
140 interference orders were separated along the optical axis. Equations 1 and 2 describe the axial separation of  
141 destructive and constructive interference orders, respectively [28,29], where  $z$  = fringe spacing,  $N$  = order,  $\lambda$  =  
142 wavelength of incident light, and  $n_m$  = refractive index of the imaging medium.

$$143 \quad z = N \left( \frac{\lambda}{2n_m} \right) \quad \text{Eq. 1}$$

$$144 \quad z = \left( N + \frac{1}{2} \right) \left( \frac{\lambda}{2n_m} \right) \quad \text{Eq. 2}$$

145 IRM images were acquired using a 10x/0.4 NA objective lens (Olympus, Japan) and the reflection signal was  
146 detected using a PMT with the detection limited to 458 nm  $\pm$  5 nm.

## 147 Reconstructing 3D Surface Curvature and Quantifying Radius of Curvature from 2D

### 148 IRM Image Data

149 All computational analyses of IRM data were performed using FIJI and Python 3.8.10 (64-bit) in a Spyder IDE  
150 5.3.2 environment on an Elitebook 840 G7 (Hewlett-Packard, USA) running a 64-bit Windows 10 Enterprise  
151 operating system with an Intel® Core™ i5-10310U 1.70GHz quad-core processor with 16 GB of 2666 MT/s DDR4  
152 RAM.

153 The IRM image data were exported as .OIB files and pre-processed using FIJI. The images were cropped to  
154 ensure the apex of the lens was centred in the image, and a median filter ( $\sigma = 2$ ) was applied to remove any high-  
155 frequency noise in the data. Images were contrast adjusted using the Contrast Limited Adaptive Histogram  
156 Equalisation (CLAHE) plugin [30] (blocksize = 127, histogram bins = 256, maximum slope = 3.00) and converted  
157 to .PNG files for analysis.

158 A custom Python pipeline [31] was created to generate 3D reconstructions of the surface of the lens specimens  
159 from 2D IRM images and calculate the median radius of curvature for each lens. Briefly, the *Calibration* script  
160 was first used to verify the correct feature detection parameters for IRM data. The position of the zeroth order  
161 minimum was taken as the centre of the lens and the radius was noted as half the width of the image. The



162 position of the intensity maxima along the radius was calculated using the *find\_peaks* function, iteratively  
163 optimising the detection thresholds for peak height, distance, and prominence to ensure that each intensity  
164 maximum was detected. A line intensity profile noting the position of each interference maxima was generated  
165 along radius ( $r$ ) (Figure 2b). The axial position of each interference maximum was calculated using Equation 1  
166 and was plotted against distance to show the curvature of the lens (Figure 2c).

167 The surface curvature of each lens was reconstructed using the *3D Reconstruction* script with the optimised  
168 setting for each lens applied from the *Calibration* script. The 3D convex surface was reconstructed by assigning  
169 an axial position, as above, to each interference order detected along each radius (360 radii measured per image).  
170 The radius of curvature ( $R$ ) value for each radius was calculated using Equation 3 [32].

$$171 \quad R = \frac{r^2 + z_N^2}{2 \cdot z_N} \quad \text{Eq. 3}$$

172 The measured radius of curvature for each lens was calculated using the *Radius Analysis* script in Python. The  
173  $R$  value for each radius was compiled into a histogram that compared the experimental measurements to the  
174 theoretical  $R$  (i.e., the manufacturer-quoted value for the lens design file used for printed). The maximum peak  
175 position from each histogram determined the  $R$  value for each lens.

176 The measured and theoretical curvature data were plotted using Prism (v8.0.2) (GraphPad Software, USA).

## 177 Tolansky Interferometry

178 Tolansky interferometry, a mode of multiple-beam interference imaging, was used as an alternative means to  
179 visualise the curved topology at the apex of the convex surface of the 3D printed lenses. The principle of  
180 Tolansky interferometry relies on two opposing highly reflective surfaces (in this case, an aluminium-coated  
181 coverglass and aluminium-coated 3D printed lens) which generate multiple reflected beams that undergo  
182 constructive and destructive interference [33–35]. The multiple beam combination acts to modify the Haidinger  
183 rings formed at the focus of the objective lens such that the nodal spacing is not altered, but the intensity  
184 distribution of the interference orders is changed. This effectively increases the axial resolution compared with  
185 IRM, such that topological features as small as 5 Å, or better, can be resolved within the interference maxima  
186 [28]. Moreover, uncoupling the reflective specimen from the reflective coverglass provides a means to translate  
187 the interference orders through the optical axis, in turn ‘scanning’ the topography of the convex lens surface in  
188 a way that IRM cannot.

189 A 3D printed lens ( $f = +19.9$  mm) and a Type 1.5 coverglass were vapour coated with a thin layer of aluminium  
190 using a thermal evaporator coating system (E306A; Edwards Vacuum, UK). Briefly, a small quantity of  
191 aluminium foil was heated on a tungsten filament under vacuum ( $1.0 \times 10^{-5}$  Pa), depositing vapourised  
192 aluminium on the surface of the lens and coverglass.

193 A custom steel objective lens collar was fabricated and fitted to a  $10 \times / 0.4$  NA objective lens (Olympus, Japan).  
194 The aluminium-coated coverglass was bonded to the top of the collar using a thin layer of epoxy resin around  
195 the circumference of the objective lens housing. An adjustment screw was included in the collar to facilitate the  
196 positioning of the collar-coverglass assembly relative to the focal length of the objective lens (approximately 3.1  
197 mm). The aluminium-coated lens specimen was suspended over the stage insert of an inverted IX81 microscope  
198 coupled to a confocal laser scanning unit (Olympus, Japan). Two glass microscope slide spacers were inserted

199 to raise the test lens so that the modified objective underneath could be focussed near to the lens surface.  
200 Tolansky interferometry was performed using the same reflection setup as in IRM experiments but employed  
201 z-scanning which uncoupled the specimen from the coverslip. Altering the relative distance between the  
202 coverglass and lens specimen resulted in axial translation of interference orders and provided a means of  
203 visualising the curved apical surface by merging images acquired at different axial positions into a single  
204 hyperstack colour-coded by depth using FIJI.

## 205 Optical Transmission Measurements

206 The percentage transmission of the printing resin was measured by comparing the mean intensity of transmitted  
207 light through resin blocks of varying thickness compared to the transmission through a single Type 1.5  
208 coverglass. Resin blocks of thickness from 1 mm to 6 mm were printed using the optimised printing parameters  
209 used for lens printing. The blocks were placed on a Type 1.5 coverglass and imaged using an IX81 inverted  
210 microscope coupled to an FV1000 confocal laser scanning unit (Olympus, Japan) configured to detect scanned  
211 transmitted light. The transmissivity of unprocessed and processed blocks (i.e., naïve and spin coated,  
212 respectively) was measured using three discrete wavelengths across the visible spectrum sequentially; a 458 nm  
213 argon laser, a 515 nm argon laser, and a 633 nm helium-neon laser (GLG3135; Showa Optronics, Japan). Images  
214 were acquired using a 4×/0.1 NA objective lens (Olympus, Japan), with dimensions 64 × 64 pixels and Kalman  
215 averaging (n = 5 frames) to minimise contributions to the image from print structures. The mean intensity across  
216 the field was measured using FIJI and compared to the optical throughput of the coverglass alone. Linear fits  
217 were conducted using Prism.

## 218 Beam Profilometry

219 An optical setup was constructed to measure the focusing ability of 3D printed lenses compared to their  
220 commercial glass counterpart. A complete parts list is included in the Supplementary Material. A 633 nm  
221 helium-neon laser source with an initial beam diameter measuring 600 μm was passed through a neutral density  
222 filter and was steered using two gimbal-mounted mirrors. The beam was expanded to a final diameter of 12.5  
223 mm using two sequential beam expanders, first through a 2.5× beam expander and then through a 7.5× beam  
224 expander. The 3D printed lens was mounted in a fixed mount (LMR1S/M; Thorlabs, USA) and a dual scanning  
225 slit beam profiler was mounted on a linear translational stage to facilitate movement along the optical axis to  
226 map the beam diameter with respect to post-lens propagation distance. Perpendicular measurements (*x* and *y*  
227 axes) of the focused beam diameter ( $1/e_2$ ) were measured at increments along the optical axis and compiled to  
228 provide a beam profile for three replicates of various planoconvex lens prescriptions. The  $1/e_2$  beam waist  
229 radius ( $w_0$ ) was calculated from the measured beam diameter and the Rayleigh Range ( $z_R$ ) for each lens was  
230 calculated using Equation 4 [36].

$$231 \quad z_R = \pi \cdot w_0^2 \cdot n / \lambda \quad \text{Eq. 4}$$

## 232 Brightfield Transmission Microscopy

233 A brightfield transmission microscopy setup was constructed to demonstrate the performance of a 3D printed  
234 lens in an imaging setup. A complete parts list is included in the Supplementary Material. A 10 mm stage

235 micrometer with 50  $\mu\text{m}$  intervals (R1L3S1P; ThorLabs, USA) and a thin section of linden tree stem (*Tilia europaea*)  
236 were imaged to assess field of view and to determine the resolution of the system. A blue light-emitting diode  
237 (LED) source ( $\lambda = 470 \text{ nm}$ ) (M470L2-C1; ThorLabs, USA) was used to illuminate the specimen. Light from the  
238 LED was brought to the specimen using a 3D printed  $f = +49.8 \text{ mm}$  planoconvex condenser lens (modelled on  
239 LA1213; ThorLabs, USA) and transmitted light was detected using a monochrome complementary metal-oxide-  
240 semiconductor (CMOS) camera (DCC3260M; ThorLabs, USA). Image acquisition was controlled via ThorCam  
241 (64-bit, v3.7.0) (ThorLabs, USA).

## 242 Results

### 243 Surface Characterisation of 3D Printed Lenses

#### 244 Measuring the Thickness of Spin Coated Resin Layers on 3D Printed Lenses

245 The thickness of the spin-coated resin layer on the printed convex lens surface was measured using 3D confocal  
246 laser scanning microscopy. The coumarin-30-spiked VidaRosa Clear resin was spin coated onto the convex  
247 surface using the same spin settings as for other lenses, creating a fluorescent resin layer (Figure 3a). A 3D  
248 confocal z-stack visualised the spin-coated layer relative to the lens surface (not fluorescent, ergo dark) (Figure  
249 3b). The mean spin coated thickness of the fluorescent resin was measured for three replicates of various lens  
250 prescriptions, with the spin coat thickness routinely ranging from 25  $\mu\text{m}$  to 45  $\mu\text{m}$  (mean thicknesses;  $f_{12.5 \text{ mm}} =$   
251 30.00  $\mu\text{m}$ ,  $f_{19.9 \text{ mm}} = 41.31 \mu\text{m}$ ,  $f_{35.0 \text{ mm}} = 42.00 \mu\text{m}$ ,  $f_{49.8 \text{ mm}} = 28.50 \mu\text{m}$ ) (Figure 3c).

#### 252 Comparing the Surface Curvature and Uniformity between 3D Printed and 253 Commercial Glass Lenses

254 The surface curvature of 3D printed lenses was first measured by conventional means using a commercial white  
255 light interferometer and a stylus profilometer. However, white light interferometry usually requires higher  
256 reflective surfaces for accurate surface measurements, and stylus profilometry is typically restricted by  
257 measuring only orthogonal straight lines along the  $x$  and  $y$  axes of the lens (Supplementary Figure 2). An  
258 alternative method was required that accurately reconstructed the transparent three-dimensional surface of the  
259 printed lenses, which provided a robust method to identify surface curvature defects that could impact optical  
260 performance.

261 Using the methods outlined in Figure 2, 2D IRM image data of printed planoconvex lens surfaces (Figure 4a)  
262 were processed to create 3D renders of the curved surface (Figure 4b). The radius of curvature was measured  
263 for each radius around the circumference of the lens and plotted as a histogram to calculate the median radius  
264 of curvature for each printed lens (Figure 4c). The radii of curvature for three replicate printed lenses of four  
265 prescriptions were compared to commercial glass lenses (Figure 4d, Supplementary Figure 3). The radius of  
266 curvature of 3D printed lenses concurred with their glass counterparts, with a slight increase due to the additive  
267 spin coating process. However, this did not hold true for longer focal length lenses, where the increased radius  
268 of curvature is more pronounced for longer focal length lenses with larger variation due to the shallower  
269 curvature. The mean radius of curvature ( $\pm$  standard deviation) measured;  $R(f_{+12.5 \text{ mm}}) = 10.76 \text{ mm} \pm 1.04 \text{ mm}$ ,  $R$   
270 ( $f_{+19.9 \text{ mm}}) = 11.31 \text{ mm} \pm 0.66 \text{ mm}$ ,  $R(f_{+35.0 \text{ mm}}) = 18.83 \text{ mm} \pm 0.40 \text{ mm}$ ,  $R(f_{+49.8 \text{ mm}}) = 31.44 \text{ mm} \pm 4.01 \text{ mm}$ . These



271 values compare to the theoretical radii of curvature for their glass counterparts;  $R(f_{+12.5\text{ mm}}) = 9.80\text{ mm}$ ,  $R(f_{+19.9}$   
272  $\text{mm}) = 10.30\text{ mm}$ ,  $R(f_{+35.0\text{ mm}}) = 18.10\text{ mm}$ ,  $R(f_{+49.8\text{ mm}}) = 25.80\text{ mm}$ .  
273 The IRM images of 3D printed lenses routinely featured a larger than expected zeroth order minimum. This  
274 suggested that the apex of the convex lenses was flat, but this was not observed in white light interferometry or  
275 stylus profilometry experiments (Supplementary Figure 2). Moreover, the curvature of the zeroth order area  
276 measured using these methods agreed with the theoretical curvature and the median radius of curvature  
277 measured by IRM. To conclude that the zeroth order area was curved, a modified z-scanning Tolansky  
278 interferometry method was employed (Figure 5a). Uncoupling of the lens specimen from the objective setup  
279 permitted translation of the interference orders over the lens surface in unison with the axial movement of the  
280 modified objective housing (Figure 5b). A maximum intensity projection from a Tolansky interferometry z-  
281 series revealed the nanoscale surface profile and the continuous curved surface of the printed lens (Figure 5c),  
282 confirming the observations using white light and stylus profilometry.

## 283 **Optical Characterisation of 3D Printed Lenses**

### 284 **Comparing the Transmissivity of 3D Printed Resin to Commercial Glass Lenses**

285 The transmissivity of the 3D print resin substrate was measured to verify the optical properties of the clear resin.  
286 The mean intensity of the transmitted light was normalised compared to a control of the same intensity of light  
287 passing through a Type 1.5 coverglass (Figure 6). Block transmissivity was increased by approximately  $2.25\times$   
288 up to greater than 90% across all tested wavelengths following spin-coating, comparable to uncoated N-BK7  
289 glass often used to manufacture glass bulk optics [37].

### 290 **Measuring the Optical Performance of 3D Printed Lenses**

291 The optical performance of four 3D printed planoconvex lens prescriptions was measured in triplicate using the  
292 setup shown in Figure 7a. The focal length approximately matched the theoretical values for all lenses, except  
293 for  $f = 49.8\text{ mm}$  (Figure 7). The measured focal lengths for each lens prescription were  $f_{12.5\text{ mm}} = 13.5\text{ mm}$ ,  $f_{19.9\text{ mm}} =$   
294  $19.0\text{ mm}$ ,  $f_{35.0\text{ mm}} = 35.0\text{ mm}$ . These data presented an error of  $f_{12.5\text{ mm}} = 8.0\%$ ,  $f_{19.9\text{ mm}} = -4.5\%$ , and  $f_{35.0\text{ mm}} = 0.0\%$   
295 compared to the focal lengths of their glass counterparts. The longer theoretical focal length lenses (i.e.,  $f_{49.8\text{ mm}}$ )  
296 did not focus the light as expected. The beam profile was elongated along the optical axis at the beam waist and  
297 was significantly displaced from the theoretical focal length, indicating the presence of spherical aberration.  
298 Each of the printed lenses were compared to the focusing performance of their commercial glass counterparts  
299 using the same profiling setup. The glass commercial lenses performed as expected, resulting in the correct focal  
300 length for each prescription.

301 The beam waist ( $w_0$ ) and Rayleigh Range were measured for each 3D printed lens and were compared to the  
302 commercial glass counterpart, except for the  $f_{49.8\text{ mm}}$  lenses due to the severity of the optical aberrations (Table 1).  
303 The mean  $w_0$  and the Rayleigh Range values for the lens prescriptions are noted in Table 1. Although the focal  
304 length of these 3D printed lenses conformed with their commercial glass counterparts, the mean  $w_0$  value was  
305 routinely larger with 3D printed lenses. Moreover, the Rayleigh range was also increased proportionally to the

306 enlarged  $w_0$ . However, 3D printed lenses were ultimately able to focus a beam to a discrete point which  
 307 demonstrated promise for implementation in optical systems.

Lens Prescription	$f_{12.5\text{ mm}}$		$f_{19.9\text{ mm}}$		$f_{35\text{ mm}}$	
	3D Printed	Commercial Glass	3D Printed	Commercial Glass	3D Printed	Commercial Glass
Mean beam waist radius ( $w_0$ ) ( $\mu\text{m}$ ) ( $\pm$ SD)	62.40 $\pm$ 21.55	16.25	25.26 $\pm$ 5.06	17.46	56.67 $\pm$ 24.10	6.39
Rayleigh range ( $z_R$ ) (mm)	19.32	1.31	3.17	1.51	15.93	0.20

308 **Table 1. Comparison of Beam Parameters in 3D printed and Glass Lenses.** The mean beam waist radius ( $w_0$ )  
 309 was measured for three different printed and glass planoconvex lens prescriptions, along with the standard  
 310 deviation (SD) for three replicate printed lenses. The Rayleigh range ( $z_R$ ) was calculated for each prescription  
 311 based on the mean  $w_0$  value.

## 312 Using 3D Printed Optics for Brightfield Transmission Imaging

313 The imaging performance of 3D printed optics was tested using the setup presented in Supplementary Figure  
 314 5, where a 3D printed lens ( $f = +49.8\text{ mm}$ ) was used as the condenser in a brightfield transmission microscope.  
 315 The setup resulted in a field of view measuring approximately 2.3 mm wide with high contrast across the full  
 316 field (Figure 8a). Moreover, brightfield transmission imaging was demonstrated using a cross section of a linden  
 317 tree stem, resolving the intricate differentiated tissue layers and structures on the order of 6  $\mu\text{m}$  (Figure 8b). This  
 318 sub-cellular resolution demonstrates the potential of 3D printed optics in biological imaging systems.

## 319 Discussion

320 We have demonstrated robust, repeatable, and accessible methods to manufacture planoconvex lenses using 3D  
 321 printing with consumer-grade instrumentation and printing materials. Moreover, we have characterised and  
 322 compared the quality of 3D printed lenses against commercially available glass counterparts of the same  
 323 prescription. A spin-coating method was employed to obviate the stepped print structure and to render the  
 324 printed lenses smooth and transparent, resulting in a thin surface coating on the order of 35  $\mu\text{m}$  thick. White  
 325 light and stylus profilometry were used to assess the surface quality of printed lenses, while confocal IRM was  
 326 used to reconstruct the surface topology of 3D printed lenses with high axial resolution and quantify the radius  
 327 of curvature, which matched that of glass lens counterparts. It is important to note that, while IRM in confocal  
 328 scanning mode provides increased contrast due to coherent illumination and the rejection of out-of-focus light  
 329 by a pinhole aperture, widefield IRM also presents feasible and accessible method to achieve similar, albeit  
 330 lower-contrast, IRM data. Tolansky multiple beam interferometry corroborated the profilometry data by  
 331 revealing the curved topology of the apex of the convex lenses. The printing and processing steps facilitated  
 332 comparable optical throughput to glass lenses, with greater than 90% transmission across the visible spectrum.

333 The optical performance of 3D printed lenses was determined by measuring their ability to focus a beam of  
334 light. Shorter focal length printed lenses were observed to have the same focal length as their glass counterparts,  
335 with moderately increased beam diameters and Rayleigh ranges. The performance of 3D printed lenses was  
336 limited in longer focal length printed lenses (i.e.,  $f = + 49.8$  mm), which exhibited extended focal lengths  
337 indicative of spherical aberration. However, our findings present 3D printing as a viable option for the  
338 manufacturing of high-quality lenses for optical instrumentation and rapid prototyping as a less costly and  
339 more accessible alternative to bulk glass optics. Finally, we demonstrated the use of 3D printed lenses for  
340 brightfield transmission imaging. Despite the aberrations observed during beam profilometry experiments  
341 using printed  $f = + 49.8$  mm lenses, the imaging results were promising. Sub-cellular spatial resolution was  
342 achieved with high contrast over a 2.3 mm field of view, showing great promise for the use of 3D printed optics  
343 in microscopy.

344 Previous studies developing 3D printed optics have used a variety of additive manufacturing methods, these  
345 with most of these typically being costly, requiring specialist equipment, or being mainly focussed on using  
346 consumer-grade printers to produce micro-optics. Fused deposition modelling, where thin glass filaments are  
347 melted, extruded and cooled into the shape of a lens, has resulted in 3D printed glass bulk optics [38]. However,  
348 the silicon dioxide substrate requires careful mixing with titanium dioxide and a complex series of drying,  
349 burnout and sintering steps performed at over 1000 °C that limit users without access to specialist equipment.  
350 Alternative filament-based methods have used CO<sub>2</sub> lasers to print transparent glass lenses from a single mode  
351 optical fibre and fused quartz filaments [39–41]. However, these molten glass methods often result in layering  
352 defects that reduce optical performance.

353 Stereolithography approaches have resulted in various techniques to manufacture resin-based optics. These  
354 methods routinely use a two-photon polymerisation-based platform to manufacture microlens arrays (MLAs)  
355 and optics, although print sizes are typically limited to only a few mm in diameter. Moreover, two-photon-  
356 based instrumentation can be prohibitively costly and presents a barrier to entry for accessible 3D printed optics.  
357 Printed lenses using, for example, a Nanoscribe printer have recently been produced [42] but they retain  
358 microstructures and layering artefacts resulting from the printing process and curvature defects that ultimately  
359 impact on their optical performance. Moreover, these techniques are often limited by their small print sizes and  
360 high materials and instrument costs, somewhat restricting their use to MLAs and other micro-optics. Recent  
361 developments in two-photon microprinting have successfully been used to fabricate bespoke micro-optics, such  
362 as 30 µm-diameter Fresnel elements for x-ray microscopy, using photopolymerising resins on a supportive  
363 silicon nitride membrane [43]. Foveated compound microlenses have also recently been produced using  
364 femtosecond direct laser writing [18], but all two-photon methods require a costly and complex ultrashort  
365 pulsed laser source. Microlens arrays have been printed using UV-induced photopolymerisation, with recent  
366 improvements seeing expansion of MLAs over large flexible substrates to improve optical performance [44] and  
367 the introduction of vibrating projection lenses during printing to smooth the surface of 3D printed micro-optics  
368 [45]. Aspheric lenses have been manufactured using UV photopolymerisation, however these specialised lenses  
369 required assembly with corrective quartz substrates and refractive index matching liquids in order to focus

370 light, but suffered from both chromatic and spherical aberration [46]. Each of these 3D printing methods are  
371 based on inaccessible and specialist equipment, which often produce 3D printed lenses that do not compare to  
372 the performance of their commercially produced glass counterparts.

373 We have described an accessible, low-cost, and reproducible method for manufacturing bespoke 3D printed  
374 lenses and a suite of characterisation methods that demonstrate their comparable performance to glass lenses.  
375 We implemented profilometry methods centred on IRM that provide high-resolution topographical  
376 reconstructions of the lens geometry, which offer alternative analysis methods for IRM and standing wave  
377 microscopy [26,47]. Using 458 nm illumination for IRM and 633 nm light for beam profiling ensured that the  
378 material properties of the printed lenses did not change during observation, as the resin absorbance peak is  
379 noted as 405 nm. The Tolansky interferometry mapping, together with white light and stylus profilometry data,  
380 suggest that the enlarged zeroth order present in IRM images was an inherent interference artefact, perhaps  
381 contributed by subtle refractive index differences between the printed layers caused by compounding exposure  
382 to light during printing [48,49]. However, despite this apparent interference artefact across all lenses, it did not  
383 impact on their optical performance. Short focal length lenses with inherently higher curvatures performed in  
384 line with their commercial glass counterparts, but longer focal length lenses on the order of  $f = 50$  mm were  
385 subject to optical aberrations. This suggests that future applications creating 3D printed multi-lens systems  
386 would better suit the inclusion of shorter focal length elements. Overall, the beam waist radius and Rayleigh  
387 range were increased compared to glass lenses, however this did not impact the ability of 3D printed lenses to  
388 effectively focus light.

389 The potential for additive manufacturing for bespoke optics, rapid design prototyping, and field diagnostics is  
390 huge. Our data show that high-quality optical elements can be produced at low cost with consumer grade  
391 equipment, totalling approximately £300. The only additional outlet would be a spin coater, which can be  
392 procured for less than £1,000. The total cost in producing a single 3D printed lens was approximately £0.11, as  
393 opposed to upwards of £50 for commercial high-grade lenses. The demonstrated optical performance of 3D  
394 printed lenses shows great promise for optical imaging and prototyping optical instrumentation. Moreover, we  
395 have shown separately that 3D printed lenses can be implemented in bioimaging applications, using both  
396 absorption and fluorescence imaging modalities [50]. The potential impact of these accessible and open  
397 manufacturing methods could also impact across low resource settings for rapid diagnostics of blood smears,  
398 for example, where 3D printing has already made significant impacts. The combination of previous 3D printed  
399 microscope chassis with 3D printed lenses would be a natural evolution to produce the first fully 3D printed  
400 optical microscope.

## 401 Conclusions

402 We present an accessible 3D printing method to manufacture high-quality optical lenses and provide  
403 characterisation methods to quantify their performance. With the prohibitive cost of bespoke bulk glass optics  
404 and difficulties in their manufacture, 3D printing offers a viable method to produce a range of lenses with a  
405 high degree of reproducibility. The quality of 3D printed lenses was determined by comparing their surface  
406 curvature, optical throughput, and ability to focus light compared to their commercial glass counterparts. Glass

407 and 3D printed lenses were observed to behave similarly for a range of short focal length prescriptions, but  
408 longer focal lengths introduced a high degree of spherical aberration. The same trend was true for the surface  
409 curvature, where highly curved lenses conformed to the radius of curvature of their glass counterparts, while  
410 the surface coat thickness was routinely on the order of 35  $\mu\text{m}$ . The transmissivity of 3D printed lenses was  
411 comparable to that of bulk N-BK7 glass across the visible spectrum. Moreover, 3D printed lenses were  
412 implemented in a brightfield transmission microscopy setup and facilitated high-quality imaging that  
413 demonstrated promise for future applications. Each of these observations concluded that 3D printing is a viable  
414 approach to reproducibly producing large volumes of high-quality optical elements that provide promise for  
415 prototyping, imaging applications, and field diagnostics.

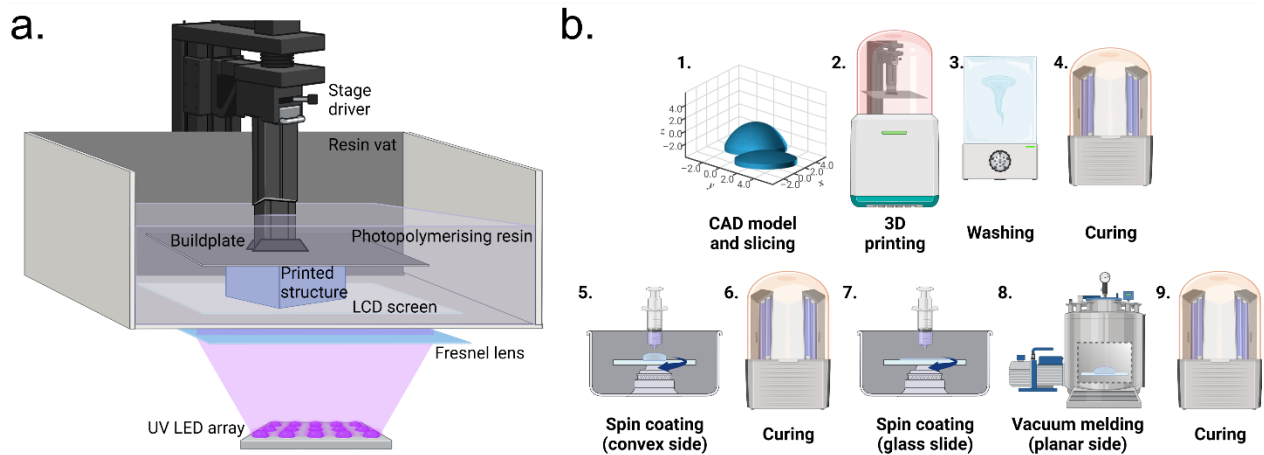
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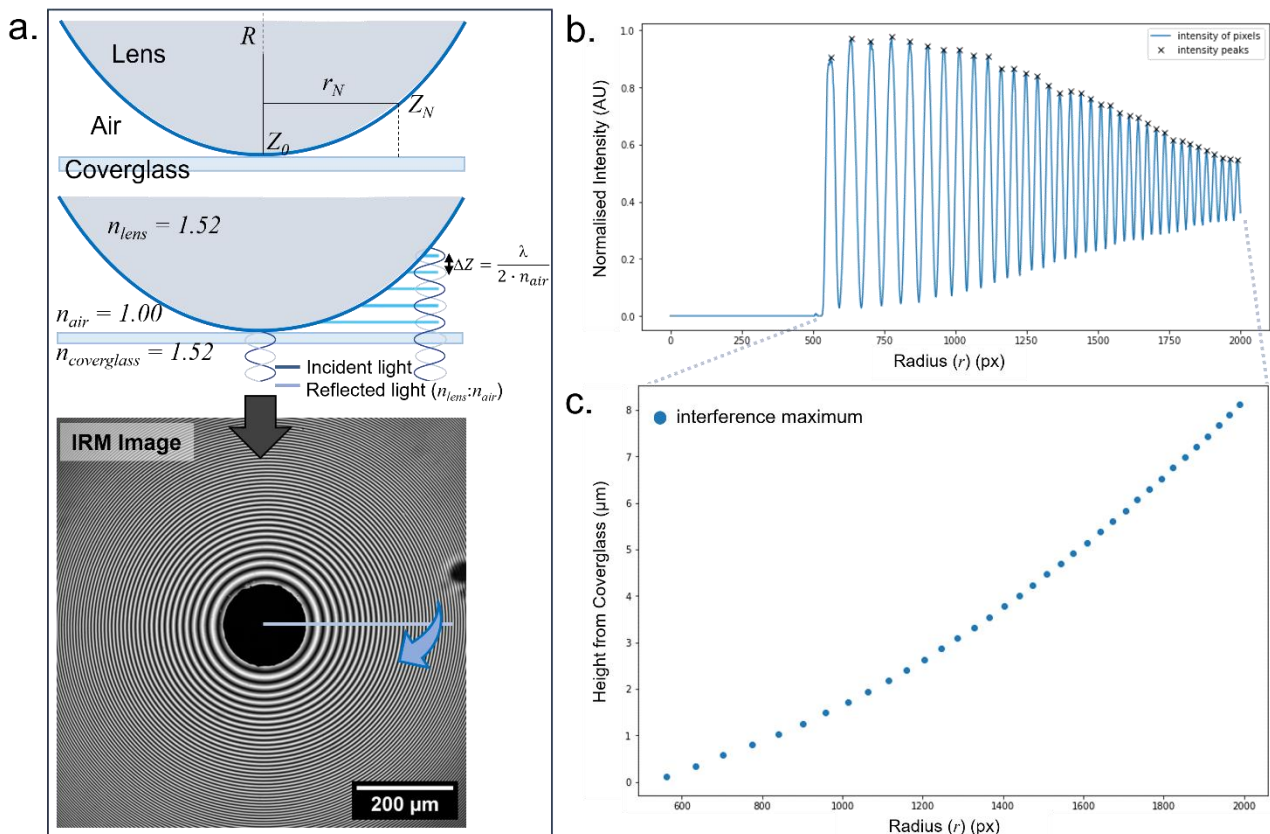


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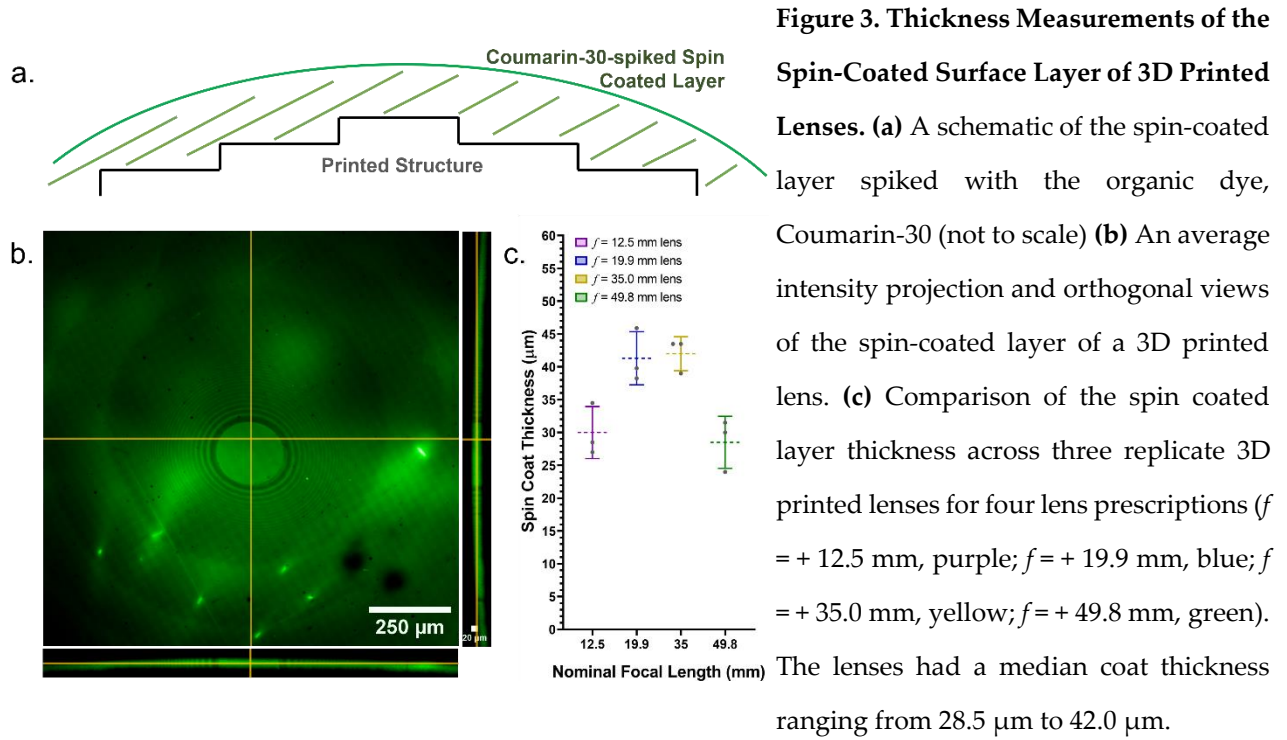
## Figures



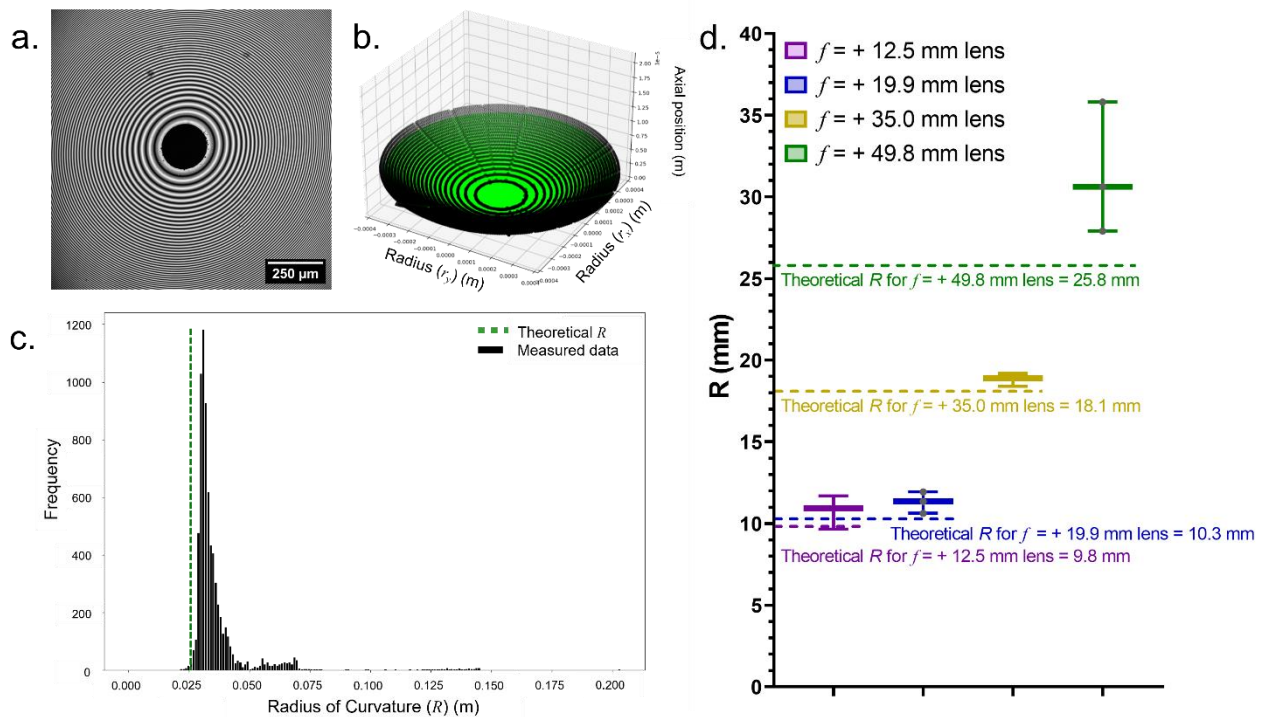
**Figure 1. Methods for Manufacturing 3D Printed Lenses.** (a) A schematic of an MSLA printer. A collimated UV light source is projected onto an LCD screen which illuminates and masks individual print layers as the buildplate is lifted from a vat of photopolymerising resin. (b) A schematic of manufacturing process for 3D printed lenses. The initial design and printing process is conducted (1-4) before a series of spin coating and curing steps to provide a smooth lens surface and increased optical quality (5-9).



**Figure 2. Reconstruction of Lens Curvature using Interference Reflection Microscopy (IRM).** (a) A schematic of the principles of IRM (not to scale). Refractive index boundaries reflect incident light and constructive/destructive interference occurs depending on the relative distance between reflective boundaries (b) An IRM image of a lens was acquired, and a line intensity profile was measured. (c) The axial position of the interference maxima (Eq. 1) was plotted against the radial peak position from the line profile.

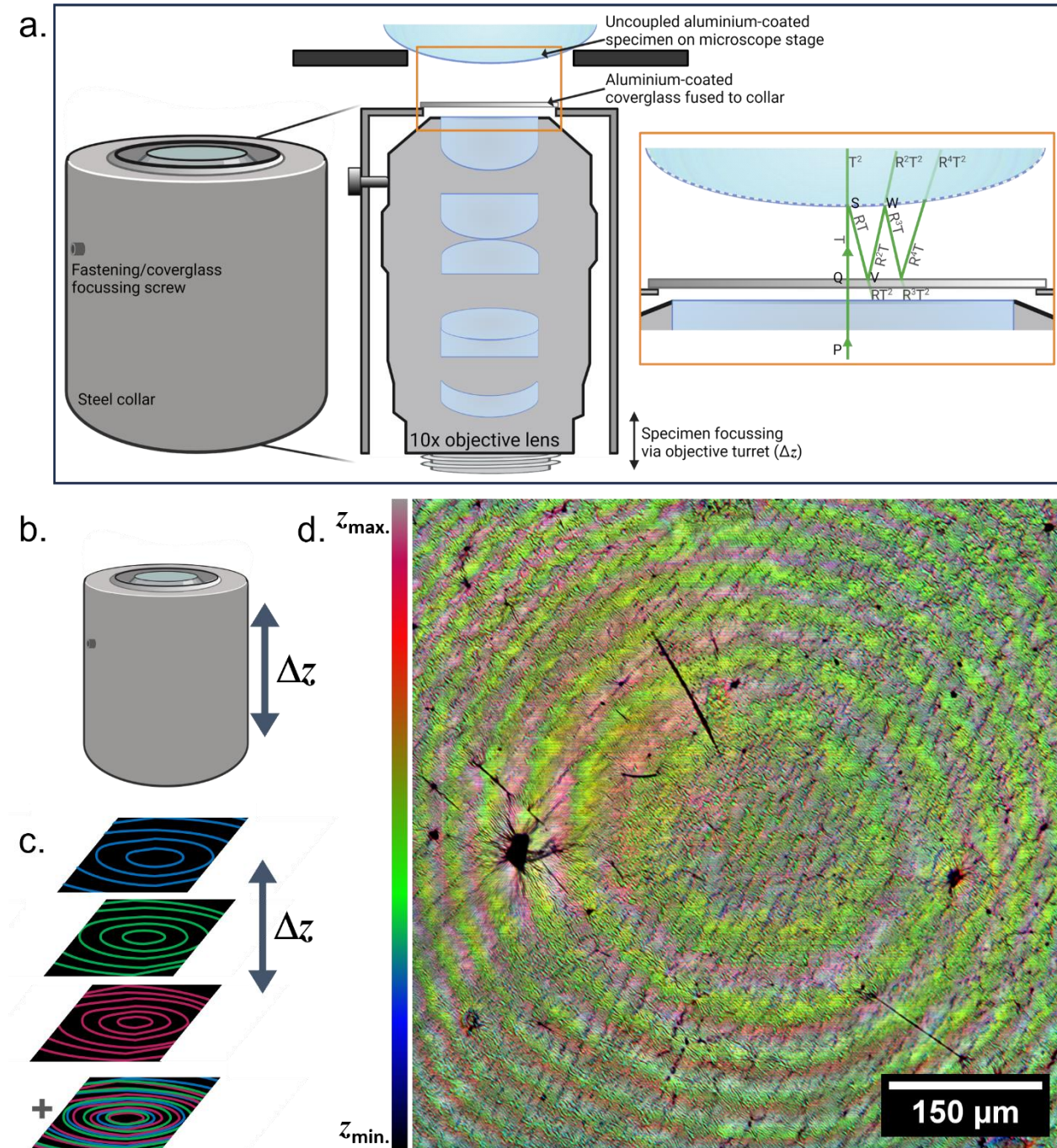


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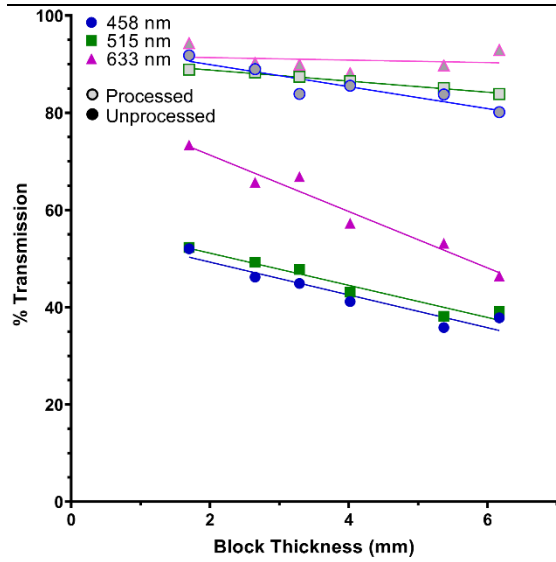


**Figure 4. Radius of Curvature Measurements of 3D Printed Lenses using IRM.** (a-b) Surface reconstruction of a 3D printed lens from 2D IRM data using method presented in Figure 2. (a) An IRM image of a 3D printed lens surface ( $\lambda = 458$  nm). (b) A 3D reconstruction of the printed lens surface (black datapoints) compared to the theoretical curvature from the print design (green). (c) A histogram of the distribution of radius of curvature ( $R$ ) values measured around the circumference of the lens (black) compared to the theoretical  $R$  (green) (See Eq. 2). (d) Measured radius of curvature values of three replicate 3D printed lenses for lens prescriptions ( $f = + 12.5$  mm, purple;  $f = + 19.9$  mm, blue;  $f = + 35.0$  mm, yellow;  $f = + 49.8$  mm, green).



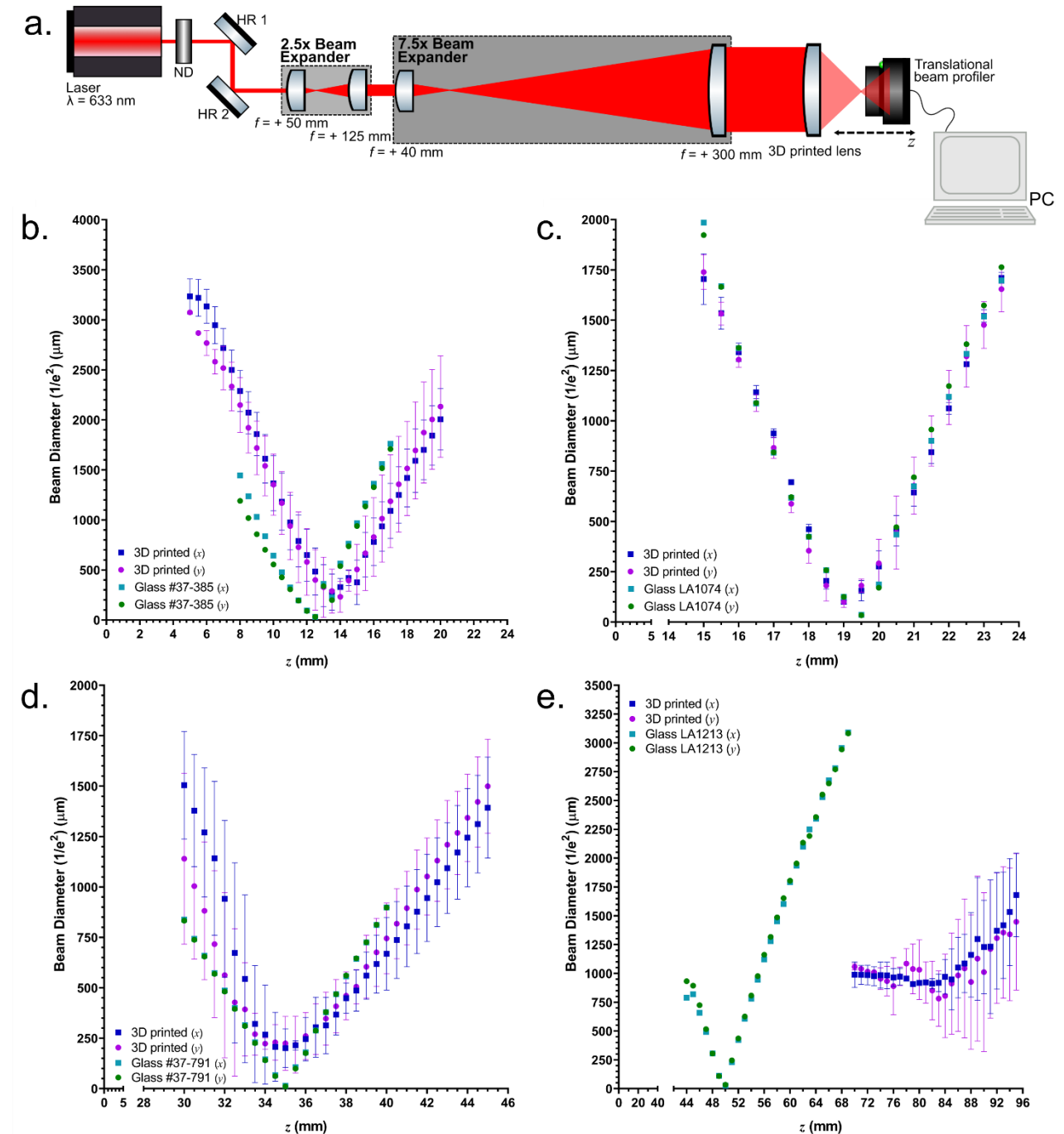


**Figure 5. Tolansky Interferometry Confirms the Apical Surface Curvature of 3D Printed Lenses.** (a) A schematic (not to scale) describing the principles of Tolansky interferometry and the design of a custom objective lens to permit 3D measurements of surface curvature of an aluminium-coated specimen. (b) Axial translation of the objective mount results in (c) translation of the interference orders, which can be colour-coded by depth and merged into a z-projection of a Tolansky interferometry acquisition. (d) The concentric interference maxima from each axial position are false coloured according to their depth and super-imposed, revealing the curved surface and nanoscale topology of the apical lens surface with higher resolution than interference optical microscopy can provide.

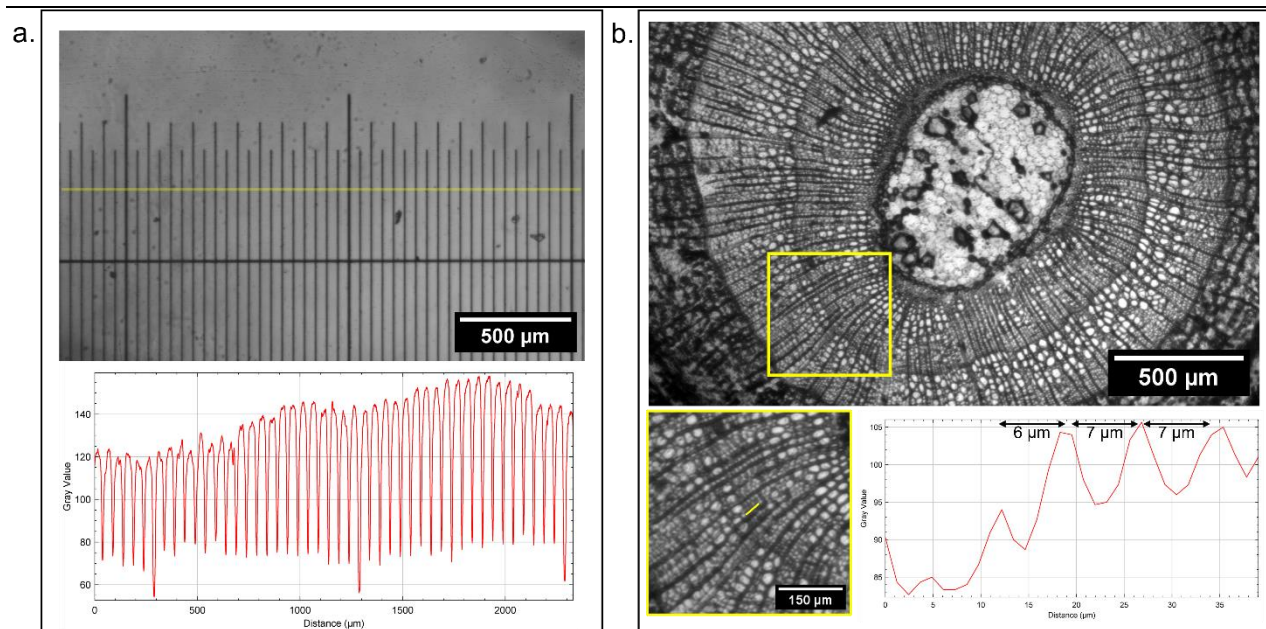


**Figure 6. Optical Throughput of Naïve and Post-processed Resin Blocks.** Naïve (i.e., unprocessed) blocks (solid markers) exhibited lower optical throughput than processed blocks (hollow markers). Optical throughput was typically higher with longer wavelengths of transmitted light, with processed blocks achieving up to 94.33% transmission at a wavelength of 633 nm.





**Figure 7. Comparing the Focusing Performance of 3D Printed Lenses.** (a) An optical setup to measure the beam profile of a 633 nm laser focused by a 3D printed lens. Two concurrent beam expanders provide a total beam expansion of 18.75 $\times$ , creating a 12.5 mm diameter beam that uses the full numerical aperture of the printed lenses. A beam profiler is translated along the optical axis to measure the focal length and metrics of the focused beam. ND = Neutral Density Filter, HR = High Reflector (Mirror), all lenses shown are commercial glass planoconvex lenses, save for the final 3D printed lens under observation. (b-e) The beam diameter along the optical axis of several 3D-printed lens prescriptions is presented. Both the  $x$  (blue) and  $y$  (purple) beam diameters are noted as a function of  $1/e^2$ . The beam profiles for lenses of theoretical focal length (a) +12.5 mm, (b) +19.9 mm, (c) +35.0 mm, and (d) +49.8 mm are presented. The focal lengths of (a), (b), and (c) concurred with their theoretical glass counterparts (presented in green for illustration in (b) and (d)) but did not agree for longer focal length lenses in (d).



**Figure 8. Brightfield Transmission Imaging Using a 3D Printed Condenser Lens.** Images demonstrating the application of a 3D printed condenser lens ( $f = + 49.8$  mm) **(a)** An image of a stage micrometer (R1L3S1P, Thorlabs, USA) measuring a field of view of approximately 2.3 mm. A line intensity profile (averaged over 5 pixels thickness) is plotted against distance, showing high contrast across the field of view. **(b)** An image of a linden tree (*Tilia europaea*) stem showing the intricate tissue layers. A magnified region of interest is presented with a yellow box, and this is digitally magnified to show individual plant cells. A line intensity profile (averaged over 5 pixels thickness) is plotted against distance and demonstrates resolution at a level that would be sufficient to record subcellular detail.

## Additional Information

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### 544 **Data and Code Availability**

545 The analysis code used to reconstruct and quantify IRM data is openly available on GitHub ([github.com/Liam-](https://github.com/Liam-M-Rooney/IRM_lenses)  
546 [M-Rooney/IRM\\_lenses](https://github.com/Liam-M-Rooney/IRM_lenses)). All data underpinning this publication are openly available from the University of  
547 Strathclyde KnowledgeBase at [doi.org/10.15129/bc849aae-198c-45f3-9e7f-2613532726af](https://doi.org/10.15129/bc849aae-198c-45f3-9e7f-2613532726af).

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### 550 **Authors' Contributions**

551 LMR: conceptualisation, data curation, analysis, investigation, methodology, supervision, validation,  
552 visualisation, manuscript preparation and review.

553 JC: conceptualisation, analysis, investigation, methodology, validation, visualisation, manuscript preparation  
554 and review.

555 BW: analysis, investigation, methodology, validation, software, visualisation, manuscript review.

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### 563 **Competing Interests**

564 We declare no competing interests.